

# Higgs production at NLO in SHERPA

Marek Schönherr

Universität Zürich

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Universität  
Zürich<sup>UZH</sup>



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FONDO NAZIONALE SVIZZERO  
SWISS NATIONAL SCIENCE FOUNDATION

# The SHERPA event generator framework

- Two multi-purpose Matrix Element (ME) generators

**AMEGIC++** JHEP02(2002)044, EPJC53(2008)501

**COMIX** JHEP12(2008)039, PRL109(2012)042001

- A Parton Shower (PS) generator

**CSSHOWER++** JHEP03(2008)038

- A multiple interaction simulation  
à la Pythia **AMISIC++** hep-ph/0601012

- A cluster fragmentation module

**AHADIC++** EPJC36(2004)381

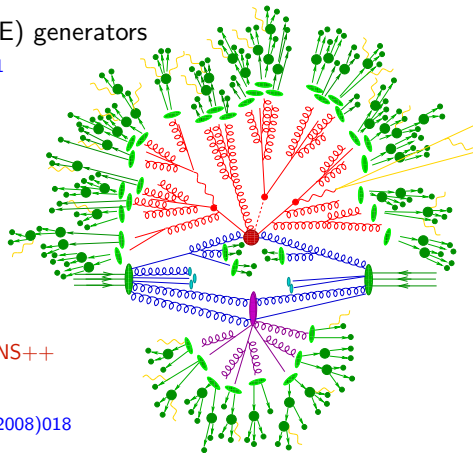
- A hadron and  $\tau$  decay package **HADRONS++**

- A higher order QED generator using  
YFS-resummation **PHOTONS++** JHEP12(2008)018

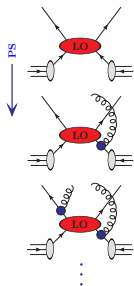
- A minimum bias simulation **SHRiMPS** to appear

**Sherpa's traditional strength is the perturbative part of the event**

MEPs (CKKW), S-Mc@NLO, MENLOPs, MEPS@NLO



# MEPS

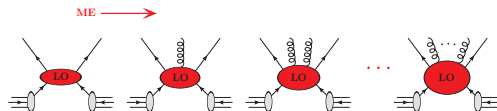


## Parton showers

resummation of (soft-)collinear limit  
 $\rightarrow$  intrajet evolution

- matrix elements (ME) and parton showers (PS) are approximations in different regions of phase space
- MEPS combines multiple LOPS – keeping either accuracy
- NLOPS elevate LOPS to NLO accuracy
- MENLOPS supplements core NLOPS with higher multiplicities LOPS

# MEPs



## Matrix elements

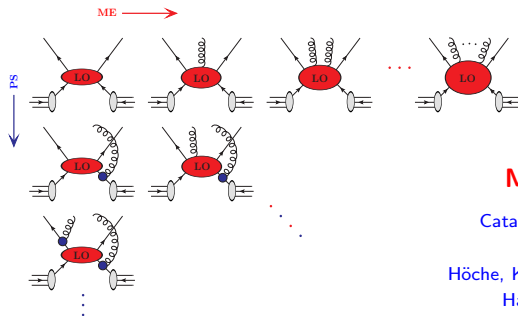
fixed-order in  $\alpha_s$

$\rightarrow$  hard wide-angle emissions

$\rightarrow$  interference terms

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# MEPS



## MEPS (CKKW, MLM)

Catani, Krauss, Kuhn, Webber JHEP11(2001)063

Lönnblad JHEP05(2002)046

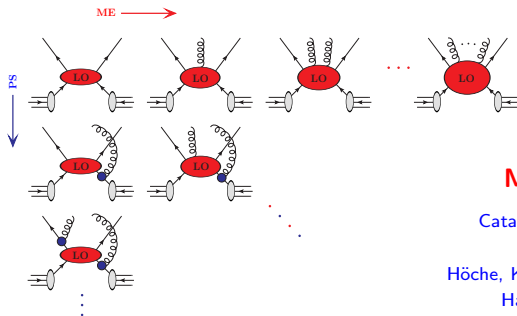
Höche, Krauss, Schumann, Siegert JHEP05(2009)053

Hamilton, Richardson, Tully JHEP11(2009)038

...

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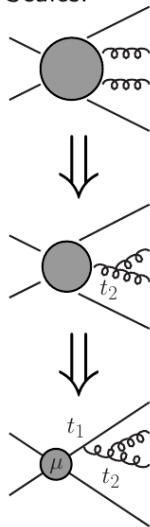
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• NLOs elevate LOPs to NLO accuracy

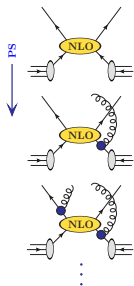
• MENLOs supplements core NLOs with higher multiplicities LOPs

$$\alpha_s^{k+n}(\mu_R) = \alpha_s^k(\mu_{\text{core}}) \alpha_s(t_1) \cdots \alpha_s(t_n)$$

Scales:



# MEPS@NLO



## NLOPS (MC@NLO, POWHEG)

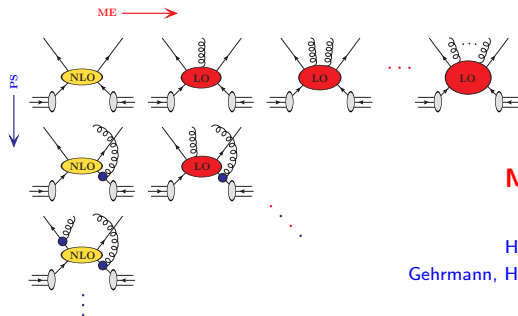
Frixione, Webber JHEP06(2002)029

Nason JHEP11(2004)040, Frixione et.al. JHEP11(2007)070

Höche, Krauss, MS, Siebert JHEP09(2012)049

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-

# MEPS@NLO



## MENLOPS

Hamilton, Nason JHEP06(2010)039

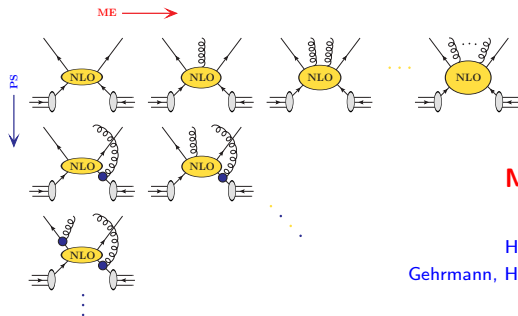
Höche, Krauss, MS, Siebert JHEP08(2011)123

Gehrmann, Höche, Krauss, MS, Siebert JHEP01(2013)144

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## MEPS@NLO

Lavesson, Lönnblad JHEP12(2008)070

Höche, Krauss, MS, Siebert JHEP04(2013)027

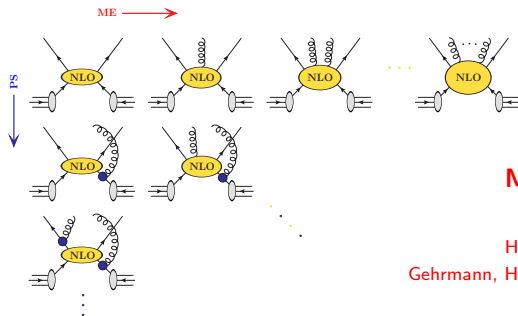
Gehrmann, Höche, Krauss, MS, Siebert JHEP01(2013)144

Lönnblad, Prestel JHEP03(2013)166

Plätzer JHEP08(2013)114

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- MEPS@NLO combines multiple NLOPS – keeping either accuracy

# MEPS@NLO



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Höche, Krauss, MS, Siebert JHEP04(2013)027

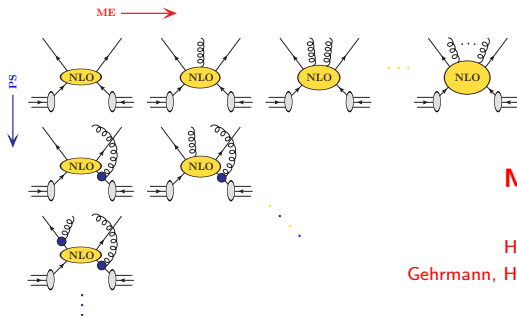
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Plätzer JHEP08(2013)114

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# MEPS@NLO



## MEPS@NLO

Lavesson, Lönnblad

Höche, Krauss, MS, Siebert

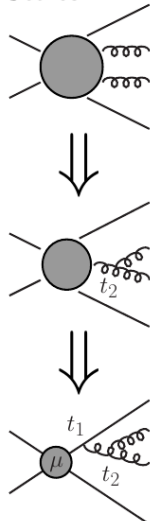
Gehrmann, Höche, Krauss, MS, Siebert

Lönnblad, Prestel

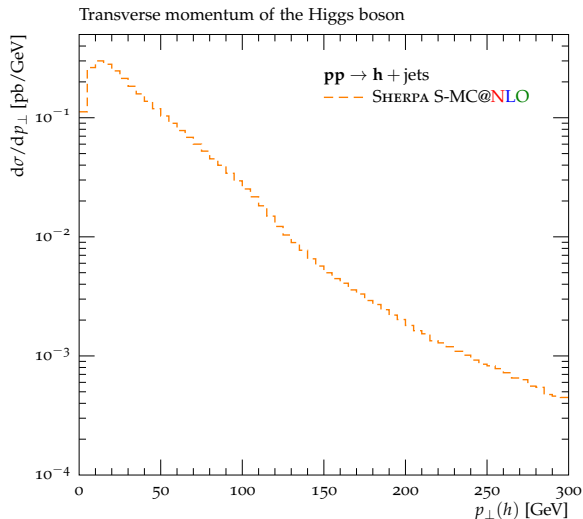
Plätzer

- matrix elements (ME) and parton showers (PS) are approx in different regions of phase space
- MEPS combines multiple LOPS – keeping either accuracy
- NLOPS elevate LOPS to NLO accuracy
- MENLOPS supplements core NLOPS with  $\alpha_s^{k+n}(\mu_R) = \alpha_s^k(\mu_{\text{core}}) \alpha_s(t_1) \cdots \alpha_s(t_n)$
- **MEPS@NLO combines multiple NLOPS**

Scales:

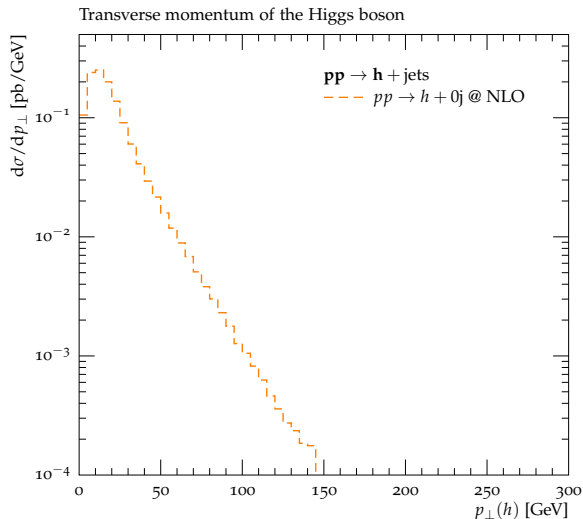


## MEPS@NLO



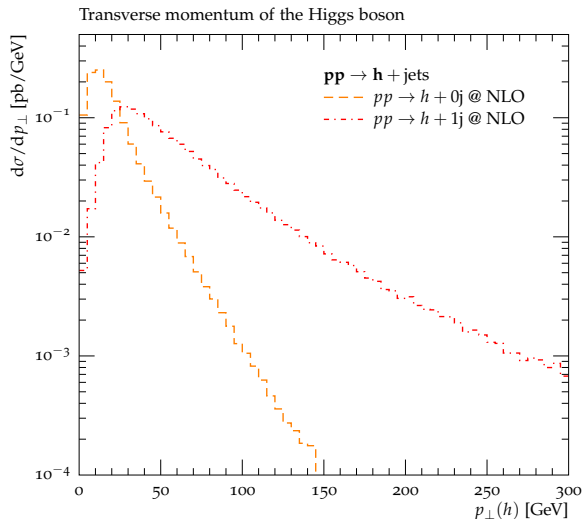
- first emission by NLOPS, restrict to  $Q_{n+1} < Q_{\text{cut}}$
- NLOPS  $pp \rightarrow h + \text{jet}$  for  $Q_{n+1} > Q_{\text{cut}}$
- restrict emission off  $pp \rightarrow h + \text{jet}$  to  $Q_{n+2} < Q_{\text{cut}}$
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- iterate
- sum all contributions
- eg.  $p_{\perp}(h) > 200$  GeV has contributions fr. multiple topologies

## MEPS@NLO



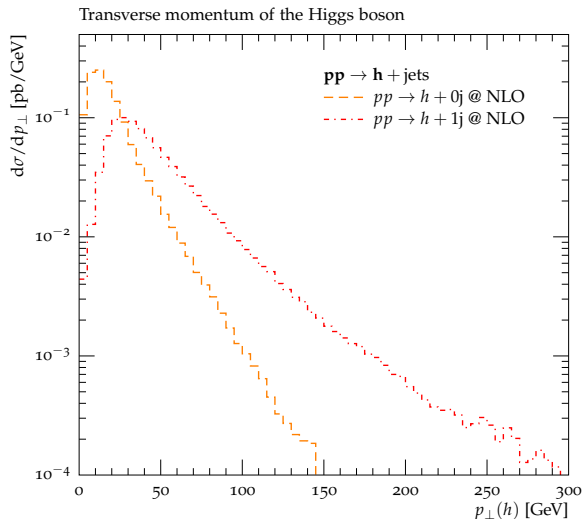
- first emission by NLOs, restrict to  $Q_{n+1} < Q_{\text{cut}}$
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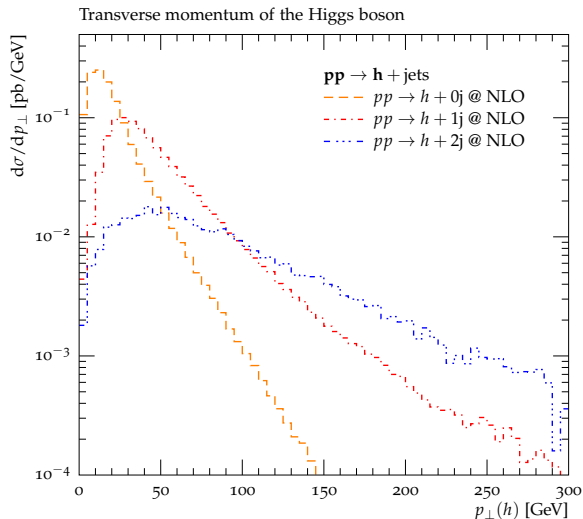
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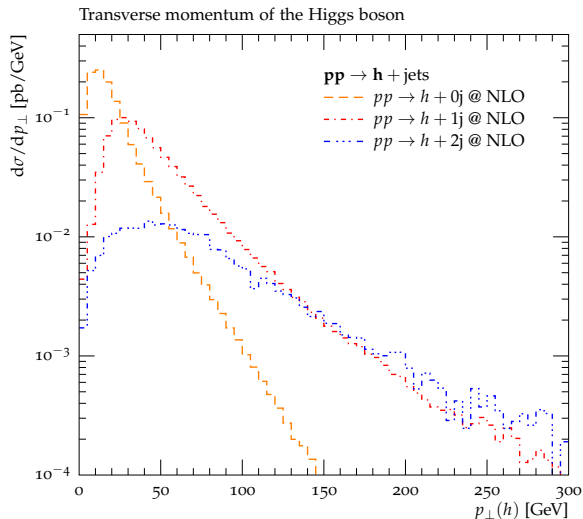
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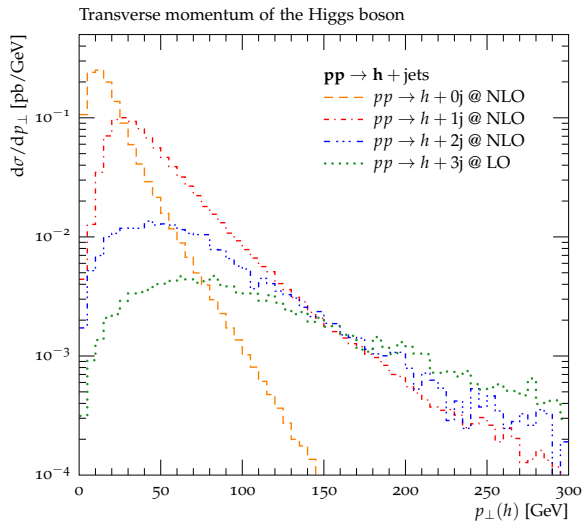


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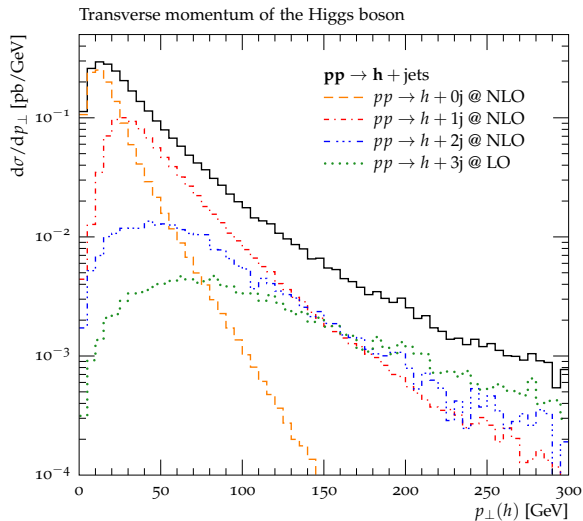
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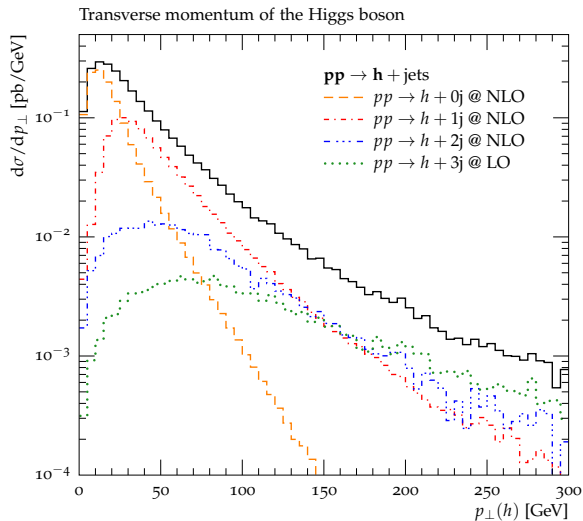
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# Parameter / Scale choices – $\mu_{R/F}$ , $\mu_Q$

$$\alpha_s^{n+k}(\mu_R^2) = \alpha_s^k(\mu_{\text{core}}^2) \alpha_s(t_1) \cdots \alpha_s(t_n) \quad \mu_{F,a/b}^2 = t_{\text{ext},a/b} \quad \mu_Q^2 = \mu_{\text{core}}^2$$

## Free choices

①  $\mu_{\text{core}}$  – scale of core process identified through clustering with inverse parton shower

②  $\mu_{R/F}$  beyond 1-loop running

- calculate with chosen  $\mu_{R/F}$

- include renormalisation and factorisation terms to shift the 1-loop running to above

$$B_n \frac{\alpha_s(\mu_R)}{\pi} \beta_0 \left( \log \frac{\mu_R}{\mu_{\text{CKKW}}} \right)^{2+n}$$

and

$$B_n \frac{\alpha_s}{2\pi} \log \frac{\mu_F}{t_{\text{ext}}} \sum_{c=q,g} \int_{x_a}^1 \frac{dz}{z} P_{ac}(z) f_c(x_a/z, \mu_F^2)$$

→ same as in UNLOPS

Lönnblad, Prestel JHEP03(2013)166, Plätzer JHEP08(2013)114

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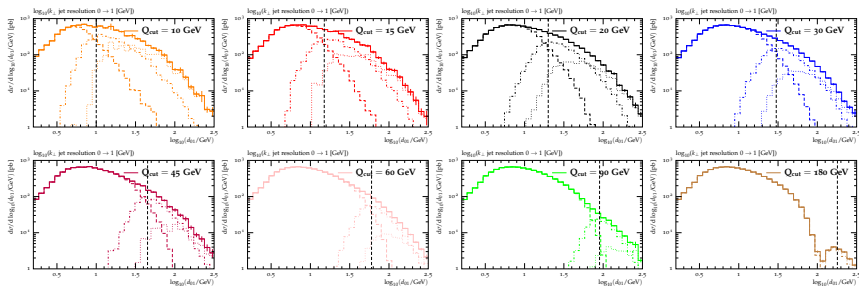
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[Lönnblad, Prestel JHEP03\(2013\)166](#), [Plätzer JHEP08\(2013\)114](#)

# Parameter / Scale choices – $Q_{\text{cut}}$

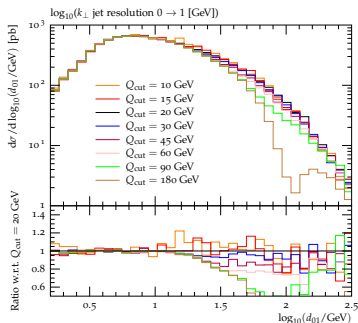
Merging cut  $Q_{\text{cut}}$  dependence ( $pp \rightarrow Z + \text{jets}$  MEPS, up to 2 in ME):



- parton shower is trusted to correctly describe emissions  $\lesssim Q_{\text{cut}}$
- changes the region where higher accuracy is used for calculation  
→ part of the uncertainty is due to degraded accuracy for large  $Q_{\text{cut}}$
- all samples are identical for  $Q < Q_{\text{cut}}^{\text{smallest}}$  and  $Q > Q_{\text{cut}}^{\text{largest}}$  by construction
- for  $Q \geq 45$  GeV shower approximation breaks down (earlier in other obs.)
- $Q_{\text{cut}}$  dependence usually small

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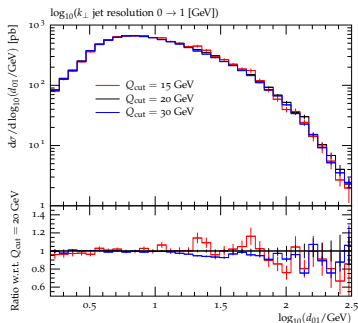


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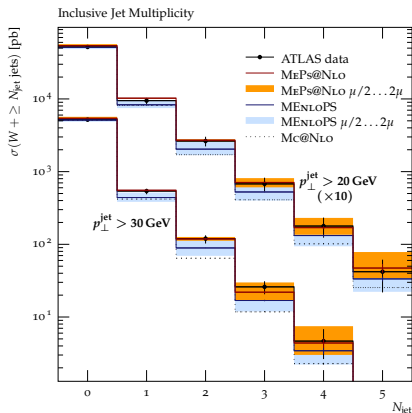
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# Recent results

## Multijet merging at NLO accuracy (MEPS@NLO)

- $pp \rightarrow W + \text{jets}$  – SHERPA+BLACKHAT    Höche, Krauss, MS, Siegert JHEP04(2013)027
- $e^+e^- \rightarrow \text{jets}$  – SHERPA+BLACKHAT  
Gehrmann, Höche, Krauss, MS, Siegert JHEP01(2013)144
- $pp \rightarrow h + \text{jets}$  – SHERPA+GOSAM/MCFM  
Höche, Krauss, MS, Siegert, contribution to YR3 arXiv:1307.1347  
Höche, Krauss, MS Phys.Rev.D90(2014)014012  
MS, Zapp, contribution to LH'13 arXiv:1405.1067
- $p\bar{p} \rightarrow t\bar{t} + \text{jets}$  – SHERPA+GOSAM/OPENLOOPS  
Höche, Huang, Luisoni, MS, Winter Phys.Rev.D88(2013)014040  
Höche, Krauss, Maierhöfer, Pozzorini, MS, Siegert arXiv:1402.6293
- $pp \rightarrow 4\ell + \text{jets}$  – SHERPA+OPENLOOPS  
Cascioli, Höche, Krauss, Maierhöfer, Pozzorini, Siegert JHEP01(2014)046
- $pp \rightarrow VH + \text{jets}, pp \rightarrow VV + \text{jets}, pp \rightarrow VVV + \text{jets}$   
– SHERPA+OPENLOOPS  
Höche, Krauss, Pozzorini, MS, Thompson, Zapp Phys.Rev.D89(2014)093015

# Results – $pp \rightarrow W + \text{jets}$

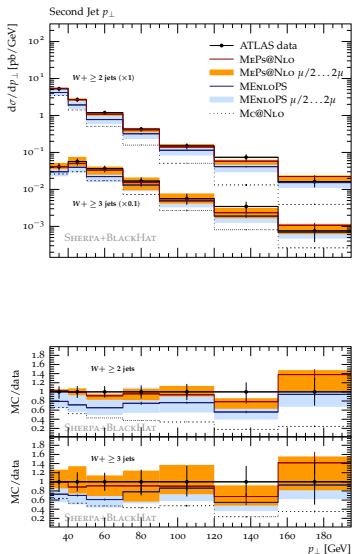
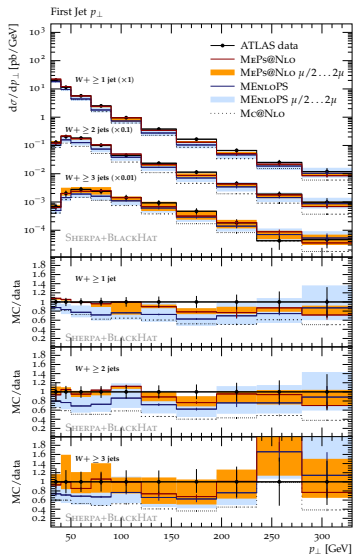


$pp \rightarrow W + \text{jets}$  (0,1,2 @ NLO; 3,4 @ LO)

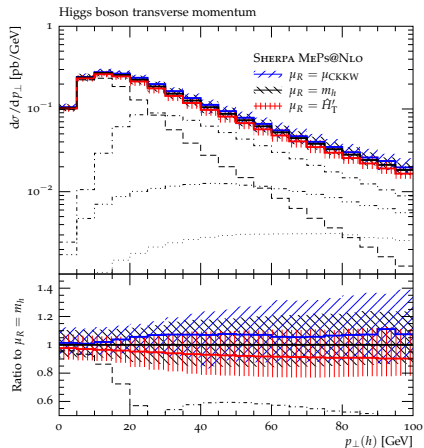
- $\mu_{R/F} \in [\frac{1}{2}, 2] \mu_{\text{def}}$   
scale uncertainty much reduced
- NLO dependence  
for  $pp \rightarrow W + 0,1,2$  jets  
LO dependence  
for  $pp \rightarrow W + 3,4$  jets
- virtual MEs from BLACKHAT
- $Q_{\text{cut}} = 30 \text{ GeV}$
- good data description

ATLAS data Phys.Rev.D85(2012)092002

# Results – $pp \rightarrow W + \text{jets}$



# Results – pp $\rightarrow$ h+jets



Höche, Krauss, MS, Phys.Rev.D90(2014)014012

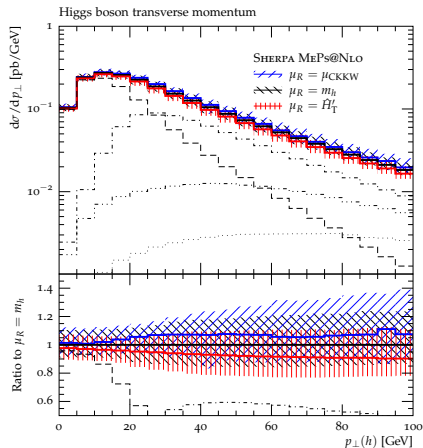
pp  $\rightarrow$  h+jets (0,1,2 @ NLO; 3 @ LO)

- $\mu_{R/F} \in [\frac{1}{2}, 2] \mu_{R/F}^{\text{def}}$
- $\mu_Q \in [\frac{1}{\sqrt{2}}, \sqrt{2}] \mu_Q^{\text{def}}$
- $Q_{\text{cut}} \in \{15, 20, 30\}$  GeV
- virtual MEs from MCFM ( $hjj$ )

use  $m_t \rightarrow \infty$  limit (EFT)

$\rightarrow$  finite  $m_t$  effects see Silvan's talk

# Results – $pp \rightarrow h + \text{jets}$



⇒ difference beyond accuracy

scale choices:  $\mu_F = \mu_Q = m_h$

①  $\mu_R = \mu_{\text{CKKW}}$

$$\alpha_s^{2+n}(\mu_{\text{CKKW}}) = \alpha_s^2(m_h) \alpha_s(t_1) \cdots \alpha_s(t_n)$$

②  $\mu_R = m_h$

③  $\mu_R = \hat{H}'_T$

need to include ren. term

$$B_n \frac{\alpha_s(\mu_R)}{\pi} \beta_0 \left( \log \frac{\mu_R}{\mu_{\text{CKKW}}} \right)^{2+n}$$

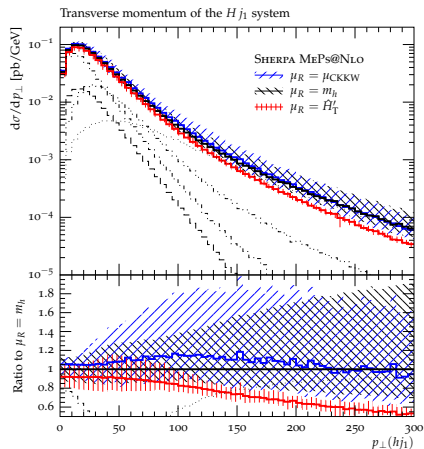
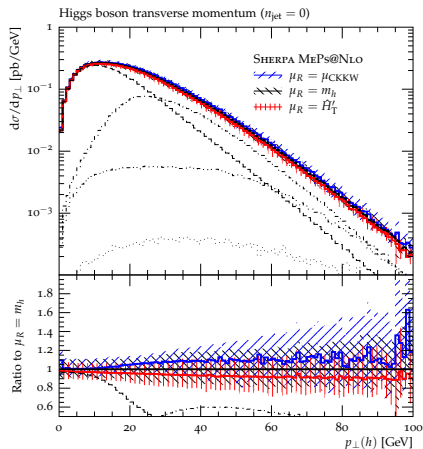
to restore 1-loop running to  $\mu_{\text{CKKW}}$   
 → otherwise PS-accuracy violated

→ same as in UNLOPS approach

Lönblad, Prestel JHEP03(2013)166

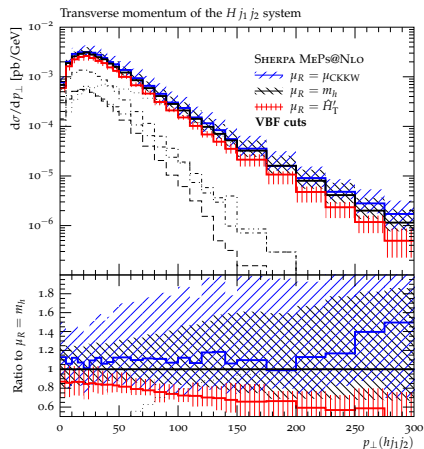
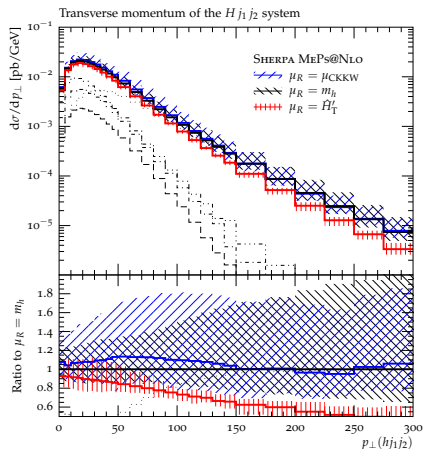
Plätzer JHEP08(2013)114

# Results – $pp \rightarrow h + \text{jets}$



- all predictions identical to MEPS@NLO accuracy
- vastly differing size of uncertainties

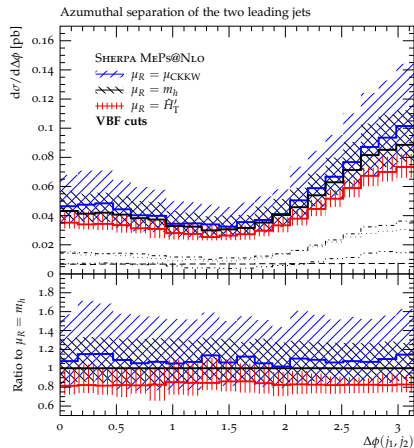
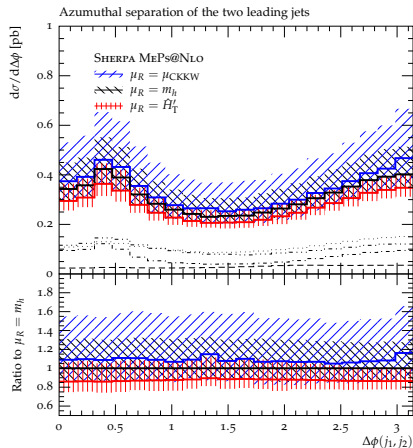
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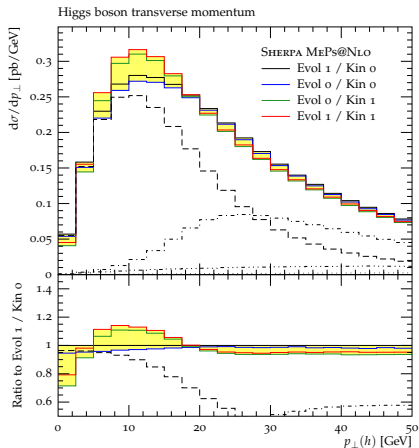


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- vastly differing size of uncertainties

# Results – pp $\rightarrow$ h+jets



## Parton shower uncertainties

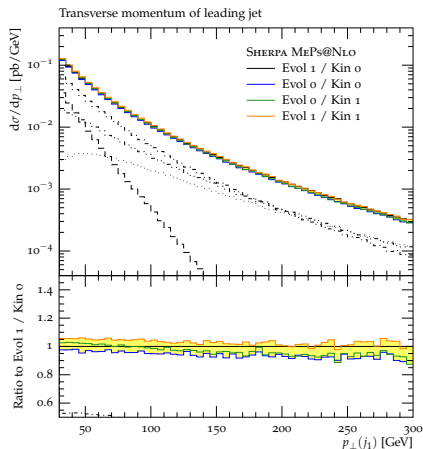
- evolution scale

		Final State
0		$2 p_i p_j \tilde{z}_{i,jk} (1 - \tilde{z}_{i,jk})$
1	$2 p_i p_j$	$\tilde{z}_{i,jk} (1 - \tilde{z}_{i,jk})$ if $i, j = g$
		$1 - \tilde{z}_{i,jk}$ if $j = g$
		$\tilde{z}_{i,jk}$ if $i = g$
		1 else
		Initial State
0		$2 p_a p_j (1 - x_{a,j,k})$
1	$2 p_a p_j$	$1 - x_{a,j,k}$ if $j = g$
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- recoil scheme

0	initial state as if final state + $\perp$ -boost <a href="#">Höhe, Schumann, Siebert Phys.Rev.D81(2010)034026</a>
1	original CS <a href="#">Catani, Seymour Nucl.Phys.B485(1997)291-419</a> <a href="#">Schumann, Krauss JHEP03(2008)038</a>
	→ similar ideas in <a href="#">Gieseke, Plätzer JHEP01(2011)024</a>

# Results – pp $\rightarrow$ h+jets



## Parton shower uncertainties

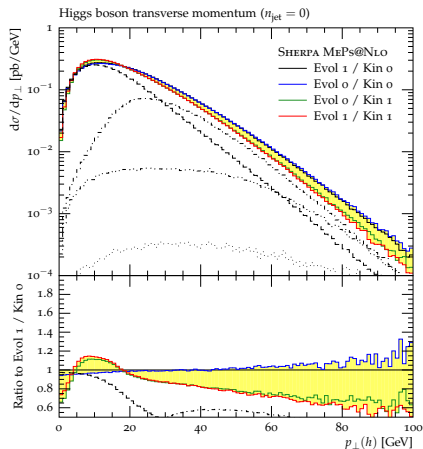
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# Results – pp $\rightarrow$ h+jets



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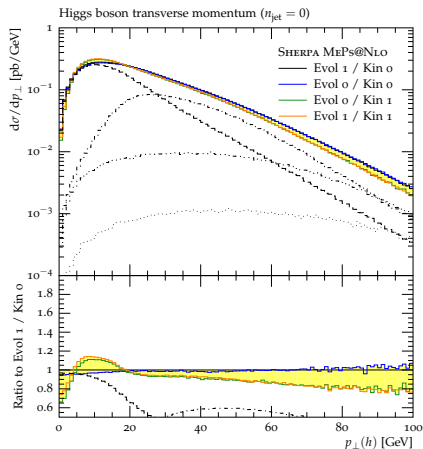
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# Results – pp $\rightarrow$ h+jets



## Parton shower uncertainties

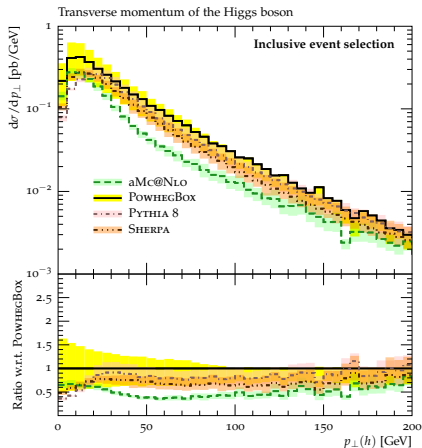
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# Les Houches comparative study – $pp \rightarrow h + \text{jets}$



LH'13 ( $h + \text{dijets}$  study) arXiv:1405.1067

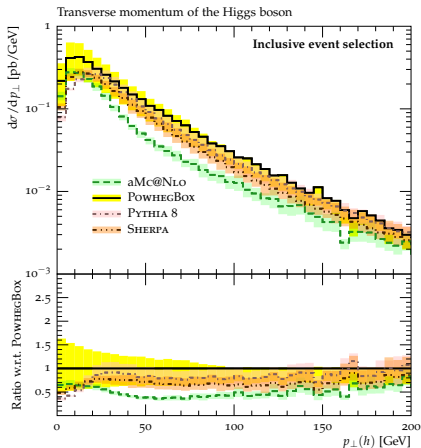
$pp \rightarrow h + \text{dijets}$  study

- HEJ (BFKL w/ ME-corr.)
- aMC@NLO (FxFx combination)
- POWHEG-BOX (HJJ-MiNLO)
- PYTHIA8 (UNLOPS merging)
- SHERPA (MEPS@NLO merging)

⇒ focus on ggF background to VBF

- two dijet-event selections  
Leading jet / Forward-backward
- two levels of cuts  
Dijet cuts / VBF cuts

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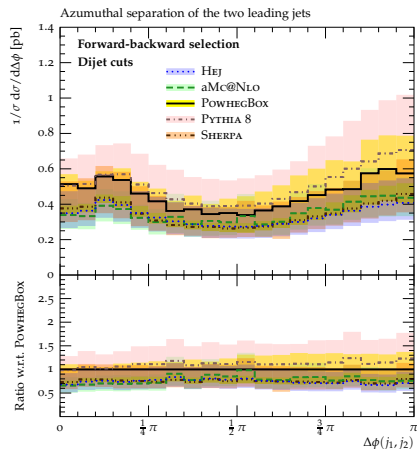
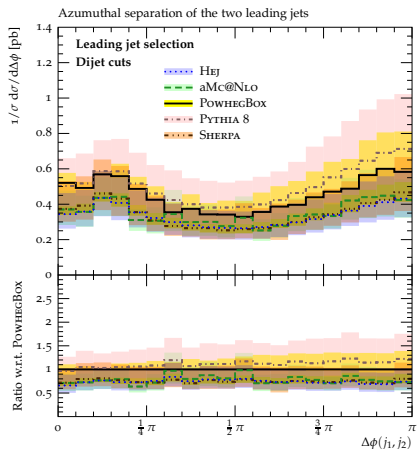
- HEJ (BFKL w/ ME-corr.) **LO**
- aMC@NLO (FxFx combination) **NLO**
- POWHEG-BOX (HJJ-MiNLO) **NLO**
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- SHERPA (MEPS@NLO merging) **NLO**

⇒ focus on ggF background to VBF

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Leading jet / Forward-backward
- two levels of cuts  
Dijet cuts / VBF cuts

- HJJ-MiNLO has no formal accuracy for inclusive observables

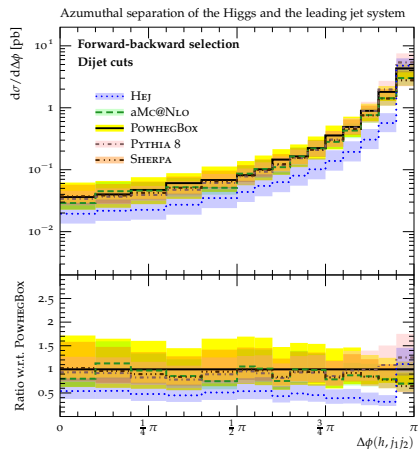
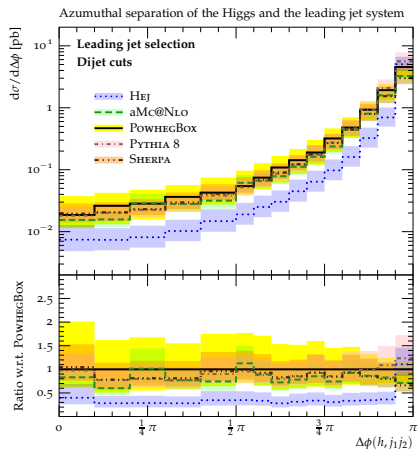
# Results – $pp \rightarrow h + \text{jets}$



- good agreement in shape between generators, different normalisations
- similar uncertainties

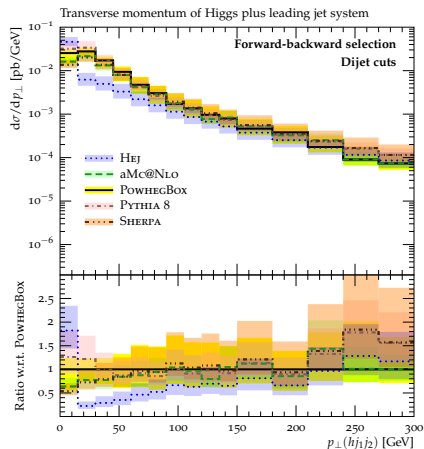
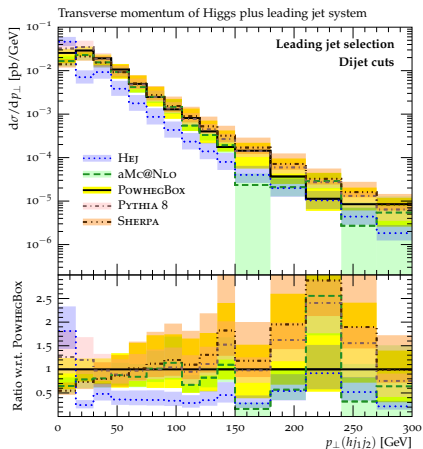


# Results – $pp \rightarrow h + \text{jets}$



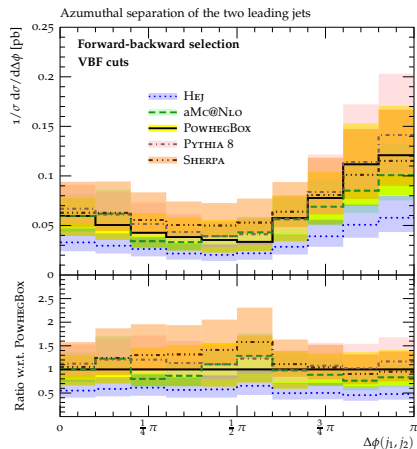
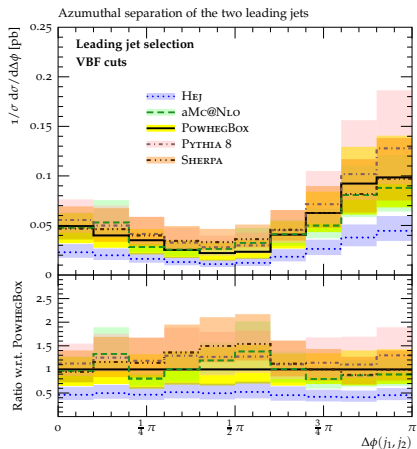
- good agreement between generators, slightly different shapes
- HEJ has less additional jet activity

# Results – $pp \rightarrow h + \text{jets}$



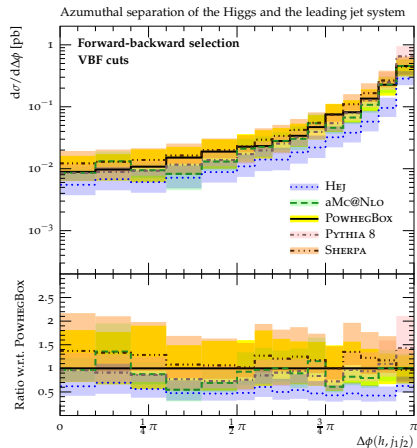
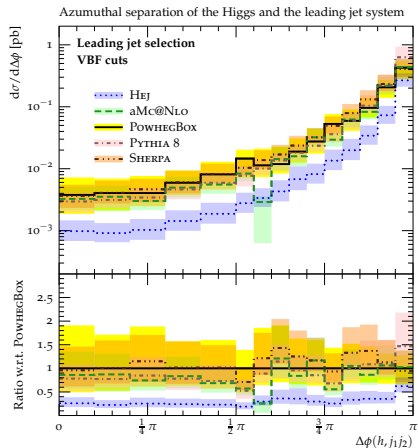
- PYTHIA8 and SHERPA have more high- $p_{\perp}$  activity
- HEJ has less additional jet activity

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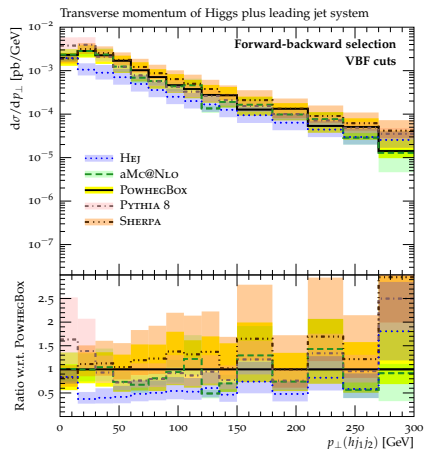
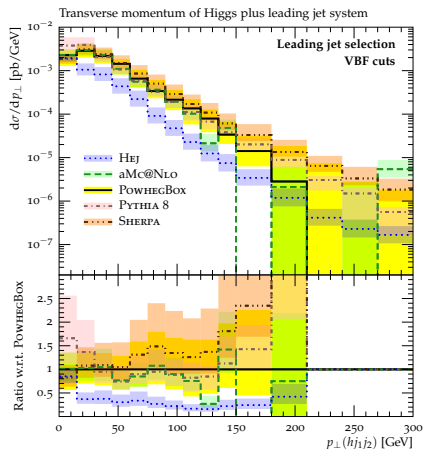
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# Results – $pp \rightarrow h + \text{jets}$



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# Results – Trilepton production

Höche, Krauss, Pozzorini, MS, Thompson, Zapp Phys.Rev.D89(2014)093015

- trilepton ( $e, \mu$ ) production analysis in  $VH$  search regions  
→ focus on theoretical uncertainties
- model all signal and background processes with consistent setup at largest available accuracy at particle level  
→ need to describe lepton isolation and jet veto efficiency simultaneously
- produce bosons on shell, model off-shell effects through Breit-Wigner smearing  
→ QCD/QED corrections to intermediate states and decay products
- most important event selection criteria

	CMS-inspired analysis	ATLAS-inspired analysis
$Z$ veto	$ m_Z - m_{SFOS}  > 25 \text{ GeV}$	no SFOS
jet veto	$p_{\perp}^{\text{jet}} < 40 \text{ GeV}$	$p_{\perp}^{\text{jet}} < 20 \text{ GeV}$

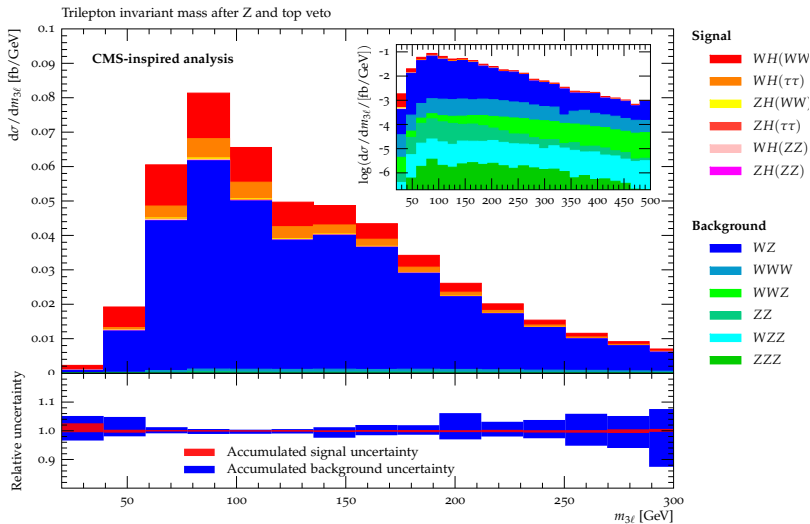
- include  $V \rightarrow \tau \rightarrow e, \mu$  decay chains and possibilities to “loose” leptons
- separate  $VVVj(j)$  from  $tVV$  and  $t\bar{t}W$  by disallowing final state  $b$ -quarks

# Results – Trilepton production

Höche, Krauss, Pozzorini, MS, Thompson, Zapp Phys.Rev.D89(2014)093015

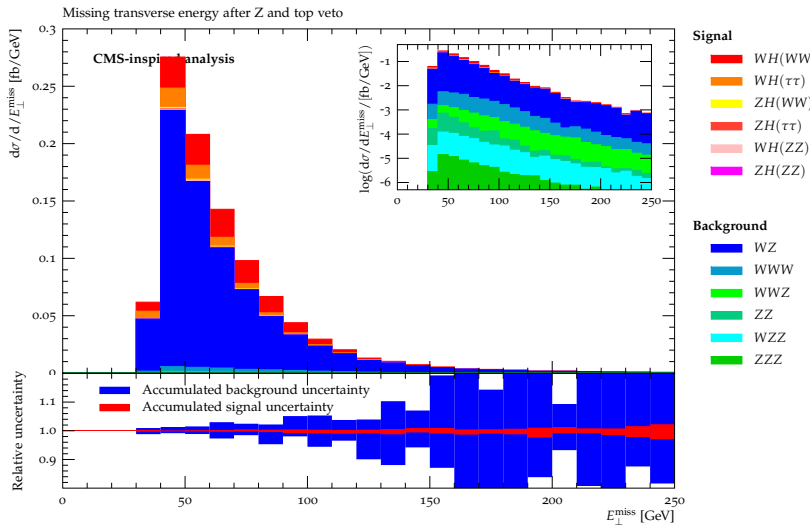
Process	Accuracy	Decays ( $\ell = e, \mu, \tau$ )
$WH + \text{jets}$	0,1j@NLO, 2j@LO	$H \rightarrow WW, W \rightarrow \ell\nu, Z \rightarrow \ell\ell, \tau \rightarrow \ell\nu\nu$ $H \rightarrow \tau\tau, W \rightarrow \ell\nu, Z \rightarrow \ell\ell, \tau \rightarrow \ell\nu\nu$
$ZH + \text{jets}$	0,1j@NLO, 2j@LO	$H \rightarrow ZZ, W \rightarrow \text{all}, Z \rightarrow \text{all}, \tau \rightarrow \text{all}$ $H \rightarrow WW, W \rightarrow \text{all}, Z \rightarrow \ell\ell, \tau \rightarrow \text{all}$ $H \rightarrow \tau\tau, Z \rightarrow \ell\ell, \tau \rightarrow \text{all}$ $H \rightarrow ZZ, Z \rightarrow \text{all}, \tau \rightarrow \text{all}$
$WZ + \text{jets}$	0,1j@NLO, 2j@LO	$W \rightarrow \ell\nu, Z \rightarrow \ell\ell, \tau \rightarrow \ell\nu\nu$
$WW + \text{jets}$	0,1j@NLO, 2j@LO	$W \rightarrow \ell\nu, \tau \rightarrow \ell\nu\nu$
$WWZ + \text{jets}$	0j@NLO, 1,2j@LO	$W \rightarrow \text{all}, Z \rightarrow \ell\ell, \tau \rightarrow \text{all}$
$ZZ + \text{jets}$	0j@NLO, 1,2j@LO	$Z \rightarrow \ell\ell, \tau \rightarrow \text{all}$
$WZZ + \text{jets}$	0j@NLO, 1,2j@LO	$W \rightarrow \text{all}, Z \rightarrow \text{all}, \tau \rightarrow \text{all}$
$ZZZ + \text{jets}$	0j@NLO, 1,2j@LO	$W \rightarrow \text{all}, Z \rightarrow \text{all}, \tau \rightarrow \text{all}$

# Results – Trilepton production

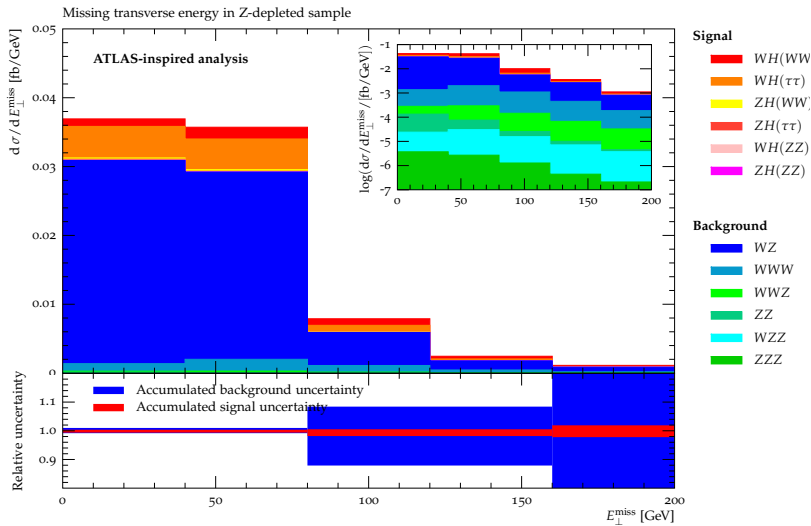




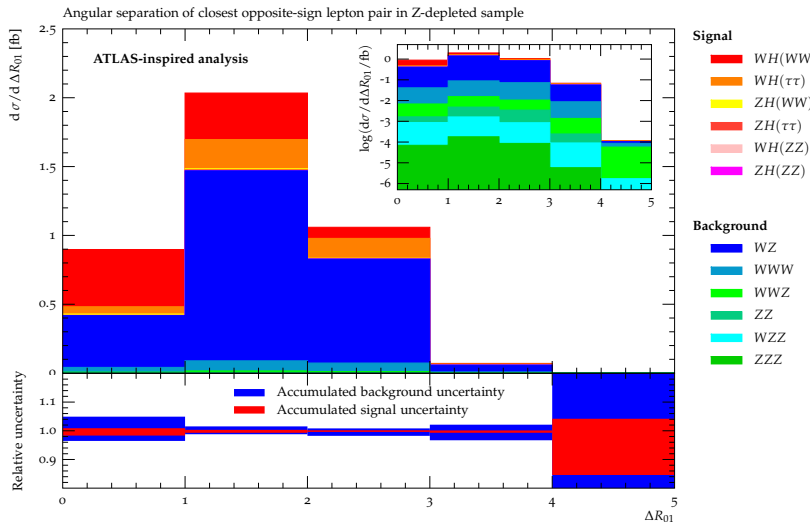
# Results – Trilepton production



# Results – Trilepton production



# Results – Trilepton production



# Conclusions

- multijet merging at NLO proceeds schematically as at LO  
→ introduce MC-counterterm to retain NLO accuracy
- preserves NLO accuracy of the ME and accuracy of the PS in resumming hierarchies of emission scales  
→ scale setting essential for recovering PS resummation  
→ core scale can be chosen freely  
→ beyond 1-loop running the scales can of course be freely chosen
- perturbative uncertainties due to  $\mu_{R/F}$ ,  $\mu_Q$  and  $Q_{\text{cut}}$  can be assessed in the fixed-order part
- intrinsic parton shower uncertainties can be partially assessed

current release SHERPA-2.1.1

<http://sherpa.hepforge.org>

Thank you for your attention!

# MEPS

Parton showers (operate in  $N_c \rightarrow \infty$  limit):

$$PS_n(t_c, t_{\max}) = \Delta_n(t_c, t_{\max}) + \int_{t_c}^{t_{\max}} dt' \mathcal{K}_n(t') \Delta_n(t', t_{\max})$$

Multijet merging at leading order:

$$d\sigma^{\text{MEPS}} = d\sigma_n^{\text{LO}} \otimes PS_n$$

- restrict the parton shower on  $2 \rightarrow n$  to emit only below  $Q_{\text{cut}}$
- arbitrary jet measure  $Q_n = Q_n(\Phi_n)$
- add the  $n + 1$  ME and its parton shower
- multiply by Sudakov wrt.  $2 \rightarrow n$  process to restore resummation
- iterate
- if  $t_n(\Phi_n) \neq Q_n(\Phi_n)$  truncated shower needed to fill gaps

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$$\begin{aligned} d\sigma^{\text{MEPs}} = & d\sigma_n^{\text{LO}} \otimes \text{PS}_n \Theta(Q_{\text{cut}} - Q_{n+1}) \\ & + d\sigma_{n+1}^{\text{LO}} \Theta(Q_{n+1} - Q_{\text{cut}}) \Delta_n(t_{n+1}, t_n) \otimes \text{PS}_{n+1} \Theta(Q_{\text{cut}} - Q_{n+2}) \\ & + d\sigma_{n+2}^{\text{LO}} \Theta(Q_{n+2} - Q_{\text{cut}}) \Delta_n(t_{n+1}, t_n) \Delta_{n+1}(t_{n+2}, t_{n+1}) \otimes \text{PS}_{n+2} \end{aligned}$$

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$$\text{PS}_n(t_c, t_{\max}) = \Delta_n(t_c, t_{\max}) + \int_{t_c}^{t_{\max}} dt' \mathcal{K}_n(t') \Delta_n(t')$$

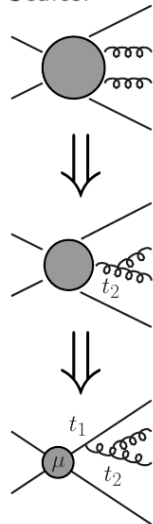
Multijet merging at leading order:

$$\begin{aligned} d\sigma^{\text{MEPS}} = & d\sigma_n^{\text{LO}} \otimes \text{PS}_n \Theta(Q_{\text{cut}} - Q_{n+1}) \\ & + d\sigma_{n+1}^{\text{LO}} \Theta(Q_{n+1} - Q_{\text{cut}}) \Delta_n(t_{n+1}, t_n) \otimes \text{PS}_{n+1} \Theta(Q_{\text{cut}} - Q_{n+2}) \\ & + d\sigma_{n+2}^{\text{LO}} \Theta(Q_{n+2} - Q_{\text{cut}}) \Delta_n(t_{n+1}, t_n) \Delta_{n+1}(t_{n+2}, t_{n+1}) \otimes \text{PS}_{n+2} \Theta(Q_{\text{cut}} - Q_{n+3}) \end{aligned}$$

- restrict the parton shower on  $2 \rightarrow n$  to emit only below  $Q_{\text{cut}}$
- arbitrary jet measure  $Q_n = Q_n(\Phi_n)$
- add the  $n + 1$  ME and its parton shower
- multiply by Sudakov wrt.  $2 \rightarrow n$  process to restore resummation
- iterate

• if  $t_n(\Phi_n) \neq Q_n(\Phi_n)$  truncated shower needed to fill gaps

Scales:



$$\alpha_s^{k+n}(\mu_R) = \alpha_s^k(\mu_{\text{core}}) \alpha_s(t_1) \cdots \alpha_s(t_n)$$

# MEPs

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Multijet merging at leading order:

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- if  $t_n(\Phi_n) \neq Q_n(\Phi_n)$  truncated shower needed to fill gaps [Nason JHEP11\(2004\)040](#)

# MEPS@NLO

Parton showers for NLOPS (need to reproduce  $N_c = 3$  singular limits for 1st em.):

$$\widetilde{\text{PS}}_n(t_c, t_{\max}) = \widetilde{\Delta}_n(t_c, t_{\max}) + \int_{t_c}^{t_{\max}} dt' \widetilde{\mathcal{K}}_n(t') \widetilde{\Delta}_n(t', t_{\max})$$

Multijet merging at next-to-leading order:

$$d\sigma^{\text{MEPS@NLO}} = d\sigma_n^{\text{NLO}} \otimes \widetilde{\text{PS}}_n$$

- NLOPS for  $2 \rightarrow n$
- add the NLOPS for  $2 \rightarrow n + 1$
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## MEPS@NLO

Parton showers for NLOPS (need to reproduce  $N_c = 3$  singular li

$$\widetilde{\text{PS}}_n(t_c, t_{\text{max}}) = \widetilde{\Delta}_n(t_c, t_{\text{max}}) + \int_{t_c}^{t_{\text{max}}} dt' \widetilde{\mathcal{K}}_n(t') \widetilde{\Delta}_n(t',$$

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$$d\sigma^{\text{MEPS@NLO}} = d\sigma_n^{\text{NLO}} \otimes \widetilde{\text{PS}}_n \Theta(Q_{\text{cut}} - Q_{n+1})$$

$$+ d\sigma_{n+1}^{\text{NLO}} \Theta(Q_{n+1} - Q_{\text{cut}}) \left( \Delta_n(t_{n+1}, t_n) - \Delta_n^{(1)} \right)$$

$$\otimes \widetilde{\text{PS}}_{n+1} \Theta(Q_{\text{cut}} - Q_{n+2})$$

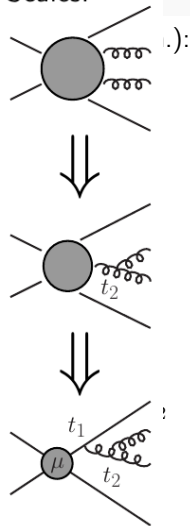
$$+ d\sigma_{n+2}^{\text{NLO}} \Theta(Q_{n+2} - Q_{\text{cut}}) \left( \Delta_n(t_{n+1}, t_n) - \Delta_n^{(1)} \right)$$

$$\times \left( \Delta_{n+1}(t_{n+2}, t_{n+1}) - \Delta_{n+1}^{(1)}(t_{n+2}, t_n) \right)$$

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- multiply by Sudakov wrt.  $2 \rightarrow n$  process to restore resummation

- remove overlap of  $\Delta_n$  and  $d\sigma_{n+1}^{\text{NLO}}$ , iteratively  $\alpha_s^{k+n}(\mu_R) = \alpha_s^k(\mu_{\text{core}}) \alpha_s(t_1) \cdots \alpha_s(t_n)$

Scales:



## MEPS@NLO

Parton showers for NLOPS (need to reproduce  $N_c = 3$  singular li

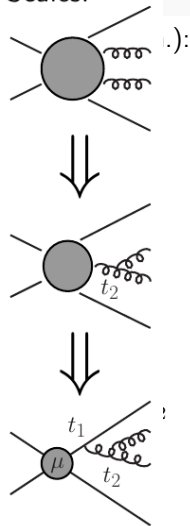
$$\widetilde{\text{PS}}_n(t_c, t_{\text{max}}) = \widetilde{\Delta}_n(t_c, t_{\text{max}}) + \int_{t_c}^{t_{\text{max}}} dt' \widetilde{\mathcal{K}}_n(t') \widetilde{\Delta}_n(t',$$

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Scales:



# MENLOPS

$$\begin{aligned}
 d\sigma^{\text{MENLOPS}} = & d\sigma_n^{\text{NLO}} \otimes \widetilde{\text{PS}}_n \Theta(Q_{\text{cut}} - Q_{n+1}) \\
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 & \quad \times \Delta_n(t_{n+1}, t_n) \Delta_{n+1}(t_{n+2}, t_{n+1}) \otimes \text{PS}_{n+2}
 \end{aligned}$$

- restrict MC@NLO expression to region  $Q < Q_{\text{cut}}$
- add in real radiation explicitly, as in MEPS
- restore logarithmic behaviour by explicit Sudakov
- local K-factor for continuity at  $Q_{\text{cut}}$

$$k_n(\Phi_{n+1}) = \frac{\bar{B}_n(\Phi_n)}{B_n(\Phi_n)} \left( 1 - \frac{H_n(\Phi_{n+1})}{R_n(\Phi_{n+1})} \right) + \frac{H_n(\Phi_{n+1})}{R_n(\Phi_{n+1})}$$

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- restore logarithmic behaviour by explicit Sudakov
- local K-factor for continuity at  $Q_{\text{cut}}$

$$k_n(\Phi_{n+1}) = \frac{\bar{B}_n(\Phi_n)}{B_n(\Phi_n)} \left( 1 - \frac{H_n(\Phi_{n+1})}{R_n(\Phi_{n+1})} \right) + \frac{H_n(\Phi_{n+1})}{R_n(\Phi_{n+1})}$$

- iterate

# MENLOPS

$$\begin{aligned}
 d\sigma^{\text{MENLOPS}} = & d\sigma_n^{\text{NLO}} \otimes \widetilde{\text{PS}}_n \Theta(Q_{\text{cut}} - Q_{n+1}) \\
 & + k_n(\Phi_{n+1}) d\sigma_{n+1}^{\text{LO}} \Theta(Q_{n+1} - Q_{\text{cut}}) \Delta_n(t_{n+1}, t_n) \\
 & \quad \otimes \text{PS}_{n+1} \Theta(Q_{\text{cut}} - Q_{n+2}) \\
 & + k_n(\Phi_{n+1}(\Phi_{n+2})) d\sigma_{n+2}^{\text{LO}} \Theta(Q_{n+2} - Q_{\text{cut}}) \\
 & \quad \times \Delta_n(t_{n+1}, t_n) \Delta_{n+1}(t_{n+2}, t_{n+1}) \otimes \text{PS}_{n+2}
 \end{aligned}$$

- restrict MC@NLO expression to region  $Q < Q_{\text{cut}}$
- add in real radiation explicitly, as in MEPS
- restore logarithmic behaviour by explicit Sudakov
- local K-factor for continuity at  $Q_{\text{cut}}$

$$k_n(\Phi_{n+1}) = \frac{\bar{B}_n(\Phi_n)}{B_n(\Phi_n)} \left( 1 - \frac{H_n(\Phi_{n+1})}{R_n(\Phi_{n+1})} \right) + \frac{H_n(\Phi_{n+1})}{R_n(\Phi_{n+1})}$$

- iterate