



τ - Physics and Lepton Flavour Violation

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Looking for LFV τ decays: LFV in the SUSY seesaw - a unique scenario

Plan: Theoretical prospects: relevant parameters and experimental implications

Learning from ν data and LFV ?

In collaboration with **M. J. Herrero**, E. Arganda and S. Antusch

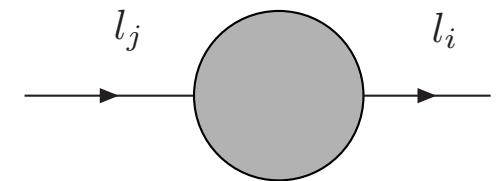
Detailed discussion in **PRD 73** (2006) 055003 (Arganda, Herrero)

and **JHEP 0611** (2006) 090 (Antusch, Arganda, Herrero, AMT)

Lepton flavour violation in charged lepton decays

Flavour violated in neutral leptons ($\nu_i \leftrightarrow \nu_j$ oscillations)

Extend SM to accommodate massive neutrinos: **SM+ ν_R**



What about **charged lepton flavour violation??**

LFV process	Present bound	Future sensitivity
$\text{BR}(\mu \rightarrow e \gamma)$	1.2×10^{-11}	1.3×10^{-13}
$\text{BR}(\tau \rightarrow e \gamma)$	1.1×10^{-7}	$\mathcal{O}(10^{-8})$
$\text{BR}(\tau \rightarrow \mu \gamma)$	6.8×10^{-8}	$\mathcal{O}(10^{-8})$
$\text{BR}(\mu \rightarrow 3e)$	1.0×10^{-12}	$\mathcal{O}(10^{-13} - 10^{-14})$
$\text{BR}(\tau \rightarrow 3e)$	2.0×10^{-7}	$\mathcal{O}(10^{-8})$
$\text{BR}(\tau \rightarrow 3\mu)$	1.9×10^{-7}	$\mathcal{O}(10^{-8})$
$\text{CR}(\mu \rightarrow e \text{ in Ti})$	4.3×10^{-12}	$\mathcal{O}(10^{-18})$

Huge experimental effort

Running & Upcoming:

BABAR, Belle, MEG

Future:

SuperKEKB (2011),

PRISM/PRIME (next decade),

Super Flavour factory (?)

SM+ ν_R : $l_j^- \rightarrow l_i^-$ transitions via $\nu_k - W^\pm$ loops are **strongly suppressed**

$$\text{BR}(l_j^- \rightarrow l_i^- \gamma)|_{\text{SM}_R} \approx (m_\nu/M_W)^2 \sim \mathcal{O}(10^{-54}) \Rightarrow \text{forever invisible..}$$

Measurement of LFV, e.g. $\text{BR}(l_j^- \rightarrow l_i^- \gamma) \Rightarrow$ **indication of New Physics!**

LFV and New Physics

- * **Seesaw mechanism** \rightsquigarrow If **Majorana** ν , a natural explanation for small m_ν
New scale - m_R ; additional **heavy singlet states** N [Leptogenesis?]

$$m_\nu \approx \frac{(Y_\nu v)^2}{m_R}$$

New parameters (masses + mixings) **unreachable (?)**
 $Y_\nu \sim \mathcal{O}(1)$ (model-dependent)

- * **Supersymmetry** \rightsquigarrow One of the most appealing **extensions of the SM!**
dark matter candidates, hierarchy problem, **testable in the near future**, ...

LFV: very sensitive to SUSY

Large Y_ν , via **SUSY** loops \Rightarrow sizable **slepton flavour mixing** ($l_i \not\leftrightarrow \tilde{l}_i, \tilde{\nu}_i$)
potentially **large LFV rates!**

If **no LFV** found \Rightarrow **Restrictions** on SUSY seesaw parameters

If **LFV measured** \Rightarrow **Important hints on SUSY seesaw parameters**

How does this interplay originate?

Minimal Supersymmetric Standard Model (MSSM) + 3 ν_R

Superfields	Boson Fields	Fermionic Partners	$SU(2)_L \times U(1)_Y \times SU(3)_c$
Gauge Multiplets			
\hat{G}	g	\tilde{g}	(1, 0, 8)
\hat{V}	W	\tilde{W}	(3, 0, 1)
\hat{B}	B	\tilde{B}	(1, 0, 1)
Matter Multiplets			
\hat{L}	$\tilde{L} = (\tilde{\nu}, \tilde{e}^-)_L^T$	$(\nu, e^-)_L^T$	(2, -1, 1)
\hat{E}^c	$\tilde{E} = \tilde{e}_R^+$	e_L^c	(1, +2, 1)
\hat{N}^c	$\tilde{N} = \tilde{\nu}_R$	ν_L	(1, 0, 1)
\hat{Q}	$\tilde{Q} = (\tilde{u}_L, \tilde{d}_L)^T$	$(u, d)_L^T$	(2, 1/3, 3)
\hat{U}^c	$\tilde{U} = \tilde{u}_R^*$	u_L^c	(1, -4/3, 3*)
\hat{D}^c	$\tilde{D} = \tilde{d}_R^*$	d_L^c	(1, 2/3, 3*)
\hat{H}_1	$(H_1^0, H_1^-)^T$	$(\tilde{H}_1^0, \tilde{H}_1^-)^T$	(2, -1, 1)
\hat{H}_2	$(H_2^+, H_2^0)^T$	$(\tilde{H}_2^+, \tilde{H}_2^0)^T$	(2, +1, 1)

Mixings: neutral Higgsinos, Binos, neutral Winos \rightarrow Neutralinos ($\tilde{\chi}^0$);
 charged Higgsinos, charged Winos \rightarrow Charginos ($\tilde{\chi}^\pm$);

$$W = Y_u \hat{H}_2 \hat{Q} \hat{U}^c + Y_d \hat{H}_1 \hat{Q} \hat{D}^c + Y_e \hat{H}_1 \hat{L} \hat{E}^c + \textcolor{red}{Y_\nu \hat{H}_2 \hat{L} \hat{N}^c + \frac{1}{2} \hat{N}^c m_N \hat{N}^c} - \mu H_1 H_2$$

And soft-breaking terms for scalars and gauginos! $\mathcal{L}^{\text{soft}} \{ M_{1/2}^i, (M_0^\phi)_{ij}, (A_0^\phi)_{ij}, B\mu \}$

SUSY Seesaw (Type I)

MSSM + 3 $\hat{\nu}_R$

$$\begin{pmatrix} \nu_{R_i} \\ \tilde{\nu}_{R_i} \end{pmatrix}$$

$$W^{\text{lepton}} = Y_e L e^c H_1 + Y_\nu L \nu^c H_2 + m_R \nu^c \nu^c$$

► Slepton terms in $\mathcal{L}^{\text{soft}}$

$$m_{\tilde{R}} = m_{\tilde{L}} = m_{\tilde{E}} = M_0 , A_\nu = A_l = A_0$$

► Enriched s-spectrum: 3 light and 3 heavy sneutrinos!

Assumptions: basis where $Y_e = Y_e^{\text{diag}}$ and $m_R = m_R^{\text{diag}}$

$$\mathbf{M}^\nu = \begin{pmatrix} 0 & m_D \\ m_D^T & m_R \end{pmatrix} \quad \begin{aligned} m_D &\rightarrow \text{Dirac mass matrix}; m_D = v_2 Y_\nu \ (\ll m_R) \\ m_R &\rightarrow \text{Heavy neutrino mass matrix - diag } (m_{N_i}) \end{aligned}$$

In the limit $m_R \gg v$, and to lowest order in $(m_D/m_R)^n$

Seesaw equation: $m_\nu = -m_D m_R^{-1} m_D^T$

$$\Rightarrow \text{find } m_D \text{ (i.e. } Y_\nu \text{)} \text{ compatible with } \left\{ \begin{array}{l} U_{\text{MNS}}^T m_\nu U_{\text{MNS}} = m_\nu^{\text{diag}} \\ U_{\text{MNS}}(\theta_{12}, \theta_{23}, \theta_{13}) \\ m_\nu^{\text{diag}}(\Delta m_{\text{sol}}^2, \Delta m_{\text{atm}}^2, \sum m_{\nu_i}) \end{array} \right.$$

SUSY Seesaw (Type I)

Solve Seesaw equation for m_D (Y_ν) \Rightarrow **Casas-Ibarra parameterisation ('01)**

Solution:

$$\textcolor{red}{m}_D = v_u \textcolor{red}{Y}_\nu = i \sqrt{m_N^{\text{diag}}} \textcolor{blue}{R} \sqrt{m_\nu^{\text{diag}}} U_{\text{MNS}}^\dagger \quad (\text{at } m_R)$$

$\textcolor{blue}{R}$ encodes possible mixings in the right-handed sector (in addition to U_{MNS})

$$\textcolor{blue}{R} = \begin{pmatrix} c_2 c_3 & -c_1 s_3 - s_1 s_2 c_3 & s_1 s_3 - c_1 s_2 c_3 \\ c_2 c_3 & c_1 c_3 - s_1 s_2 s_3 & -s_1 c_3 - c_1 s_2 s_3 \\ s_2 & s_1 c_2 & c_1 c_2 \end{pmatrix} \quad \begin{aligned} R^T R &= 1 \\ c_i &= \cos \theta_i ; s_i = \sin \theta_i \\ \theta_i &\text{ complex angles} \end{aligned}$$

Parameterise Y_ν while in agreement with low-energy ν -data

Y_ν strongly reflects the chosen values of θ_i , m_N and θ_{13} !

Notice that **additional constraints** on Y_ν can be obtained from

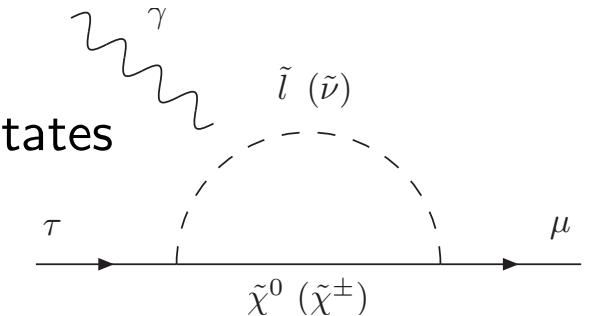
- **BAU (leptogenesis)** constraints on R -angles and on m_{N_1}
- Potential **contributions** to e , μ and τ **electric dipole moments**

Origin of large LFV in SUSY Seesaw

Focus on the **SUSY contributions** to $\tau \rightarrow \mu\gamma$

Like in SM, SUSY LFV relies on **misalignment** of eigenstates

\Rightarrow **How does** $\tau \not\leftrightarrow \tilde{\tau}, \tilde{\nu}_\tau; \mu \not\leftrightarrow \tilde{\mu}, \tilde{\nu}_\mu;$ **originate?**



LFV originates from large size and non-trivial structure of Y_ν !

Borzumati, Masiero (86); Hisano et al (96);...

- ▶ If Majorana ν s, Y_ν can be $\mathcal{O}(1)$
- ▶ To accommodate **data on ν -oscillations**, Y_ν cannot be **diagonal**!
- ▶ Even for **universal** soft-breaking terms @ M_{GUT} ,

RGE running of Y_ν from M_{GUT} down to **Seesaw scale m_R**

induces **flavour-violating** terms in slepton soft-breaking masses

$M_{LLij}^2 \neq 0, M_{LRij}^2 \neq 0$, etc \Rightarrow **slepton-lepton misalignment**

Flavour violation in the charged slepton sector!

Large LFV decay rates ($l_j \rightarrow l_i \gamma, l_j \rightarrow 3 l_i$) **within experimental reach!**

On the computation of LFV observables

All processes: **Full one-loop computation** (full-vertex, physical eigenstates)

All relevant one-loop SUSY diagrams included

- **Radiative decays $l_j \rightarrow l_i \gamma$:** $\text{BR}(l_j \rightarrow l_i \gamma) = \frac{e^2}{16\pi} \frac{m_{l_j}^5}{\Gamma_{l_j}} (|A_2^L|^2 + |A_2^R|^2)$

Loops: chargino-sneutrino and charged slepton-neutralino

- **Decays into 3 leptons $l_j \rightarrow 3l_i$:** photon-, Z -, Higgs-penguins; box diagrams

$$\text{Dominant photon-penguin: } \text{BR}(l_j \rightarrow 3l_i) = \frac{\alpha}{3\pi} \left(\log \frac{m_{l_j}^2}{m_{l_i}^2} - \frac{11}{4} \right) \times \text{BR}(l_j \rightarrow l_i \gamma)$$

Useful analytical insight: **Leading Log Approximation & MIA**

$$\begin{array}{ll} \cancel{l} & \left\{ \begin{array}{l} (\tilde{M}_{LL})_{ij}^2 \propto (3M_0^2 + A_0^2) (Y_\nu^* L Y_\nu^T)_{ij} \\ (\tilde{M}_{LR})_{ij}^2 \propto A_0 v_1 Y_i^e (Y_\nu^* L Y_\nu^T)_{ij} \\ (\tilde{M}_{RR})_{ij}^2 = 0; \quad L_i = \log(\frac{M_X}{m_{N_i}}) \end{array} \right. \\ \cancel{\tilde{l}} & \left. \begin{array}{l} \text{BRs} \left\{ \begin{array}{l} \tan^2 \beta \uparrow \\ |m_{N_3} \log m_{N_3}|^2 \downarrow \end{array} \right. \end{array} \right. \end{array}$$

$$\text{LLog \& MIA} \Rightarrow \text{BR}(\tau \rightarrow \mu \gamma) \approx \frac{\alpha^3 \tan^2 \beta}{G_F^2 m_{\text{SUSY}}^8} \left| \frac{1}{8\pi^2} (3M_0^2 + A_0^2) (Y_\nu^\dagger L Y_\nu)_{32} \right|^2$$

Not always good approximations...

Our approach

- Work in **MSSM + $3\hat{\nu}_R$** , universal conditions at $M_{\text{GUT}} \approx 2 \times 10^{16}$ GeV

CMSSM: SPS benchmark points $\text{SPS}_X(M_0, M_{1/2}, A_0, \tan \beta, \text{sign } \mu)$

- Compatibility with neutrino data (assume **hierarchical heavy neutrinos**)

$$m_{\nu_i} \left\{ \begin{array}{l} \Delta m_{\text{sol}}^2 = 8 \times 10^{-5} \text{ eV}^2 \\ \Delta m_{\text{atm}}^2 = 2 \times 10^{-3} \text{ eV}^2 \\ \text{Hierarchical } \mathbf{m}_{\nu_i} \end{array} \right. \quad U_{\text{MNS}} \left\{ \begin{array}{l} \theta_{12} = 30^\circ; \theta_{23} = 45^\circ \\ \delta = \phi_1 = \phi_2 = 0 \\ \text{Impact of } 0 \lesssim \theta_{13} \lesssim 10^\circ \end{array} \right.$$

- Require “viable” BAU: $n_B/n_\gamma \in [10^{-10}, 10^{-9}]$ (WMAP $\sim 6.1 \times 10^{-10}$, '06)
- Compatibility $\text{EDM}_{e\mu\tau} \lesssim (6.9 \times 10^{-28}, 3.7 \times 10^{-19}, 4.5 \times 10^{-17}) \text{ e.cm}$
- Theoretical **predictions for LFV decay rates**:

SPheno 2.2.2: Numerical integration of two-loop RGEs from M_X to M_Z

Additionally: Neutrino RGEs, LFV decays (full 1-loop), Leptogenesis, EDMs

- Compare with present and future LFV bounds!

The Constrained MSSM

Assume (well-motivated) relations between soft-breaking parameters

universality of soft gaugino masses: $M_i = M_{1/2}$

universality of soft breaking scalar masses: $(M_0^{\tilde{\phi}})_{ij} = M_0$

universality of soft breaking trilinear couplings: $(A_0^{\tilde{\phi}})_{ij} = A_0$

at a common scale, gauge-coupling unification scale, $M_X \approx 10^{16}$ GeV

CMSSM (mSUGRA-like): $M_{1/2}, M_0, A_0, \tan\beta, \text{sign}(\mu)$

Predictions for the low-energy spectrum: solve RGEs !

Further simplify: consider benchmark points (Snowmass Points and Slopes):

SPS 1a, 1b	typical mSUGRA points, bulk cosmological region, med-large $\tan\beta$
SPS 2	light gauginos, focus point region for relic density
SPS 3	small $\tilde{l} - \tilde{\chi}^0$ mass difference, coannihilation region
SPS 4, 5	extreme $\tan\beta$: high and low

Similar analysis for any other (reconstructed) SUSY model

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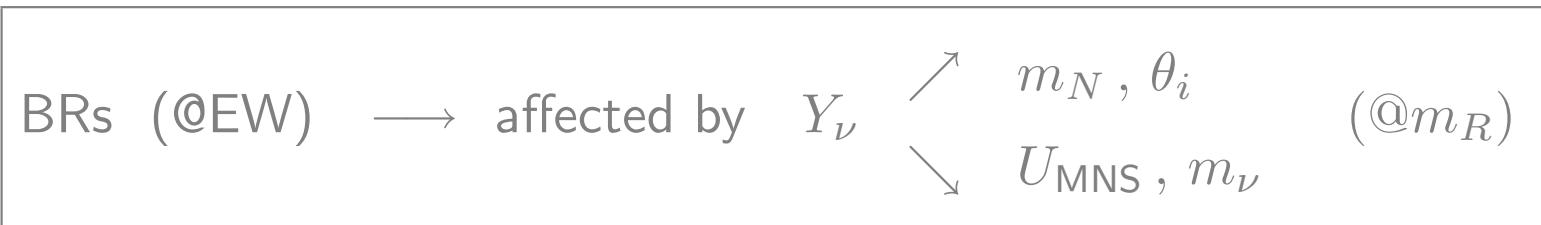
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- Compare with present and future LFV bounds!

Predictions for LFV within the type I SUSY Seesaw

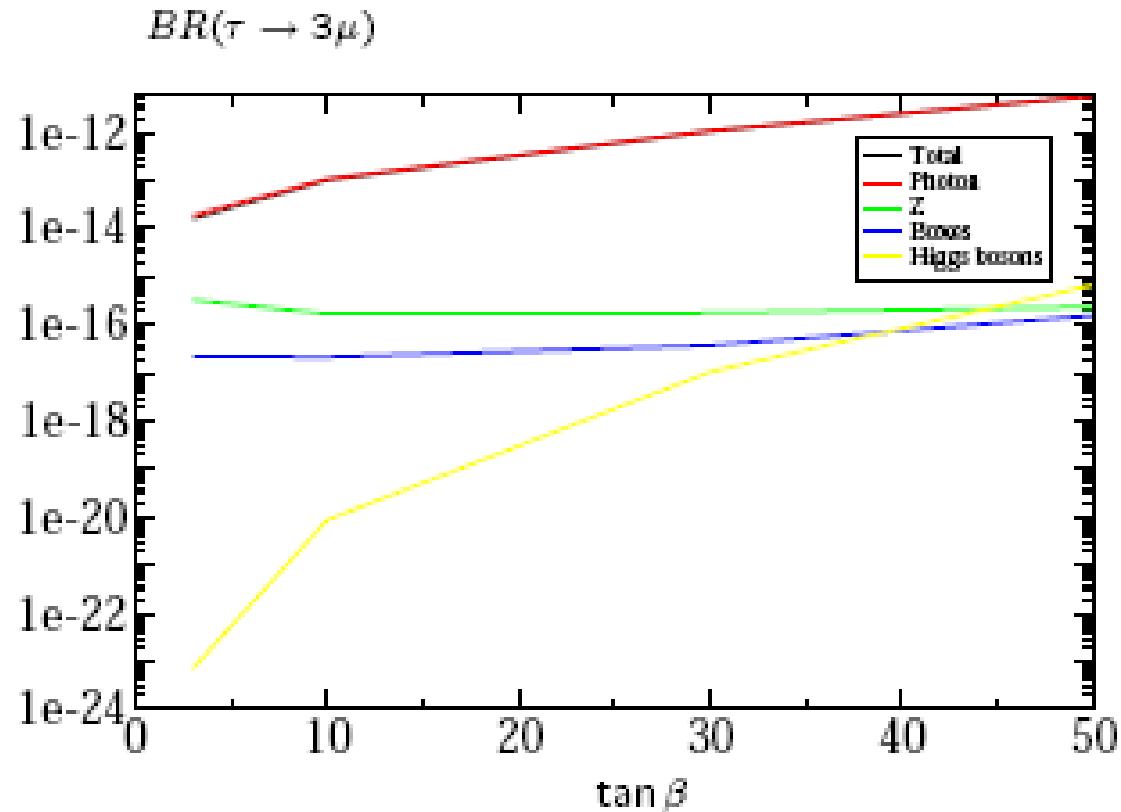
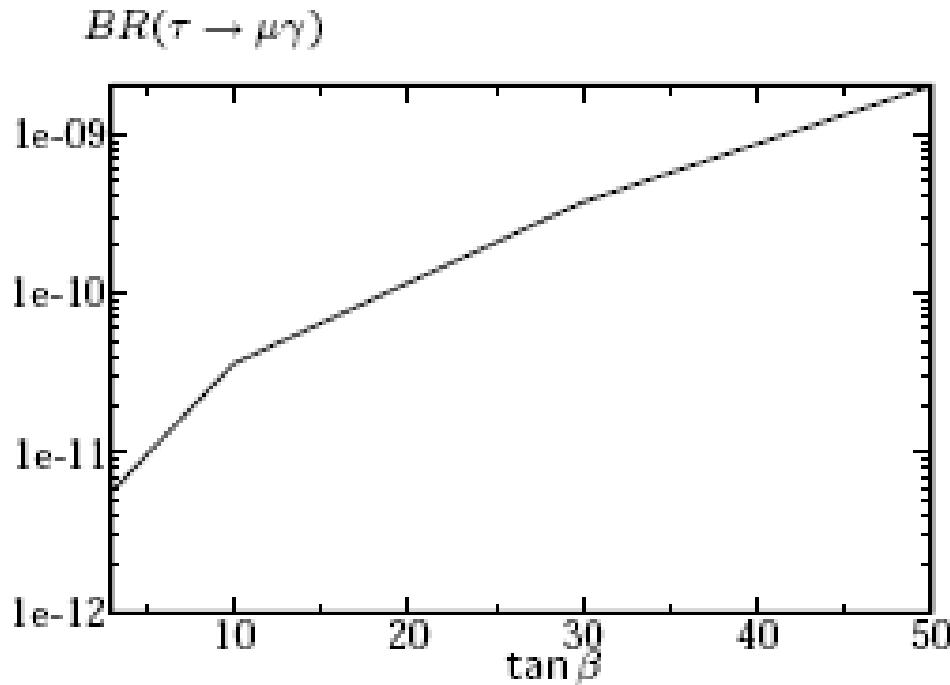
- Dependence on CMSSM parameters
 - Dependence on the size of the Yukawa couplings
 - Impact of a measurement of θ_{13}



CMSSM parameters: $\tan \beta$

$M_0 = 400 \text{ GeV}$, $M_{1/2} = 300 \text{ GeV}$, $A_0 = 0$, $m_N = 10^{14} \text{ GeV}$ (degen.), $\theta_i = 0$

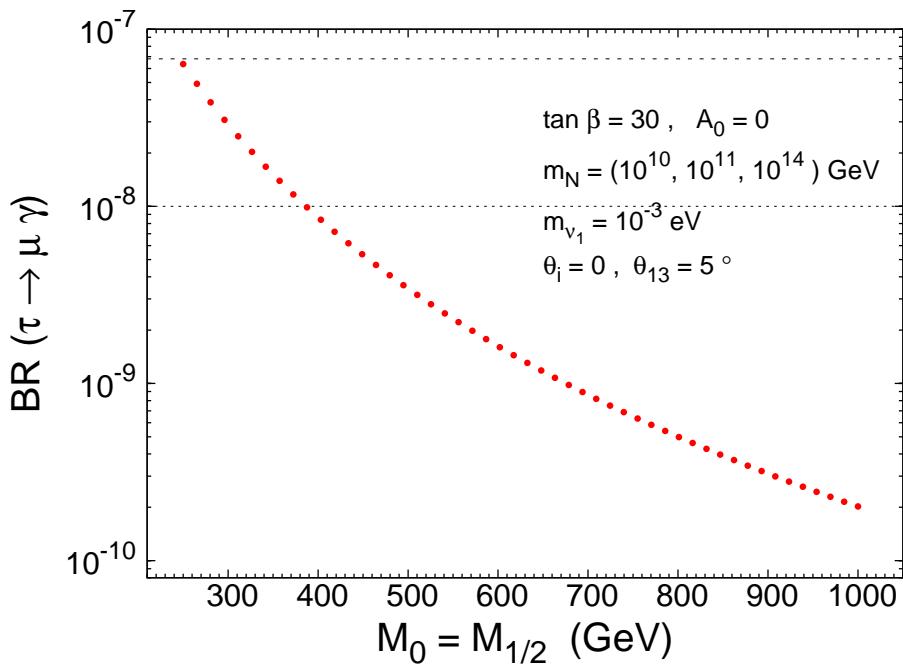
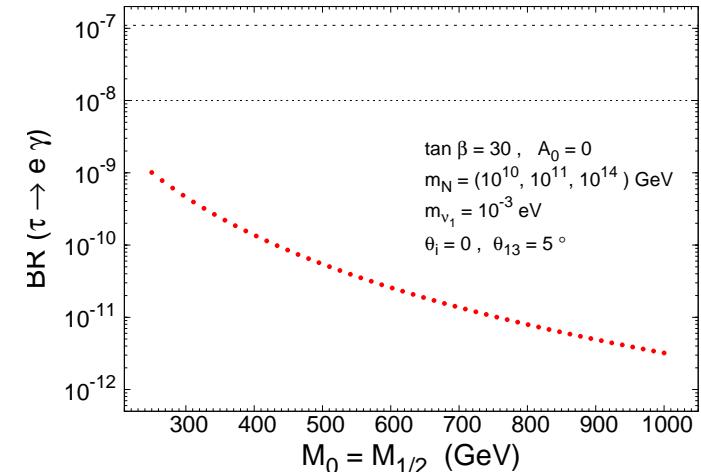
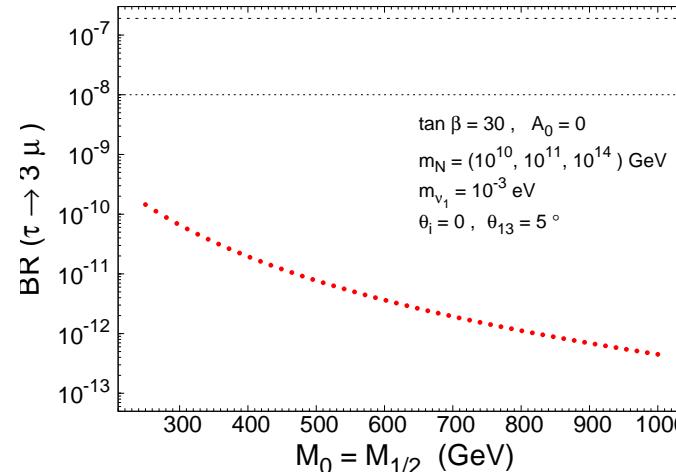
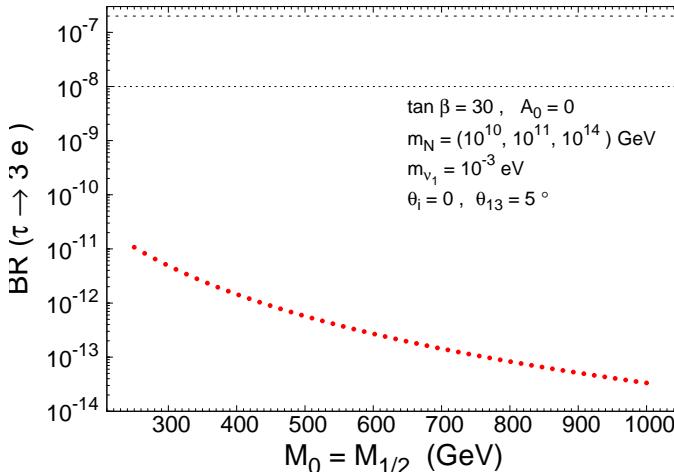
From Arganda and Herrero, PRD 73 (2006)



- ★ BRs dependence on $\tan \beta$: $\text{BR}(\tau \rightarrow 3l_i)$, $\text{BR}(\tau \rightarrow l_j \gamma)$ grow as $\tan^2 \beta$
- ★ $\text{BR}(\tau \rightarrow \mu\mu\mu)/\text{BR}(\tau \rightarrow \mu\gamma) \approx 1/440$
 \Rightarrow reflects **strong γ -penguin dominance** (all $\tan \beta$)
- ★ Degenerate N_i : typically **lower BRs** than **hierarchical N_i**

CMSSM parameters: M_0 and $M_{1/2}$

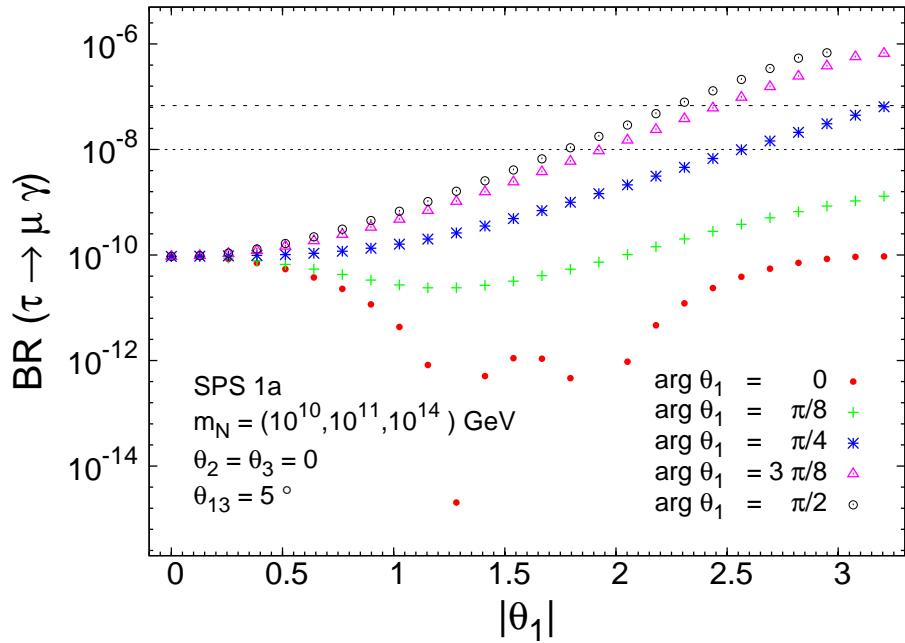
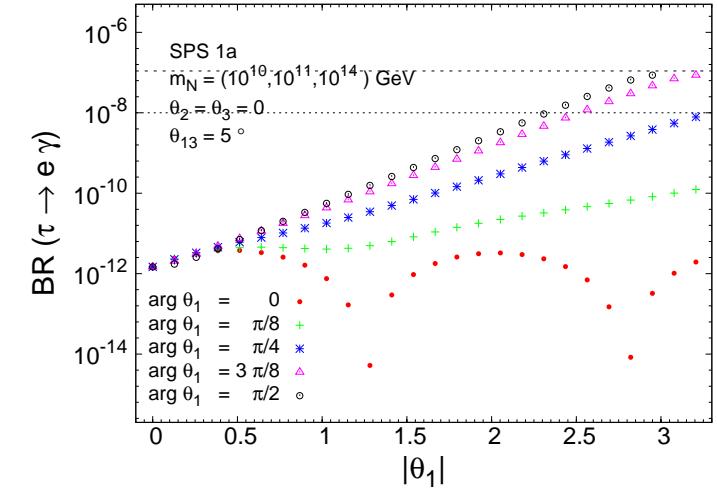
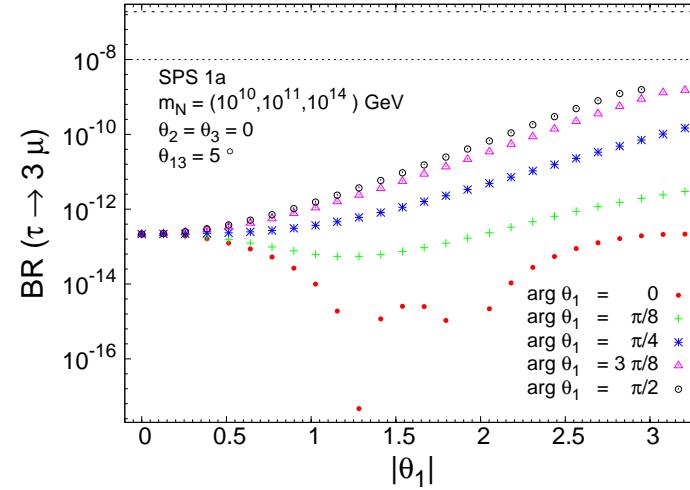
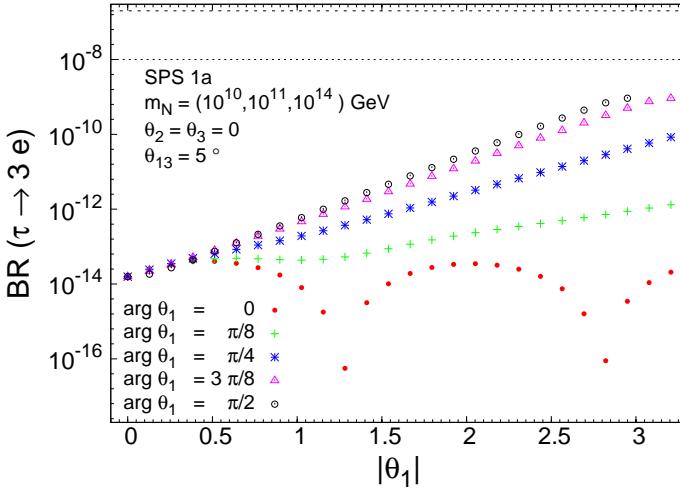
$M_0 = M_{1/2}$, $\tan \beta = 30$, $A_0 = 0$, $m_N = (10^{10}, 10^{11}, 10^{14})$ GeV, $\theta_{13} = 5^\circ$, $\theta_i = 0$



- ★ Important dependence on M_0 and $M_{1/2}$
- ★ Larger BRs obtained for low M_0 ($M_{1/2}$)
- ★ Comparative analysis:
 $\tau \rightarrow \mu\gamma$ most challenging channel
 $\tan \beta = 30$, $\text{BR}(\tau \rightarrow \mu\gamma)$ within exp. reach!
 \Rightarrow Disfavouring M_0 ($M_{1/2}$) $\lesssim 300$ GeV

Large Yukawa couplings: role of θ_1

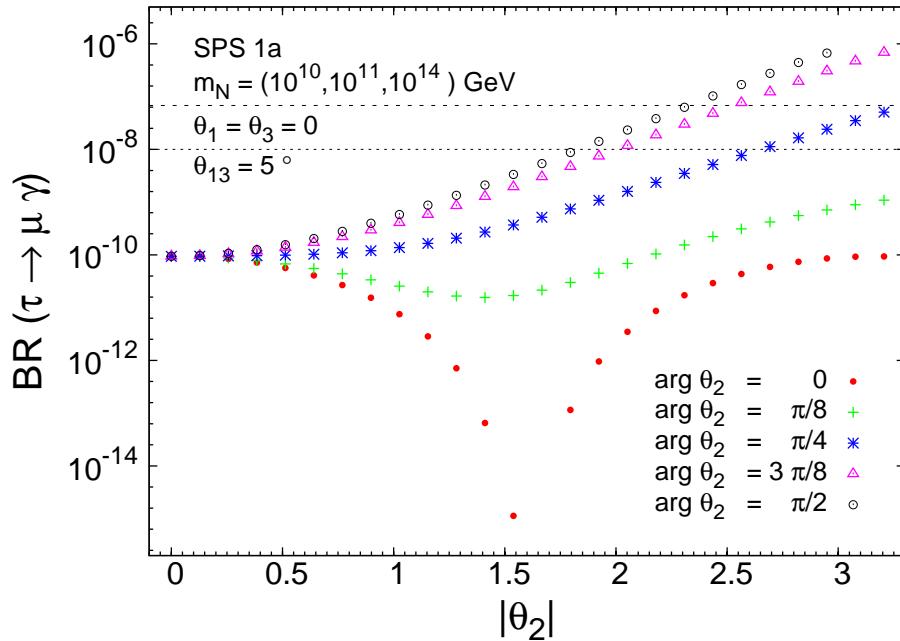
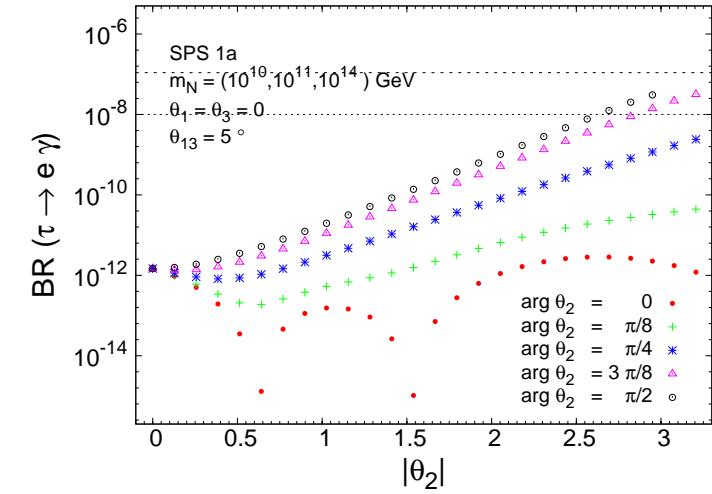
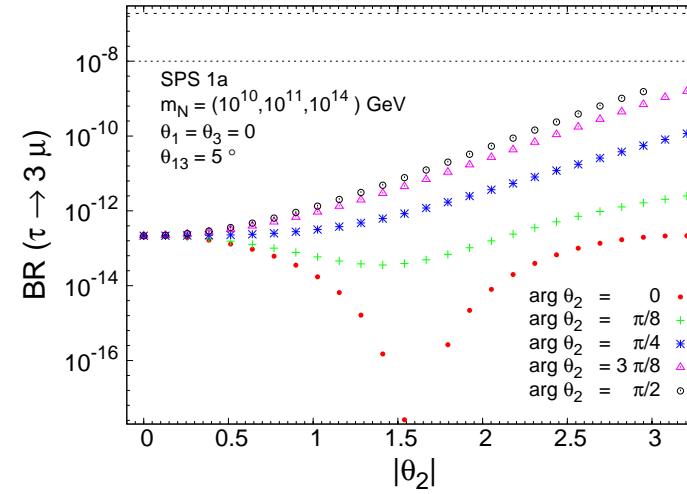
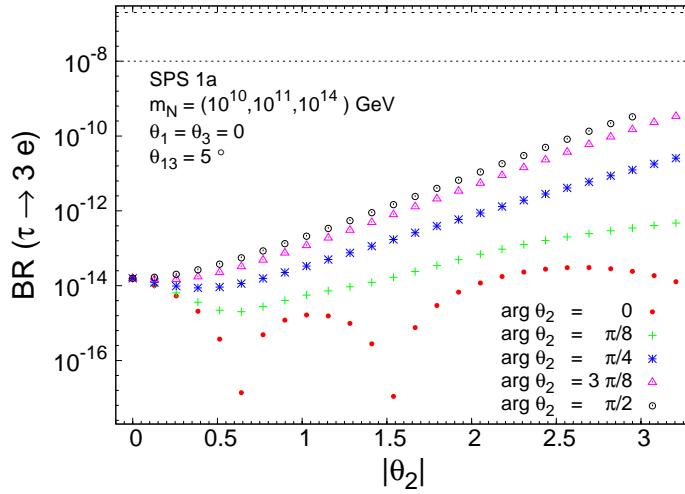
SPS 1a ($\tan \beta = 10$), $m_N = (10^{10}, 10^{11}, 10^{14})$ GeV, $\theta_{13} = 5^\circ$, $\theta_2 = \theta_3 = 0$



- ★ **R-matrix complex angles strongly affect Y_ν**
- arg $\theta_1 = 0, \pi/8, \pi/4, 3\pi/8, \pi/2$
- ★ **BRs closely follow pattern of $Y_\nu(\theta_1)$**
- ★ **Large, complex θ_1 can induce $\mathcal{O}(10^4)$ enhancement in BRs**
- ⇒ For **SPS 1a**, $\text{BR}(\tau \rightarrow \mu\gamma)$ already constrains values of θ_1

Large Yukawa couplings: role of θ_2

SPS 1a ($\tan\beta = 10$), $m_N = (10^{10}, 10^{11}, 10^{14})$ GeV, $\theta_{13} = 5^\circ$, $\theta_1 = \theta_3 = 0$



★ Similar role played by θ_2

$\arg \theta_2 = 0, \pi/8, \pi/4, 3\pi/8, \pi/2$

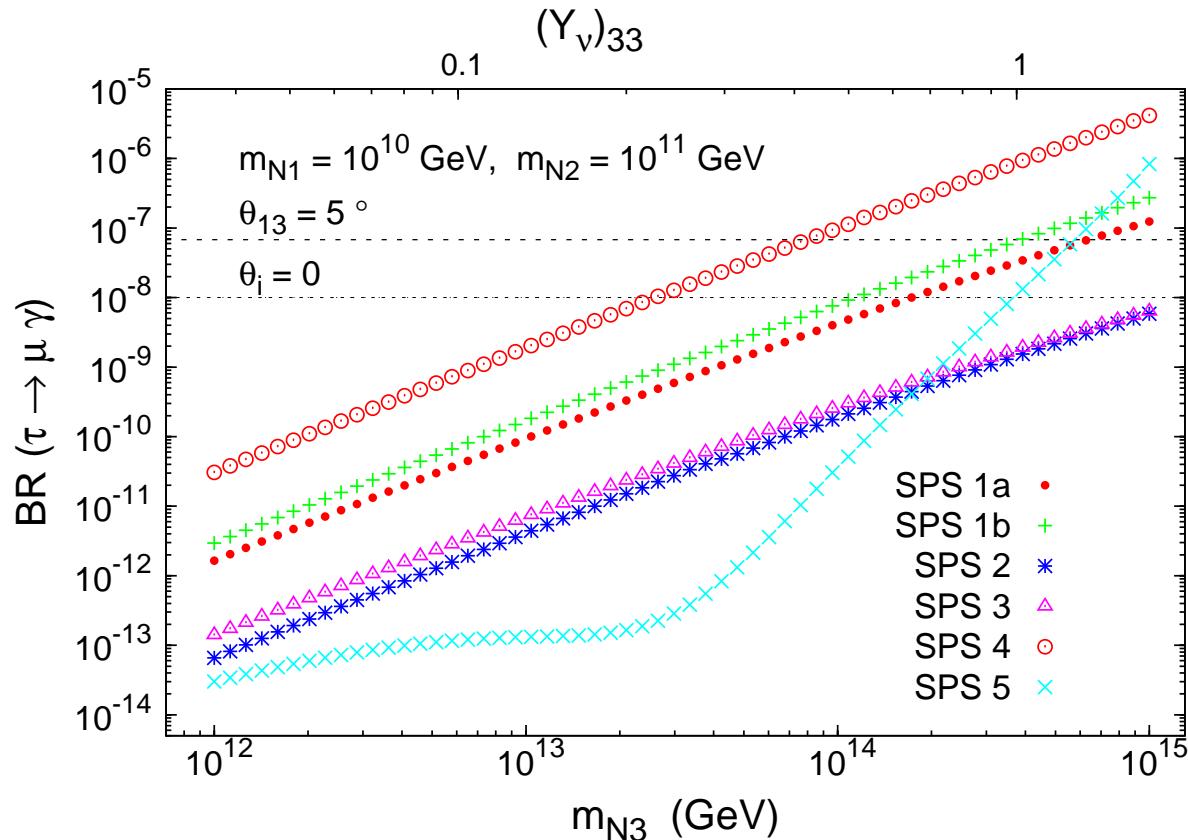
★ $\text{BR}(\tau \rightarrow \mu\gamma)$ also **constrains** values of θ_2

★ **Baryon Asymmetry of the Universe:**

Leptogenesis requires $\theta_1, \theta_2 (\arg \theta_{1,2}) \neq 0 \dots$

θ_3 does not play a relevant role for BRs

Large Yukawa couplings: m_{N_3}



★ SUSY seesaw scenario \Rightarrow **indirect bounds** via BRs:

$$\begin{cases} m_{N_3} \lesssim 10^{14} \text{ GeV (SPS 4)} \\ m_{N_3} \lesssim 5 \times 10^{14} \text{ GeV (1a,1b)} \end{cases}$$

► **SPS 1a, SPS 1b, SPS 2, SPS 3, SPS 4** - Good agreement with LLog

$$\text{BR} \propto \tan^2 \beta: \text{BR}_4 > \text{BR}_{1b} \gtrsim \text{BR}_{1a} > \text{BR}_3 \gtrsim \text{BR}_2 > \text{BR}_5 \quad \text{BR} \propto |m_{N_3} \log m_{N_3}|^2$$

► **SPS 5**: LLog not good approximation!

★ m_{N_3} drives Y_ν :

hierarchical N_i , $m_D \propto \sqrt{m_N^{\text{diag}}}$

strongly affects $\text{BR}(\tau \rightarrow \mu \gamma)$
and all other BRs!

★ Without predictive framework

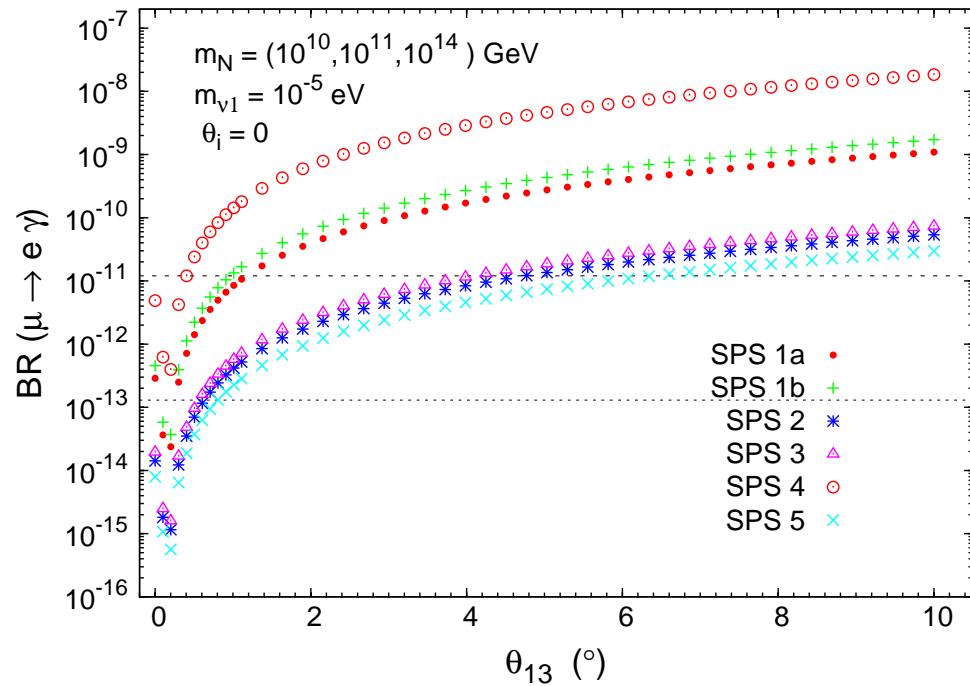
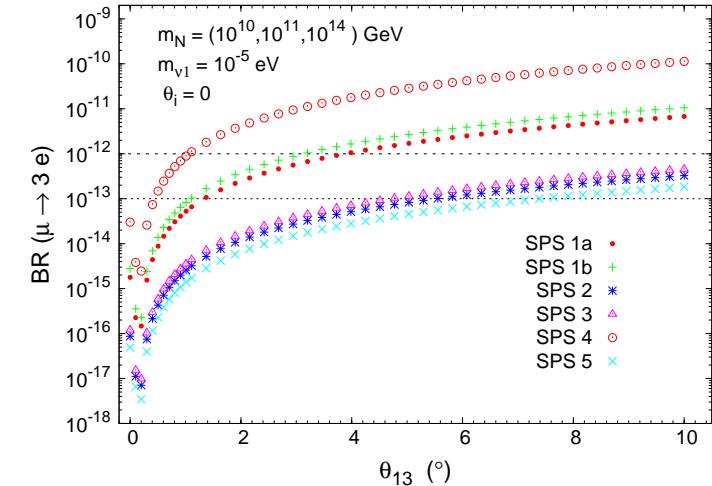
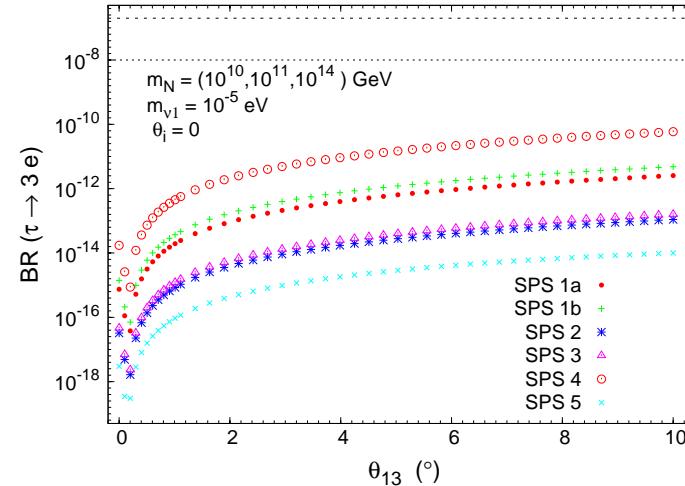
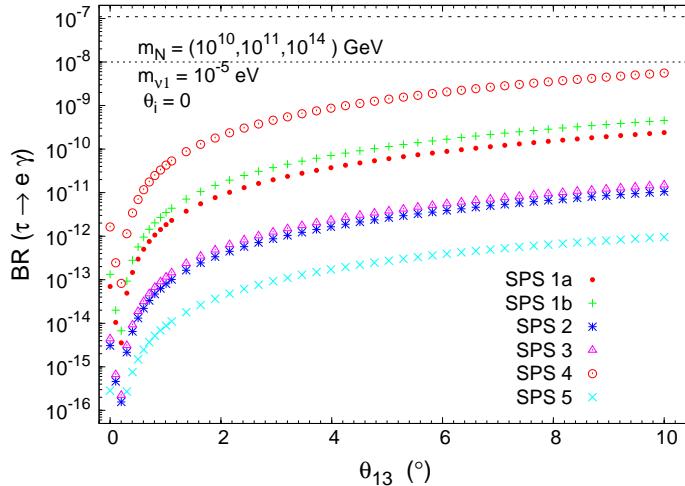
$$m_{N_3} \in [10^{10}, 10^{15}] \text{ GeV}$$

(indirect lower bound - Leptogenesis)

Useful hints on m_{N_3} from potential BR and θ_{13} measurement!!

Sensitivity to θ_{13} : τ and μ decays

SPS X, $m_N = (10^{10}, 10^{11}, 10^{14})$ GeV, $\theta_i = 0$

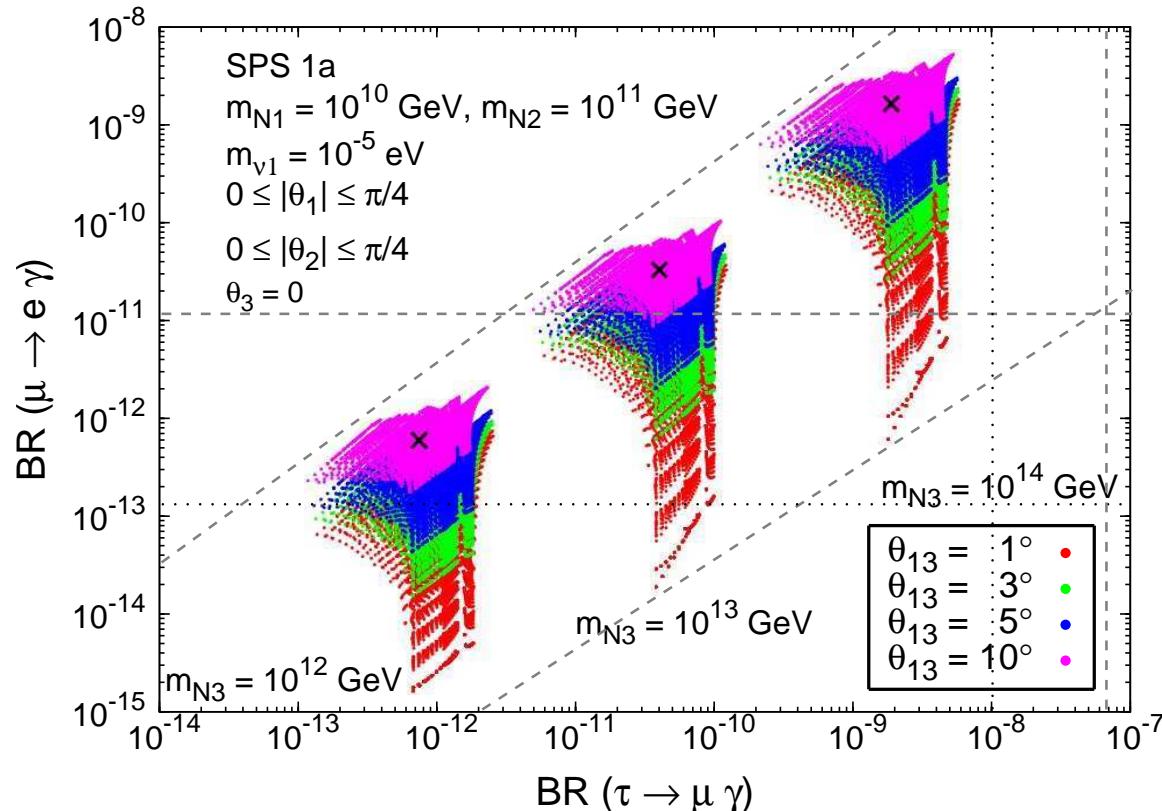


- * $\mu, \tau \rightarrow e\gamma$ and $\mu, \tau \rightarrow 3e$ exhibit large sensitivity to θ_{13} (for small θ_i)
- * $\tau \rightarrow \mu\gamma$ and $\tau \rightarrow 3\mu$: no θ_{13} dependence
- * Measurement of θ_{13} is crucial!
- Maximal impact for $\mu \rightarrow e\gamma$
- Experimentally: SUSY & Seesaw scenarios, $\theta_{13} \gtrsim 2^{\circ}$ within MEG reach

Impact of θ_{13} on LFV observables: hints on m_{N_3}

Experimental **measurement of θ_{13}** in a near future! SUSY @ LHC??

Correlations of $\text{BR}(\mu \rightarrow e \gamma)$ and $\text{BR}(\tau \rightarrow \mu \gamma)$ \Rightarrow Optimises impact of θ_{13}^{exp}



- ★ Other m_{N_3} within “corridor”
- ★ Conclude on impact of θ_{13} measurement on allowed/excluded m_{N_3}
- ★ $\text{BR}(\mu \rightarrow e \gamma)|_{\text{06}} + \theta_{13} = 1^\circ$
 $\Rightarrow m_{N_3} < 10^{14} \text{ GeV}$
- ★ $\text{BR}(\mu \rightarrow e \gamma)|_{\text{MEG}} + \theta_{13} = 1^\circ$
 $\Rightarrow m_{N_3} < 3 \times 10^{12} \text{ GeV}$

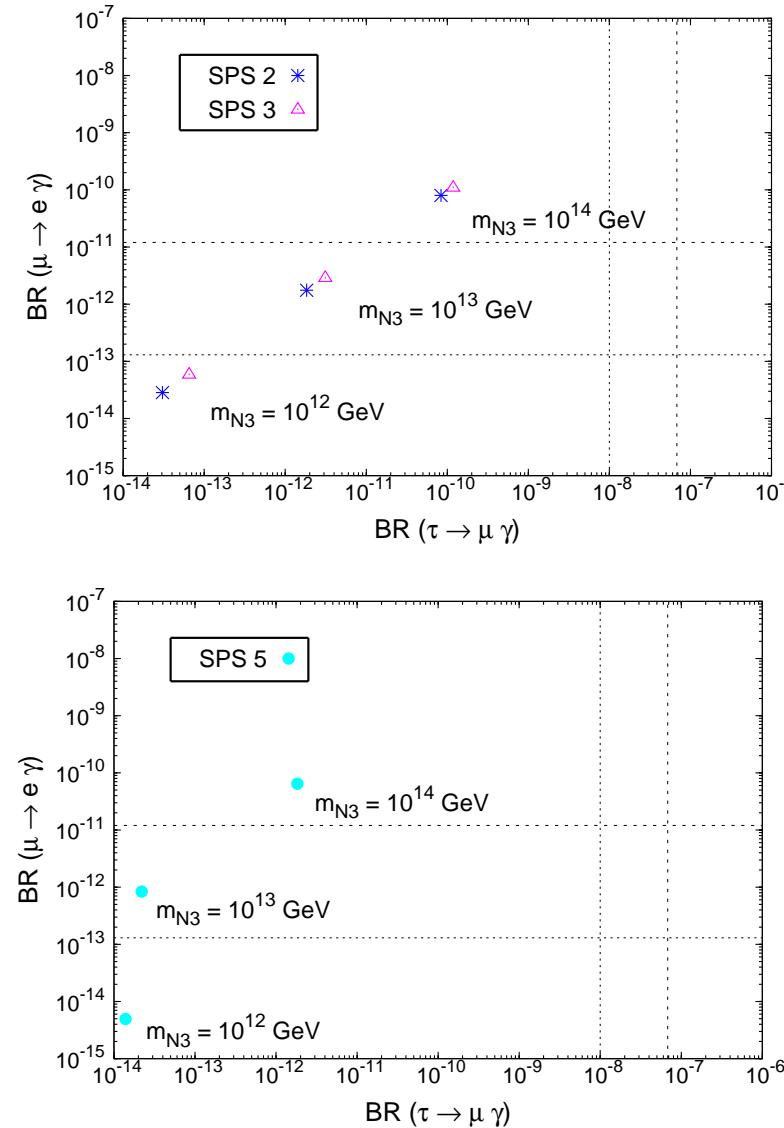
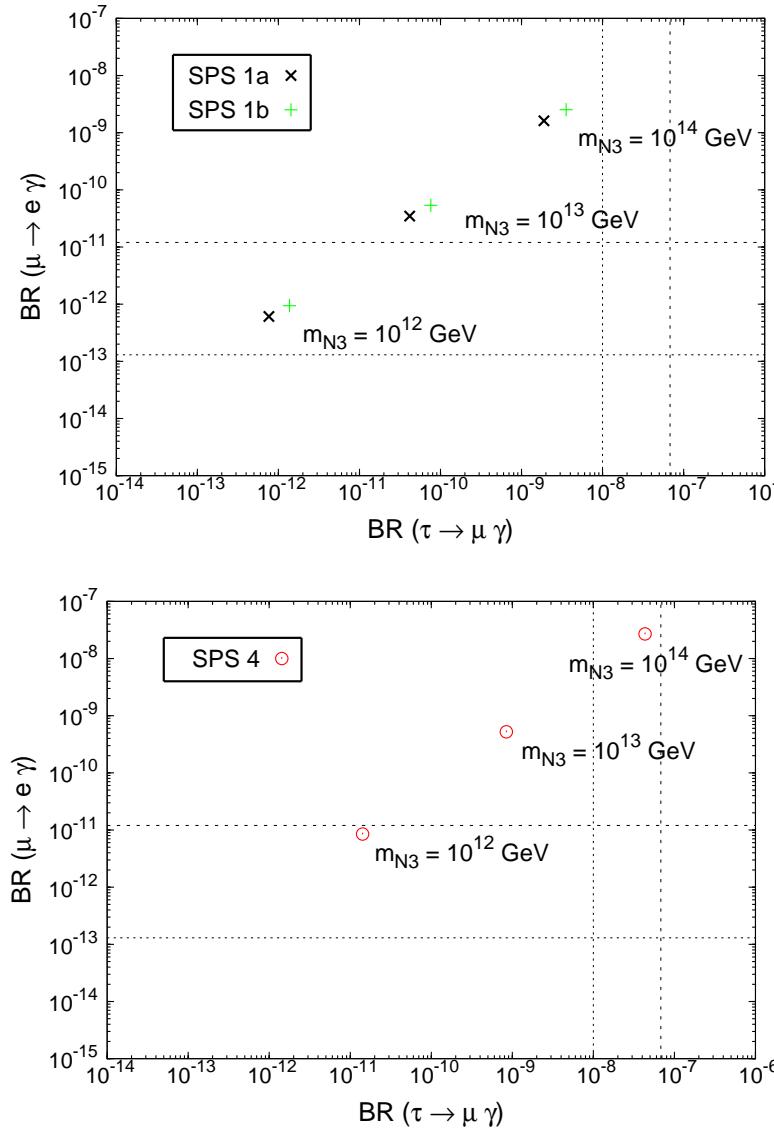
Discovery of SUSY (LHC, Tevatron)

Measurement of θ_{13}

Observation of LFV τ, μ decays

} \Rightarrow Window to higher scales!
Insight into Seesaw parameters

Impact of θ_{13} on LFV observables: hints on m_{N_3}



★ Similar conclusions for other **SPS** points: **SPS 4 most stringent**

$$\theta_{13} \approx 10^\circ \Rightarrow m_{N_3} \lesssim 2 \times 10^{12} \text{ GeV}$$

► Present bounds from $\tau \rightarrow \mu \gamma$, $\mu \rightarrow e \gamma$ already exclude $m_{N_3} \gtrsim 10^{14} \text{ GeV}$

Final remarks

► Framework:

Lepton flavour violation may provide first **evidence** of New Physics!

LFV in SUSY Seesaw: obtain large $\text{BR}(l_j \rightarrow l_i \gamma)$ and $\text{BR}(l_j \rightarrow 3 l_i)$

Consistency with low-energy ν -data, **EDMs** (and BAU via thermal leptogenesis)

Studied **BR dependence** on θ_{13} and other key parameters (m_{N_i} , θ_i , $\tan \beta$)

► Results:

$\tau \rightarrow \mu \gamma$, $\mu \rightarrow e \gamma$ and $\mu \rightarrow 3e$ within present experimental reach

for SPS 4 and $m_{N_3} = 10^{14}$ GeV

For $R \approx 1$, important **impact** of θ_{13} (factor 10^5 for $\theta_{13} \in [0, 10^\circ]$)

Sensitivity manifest in $\mu \rightarrow e \gamma$, $\mu \rightarrow 3e$, $\tau \rightarrow e \gamma$ and $\tau \rightarrow 3e$

Possible **interplay** of future **experimental measurements** in

hinting towards **SUSY** and **Seesaw** parameters via **LFV!**

If SUSY is present at low scale ($\lesssim 1$ TeV)

If ν masses are originated by a type-I Seesaw mechanism

**LFV τ and μ decays offer a unique laboratory
that together with low-energy ν data can provide
insight into the heavy neutrino sector,
and into Seesaw parameters**

Additional transparencies

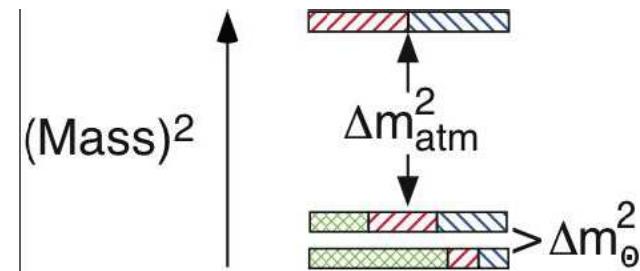
Implications of ν -oscillations

- Neutrinos are massive (very light); large mixing in the lepton sector
- ★ Experimentally (June '06): three mass eigenstates (possibly $m_{\nu_1} = 0$)

$$\Delta m_{\text{sol}}^2 = m_{\nu_2}^2 - m_{\nu_1}^2 \approx 8 \times 10^{-5} \text{ eV}^2$$

$$\Delta m_{\text{atm}}^2 = m_{\nu_3}^2 - m_{\nu_1}^2 \approx 2 \times 10^{-3} \text{ eV}^2$$

$$\sum m_{\nu_i} \lesssim 0.7 \text{ eV (Cosmological)}$$



also Tritium β -decays and $0\nu2\beta$ decays

- ★ And three mixing angles (and CP violating phases??): U_{MNS}

$$\left. \begin{array}{l} \sin^2 \theta_{12} \approx 0.3 \\ \sin^2 \theta_{23} \approx 0.5 \\ \sin^2 \theta_{13} \leq 0.04 \end{array} \right\} \begin{pmatrix} c_{12} c_{13} & s_{12} c_{13} & s_{13} e^{-i\delta} \\ -s_{12} c_{23} - c_{12} s_{23} s_{13} e^{i\delta} & c_{12} c_{23} - s_{12} s_{23} s_{13} e^{i\delta} & s_{23} c_{13} \\ s_{12} s_{23} - c_{12} c_{23} s_{13} e^{i\delta} & -c_{12} s_{23} - s_{12} c_{23} s_{13} e^{i\delta} & c_{23} c_{13} \end{pmatrix} \cdot P(\phi_{1,2})$$

- Extend SM! Add right-handed neutrinos ν_R ! Dirac or Majorana ν ??

Neutrino masses

- ▶ **Dirac ν :** Only “standard” Dirac neutrino mass term: $Y_\nu H \bar{\nu}_L \nu_R$
 EWBS: $m_\nu = Y_\nu v$; $m_\nu \ll v \Rightarrow Y_\nu \sim \mathcal{O}(10^{-11})$ —> **Naturalness** problem

- ▶ **Majorana $\nu \longleftrightarrow$ Seesaw mechanism**

Type-I Seesaw: [Minkowski; Gell-Mann, Ramond, Slansky; Yanagida; Glashow;
 Mohapatra, Senjanović]

New scale - m_R ; additional **heavy singlet states N** $\langle m_N \rangle \sim \mathcal{O}(m_R)$

$$\mathcal{L}_R^{\text{Maj}} \sim m_R \bar{\nu}_R^c \nu_R$$

$$m_\nu \approx \frac{(Y_\nu v)^2}{m_R}$$

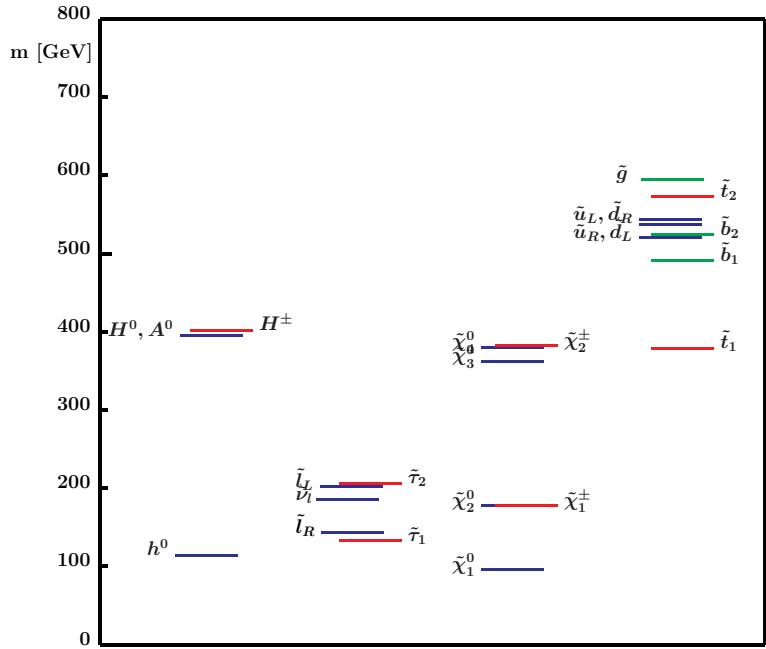
 $(Y_\nu \sim \mathcal{O}(1), \text{ model-dependent})$

New **parameters** (masses + complex mixings) - Experimentally **unreachable**
Opens the possibility of generating **Baryogenesis via Leptogenesis**

Type-II Seesaw: [Barbieri et al; Marshak, Mohapatra; Cheng, Li; Magg, Wetterich;
 Schechter, Valle; Mohapatra, Senjanović]

$$\mathcal{L}_L^{\text{Maj}} \sim Y_\nu^\Delta \Delta_H \bar{l}_L^c l_L \quad \leftrightarrow \quad \text{triplet Seesaw} \quad \text{We will not pursue this here..}$$

SPS points



SPS 1a

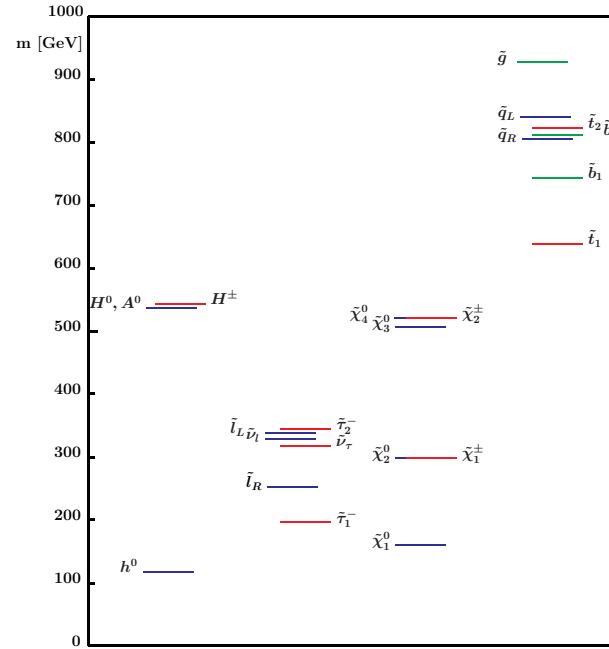
$$M_0 = 100 \text{ GeV}$$

$$M_{1/2} = 250 \text{ GeV}$$

$$A_0 = -100 \text{ GeV}$$

$$\tan \beta = 10$$

$$\mu > 0$$



SPS 1b

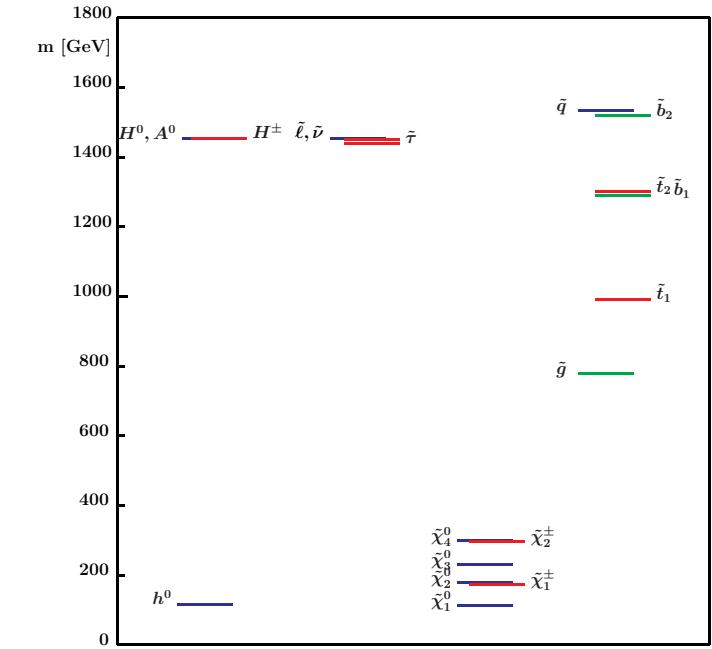
$$M_0 = 200 \text{ GeV}$$

$$M_{1/2} = 400 \text{ GeV}$$

$$A_0 = 0 \text{ GeV}$$

$$\tan \beta = 30$$

$$\mu > 0$$



SPS 2

$$M_0 = 1450 \text{ GeV}$$

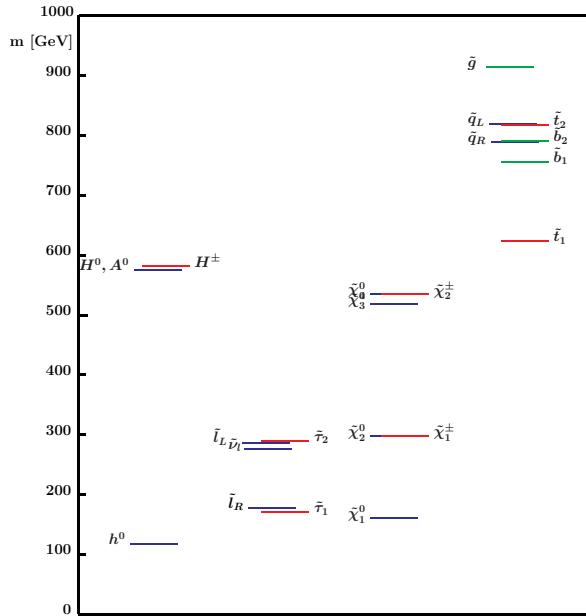
$$M_{1/2} = 300 \text{ GeV}$$

$$A_0 = 0 \text{ GeV}$$

$$\tan \beta = 10$$

$$\mu > 0$$

SPS points



SPS 3

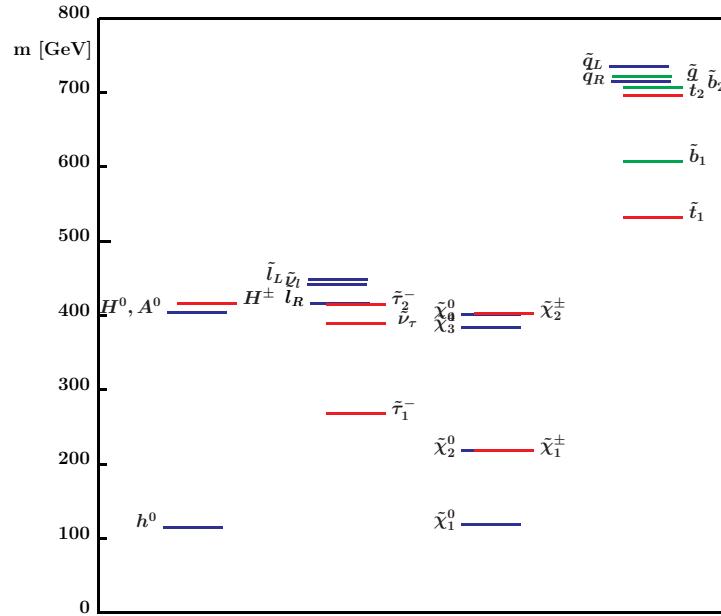
$$M_0 = 90 \text{ GeV}$$

$$M_{1/2} = 300 \text{ GeV}$$

$$A_0 = 0 \text{ GeV}$$

$$\tan \beta = 10$$

$$\mu > 0$$



SPS 4

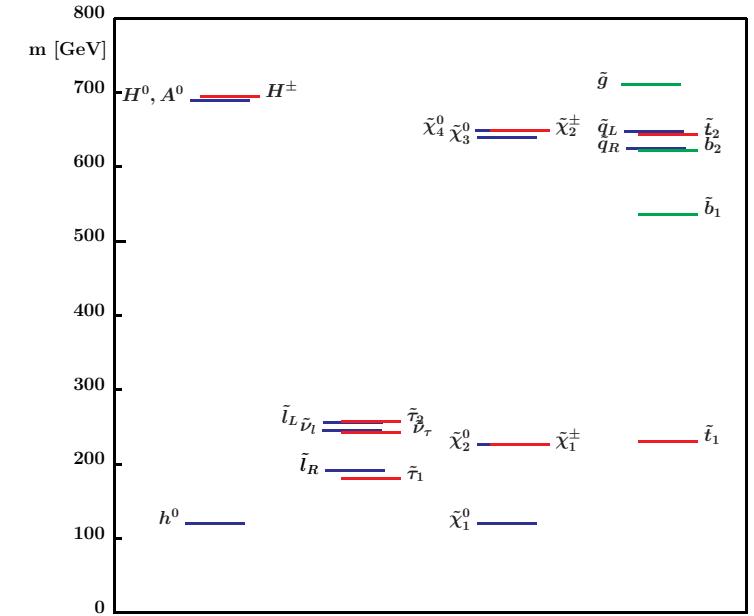
$$M_0 = 400 \text{ GeV}$$

$$M_{1/2} = 300 \text{ GeV}$$

$$A_0 = 0 \text{ GeV}$$

$$\tan \beta = 50$$

$$\mu > 0$$



SPS 5

$$M_0 = 150 \text{ GeV}$$

$$M_{1/2} = 300 \text{ GeV}$$

$$A_0 = -1000 \text{ GeV}$$

$$\tan \beta = 5$$

$$\mu > 0$$

BAU via Thermal Leptogenesis

Seesaw mechanism: $\left\{ \begin{array}{l} \text{smallness of neutrino masses} \\ \text{Baryogenesis via Leptogenesis} \end{array} \right.$

Lepton asymmetry dynamically generated via **out-of-equilibrium**,

CP violating decay of lightest RH neutrino N_1 [Fukujita,Yanagida ('86)]

- ▶ Correct treatment: **flavour matters!** (CP violation from U_{MNS} and R !)
[Barbieri et al, Abada et al, ...]
- ▶ **Here:** “one-flavour” *approximation*... (to be improved! CP only from R)

Lepton asymmetry: $Y_{\text{B-L}}^{\text{MSSM}} \approx -\eta \left[\frac{1}{2}(\varepsilon_1 + \tilde{\varepsilon}_1) Y_{N_1}^{\text{eq}} + \frac{1}{2}(\varepsilon_{\tilde{1}} + \tilde{\varepsilon}_{\tilde{1}}) Y_{\tilde{N}_1}^{\text{eq}} \right]$
 \Downarrow efficiency factor \nearrow \searrow decay asymmetry $\Rightarrow (R \in \text{ })$

Baryon asymmetry: $Y_B = \alpha(N_H) Y_{\text{B-L}}$

(partial sphaleron conversion [Kuzmin,Rubakov,Shaposhnikov ('85)])

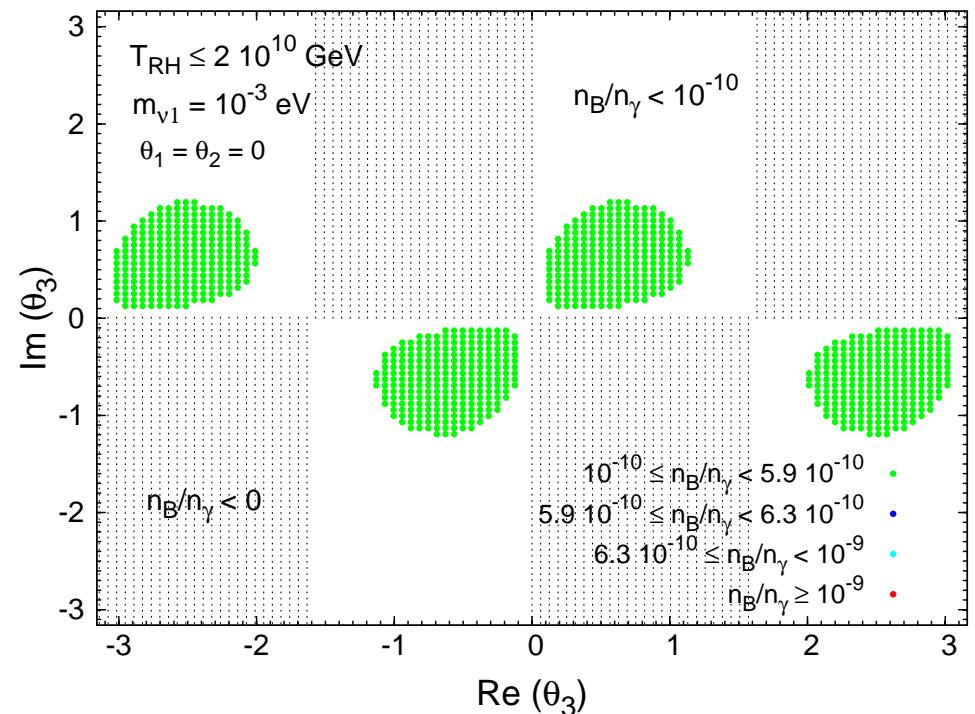
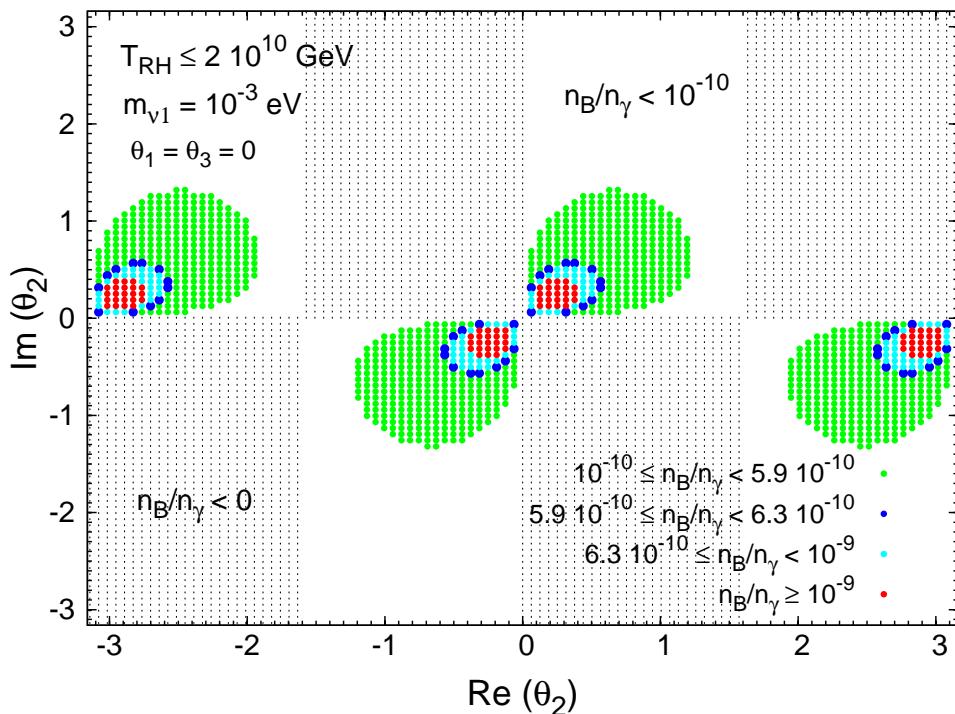
BAU in agreement with WMAP ('06): $n_B/n\gamma \approx (6.10 \pm 0.21) \times 10^{-10}$

First constraints on Seesaw parameters

► BAU constraints on *R*-angles:

(no U_{MNS} - one-flavour approx.)

$\left. \begin{array}{l} |\theta_2|, |\theta_3| : \text{small, but non-zero} \\ |\theta_1| : \text{generically not constrained} \end{array} \right\}$



Notice: different favoured regions in **flavour dependent treatment!** [Antusch, AT ('06)]

First constraints on Seesaw parameters

- **BAU constraints on R -angles:** $\left\{ \begin{array}{l} |\theta_2|, |\theta_3| : \text{small, but non-zero} \\ |\theta_1| : \text{generically not constrained} \end{array} \right.$
(no U_{MNS} - one-flavour approx.)

- Additional **BAU constraints** on m_{N_1} : $m_{N_1} \sim \mathcal{O}(T_{\text{RH}}) \sim 10^{10} \text{ GeV}$

Dramatic efficiency loss for $m_{N_1} \gg T_{\text{RH}}$ [Giudice et al ('03)]

BBN gravitino problem: constraints on T_{RH} depending on $m_{3/2}$

Gravitino decay and non-thermal LSP production: $T_{\text{RH}} \lesssim 2 \times 10^{10} \text{ GeV}$
 $(m_{\text{LSP}} \sim 100 \text{ GeV})$

- Contributions to **EDMs**:

Even for CP-conserving U_{MNS} , CP violating phases in R , thus **complex** Y_ν

Potential **contributions** to e , μ and τ electric dipole moments

Ensure **compatibility** with current **bounds**:

$$\text{EDM}_{e\mu\tau} \lesssim (6.9 \times 10^{-28}, 3.7 \times 10^{-19}, 4.5 \times 10^{-17}) \text{ e.cm}$$