Conformal Bootstrap a new way to compute in Quantum Field Theory

Slava Rychkov



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Problem

Find an efficient algorithm to solve strongly coupled QFTs

Lattice field theory?

• Benchmark problem: Light hadron spectrum at % level



$$N_f = 2 + 1$$

Durr et al, Science 322(2008) 1224

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• Cost: \approx 1 supercomputer-year \approx 100,000 single core-years

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aka fixed points of RG flow

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Studying these fixed points is proving more difficult on the lattice than the ordinary QCD

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 "quantum spin liquid" state in condensed matter
 E.g. N_f = 4 for *Herbertsmithite*. ZnCu₃(OH)₆Cl₂





119°×

CFTs in Euclidean d=3

$$\mathcal{L} = (\partial \phi)^2 + m^2 \phi^2 + \lambda \phi^4$$

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describes - Curie point of uniaxial magnets in 3d - liquid-vapor phase transitions

Plan

1. CFT generalities

2. Conformal Bootstrap

CFT kinematics i.e. everything you get for free

• Observables in CFT are correlation functions (no S-matrix)

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Example: each CFT contains a stress tensor operator $T_{\mu\nu}$, conserved and traceless, of dimension Δ =d

Operators transform covariantly under

Conformal transformations $x \to x' = f(x)$

E.g. for scalars:

$$A(x') = \lambda(x)^{-\Delta} A(x)$$

$$\lambda = |\partial x'/\partial x|^{1/d}$$
 local scale factor

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Poincaré + dilatations + special conformal transformations $\delta_a x_\mu = 2(a.x)x^\mu - x^2 a^\mu$

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 $\langle A_1(x_1)A_2(x_2)A_3(x_3)\rangle = f_{123}|x_{12}|^{\Delta_3 - \Delta_1 - \Delta_2}|x_{13}|^{\dots}|x_{23}|^{\dots}$

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CFT data

to find these = to solve the CFT

Example: CFT data from experiment



Lipa et al, Phys.Rev.Lett. 76,944(1996)

 λ -point of He-4 space shuttle measurement

Reached T-T_c~nK
$$\Rightarrow \xi/a\sim 10^6$$



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Common disadvantages:

- conformal symmetry left unused
- 3pt function coefficients f_{ijk} much harder to get than Δ_i (never done in practice)

Conformal Bootstrap

Original idea Ferrara, Gatto, Grillo 1973; Polyakov 1974

Concrete realization in 2d Belavin, Polyakov, Zamolodchikov 1983

First results in 4d (motivated by BSM) Rattazzi, S.R., Tonni, Vichi 2008 (Using 4d conformal blocks by Dolan,Osborn 2000)

First results in 3d (applications to Ising-3) El-Showk, Paulos, Poland, S.R., Simmons-Duffin, Vichi 2012

Put a Lorentz-inv theory on $S^1 \times S^1$









2d CFT: $E_n(L) = \frac{2\pi}{L} (\Delta_n - c/12) \implies \text{constraint on } \Delta$'s Cardy 1986
Consider 4-pt function:



Consider 4-pt function: sphere S A_1 A_1 A_4 A_4 A_4 A_2 A_3













• States:
$$E_n \leftrightarrow \Delta_n$$
 $|n\rangle \leftrightarrow A_n(0)$

n

states on S

• Matrix elements:

 $\langle A_1 A_2 | n \rangle = f_{12n} \times (\text{kinematical factor})$

Get

$$\langle A_1 A_2 A_3 A_4 \rangle = \sum_n f_{12n} f_{34n} G_n(x_1, x_2, x_3, x_4)$$





conformal bootstrap eqn



conformal bootstrap eqn

1. *functional* constraint on CFT data - must be satisfied for every

*X*₁,...,*X*₄

2. one equation for every 4-point function

Concrete application: 3d Ising model CFT

Many other applications exist:

- Wilson-Fisher fixed point in 2<d<4
- O(N) models in d=3

...

- SUSY conformal theories in d=2,3,4,6 dimensions

Local operators

3d Ising model CFT can be def'd as IR fixed point of

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These are just `names', because the operators are strongly renormalized.

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E.g. $[\phi^2]=1$ in the UV (free theory) becomes $[\phi^2] \approx 1.41$ in the IR (CFT) In bootstrap, ϕ -content of operators is irrelevant (only their dimension, spin, and Z₂ count) In bootstrap, ϕ -content of operators is irrelevant (only their dimension, spin, and Z₂ count)

Very important experimental fact: 3d Ising CFT contains 2 and only 2 relevant (i.e. $\Delta < 3$) scalar operators:

- one Z2-odd, **call it** σ (it's coming from ϕ in the UV) - one Z2-even, **call it** ϵ (coming from ϕ^2) In bootstrap, ϕ -content of operators is irrelevant (only their dimension, spin, and Z₂ count)

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Coupling $\langle \sigma \sigma \varepsilon \rangle$ allowed by Z2 symmetry $\Rightarrow f_{\sigma\sigma\varepsilon} \neq 0$

More generally, consider OPE $\sigma \, x \, \sigma$

It contains all operators A such that $f_{\sigma\sigma A} \neq 0$ (infinitely many per each spin)

$$\sigma \times \sigma = \mathbf{1} + \varepsilon + \varepsilon' + \dots \quad (\ell = 0)$$
$$+ T_{\mu\nu} + T'_{\mu\nu} + \dots \quad (\ell = 2)$$
$$+ (\ell = 4, 6, \dots)$$

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To solve the theory we must find dimensions and f's for all these operators

Bootstrap Oracle

The algorithm explores the space of CFT data and rigorously rules out the impossible sets



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$(\Delta\sigma, \Delta\varepsilon)$ -oracle using $<\sigma\sigma\sigma\sigma\sigma>$

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Zoom:

Conjecture: 3d Ising minimizes central charge known fact in d=2

From minimum localization:

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All other operators and their f's become fixed at the minimum:

$$\Delta_{\epsilon} = 1.41267(13) \qquad f_{\sigma\sigma\epsilon}^2 = 1.10636(9) \\ \Delta_{\epsilon'} = 3.8303(18) \qquad f_{\sigma\sigma\epsilon'}^2 = 0.002810(6)$$

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These are **world's most precise** determinations of 3d Ising model CFT data

$(\Delta\sigma, \Delta\varepsilon)$ -oracle using $<\sigma\sigma\sigma\sigma\sigma>$, $<\varepsilon\varepsilon\varepsilon\varepsilon>$, $<\sigma\sigma\varepsilon\varepsilon>$ Kos,Poland,Simmons-Duffin 2014



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Strongly coupled CFT computation with rigorous error bars

Recall bootstrap:

$$\sum_{n} f_{12n} f_{34n} \times (...) = 2 \leftrightarrow 4$$
Quadratic equation for f's

But if identical operators :

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Is intersection nonempty?

Linear programming Dantzig 1947

Crucial: $p_n \ge 0$ is a convex condition



Several correlators <0000>, <2222>, <0022>

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non sign-definite

$$\mathbb{P}_{n} = \begin{pmatrix} f_{\sigma\sigma n}^{2} & f_{\sigma\sigma n} f_{\epsilon\epsilon n} \\ f_{\sigma\sigma n} f_{\epsilon\epsilon n} & f_{\epsilon\epsilon n}^{2} \end{pmatrix} \succeq 0$$

positive semidefinite $(\lambda_1 = 0, \lambda_2 = f_{\sigma\sigma n}^2 + f_{\epsilon\epsilon n}^2)$

Several correlators <0000>, <222>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <2002>, <20

$$\sum_{n} f_{\sigma\sigma n}^{2} \times (...) = 2 \leftrightarrow 4 \qquad \sum_{n} f_{\epsilon\epsilon n}^{2} \times (...) = 2 \leftrightarrow 4$$
$$\sum_{n} f_{\sigma\sigma n} f_{\epsilon\epsilon n} \times (...) = \sum_{n} f_{\sigma\epsilon n}^{2} \times (...)$$
non sign-definite
$$\mathbb{P}_{n} = \begin{pmatrix} f_{\sigma\sigma n}^{2} & f_{\sigma\sigma n} f_{\epsilon\epsilon n} \\ f_{\sigma\sigma n} f_{\epsilon\epsilon n} & f_{\epsilon\epsilon n}^{2} \end{pmatrix} \succeq 0$$
positive semidefinite $(\lambda_{1} = 0, \lambda_{2} = f_{\sigma\sigma n}^{2} + f_{\epsilon\epsilon n}^{2})$ Also a convex condition!

Linear programming \rightarrow Semidefinite programming

Conclusions and future

• conformal bootstrap works in any d

for some models (Ising-3, O(N)) even better than expected
lots of mileage out of a few constraints

Short-term: go through the list of known CFTs

Long-term: classify CFTs with a small number of low dimension operators

Backup

RG vs Monte Carlo and experiment

Ising-3		Monte Carlo Hasenbusch 2010 (20 CPU-years)	RG Guida, Zinn-Justin 1998 (5-7 loops + Borel)
	Δ(φ)	0.51814(5)	0.51675(125)
	Δ(φ²)	1.41275(25)	1.41370(330)

agrees, with O(10) larger errors

O(2) d=3		Experiment Lipa et al 1996	MC+HT Campostrini et al 2006	RG Guida, Zinn-Justin 1998
	Δ(φ²)	1.5094(2)	1.5112(2)	1.5081(33)
		RG inconclusiv		

QCD conformal window

SU(3) massless gauge theory in d=3+1 with Nf fermions

Flows to a conformal IR fixed point for $N_F^* \leq N_F \leq 16$

