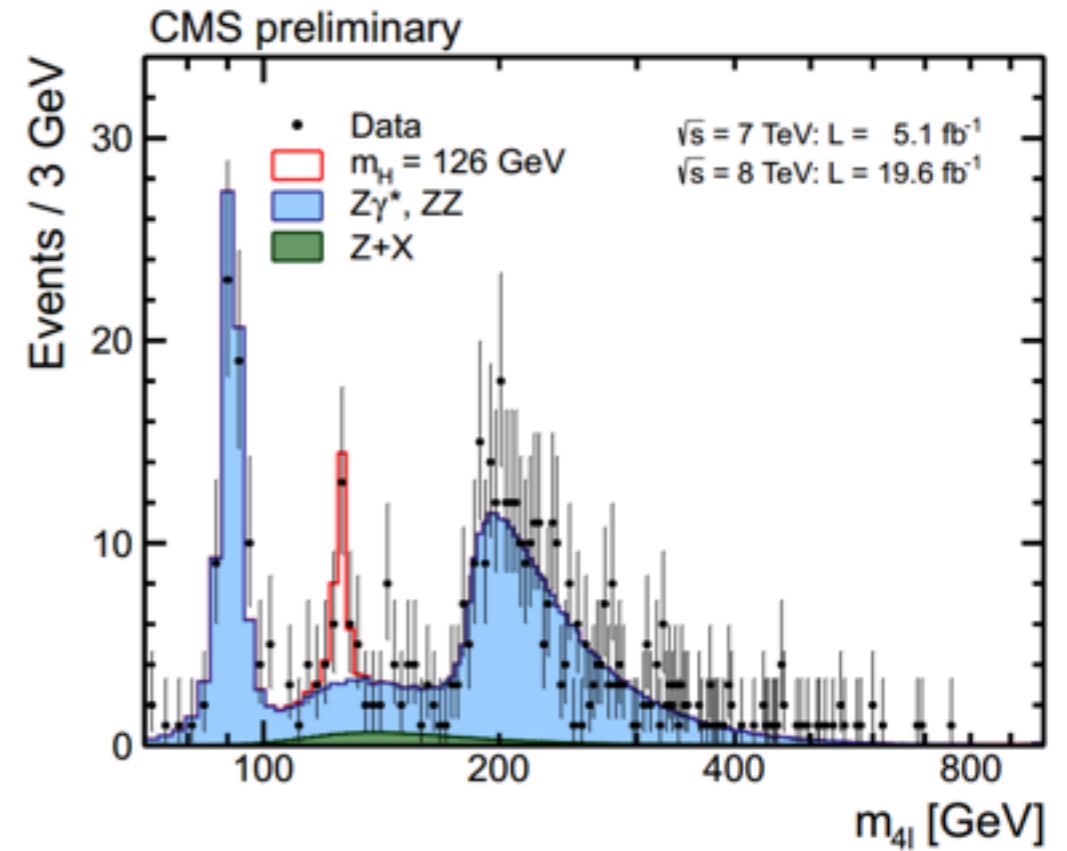
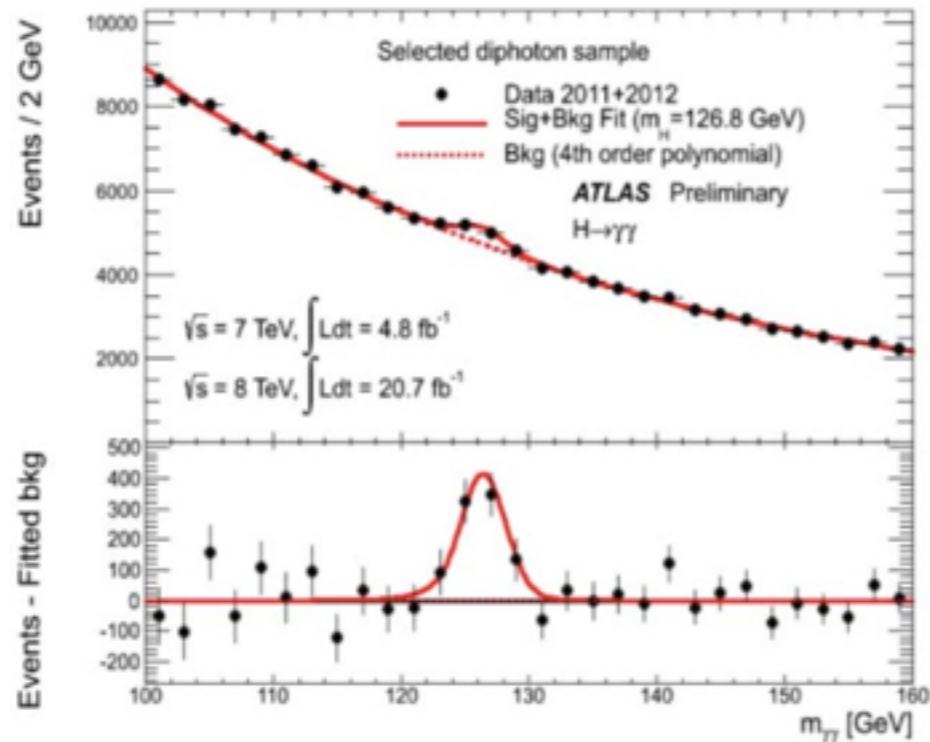


Higgs and Dark Matter Phenomenology at Colliders

Michael Spannowsky

IPPP, Durham University

Combined results for each experiment



- Huge **international** and **intergenerational success!**
- First observed in clean final states: photons, ZZ, WW
- Now more channels, e.g. taus
- In absence of other resonances Higgs is window to new physics



Results from Run 1: 'The End of the Beginning'

Mass:

ATLAS	CMS (new ZZ(4l) not used)
$125.5 \pm 0.2 \text{ (stat)}^{+0.5}_{-0.6} \text{ (syst) GeV}$	$125.7 \pm 0.3 \text{ (stat)} \pm 0.3 \text{ (syst) GeV}$

Spin

▶ Tested spin-1 and 0^- excluded with 1-CLs > 0.99%

CP

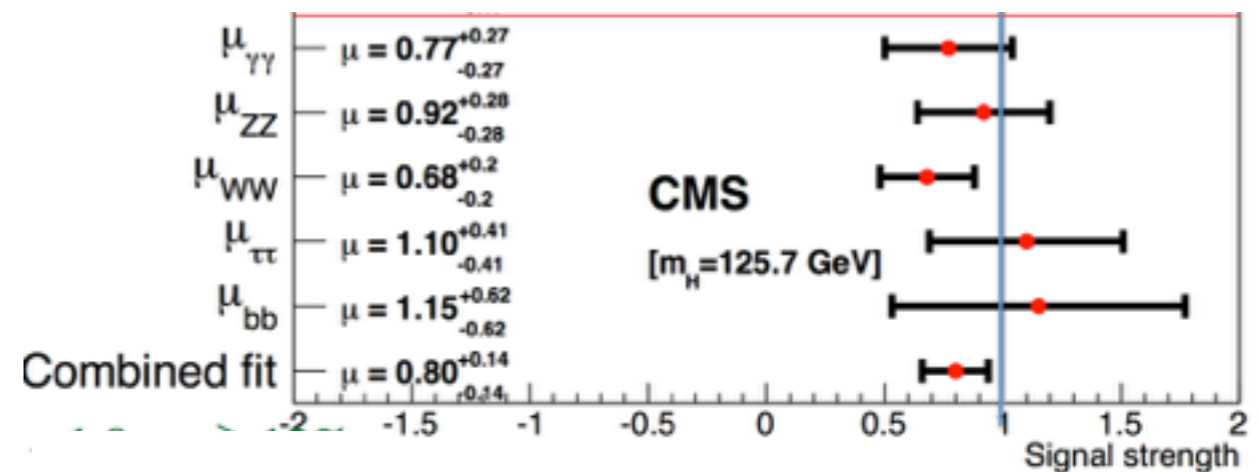
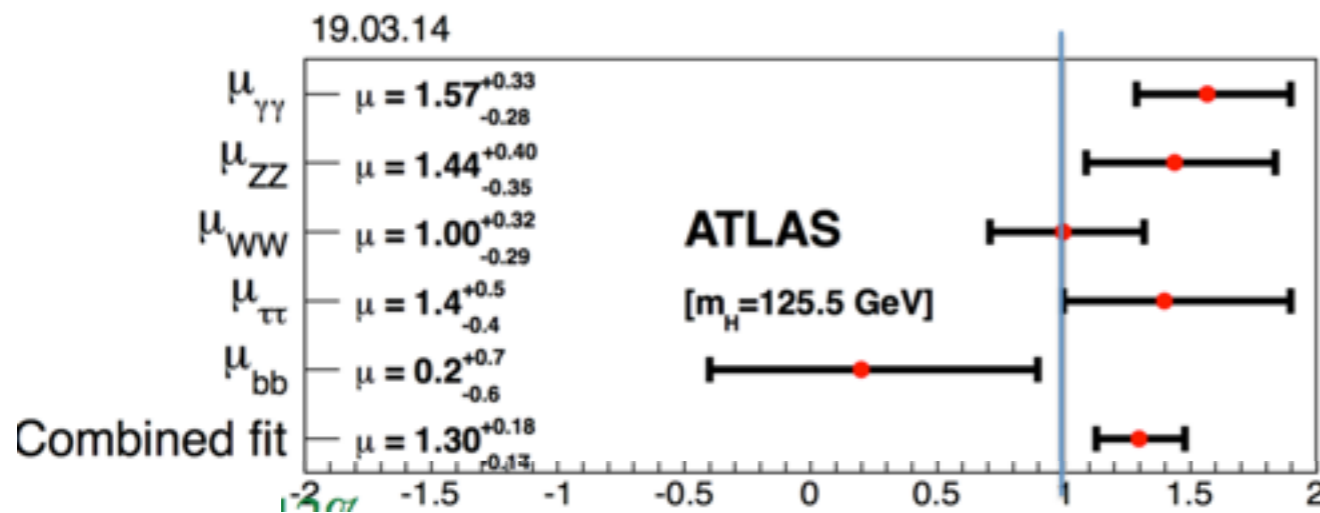
▶ Tested Spin-2 models excluded with 1-CLs > 0.95%

width

▶ Combine 4l and 2l2n decay channels.

▶ **Observed (expected) 95%CL limits: $\Gamma < 8.5(4.2) \times \Gamma_{SM}$**

Couplings



Organisation

YETI

BUSSTEPP

RAL

S.Ex. Fellowships

ITNs:

Workshops

HXWG

Width

HEJ

MCNet

Joint Exp-Theory

Higgs couplings

Sherpa

HiggsTools

Postdocs

ERC grant

Herwig++

IPPP

Portals

SUSY

Jet substructure
tools

Composite Higgs

BlackHat

VBFNLO

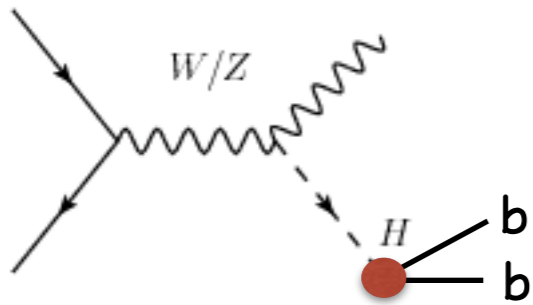
Simplified Models

EFT vs non-EFT

Support

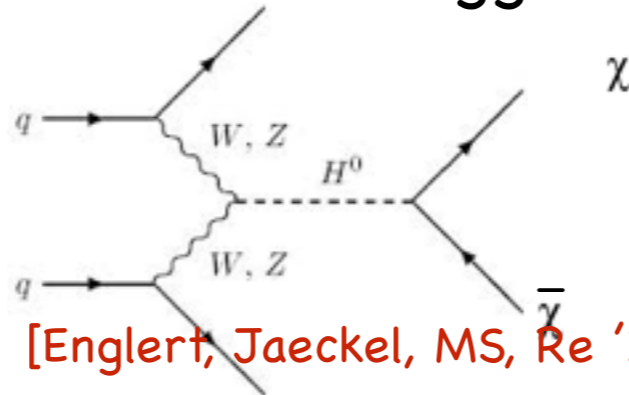
Interpretation

Higgs-bottom coupling



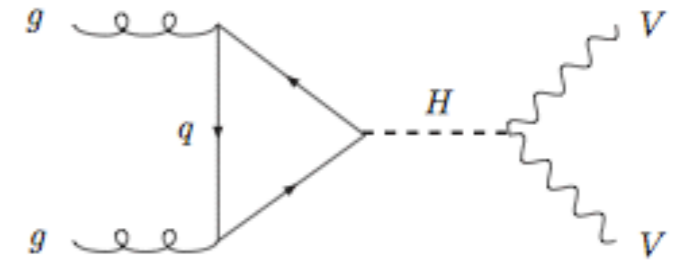
[Soper, MS '10]
[Soper, MS '11]

Invisible Higgs



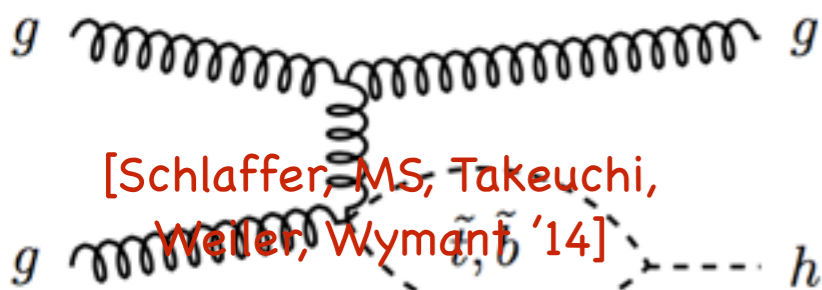
[Englert, Jaeckel, MS, Re '11]
[Englert, MS, Wymant '13]
[Bernacek, Plehn, Schichtel, Tattersall '14]

Off-shell Higgs (Width)



[Englert, McCullough, MS '15]
[Englert, MS '14]
[Englert, Soreq, MS '14]

ct vs cg in H+jet

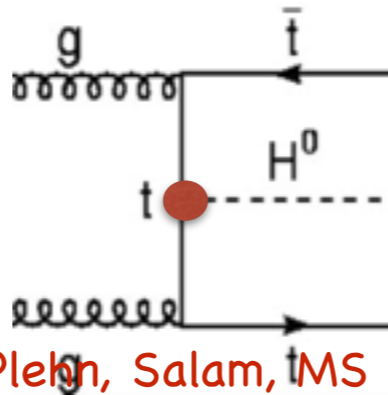


[Schlaffer, MS, Takeuchi, Weiler, Wymant '14]

[Buschmann, Englert, Goncalves, Plehn, MS '14]

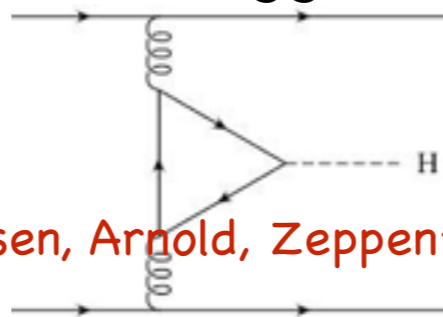
[Buschmann, Goncalves, Kutimalai, Krauss, Plehn, '14]

Higgs-top coupling



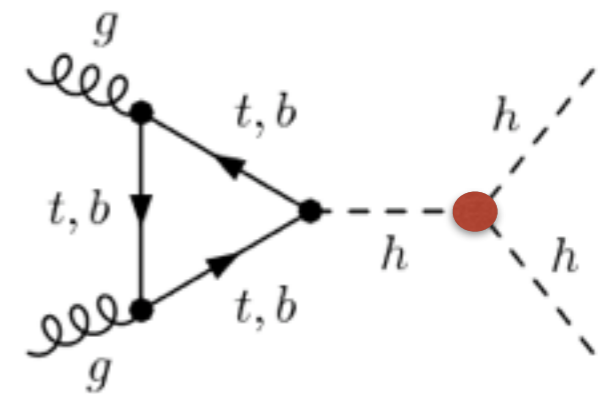
[Plehn, Salam, MS '09]
[Artoisenet, de Aquino, Maltoni, Mattelaer '09]

CP Higgs



[Andersen, Arnold, Zeppenfeld '10]
[Dolan, Harris, Jankowiak, MS '14]

Higgs selfcoupling



[Dolan, Englert, MS '13 '14]
[Barr, Dolan, Englert, MS '14]
[Ferreira, Papaefstathiou, MS '14]
[Dolan, Englert, Greiner, MS '14]
[Englert, Krauss, MS, Thompson '14]

Constraining the Higgs width at the LHC?

- alternative method using interference effects directly see [Dixon, Li '13]

Constraining the Higgs boson width with ZZ production at the LHC

Fabrizio Caola^{1,*} and Kirill Melnikov^{1,†}

¹*Department of Physics and Astronomy, Johns Hopkins University, Baltimore, USA*

We point out that existing measurements of $pp \rightarrow ZZ$ cross-section at the LHC in a broad range of ZZ invariant masses allow one to derive a model-independent upper bound on the Higgs boson width, thanks to strongly enhanced off-shell Higgs contribution. Using CMS data and considering events in the interval of ZZ invariant masses from 100 to 800 GeV, we find $\Gamma_H \leq 38.8 \Gamma_H^{\text{SM}} \approx 163$ MeV, at the 95% confidence level. Restricting ZZ invariant masses to $M_{ZZ} \geq 300$ GeV range, we estimate that this bound can be improved to $\Gamma_H \leq 21 \Gamma_H^{\text{SM}} \approx 88$ MeV.



[Caola, Melnikov PRD 88]

Measurement done in CMS-PAS-HIG-14-002 and presented at Moriond '14

By now ATLAS has performed same measurement



ATLAS

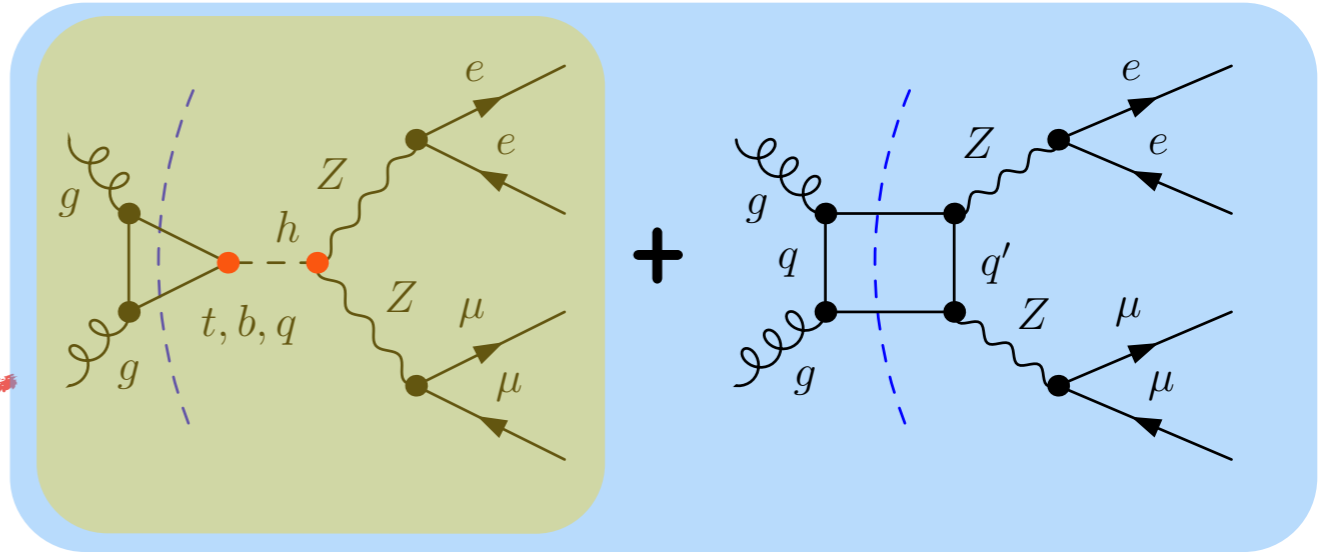
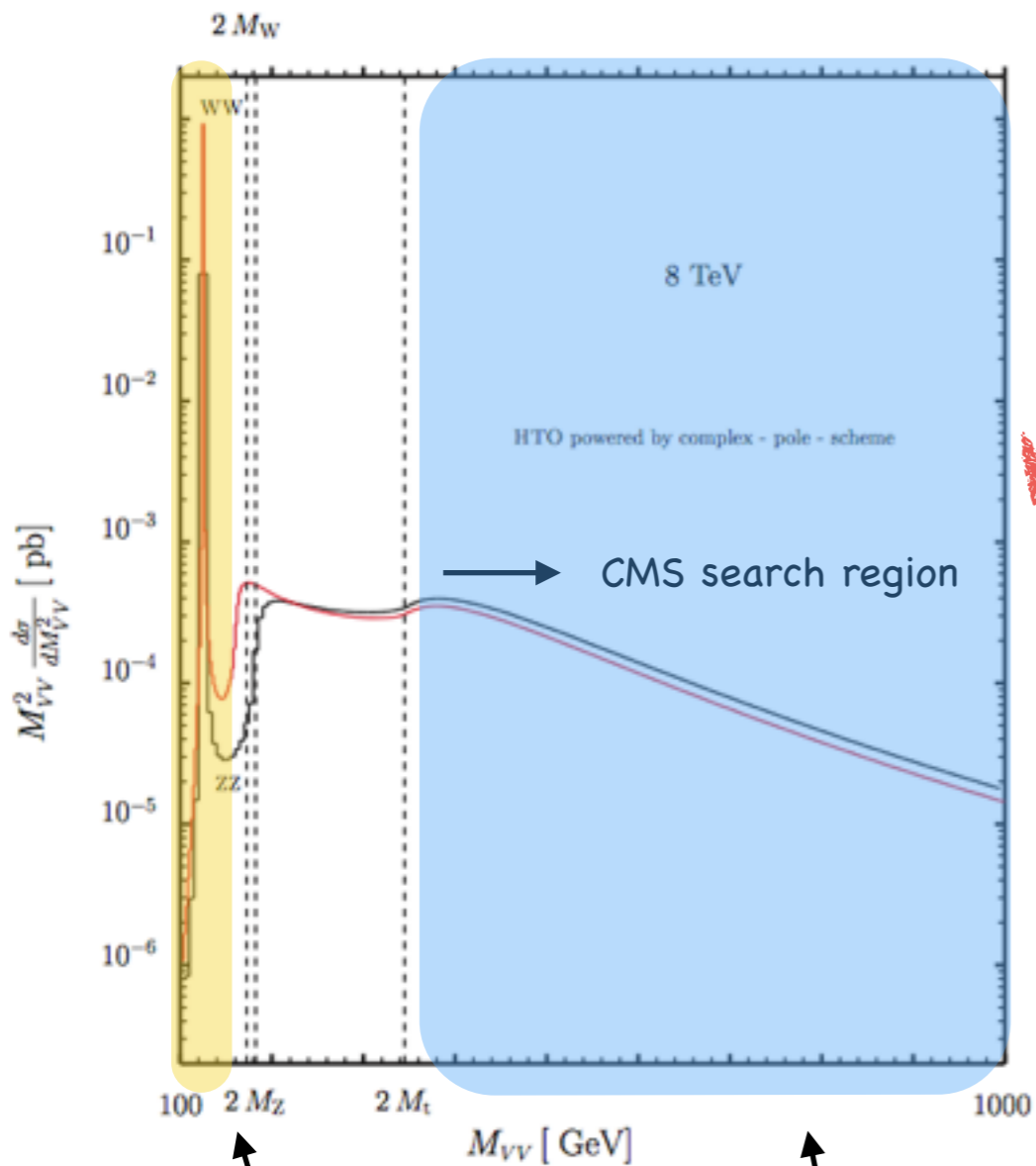
LIFE AND PHYSICS
JON BUTTERWORTH
HOSTED BY THE GUARDIAN

[WIKI.](#)

How wide is a Higgs?

In accord with Heisenberg's uncertainty principle, short-lived particles have uncertain mass. So the Higgs boson, which gives mass to other particles, is uncertain about its own mass. New results from the CMS experiment at the CERN LHC have started to tell us how uncertain

CMS Measurement



I. Count events in on-shell region

→ fix signal strength $\mu_{i,j} = \sigma_{H,i} \times BR_j \sim \frac{g_{ggH} g_{HZZ}}{\Gamma_H}$

II. measure $g_{ggH}^2 g_{HZZ}^2$ in off-shell region using angular correlations of 4l decay products

III. insert off-shell coupling measurement in on-shell signal strength to bound width

$$\sigma_{gg \rightarrow H \rightarrow ZZ}^{\text{on-peak}} \sim \frac{g_{ggH}^2 g_{HZZ}^2}{\Gamma_H}$$

$$\sigma_{gg \rightarrow H \rightarrow ZZ}^{\text{off-peak}} \sim g_{ggH}^2 g_{HZZ}^2$$

Obs.(exp.) @95% C.L:

$$\Gamma_H < 4.2 (8.5) \Gamma_H^{\text{SM}}$$

$$\Gamma_H < 17.4 (35.3) \text{ MeV}$$



[Kauer, Passarino 2011]

Unfortunately, method has loop-holes:

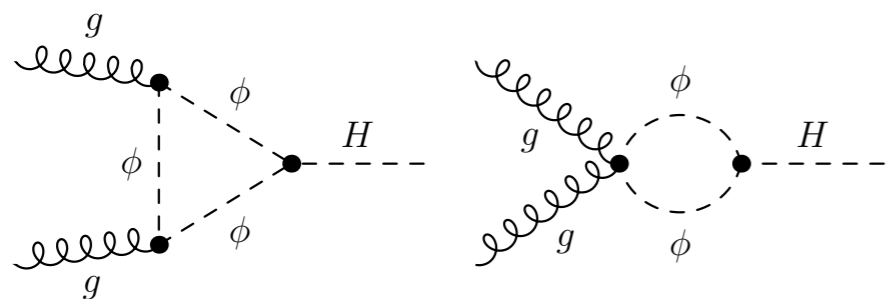
[Englert, MS '14]

[Logan, '15]

- In SM couplings in on-shell and off-shell region intimately related
- Direct correlation of on-shell $g_{ggh}^2 g_{hZZ}^2$ and off-shell $g_{ggh}^2(\sqrt{s}) g_{hZZ}^2(\sqrt{s})$ necessary ingredient for width measurement -> **can be broken by BSM effects**

Example: Higgs-portal (toy) model

$$\mathcal{L}_\phi = |D_\mu \phi|^2 - \tilde{m}_\phi^2 |\phi|^2 - \lambda |\phi|^2 |H|^2 + \dots$$



scalar only charged under $SU(3)_C$

$$m_\phi^2 = \tilde{m}_\phi^2 + \lambda v^2 \quad \text{free parameter}$$

off-shell CS

$$g_{ggh}(m_h) > g_{ggh,SM} \rightarrow \Gamma > \Gamma_{SM} \quad \text{for } \mu \sim 1$$

m_ϕ	μ (h peak)	Γ_h/Γ_h^{SM}	$\bar{\sigma}/\bar{\sigma}^{SM}$ [$m(4\ell) \geq 330$ GeV] ^a
70 GeV	$\simeq 1.0$	$\simeq 5$	-2%
170 GeV	$\simeq 1.0$	$\simeq 4.7$	+80%
170 GeV	$\simeq 1.0$	$\simeq 1.7$	+6%

^aWe impose the cut set used by CMS [17] without the MELA cut [34].

Despite increased on-shell coupling (and Higgs width) negligible contribution in off-shell region

- Here only simplest toy model - thus question: WHEN IS WIDTH INTERPRETATION VALID

- Width interpretation ONLY interesting if model-independent

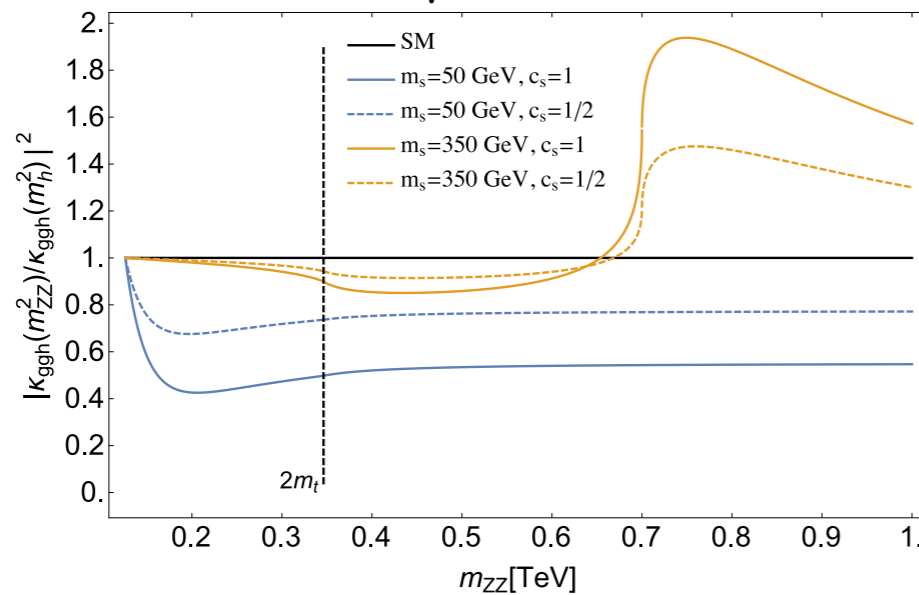
Within a model width is fixed (not free parameter of theory), result of QFT

Width measurement is result of global coupling fit.

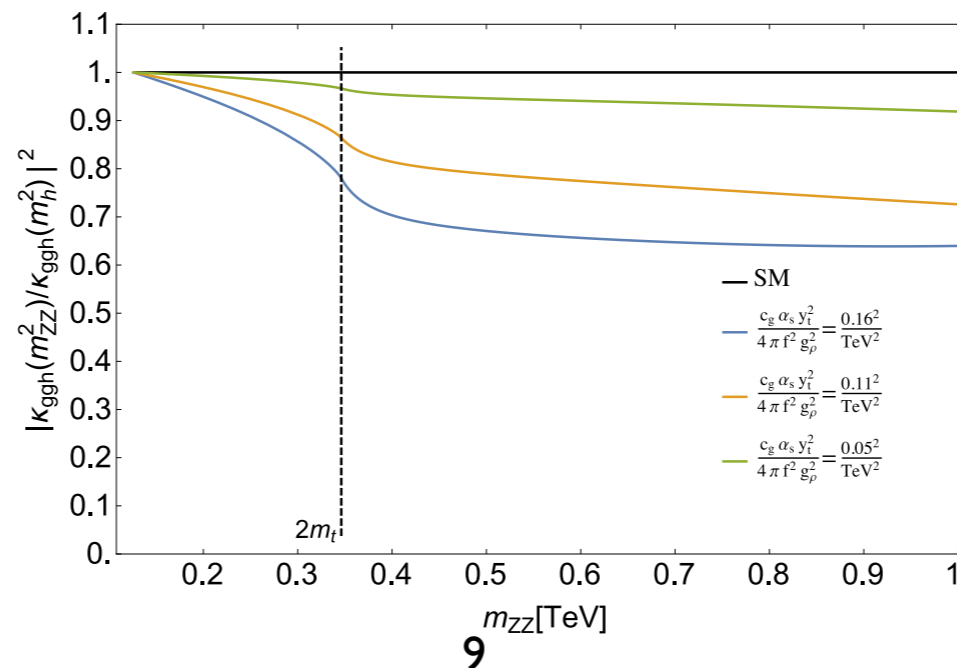
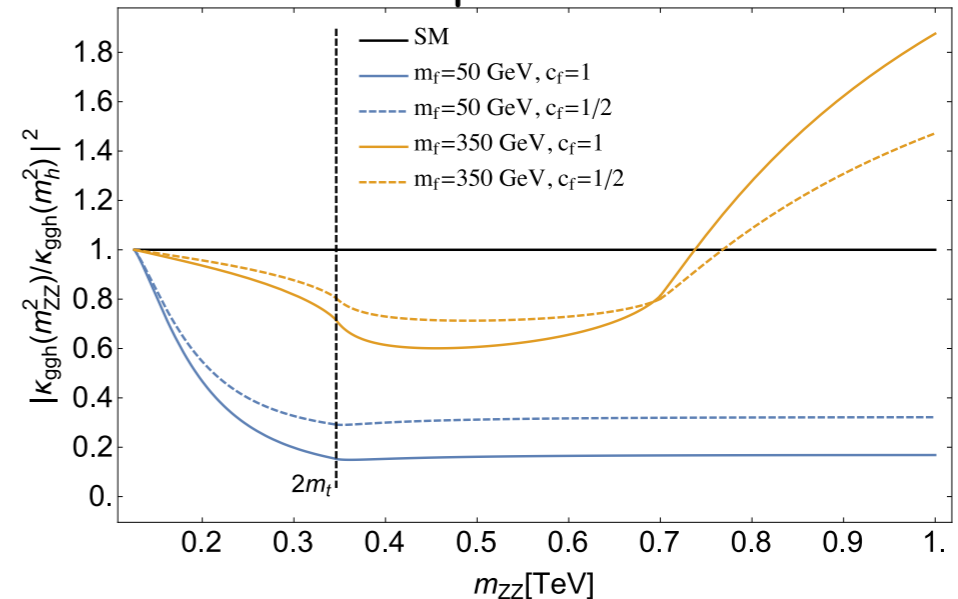
- But for classes of models a width interpretation is valid: [Englert, Soreq, MS '14]

Necessary condition: $R(m_{ZZ}^2) = \frac{ggH(m_{ZZ}^2)/ggH_{SM}(m_{ZZ}^2)}{ggH(m_H^2)/ggH_{SM}(m_H^2)} \simeq 1$ complex valued double ratio

color triplet scalar



color triplet fermions



effective GGHH coupling

- Way to close loophole of Caola-Melnikov method by using WBF process:

Same tree-level coupling in production and decay

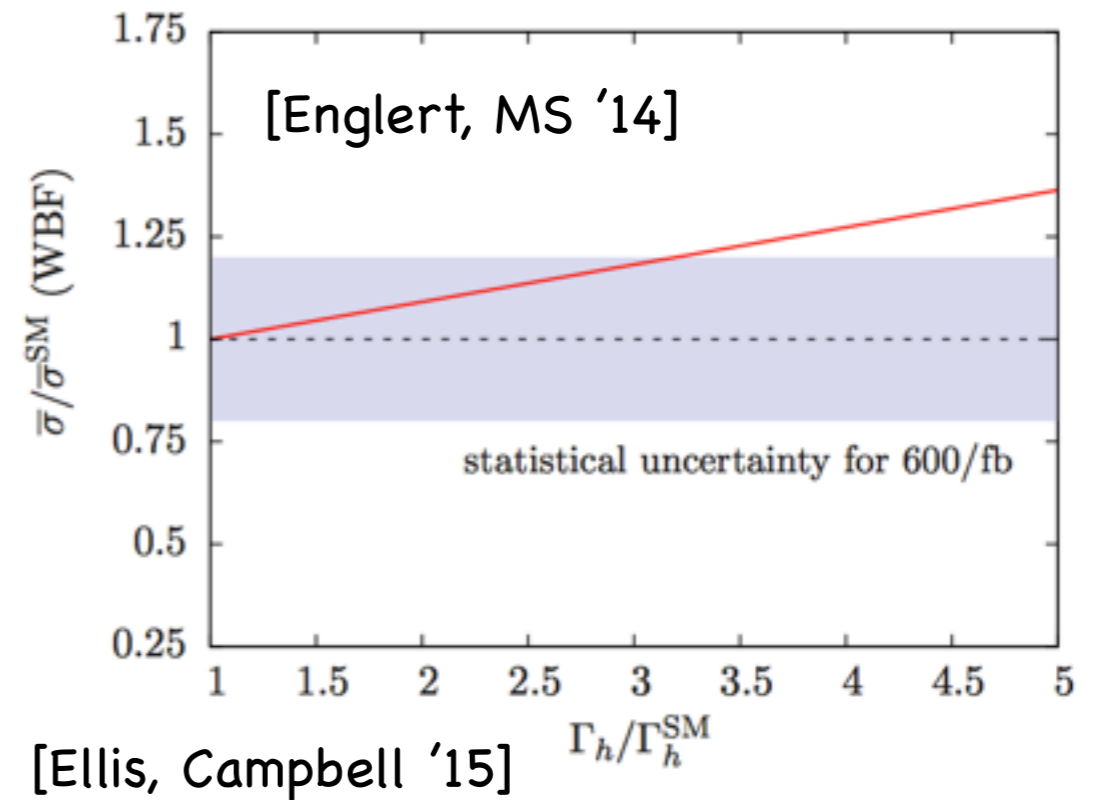
T parameter links WWH and ZZH

$$\mathcal{L}_{HD} = F_{HD} \text{tr} \left[\mathbf{H}^\dagger \mathbf{H} - \frac{v^2}{4} \right] \cdot \text{tr} \left[(\mathbf{D}_\mu \mathbf{H})^\dagger (\mathbf{D}^\mu \mathbf{H}) \right]$$

assumed

$$hW_\mu^+ W_\nu^- : igM_W g_{\mu\nu} \frac{v^2 F_{HD}}{2}$$

$$hZ_\mu Z_\nu : ig \frac{M_W}{\cos^2 \theta_W} g_{\mu\nu} \frac{v^2 F_{HD}}{2}$$

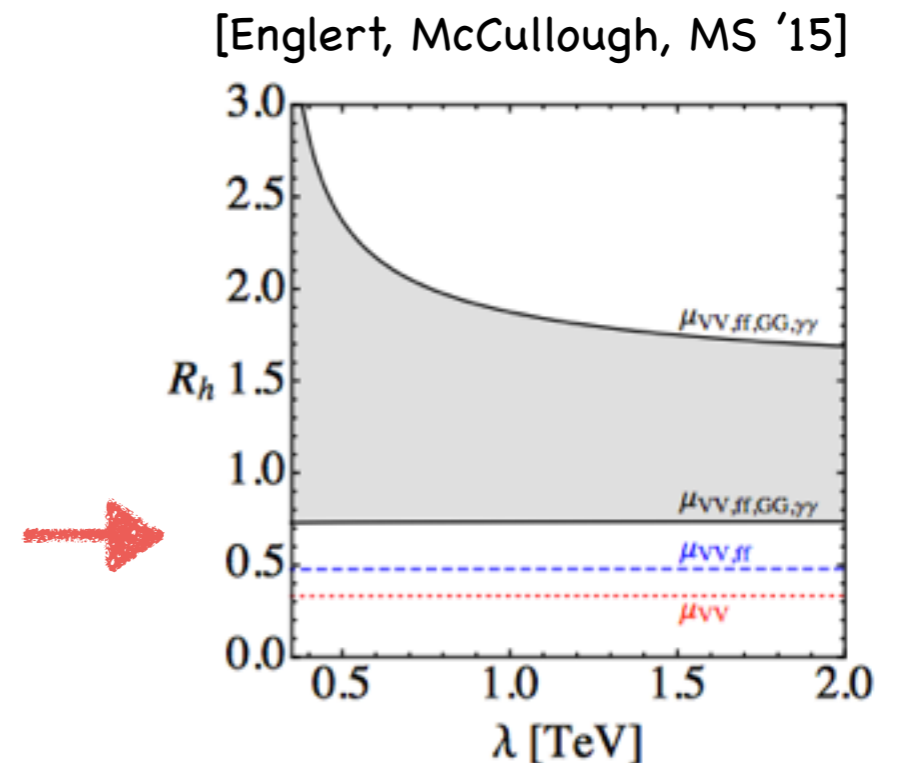


- Use LEP as off-shell Higgs factory:

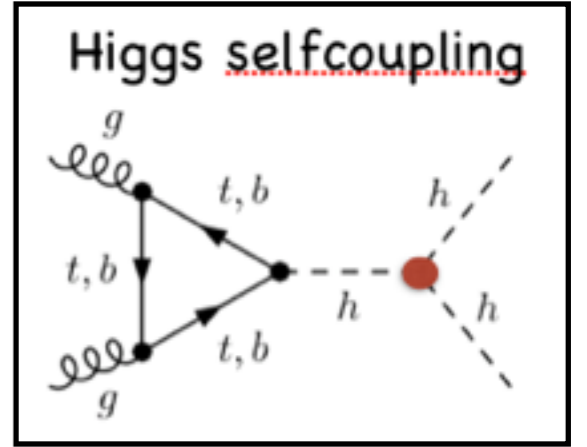
Why off-shell LHC coupling measurement?

LEP has already performed precise Higgs coupling measurements!

Use LEP coupling measurements and plug into LHC signal strength measurement

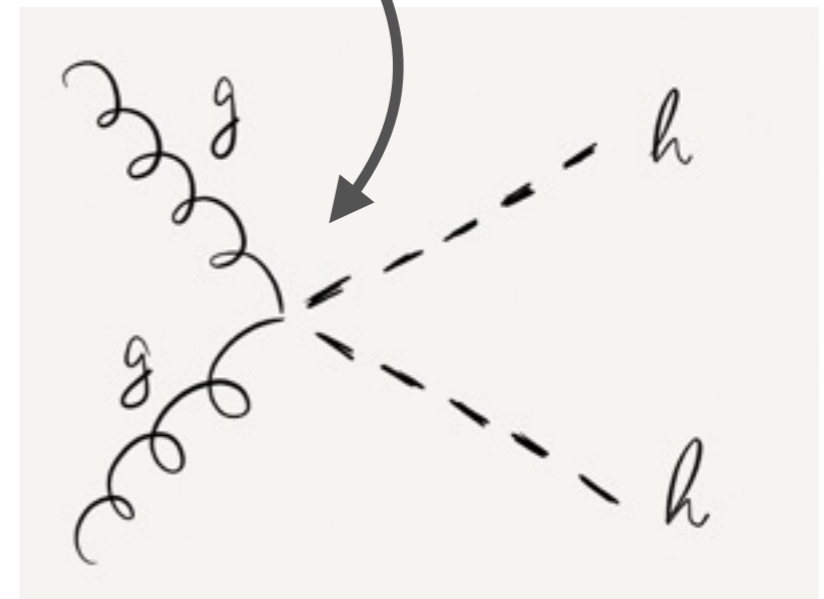
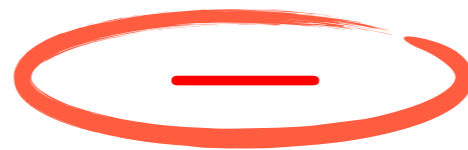
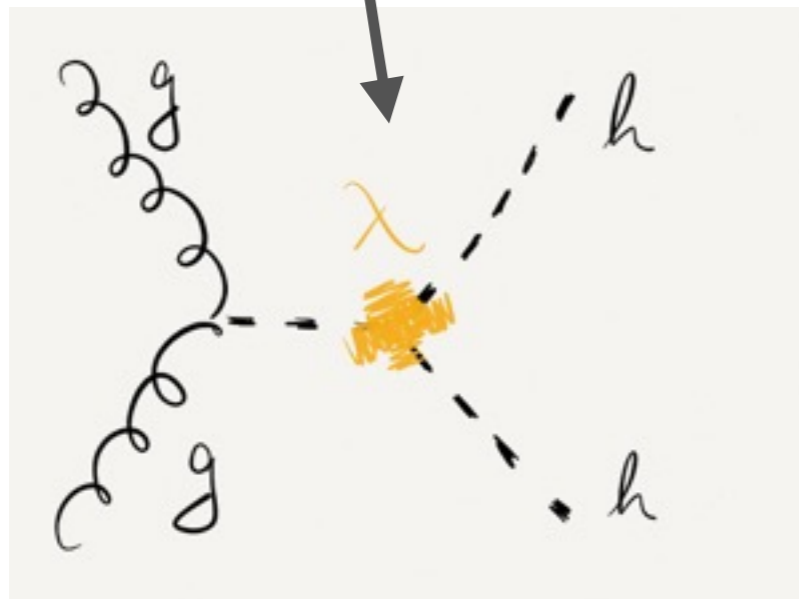


Higgs self-coupling measurements

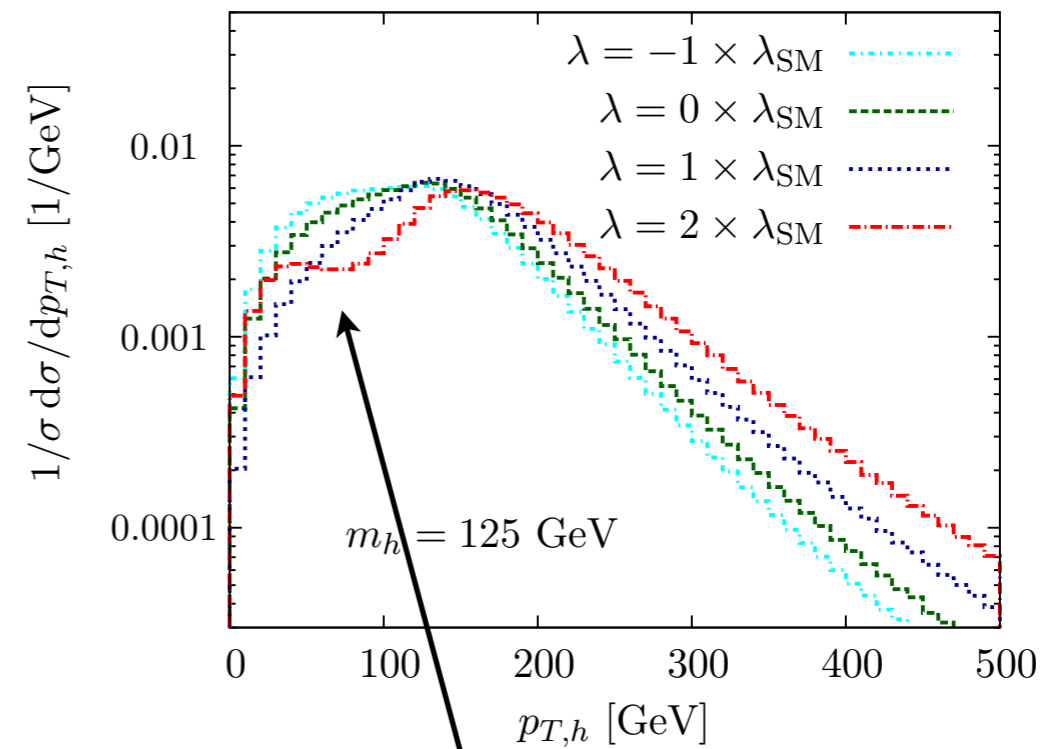
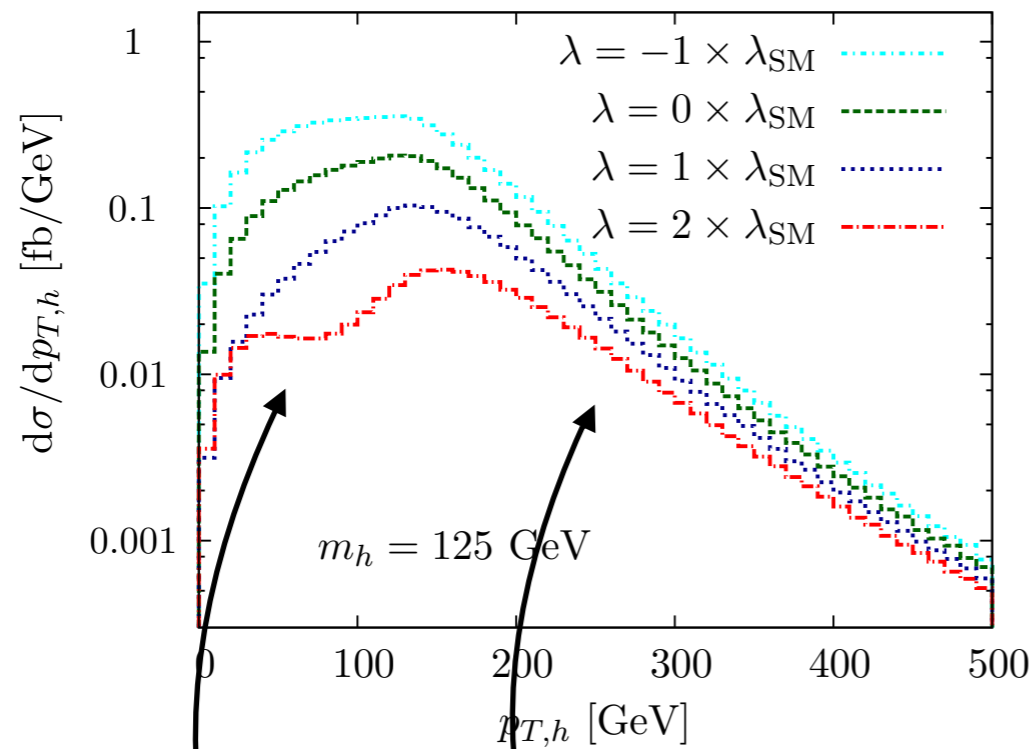


$$\begin{aligned}
 -\mathcal{L} \supset & \frac{1}{2} m_h^2 h^2 + \sqrt{\frac{\eta}{2}} m_h h^3 + \frac{\eta}{4} h^4 \longrightarrow \text{Potential needs at least diHiggs production!} \\
 & - g m_V V^2 h - \frac{m_f}{v} \bar{f} f h \\
 & - \frac{\alpha_s}{12\pi} G_{\mu\nu}^a G^{a\mu\nu} \log(1 + h/v) \\
 & = - \frac{\alpha_s}{12\pi v} G_{\mu\nu}^a G^{a\mu\nu} h + \frac{\alpha_s}{24\pi v^2} G_{\mu\nu}^a G^{a\mu\nu} h^2 + \dots
 \end{aligned}$$

$= \lambda_{\text{SM}} = g^2 m_h^2 / m_W^2$

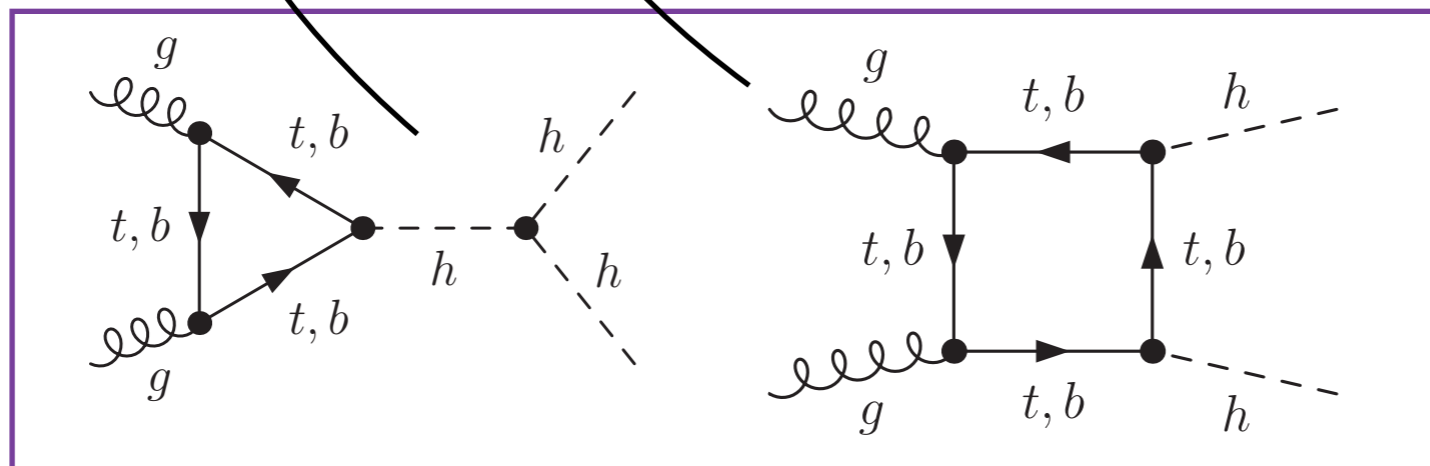
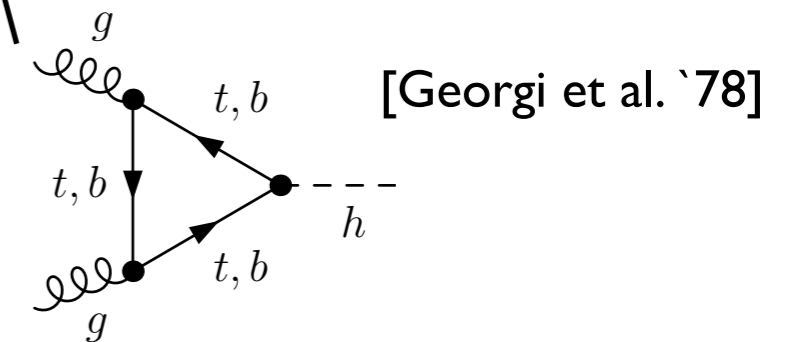


Higgs selfcoupling in HH+X



has maximum contribution for

$$s = (p_{h,1} + p_{h,2})^2 = 4m_t^2$$



Total cross section: [HXSWG]

$$pp \rightarrow HH(\lambda = 1) = 33.86 \text{ fb}$$

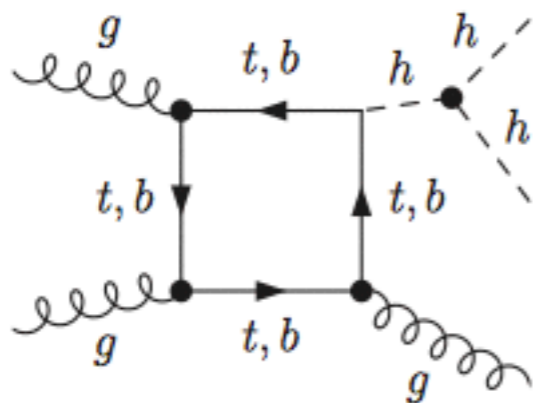
$$pp \rightarrow HH(\lambda = 0) = 71.01 \text{ fb}$$

$$pp \rightarrow HH(\lambda = 2) = 15.85 \text{ fb}$$

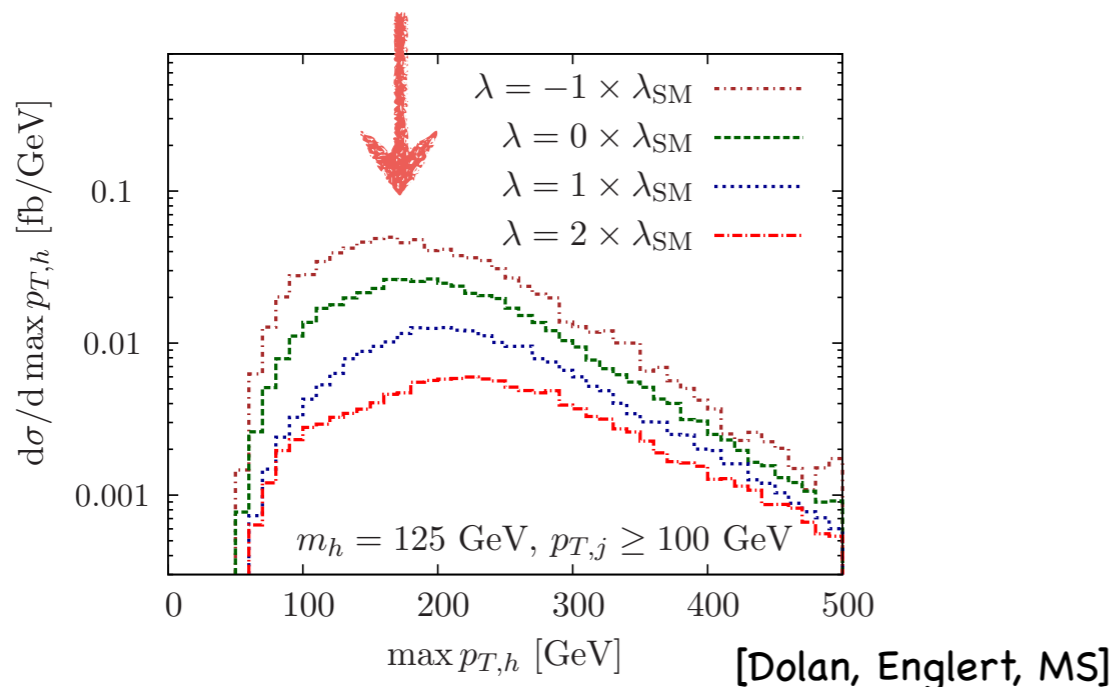
Decay	Issues	Expectation 3000 ifb	References
$b\bar{b}\gamma\gamma$	<ul style="list-style-type: none"> • Signal small • BKG large & difficult to asses • Simple reconst. 	$S/B \simeq 1/3$ $S/\sqrt{B} \simeq 2.5$	[Baur, Plehn, Rainwater] [Yao 1308.6302] [Baglio et al. JHEP 1304] [Azatov et al. JHEP '15]
$b\bar{b}\tau^+\tau^-$	<ul style="list-style-type: none"> • tau rec tough • largest bkg tt • Boost+MT2 might help 	differ a lot $S/B \simeq 1/5$ $S/\sqrt{B} \simeq 5$	[Dolan, Englert, MS] [Barr, Dolan, Englert, MS] [Baglio et al. JHEP 1304]
$b\bar{b}W^+W^-$	<ul style="list-style-type: none"> • looks like tt • Need semilep. W to rec. two H • Boost + BDT proposed 	differ a lot best case: $S/B \simeq 1.5$ $S/\sqrt{B} \simeq 8.2$	[Dolan, Englert, MS] [Baglio et al. JHEP 1304] [Papaefstathiou, Yang, Zurita 1209.1489]
$b\bar{b}b\bar{b}$	<ul style="list-style-type: none"> • Trigger issue (high pT kill signal) • 4b background large difficult with MC • Subjets might help 	$S/B \simeq 0.02$ $S/\sqrt{B} \leq 2.0$	[Dolan, Englert, MS] [Ferreira de Lima, Papaefstathiou, MS] [Wardrope et al, 1410.2794]
others	<ul style="list-style-type: none"> • Many taus/W not clear if 2 Higgs • Zs, photons no rate 		

More jets more fun:

First to calculate HH+jet and HH+2jets beyond effective theory



keep m_{HH} small \rightarrow retain sensitivity for high- p_T Higgs

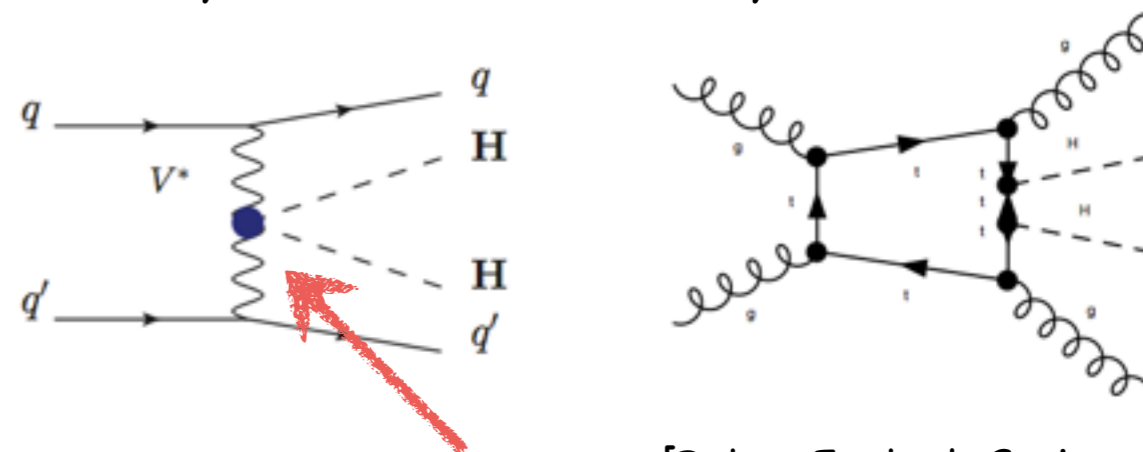


Additional jet can help to suppress backgrounds:

$$hhj \rightarrow b\bar{b}\tau^+\tau^-j \quad S/B \sim 3/2$$

$$hh \rightarrow b\bar{b}\tau^+\tau^- \quad S/B \sim 1/2$$

Unfortunately rate small

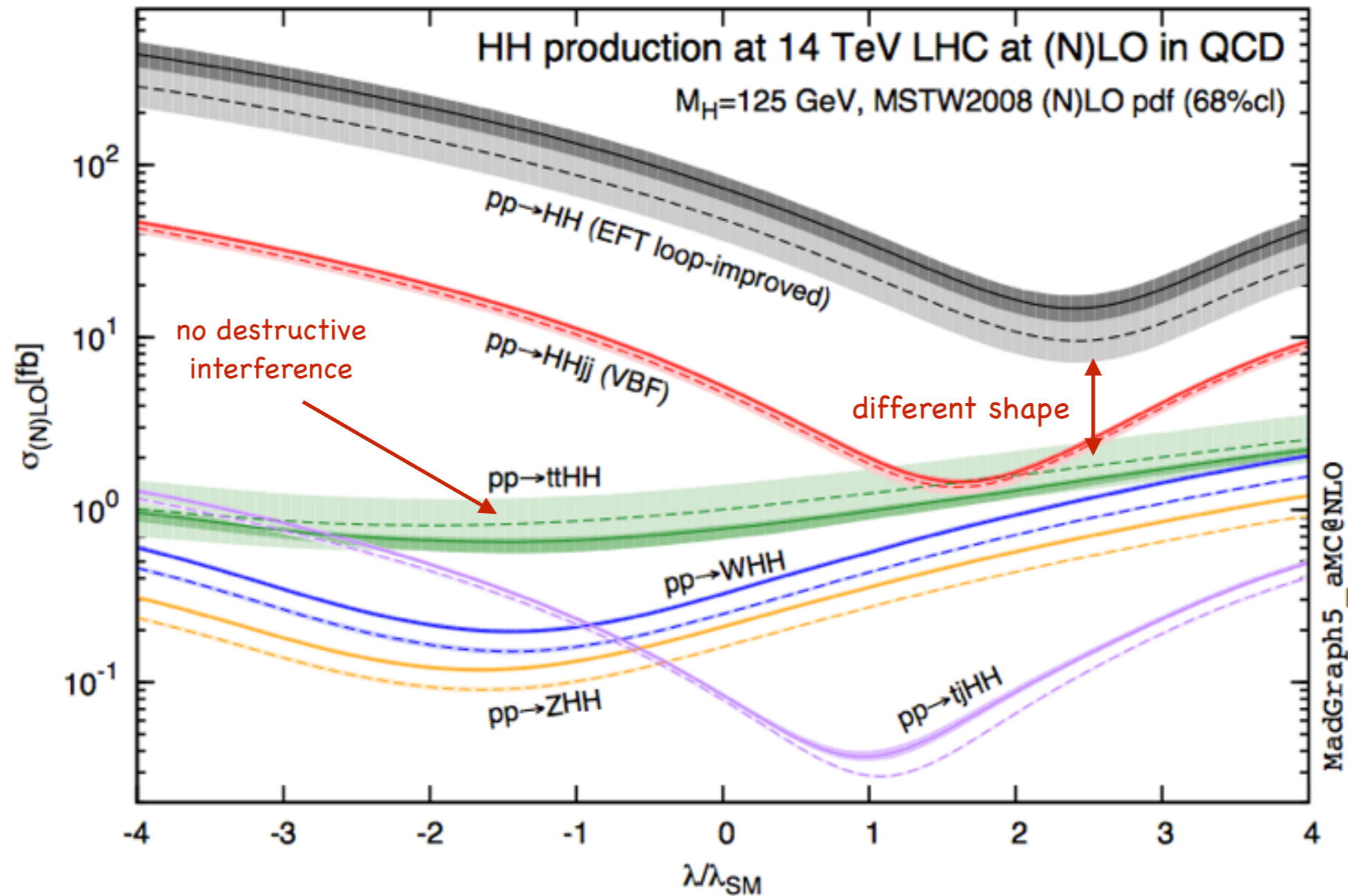


[Dolan, Englert, Greiner, MS]

- Want to study VVHH
Directly related to long. gauge boson scattering
Though strongly modified in comp. Higgs models
- In SM fixed: $g_{WWhh} = e^2/(2s_w^2)$
 $g_{ZZhh} = e^2/(2c_w^2s_w^2)$
- Unfortunately gluon fusion dominating over WBF
Usual WBF cuts, e.g. central jet vetos not applicable.
- WBF only measurable for large enhancement of SM coupling value

Higgs selfcoupling in ttHH

[Frederix, Frixione, Hirschi, Maltoni, Mattelaer, Torielli, Vryonidou, Zaro '14]



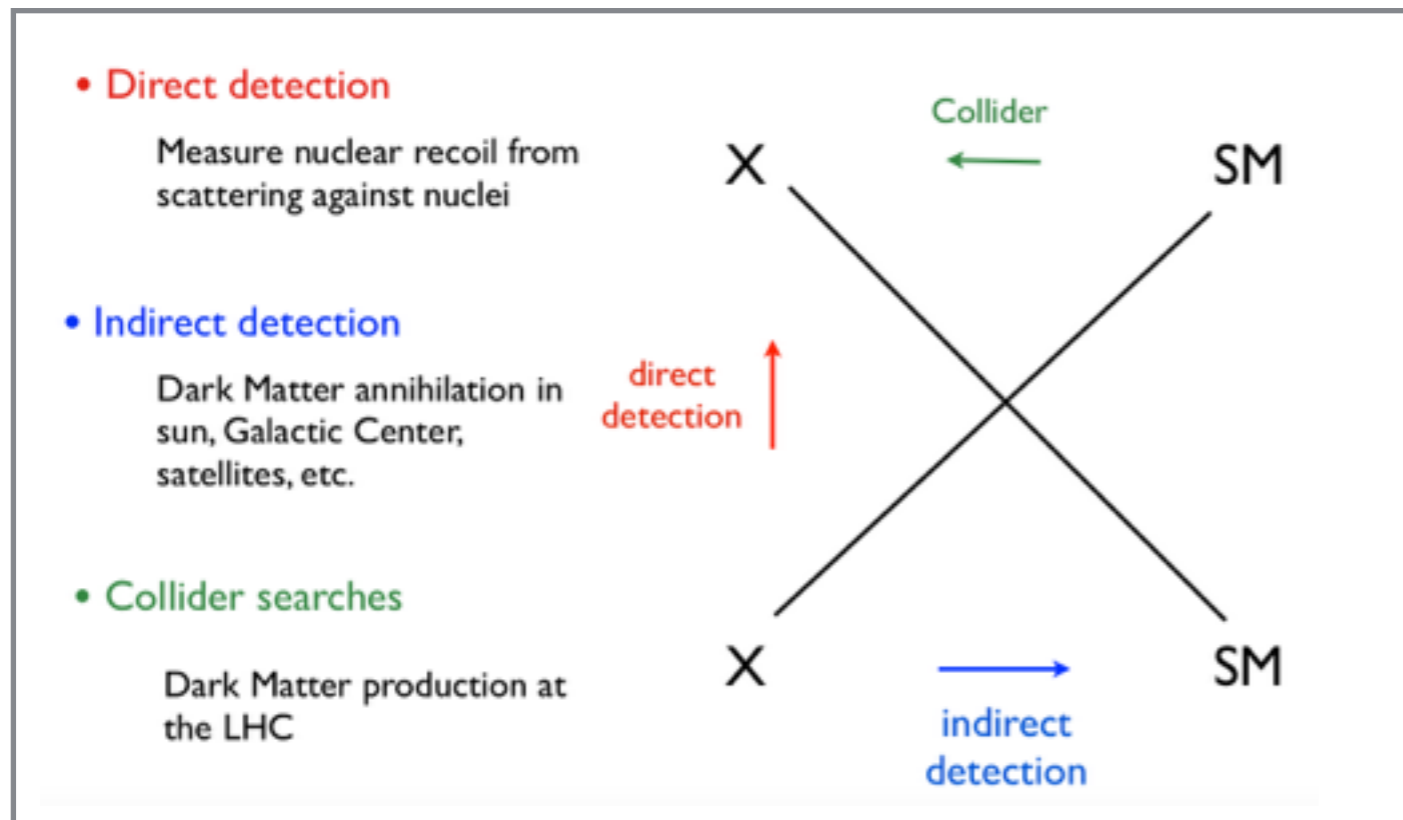
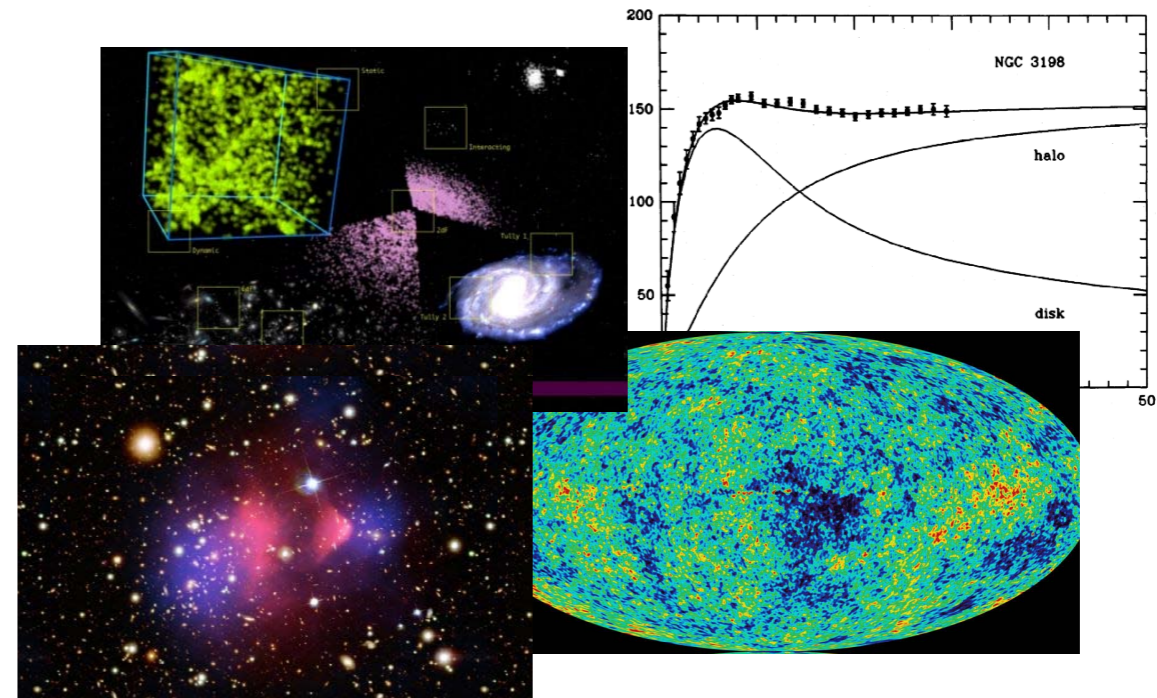
Enhanced selfcoupling can be measured in tthh

[Englert, Krauss, MS, Thompson '14]

Higgs/Scalars and their Dark Matter relation

Evidence for Dark Matter overwhelming:

- Spiral Galaxy rotation curves
- Gravitational lensing
- Acoustic peaks

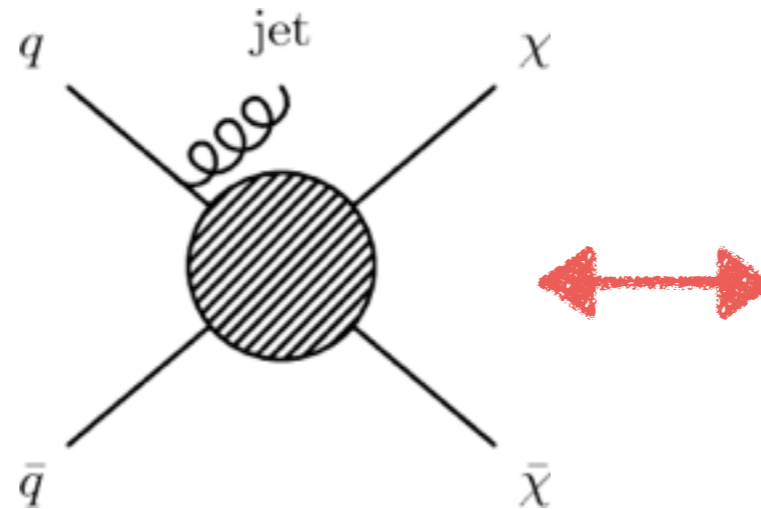


Several ways to look for Dark Matter

Which way more sensitive depends mostly on nature of mediator

Effective theory approach:

- Parametrise interactions in terms of eff. operator
- Simplest way of capturing interactions



➔ Used to be preferred choice of experiments to present results

- However, only valid if interaction not resolved

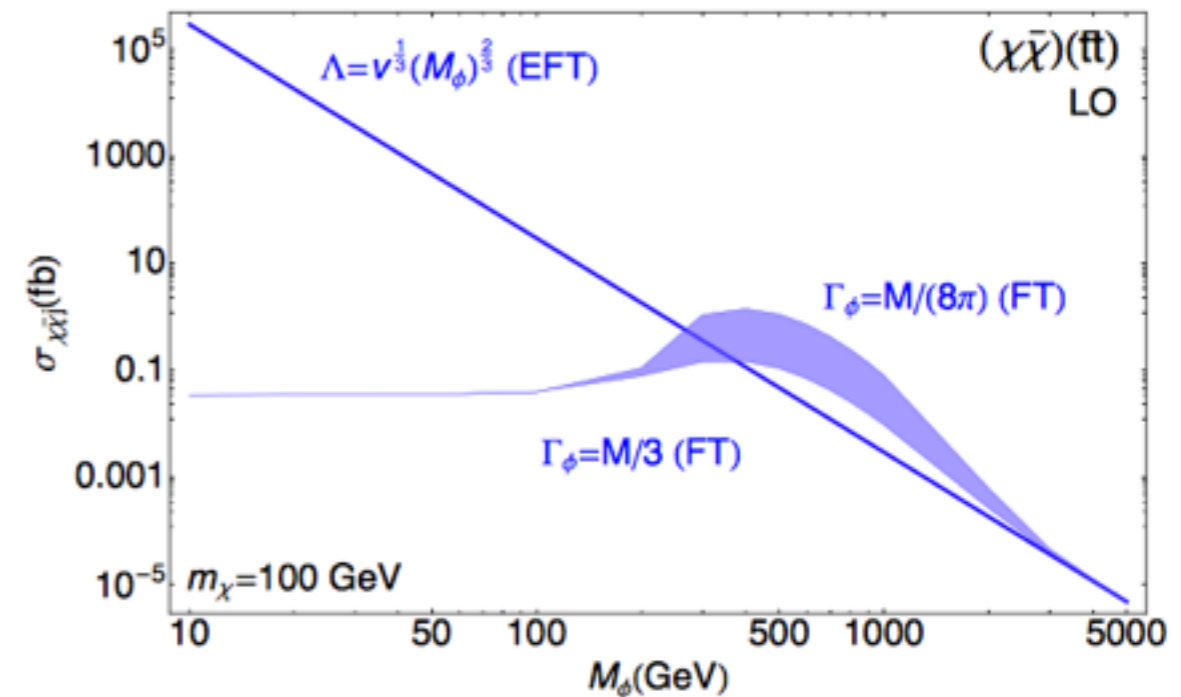
Name	Operator	Coefficient
D1	$\bar{\chi}\chi\bar{q}q$	m_q/M_*^3
D2	$\bar{\chi}\gamma^5\chi\bar{q}q$	im_q/M_*^3
D3	$\bar{\chi}\chi\bar{q}\gamma^5q$	im_q/M_*^3
D4	$\bar{\chi}\gamma^5\chi\bar{q}\gamma^5q$	m_q/M_*^3
D5	$\bar{\chi}\gamma^\mu\chi\bar{q}\gamma_\mu q$	$1/M_*^2$
D6	$\bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu q$	$1/M_*^2$
D7	$\bar{\chi}\gamma^\mu\chi\bar{q}\gamma_\mu\gamma^5q$	$1/M_*^2$
D8	$\bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu\gamma^5q$	$1/M_*^2$
D9	$\bar{\chi}\sigma^{\mu\nu}\chi\bar{q}\sigma_{\mu\nu}q$	$1/M_*^2$
D10	$\bar{\chi}\sigma_{\mu\nu}\gamma^5\chi\bar{q}\sigma_{\alpha\beta}q$	i/M_*^2
D11	$\bar{\chi}\chi G_{\mu\nu}G^{\mu\nu}$	$\alpha_s/4M_*^3$
D12	$\bar{\chi}\gamma^5\chi G_{\mu\nu}G^{\mu\nu}$	$i\alpha_s/4M_*^3$
D13	$\bar{\chi}\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$	$i\alpha_s/4M_*^3$
D14	$\bar{\chi}\gamma^5\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$	$\alpha_s/4M_*^3$

Going beyond:

- At colliders momentum transfer too large for EFT approach

➔ Need simplified models

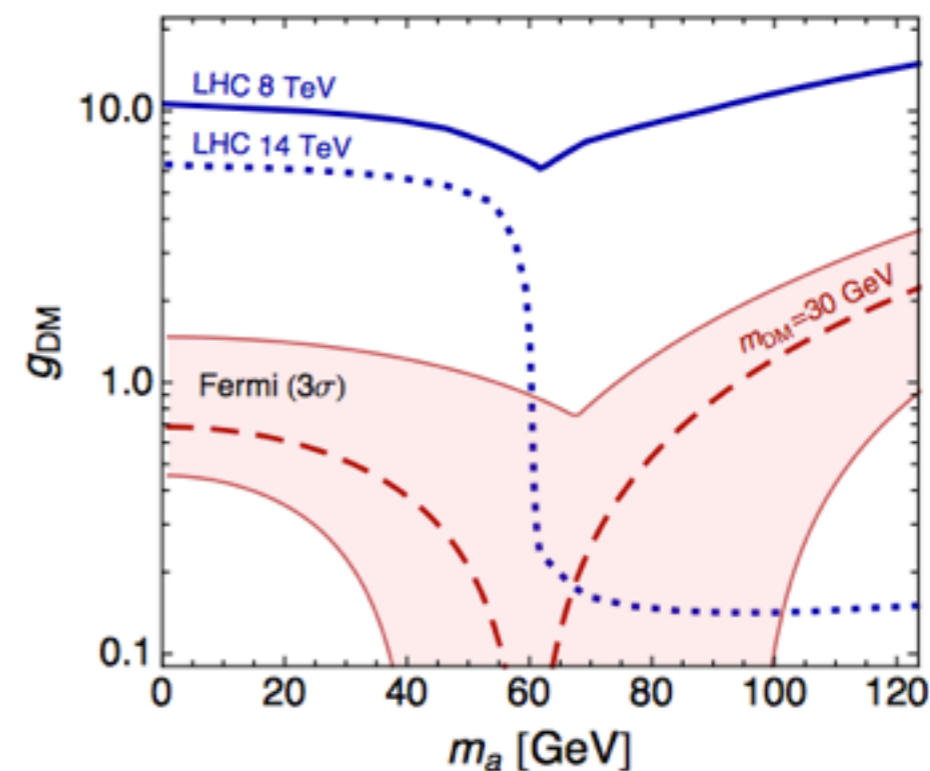
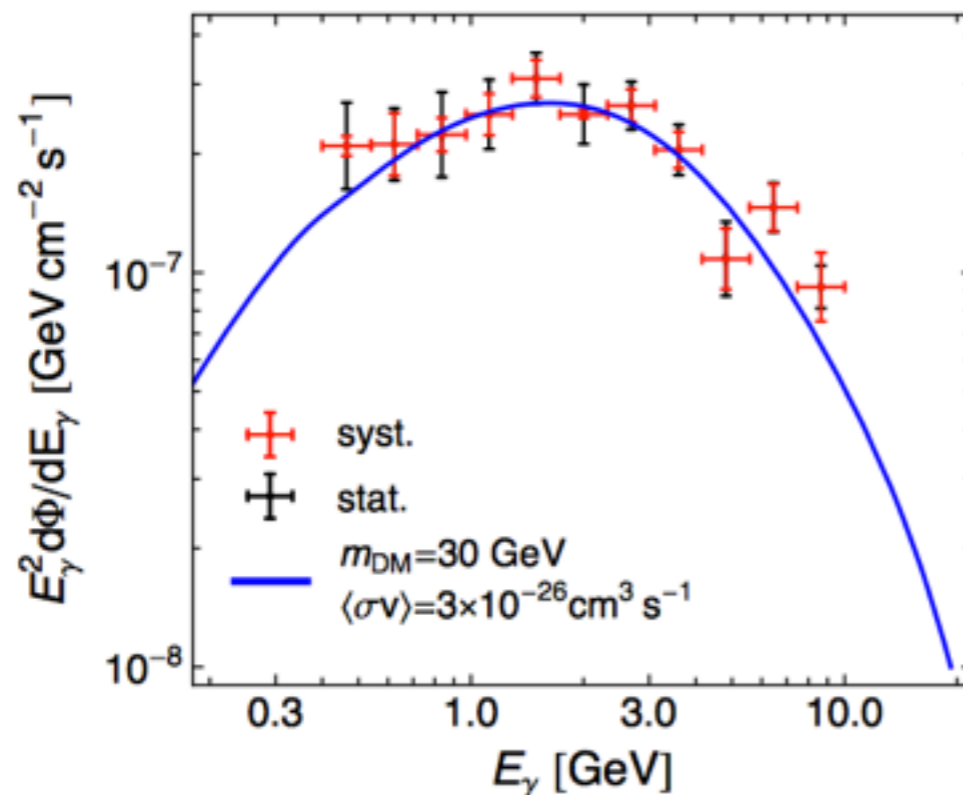
[Fox, Williams '12] [Buchmueller, Dolan, McCabe '13]



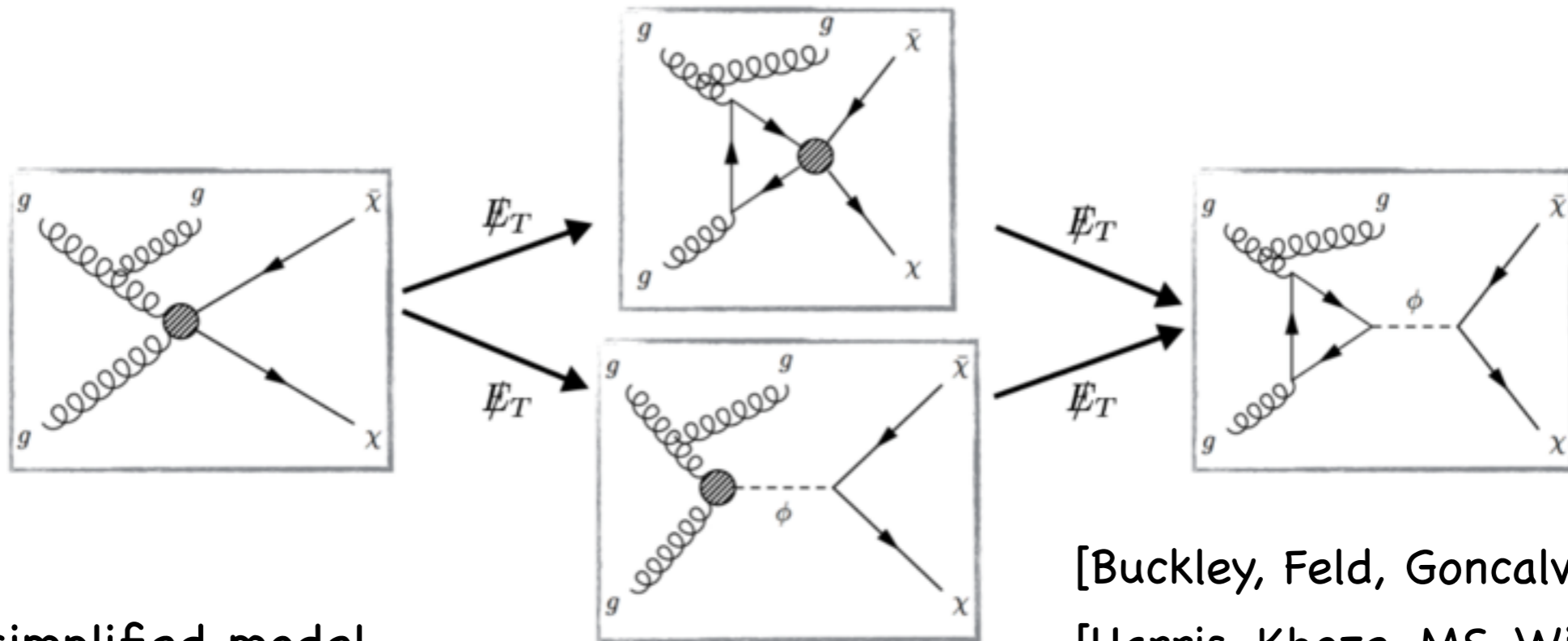
Dark matter could interact with SM via scalar mediator

- CP-even scalar, e.g. Higgs portal, or CP-odd scalar
- Dark Matter interacting via CP-odd scalar difficult to find
 - ➔ Direct detection interaction velocity suppressed
 - ➔ Difficult to produce at colliders, e.g. LEP
 - ➔ But might give visible signal in indirect detection
"Coy Dark Matter" and can fit GC excess

[Boehm, Dolan,
McCabe, MS,
Wallace '14]



Searching scalar DM-mediators in mono-jets



simplified model

[Buckley, Feld, Goncalves '14]

[Harris, Khoze, MS, Williams '14]

$$\mathcal{L}_{\text{pseudo-scalar}} \supset -\frac{1}{2}m_{\text{MED}}^2 P^2 - g_{\text{DM}} P \bar{\chi} \gamma^5 \chi - g_{SM}^t P \bar{t} \gamma^5 t - g_{SM}^b P \bar{b} \gamma^5 b$$

$$\mathcal{L}_{\text{scalar}} \supset -\frac{1}{2}m_{\text{MED}}^2 S^2 - g_{\text{DM}} S \bar{\chi} \chi - g_{SM}^t S \bar{t} t - g_{SM}^b S \bar{b} b$$

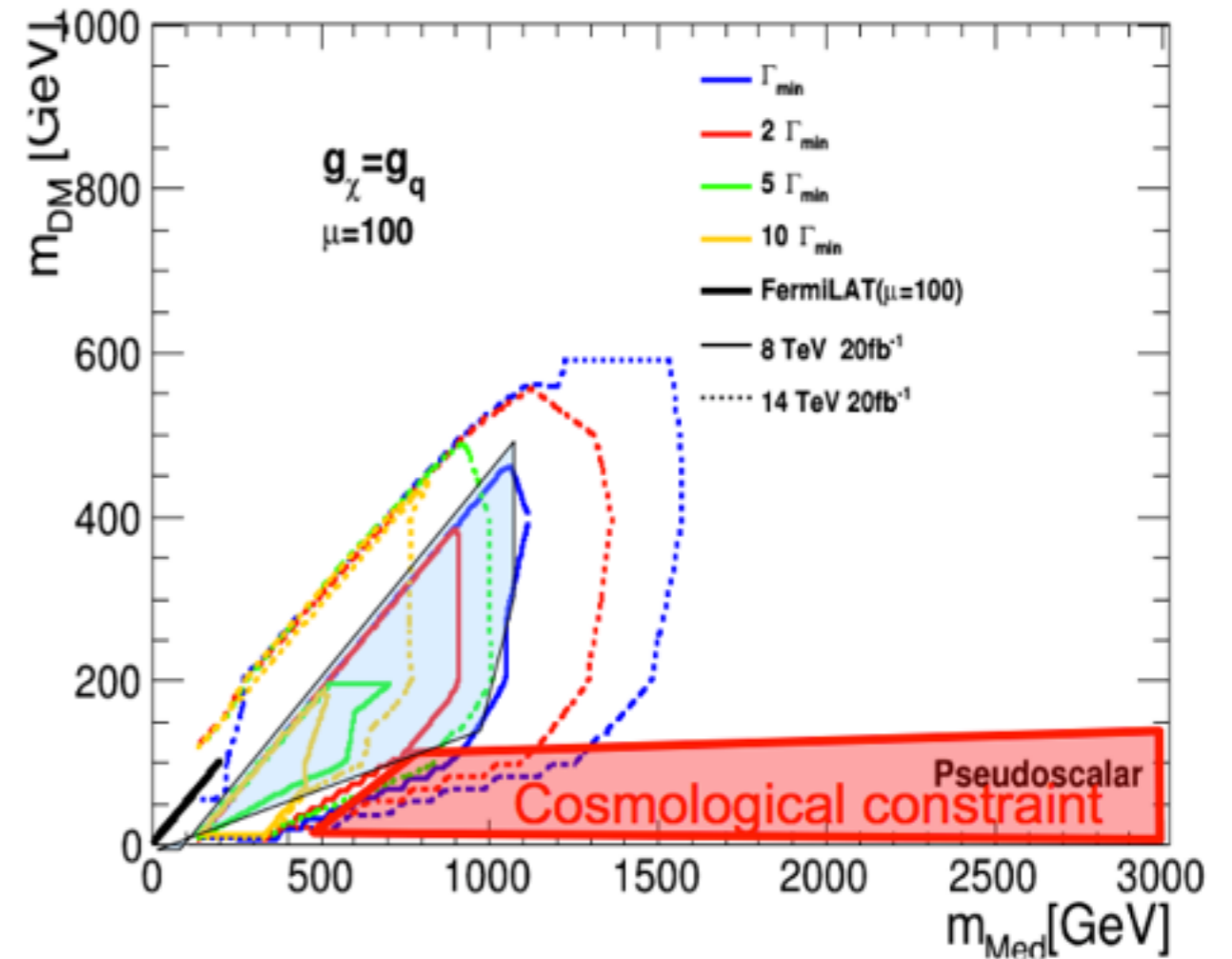
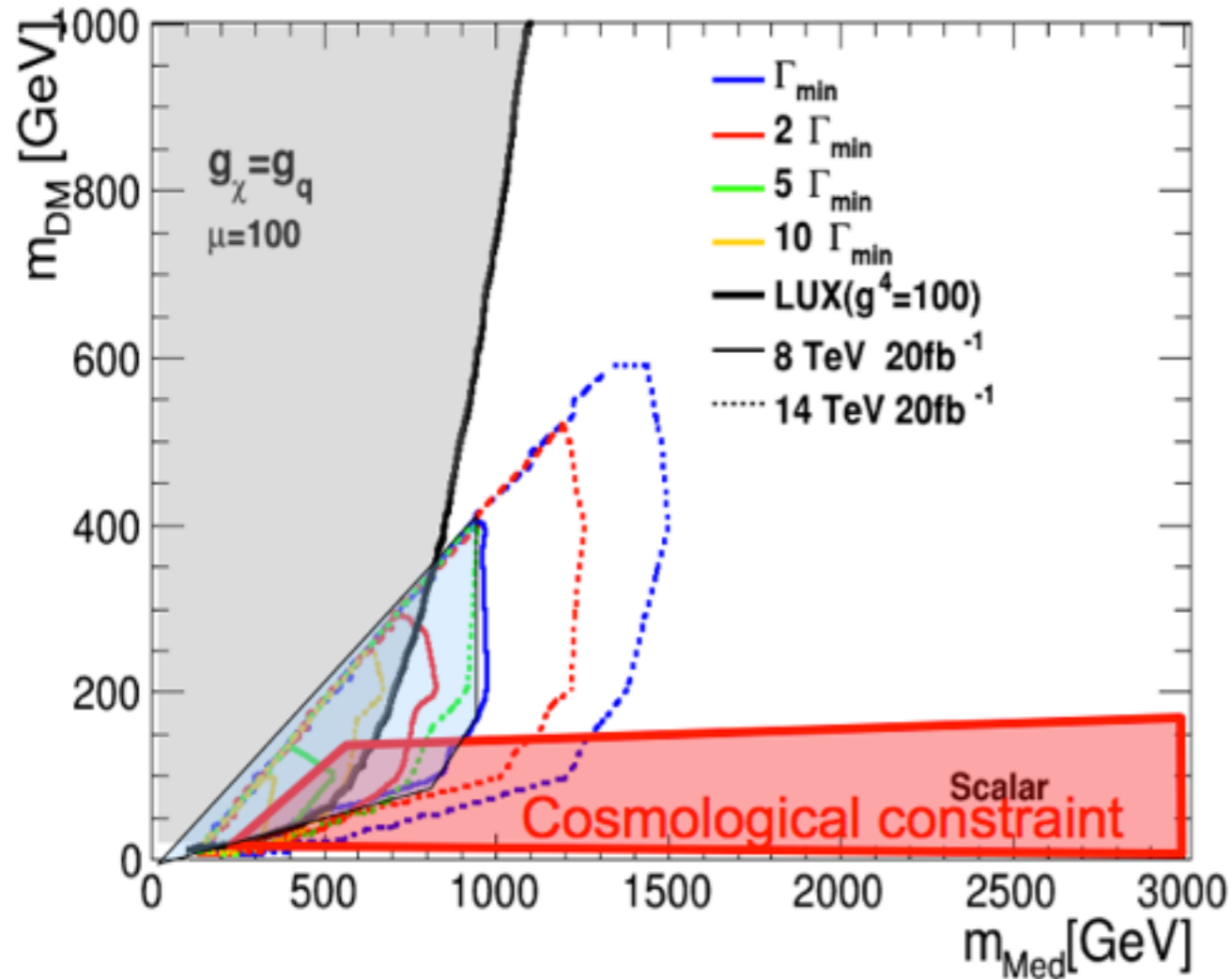
4 relevant parameters for phenomenology

1. mediator mass m_{MED}

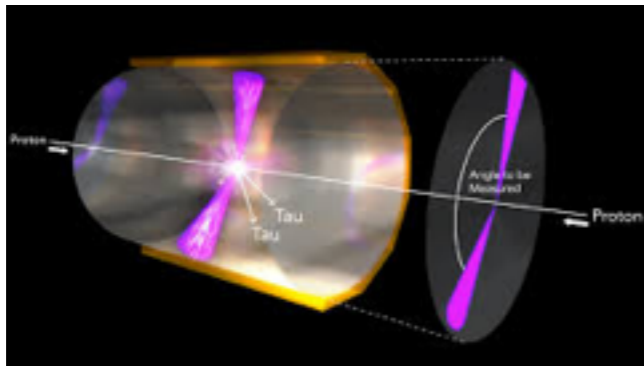
2. mediator width Γ_{MED}

3. dark matter mass m_{DM}

4. effective coupling parameter $g_q \cdot g_\chi$



- For light Dark Matter and heavy mediators the LHC can provide complementary information to DD and ID experiments
- A joint effort of all possible ways to look for (coy) Dark Matter is needed to maximize our chances to find it



Summary



- Scalar sector most interesting for coming years, i.e. strong crosstalk between different experiments
- Higgs and Dark Matter Phenomenology hot topics for upcoming runs
- IPPP has many staff members involved in this line of research
- For a successful program close collaboration between experimentalists and theorists essential