

FUTURE FLAVOUR PHYSICS, Abingdon, June 21-22, 2007

NEW PHYSICS AND FLAVOUR

Antonio Masiero
Univ. of Padua
And
INFN, Padua

A FUTURE FOR FLAVOR PHYSICS IN OUR SEARCH BEYOND THE SM?

- The traditional **competition** between direct and indirect (FCNC, CPV) searches to establish who is going **to see the new physics first** is no longer the priority, rather
- **COMPLEMENTARITY** between direct and indirect searches for New Physics is the key-word
- Twofold meaning of such complementarity:
 - i) **synergy in “reconstructing” the “fundamental theory”** staying behind the signatures of NP;
 - ii) **coverage of complementary areas of the NP parameter space** (ex.: multi-TeV SUSY physics)

WHY TO GO BEYOND THE SM

“OBSERVATIONAL” REASONS

•HIGH ENERGY PHYSICS

NO (but $A_{FB}^{Z \rightarrow bb}$)

•FCNC, $CP \neq$

NO (but $b \rightarrow s\bar{q}q$ penguin, V_{ub} ...)

•HIGH PRECISION LOW-EN.

NO (but $(g-2)_\mu$...)

•NEUTRINO PHYSICS

YES $m_\nu \neq 0, \theta_\nu \neq 0$

•COSMO - PARTICLE PHYSICS

YES (DM, ΔB_{cosm} , INFLAT., DE)

THEORETICAL REASONS

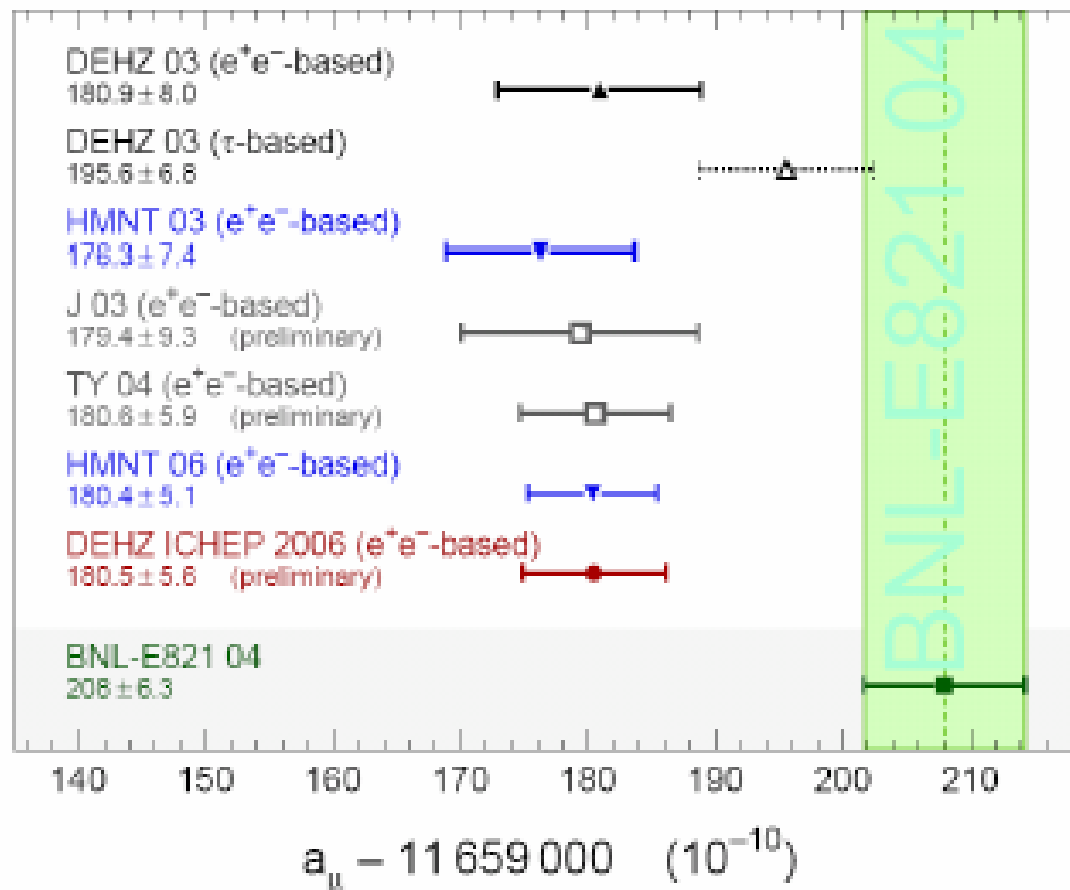
•INTRINSIC INCONSISTENCY OF SM AS QFT

NO (spont. broken gauge theory without anomalies)

•NO ANSWER TO QUESTIONS THAT “WE” CONSIDER “FUNDAMENTAL” QUESTIONS TO BE ANSWERED BY A “FUNDAMENTAL” THEORY

YES (hierarchy, unification, flavor)

Status of $g_{\mu}-2$



Whereas τ based prediction agrees with the measurement within 1σ
all recent $e+e-$ based predictions have a deviation with data at over 3σ

Present “Observational” Evidence for New Physics

- NEUTRINO MASSES 
- DARK MATTER 
- MATTER-ANTIMATTER ASYMMETRY 
- INFLATION 

THE FATE OF LEPTON NUMBER

L VIOLATED

L CONSERVED

ν Majorana ferm.

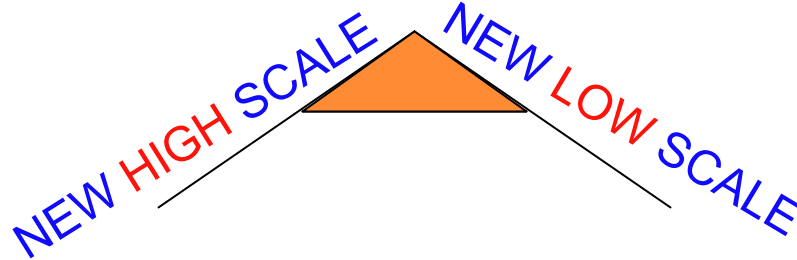
ν Dirac ferm.
(dull option)

SMALLNESS of m_ν

$$h \bar{\nu}_L H \nu_R \rightarrow m_\nu = h \langle H \rangle \quad M_\nu < 1 \text{ eV} \rightarrow h < 10^{-6}$$

EXTRA-DIM. ν_R in the bulk: small overlap?

PRESENCE OF A NEW PHYSICAL MASS SCALE



SEE - SAW MECHAN.

MAJORON MODELS

Minkowski; Yanagida; Gell-Mann, Ramond, Slansky; Mohapatra-Senjanovic

Gelmini, Roncadelli

ν_R ENLARGEMENT OF THE FERMIONIC SPECTRUM

Δ ENLARGEMENT OF THE HIGGS SCALAR SECTOR

$$M \nu_R \nu_R + h \bar{\nu}_L \phi \nu_R$$

$$h \bar{\nu}_L \nu_L \Delta$$

$$m_\nu = h \langle \Delta \rangle$$

$$\begin{matrix} \nu_L & \sim O & \nu_R \\ \nu_R & h \langle \Phi \rangle & M \end{matrix}$$

LR Models?

N.B.: EXCLUDED BY LEP!

STABLE ELW. SCALE WIMPs from PARTICLE PHYSICS

	SUSY (x^μ, θ)	EXTRA DIM. (x^μ, j^i)	LITTLE HIGGS. SM part + new part
1) ENLARGEMENT OF THE SM	Anticomm. Coord.	New bosonic Coord.	to cancel Λ^2 at 1-Loop
2) SELECTION RULE	<u>R-PARITY LSP</u>	<u>KK-PARITY LKP</u>	<u>T-PARITY LTP</u>
→ DISCRETE SYMM.	Neutralino spin 1/2	spin1	spin0
→ STABLE NEW PART.	m_{LSP}	m_{LKP}	m_{LTP}
3) FIND REGION (S) PARAM. SPACE WHERE THE "L" NEW PART. IS NEUTRAL + $\Omega_L h^2$ OK	~100 - 200 GeV *	~600 - 800 GeV	~400 - 800 GeV

* But abandoning gaugino-masss unif. → Possible to have m_{LSP} down to 7 GeV

ELW. SYMM. BREAKING STABILIZATION VS. FLAVOR PROTECTION: THE SCALE TENSION

$$M(B_d - \bar{B}_d) \sim c_{\text{SM}} \frac{(y_t V_{tb}^* V_{td})^2}{16 \pi^2 M_W^2} + c_{\text{new}} \frac{1}{\Lambda^2}$$

If $c_{\text{new}} \sim c_{\text{SM}} \sim 1$

Isidori

$\Lambda > 10^4 \text{ TeV}$ for $O^{(6)} \sim (\bar{s} d)^2$
 [$K^0 - \bar{K}^0$ mixing]

$\Lambda > 10^3 \text{ TeV}$ for $O^{(6)} \sim (\bar{b} d)^2$
 [$B^0 - \bar{B}^0$ mixing]

UV SM COMPLETION TO STABILIZE THE ELW.
 SYMM. BREAKING: $\Lambda_{\text{UV}} \sim O(1 \text{ TeV})$

FLAVOR BLINDNESS OF THE NP AT THE ELW. SCALE?

- **THREE DECADES OF FLAVOR TESTS** (Redundant determination of the UT triangle \longrightarrow verification of the SM, theoretically and experimentally “high precision” FCNC tests, ex. $b \longrightarrow s + \gamma$, CP violating flavor conserving and flavor changing tests, lepton flavor violating (LFV) processes, ...) clearly state that:
- A) in the **HADRONIC SECTOR** the CKM flavor pattern of the SM represents the main bulk of the flavor structure and of CP violation;
- B) in the **LEPTONIC SECTOR**: although neutrino flavors exhibit large admixtures, LFV, i.e. non – conservation of individual lepton flavor numbers in FCNC transitions among charged leptons, is extremely small: once again the SM is right (to first approximation) predicting negligibly small LFV

FROM DETERMINATION TO VERIFICATION OF THE CKM PATTERN FOR HADRONIC FLAVOR DESCRIPTION

$$|V_{us}| \equiv \lambda, \quad |V_{cb}|, \quad R_b, \quad \gamma, \quad \text{TREE LEVEL}$$

$$|V_{us}| \equiv \lambda, \quad |V_{cb}|, \quad R_t, \quad \beta. \quad \text{ONE - LOOP}$$

$$R_b \equiv \frac{|V_{ud}V_{ub}^*|}{|V_{cd}V_{cb}^*|} = \sqrt{\bar{\varrho}^2 + \bar{\eta}^2} = \left(1 - \frac{\lambda^2}{2}\right) \frac{1}{\lambda} \left| \frac{V_{ub}}{V_{cb}} \right|$$

$$R_t \equiv \frac{|V_{td}V_{tb}^*|}{|V_{cd}V_{cb}^*|} = \sqrt{(1 - \bar{\varrho})^2 + \bar{\eta}^2} = \frac{1}{\lambda} \left| \frac{V_{td}}{V_{cb}} \right|.$$

$$R_b = \sqrt{1 + R_t^2 - 2R_t \cos \beta}, \quad \cot \gamma = \frac{1 - R_t \cos \beta}{R_t \sin \beta}, \quad \text{A. BURAS et al.}$$

Reference Unitarity Triangle and UUT (CMFV)

$$(R_b)_{\text{CMFV}} = 0.363 \pm 0.016$$

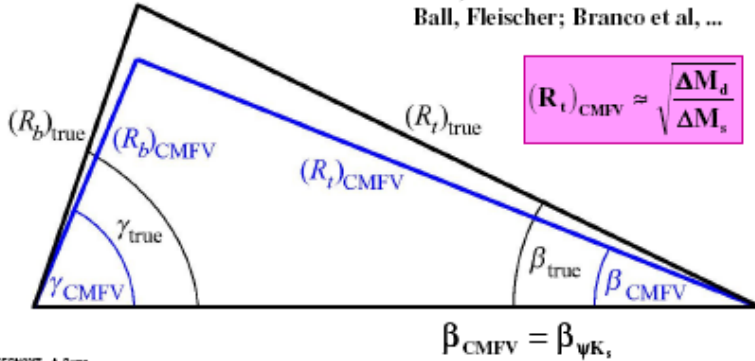
$$(R_b)_{\text{true}} = 0.428 \pm 0.027$$

$$\gamma_{\text{true}} = (82 \pm 20)^\circ$$

$$\sin 2\beta_{\text{CMFV}} = 0.675 \pm 0.026$$

$$\sin 2\beta_{\text{true}} = 0.749 \pm 0.063 \quad \text{"true" = RUT}$$

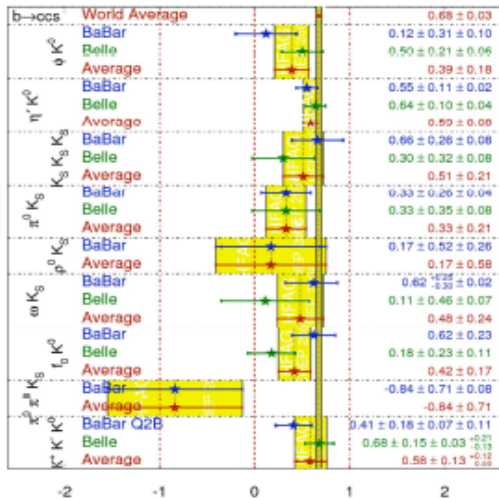
Blanke, AJB, Guadagnoli, Tarantino
 Ufit, CKMfit
 Ball, Fleischer; Branco et al, ...



"sin 2β Problem"

Is this a |V_{ub}| Problem?

Preliminary
 $\sin(2\beta^{\text{eff}}) \equiv \sin(2\phi_1^{\text{eff}})$ HFAG
 ICHP 2006
 PRELIMINARY




second

"sin2β Problem"

Single channels understood?
 Allowed to take the avg.?


Is CP violation entirely due to the KM mechanism? Y.Nir

For CPV in FLAVOR CHANGING* PROCESSES it is VERY LIKELY** that the KM mechanism represents the MAIN SOURCE***

- *FC CPV : as for flavor conserving CPV there could be new phases different from the CKM phase (importance of testing EDMs!)
- **VERY LIKELY: the alternative is to invoke some rather puzzling coincidence (e.g., it could be that $\sin 2\beta$ is not that predicted by the SM , but $H_{SM} + H_{NP}$ in the B_d - \bar{B}_d mixing has the same phase as that predicted by the SM alone or it could be that the phase of the NP contribution is just the same as the SM phase)
- *** MAIN SOURCE : Since $S_{\psi K}$ is measured with an accuracy ~ 0.04 , while the SM accuracy in predicting $\sin 2\beta$ is ~ 0.2  still possible to have

$$H_{NP} \leq 20\% H_{SM} \text{ in } B_d - \bar{B}_d \text{ mixing}$$

□ What to make of this triumph of the CKM pattern in flavor tests?

New Physics at the Elw.
Scale is Flavor Blind
CKM exhausts the flavor
changing pattern at the elw.
Scale 

MINIMAL FLAVOR
VIOLATION

MFV : Flavor originates only
from the SM Yukawa coupl.

New Physics introduces
NEW FLAVOR SOURCES in
addition to the CKM pattern.
They give rise to
contributions which are
<20% in the “flavor
observables” which have
already been observed!

What a SuperB can do in testing CMFV

L. Silvestrini at SuperB IV

Minimal Flavour Violation

In MFV models with one Higgs doublet or low/moderate $\tan\beta$ the NP contribution is a shift of the Inami-Lim function associated to top box diagrams

$$S_0(x_t) \rightarrow S_0(x_t) + \delta S_0(x_t)$$

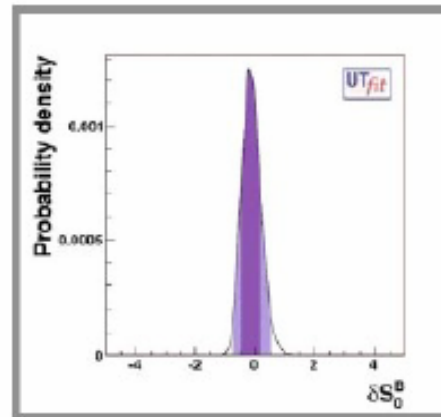
$$\delta S_0(x_t) = 4a \left(\frac{\Lambda_0}{\Lambda} \right)^2$$

$$\Lambda_0 = \frac{\lambda_t \sin^2 \theta_W M_W}{\alpha} \simeq 2.4 \text{ TeV}$$

(D'Ambrosio et al., hep-ph/0207036)

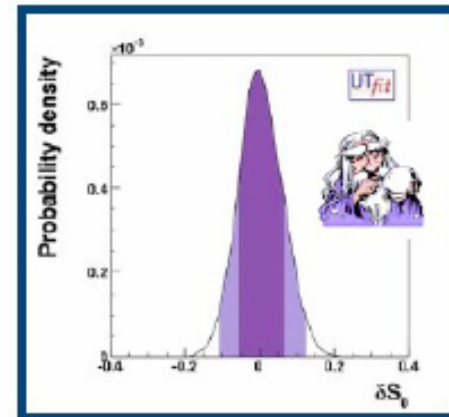
$$\delta S_0^B = \delta S_0^K$$

The "worst" case:
we still probe
virtual particles
with masses up to
 $\sim 12 M_W \sim 1 \text{ TeV}$



$$\delta S_0 = -0.16 \pm 0.32$$

$$\Lambda > 5.5 \text{ TeV @95\%}$$

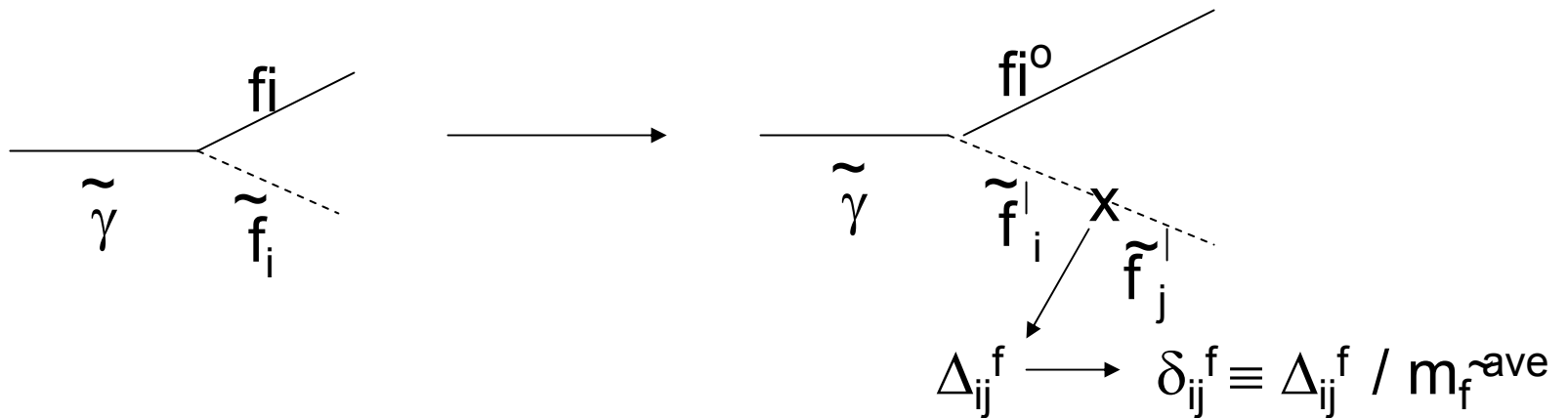


$$\delta S_0 = 0.004 \pm 0.059$$

$$\Lambda > 28 \text{ TeV @95\%}$$

SCKM basis

SUPER CKM: basis in the LOW - ENERGY phenomenology where through a rotation of the whole superfield (fermion + sfermion) one obtains DIAGONAL Yukawa COUPL. for the corresponding fermion field



Unless m_f and $m_{\tilde{f}}$ are aligned, f^0 is not a mass eigenstate

Hall, Kostelecki, Raby

BOUNDS ON THE HADRONIC FCNC: 1 - 3 DOWN GENERATION

	$ \Re(\delta_{13}^d)_{LL} $		$ \Re(\delta_{13}^d)_{LL=RR} $	
x	TREE	NLO	TREE	NLO
0.25	4.9×10^{-2}	6.2×10^{-2}	3.1×10^{-2}	1.9×10^{-2}
1.0	1.1×10^{-1}	1.4×10^{-1}	3.4×10^{-2}	2.1×10^{-2}
4.0	6.0×10^{-1}	7.0×10^{-1}	4.7×10^{-2}	2.8×10^{-2}
	$ \Im(\delta_{13}^d)_{LL} $		$ \Im(\delta_{13}^d)_{LL=RR} $	
x	TREE	NLO	TREE	NLO
0.25	1.1×10^{-1}	1.3×10^{-1}	1.3×10^{-2}	8.0×10^{-3}
1.0	2.6×10^{-1}	3.0×10^{-1}	1.5×10^{-2}	9.0×10^{-3}
4.0	2.6×10^{-1}	3.4×10^{-1}	2.0×10^{-2}	1.2×10^{-2}
	$ \Re(\delta_{13}^d)_{LR} $		$ \Re(\delta_{13}^d)_{LR=RL} $	
x	TREE	NLO	TREE	NLO
0.25	3.4×10^{-2}	3.0×10^{-2}	3.8×10^{-2}	2.6×10^{-2}
1.0	3.9×10^{-2}	3.3×10^{-2}	8.3×10^{-2}	5.2×10^{-2}
4.0	5.3×10^{-2}	4.5×10^{-2}	1.2×10^{-1}	—
	$ \Im(\delta_{13}^d)_{LR} $		$ \Im(\delta_{13}^d)_{LR=RL} $	
x	TREE	NLO	TREE	NLO
0.25	7.6×10^{-2}	6.6×10^{-2}	1.5×10^{-2}	9.0×10^{-3}
1.0	8.7×10^{-2}	7.4×10^{-2}	3.6×10^{-2}	2.3×10^{-2}
4.0	1.2×10^{-1}	1.0×10^{-1}	2.7×10^{-1}	—

SuperB vs. LHC Sensitivity Reach in testing Λ_{SUSY}

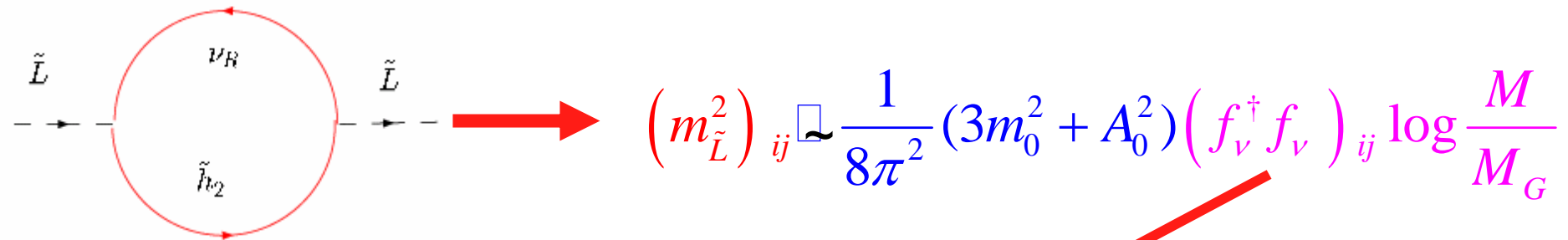
	superB	general MSSM	high-scale MFV
$ \left(\delta_{13}^d\right)_{LL} (LL \gg RR)$	$1.8 \cdot 10^{-2} \frac{m_{\tilde{q}}}{(350\text{GeV})}$	1	$\sim 10^{-3} \frac{(350\text{GeV})^2}{m_{\tilde{q}}^2}$
$ \left(\delta_{13}^d\right)_{LL} (LL \sim RR)$	$1.3 \cdot 10^{-3} \frac{m_{\tilde{q}}}{(350\text{GeV})}$	1	—
$ \left(\delta_{13}^d\right)_{LR} $	$3.3 \cdot 10^{-3} \frac{m_{\tilde{q}}}{(350\text{GeV})}$	$\sim 10^{-1} \tan \beta \frac{(350\text{GeV})}{m_{\tilde{q}}}$	$\sim 10^{-4} \tan \beta \frac{(350\text{GeV})^3}{m_{\tilde{q}}^3}$
$ \left(\delta_{23}^d\right)_{LR} $	$1.0 \cdot 10^{-3} \frac{m_{\tilde{q}}}{(350\text{GeV})}$	$\sim 10^{-1} \tan \beta \frac{(350\text{GeV})}{m_{\tilde{q}}}$	$\sim 10^{-3} \tan \beta \frac{(350\text{GeV})^3}{m_{\tilde{q}}^3}$

SuperB can probe MFV (with small-moderate $\tan\beta$) for TeV squarks; for a generic non-MFV MSSM \longrightarrow sensitivity to squark masses > 100 TeV ! L. Silvestrini

SUSY SEESAW: Flavor universal SUSY breaking and yet large lepton flavor violation

Borzumati, A. M. 1986 (after discussions with W. Marciano and A. Sanda)

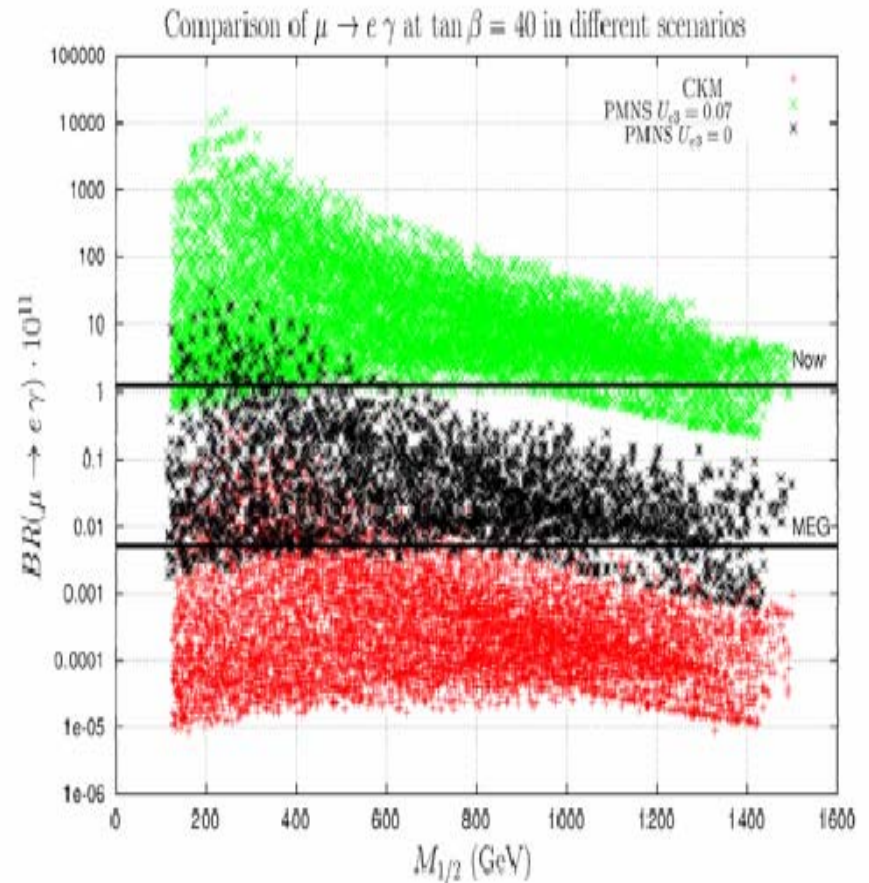
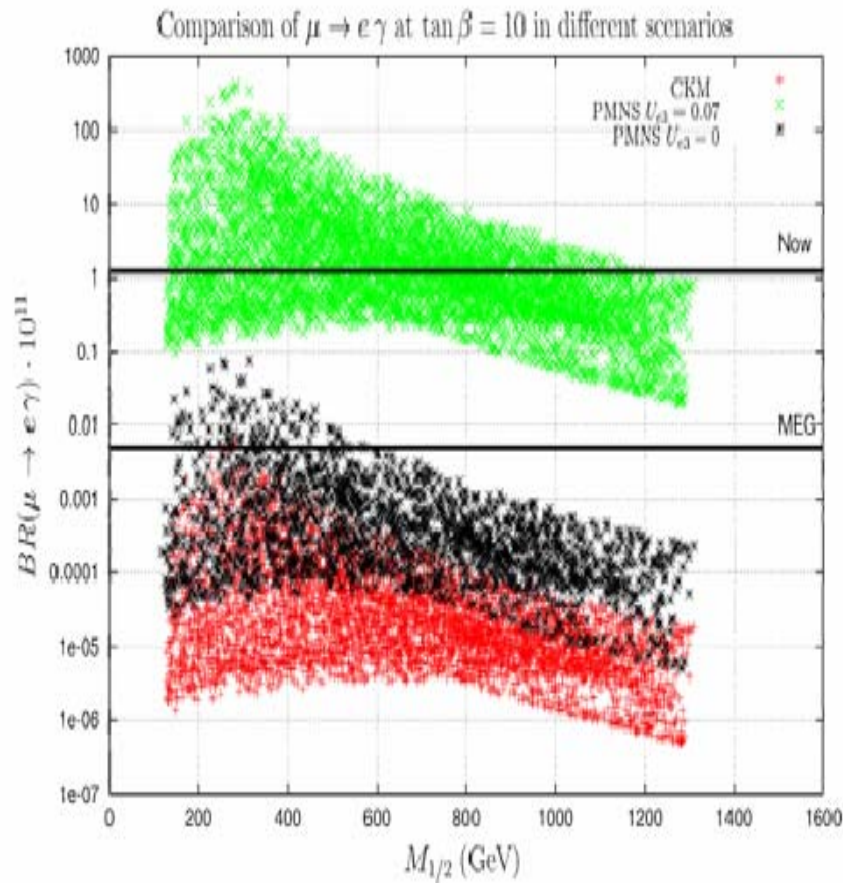
$$L = f_l \bar{e}_R L h_1 + f_\nu \bar{\nu}_R L h_2 + M \nu_R \nu_R$$



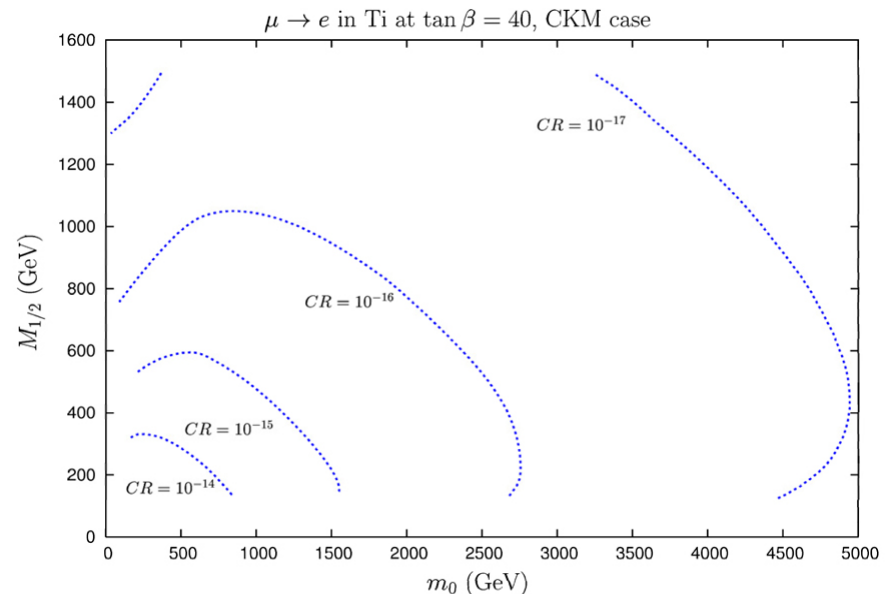
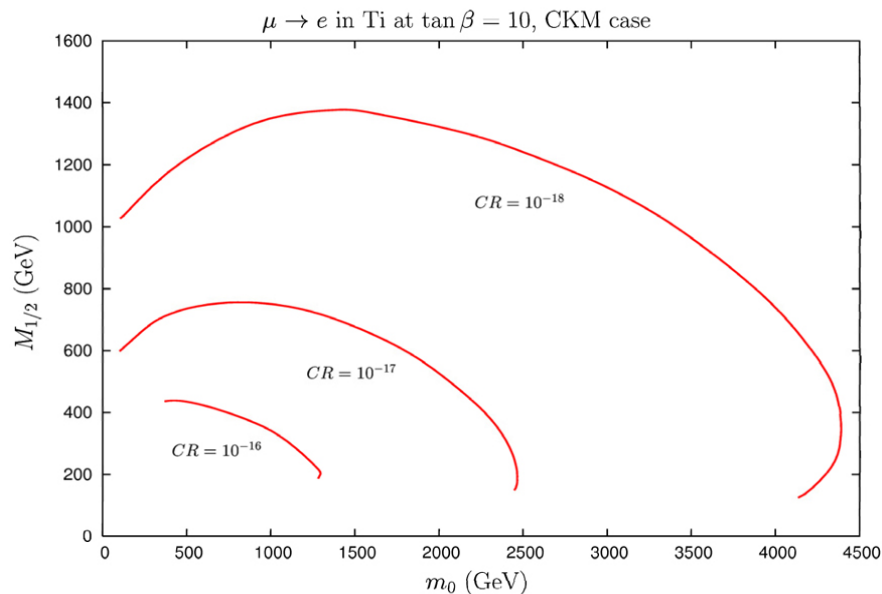
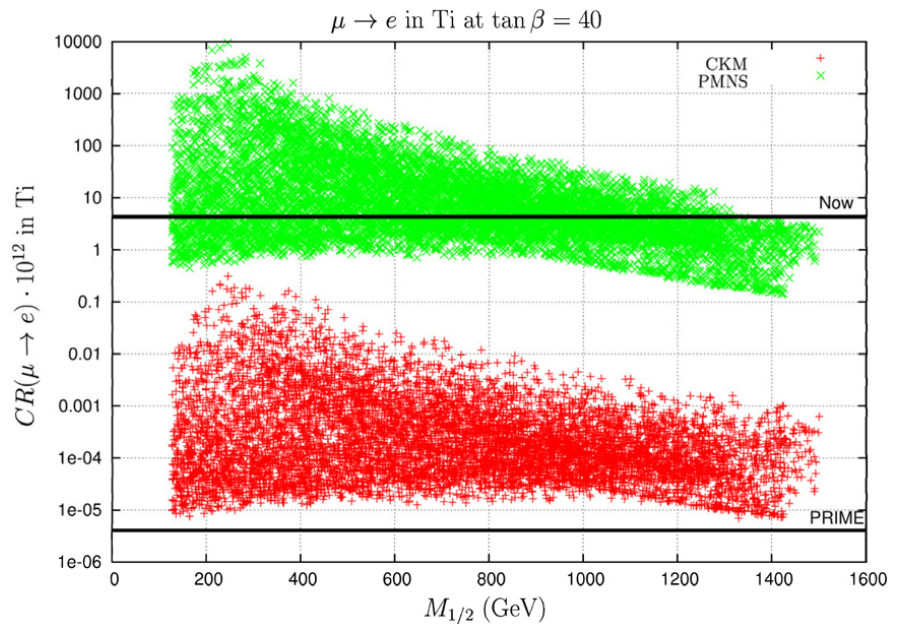
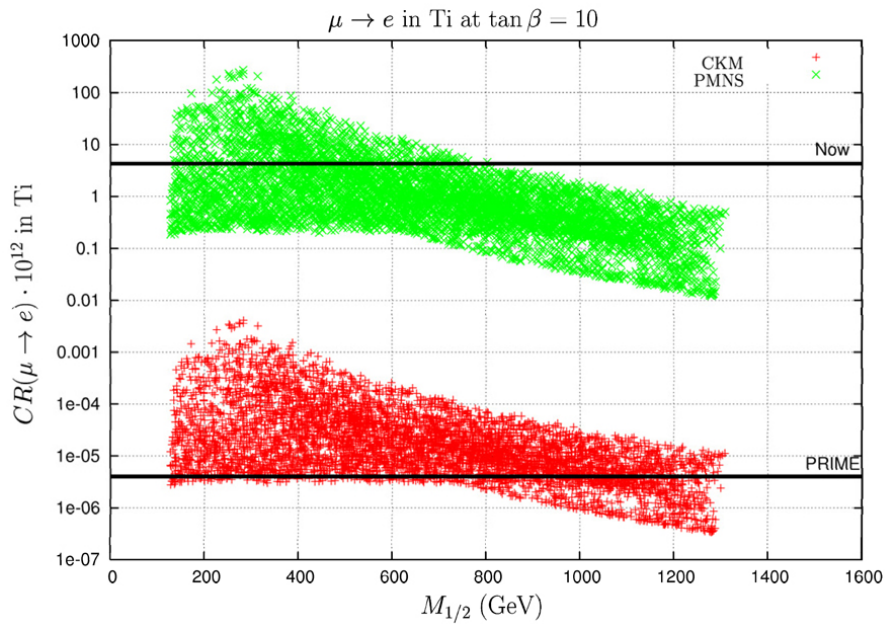
Non-diagonality of the slepton mass matrix in the basis of diagonal lepton mass matrix depends on the unitary matrix U which diagonalizes $(f_\nu^\dagger f_\nu)$

$\mu \rightarrow e + \gamma$ in SUSYGUT: past and future

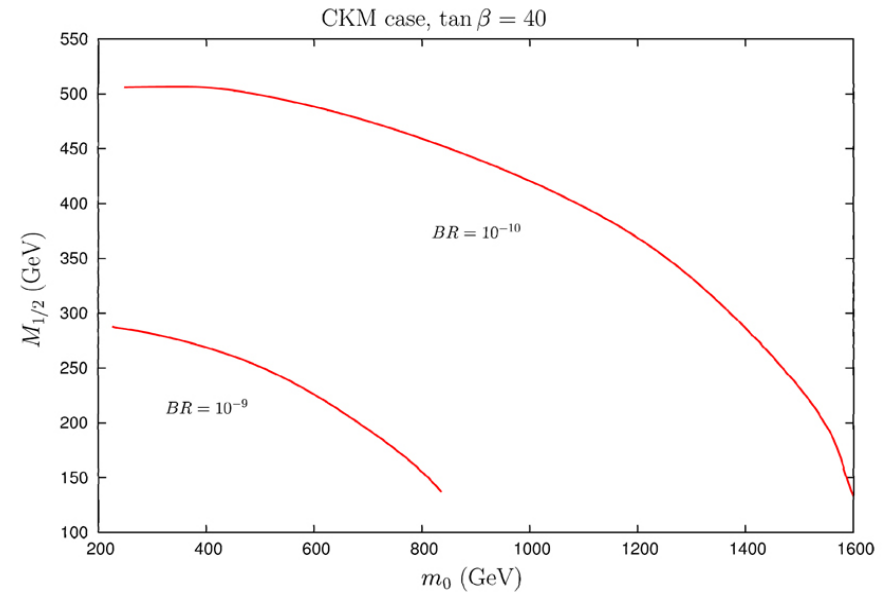
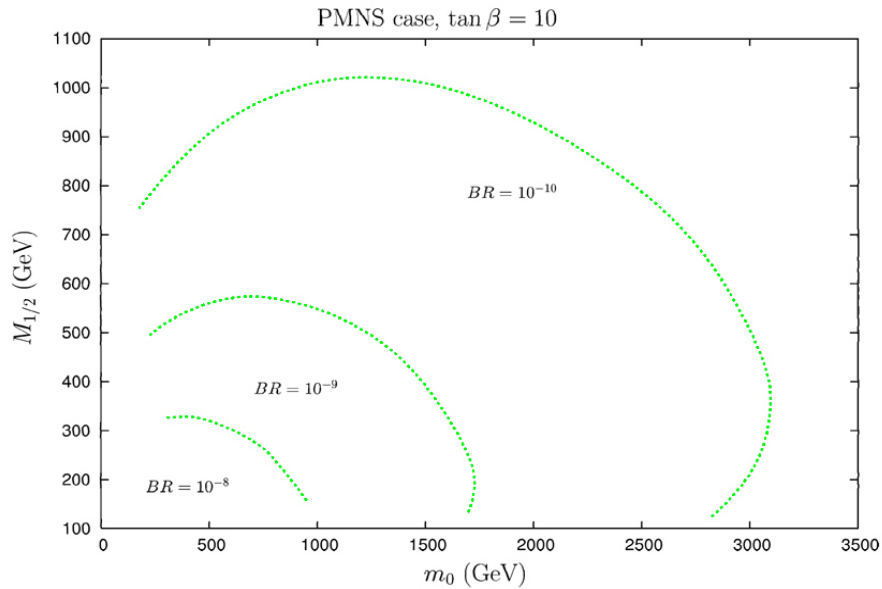
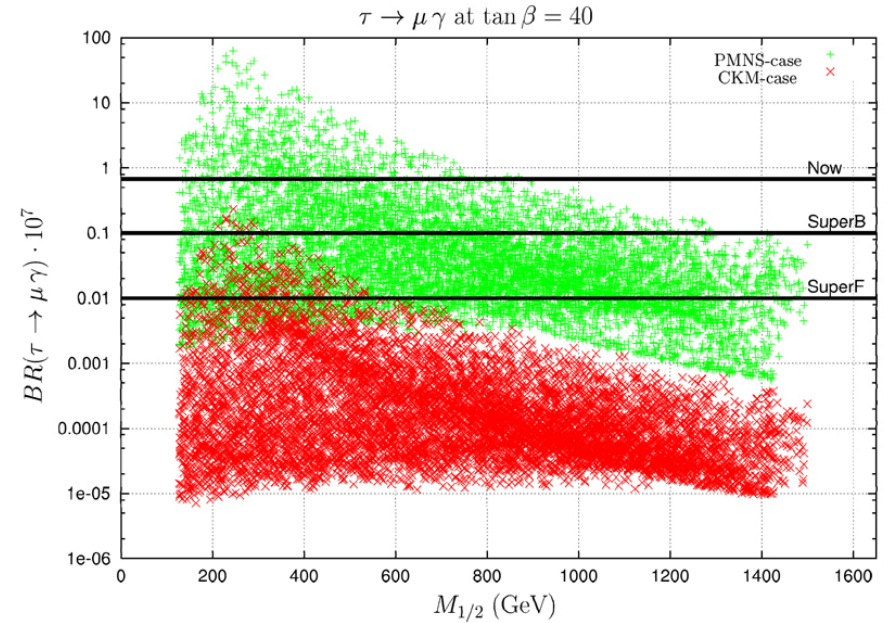
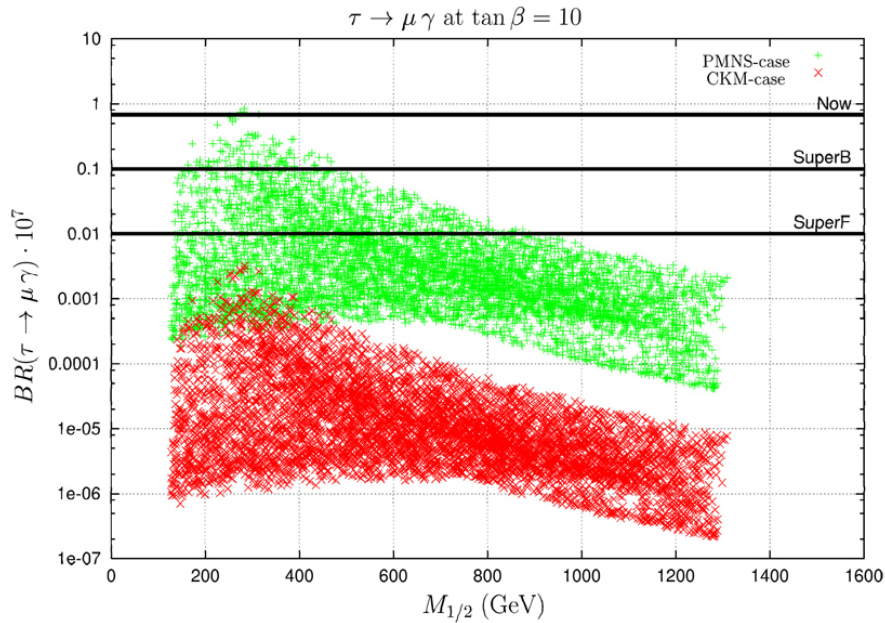
$\mu \rightarrow e \gamma$ in the $U_{e3} = 0$ PMNS case



$\mu \rightarrow e$ in Ti and **PRISM/PRIME** conversion experiment



$\tau \rightarrow \mu \gamma$ and the **Super B** (and **Flavour**) factories



LFV \longleftrightarrow LHC SENSITIVITIES IN PROBING THE SUSY PARAM. SPACE

TABLE IX: Reach in $(m_0, m_{\tilde{g}})$ of the present and planned experiment from their $\tau \rightarrow \mu \gamma$ sensitivity.

Exp.	PMNS		CKM	
	$t_\beta = 40$	$t_\beta = 10$	$t_\beta = 40$	$t_\beta = 10$
BaBar, Belle	1.2 TeV	no	no	no
SuperKEKB	2 TeV	0.9 TeV	no	no
Super Flavour ^a	2.8 TeV	1.5 TeV	0.9 TeV	no

^aPost-LHC era proposed/discussed experiment

DEVIATION from $\mu - e$ UNIVERSALITY

A.M., Paradisi, Petronzio

- Denoting by $\Delta r_{NP}^{e-\mu}$ the deviation from $\mu - e$ universality in $R_{K,\pi}$ due to new physics, i.e.:

$$R_{K,\pi} = R_{K,\pi}^{SM} \left(1 + \Delta r_{K,\pi NP}^{e-\mu} \right),$$

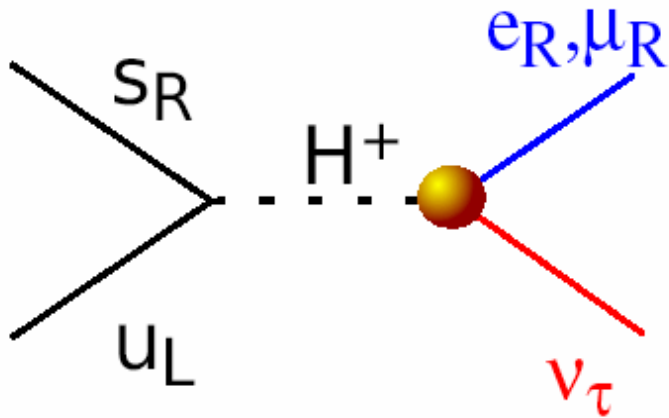
- we get at the 2σ level:

$$-0.063 \leq \Delta r_{K NP}^{e-\mu} \leq 0.017 \quad \text{NA48/2}$$

$$-0.0107 \leq \Delta r_{\pi NP}^{e-\mu} \leq 0.0022 \quad \text{PDG}$$

H mediated LFV SUSY contributions to R_K

$$R_K^{LFV} = \frac{\sum_i K \rightarrow e\nu_i}{\sum_i K \rightarrow \mu\nu_i} \simeq \frac{\Gamma_{SM}(K \rightarrow e\nu_e) + \Gamma(K \rightarrow e\nu_\tau)}{\Gamma_{SM}(K \rightarrow \mu\nu_\mu)}, \quad i = e, \mu, \tau$$



$$eH^\pm \nu_\tau \rightarrow \frac{g_2}{\sqrt{2}} \frac{m_\tau}{M_W} \Delta_R^{31} \tan^2 \beta$$

$$\Delta_R^{31} \sim \frac{\alpha_2}{4\pi} \delta_{RR}^{31}$$

$$\Delta_R^{31} \sim 5 \cdot 10^{-4} \quad t_\beta = 40 \quad M_{H^\pm} = 500 \text{ GeV}$$

$$\Delta r_K^{e-\mu} \simeq \left(\frac{m_K^4}{M_{H^\pm}^4} \right) \left(\frac{m_\tau^2}{m_e^2} \right) |\Delta_R^{31}|^2 \tan^6 \beta \approx 10^{-2}$$

Extension to B \rightarrow $l\nu$ deviation from universality
Isidori, Paradisi

Large ν mixing \leftrightarrow large b-s transitions in SUSY GUTs

In SU(5) $d_R \longleftrightarrow l_L$ connection in the 5-plet
Large $(\Delta^l_{23})_{LL}$ induced by large f_ν of $O(f_{\text{top}})$
is accompanied by large $(\Delta^d_{23})_{RR}$

In **SU(5)** assume large f_ν (Moroi)

In **SO(10)** f_ν large because of an underlying Pati-Salam symmetry

(Darwin Chang, A.M., Murayama)

See also: Akama, Kiyo, Komine, Moroi; Hisano, Moroi, Tobe, Yamaguchi, Yanagida; Hisano, Nomura; Kitano, Koike, Komine, Okada

FCNC HADRON-LEPTON CONNECTION IN SUSYGUT

If



soft SUSY breaking terms arise
at a scale $> M_{GUT}$, they have to respect
the underlying quark-lepton GU symmetry



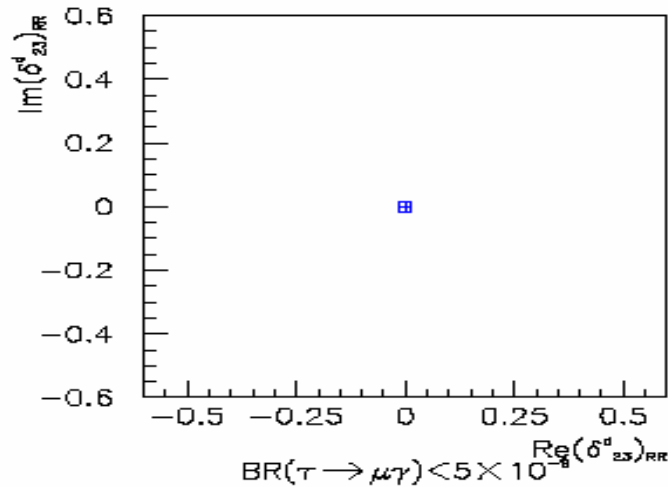
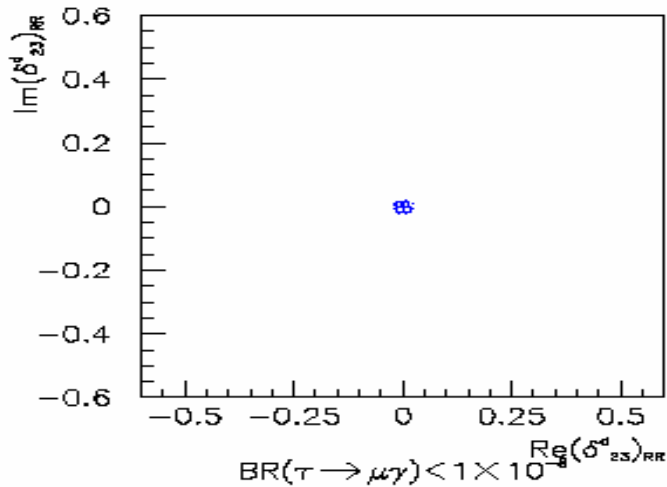
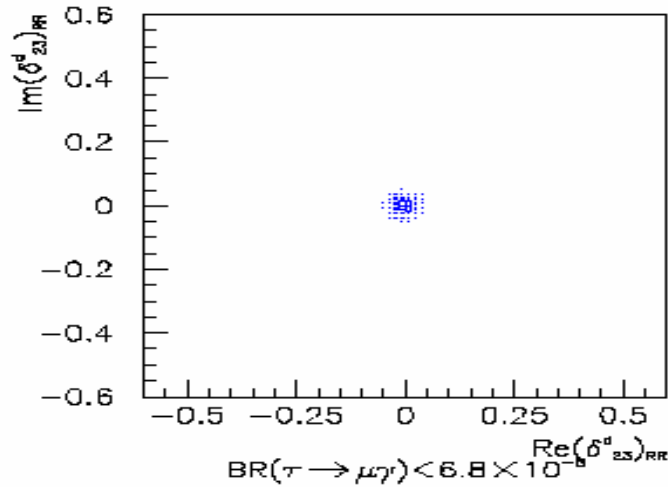
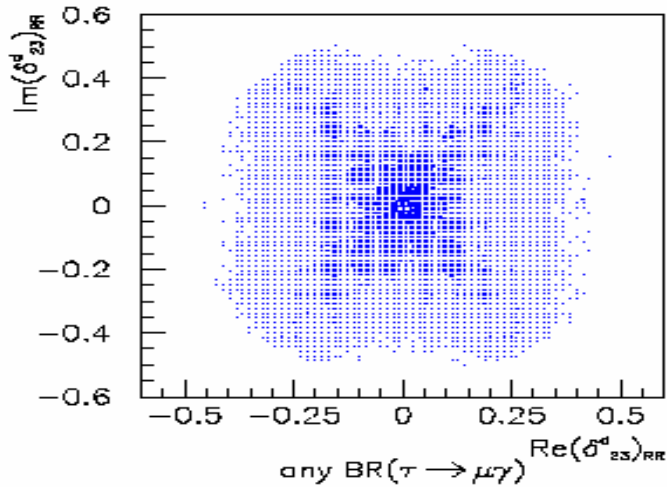
constraints on δ^{quark} from LFV and
constraints on δ^{lepton} from hadronic FCNC

Ciuchini, A.M., Silvestrini, Vempati, Vives PRL

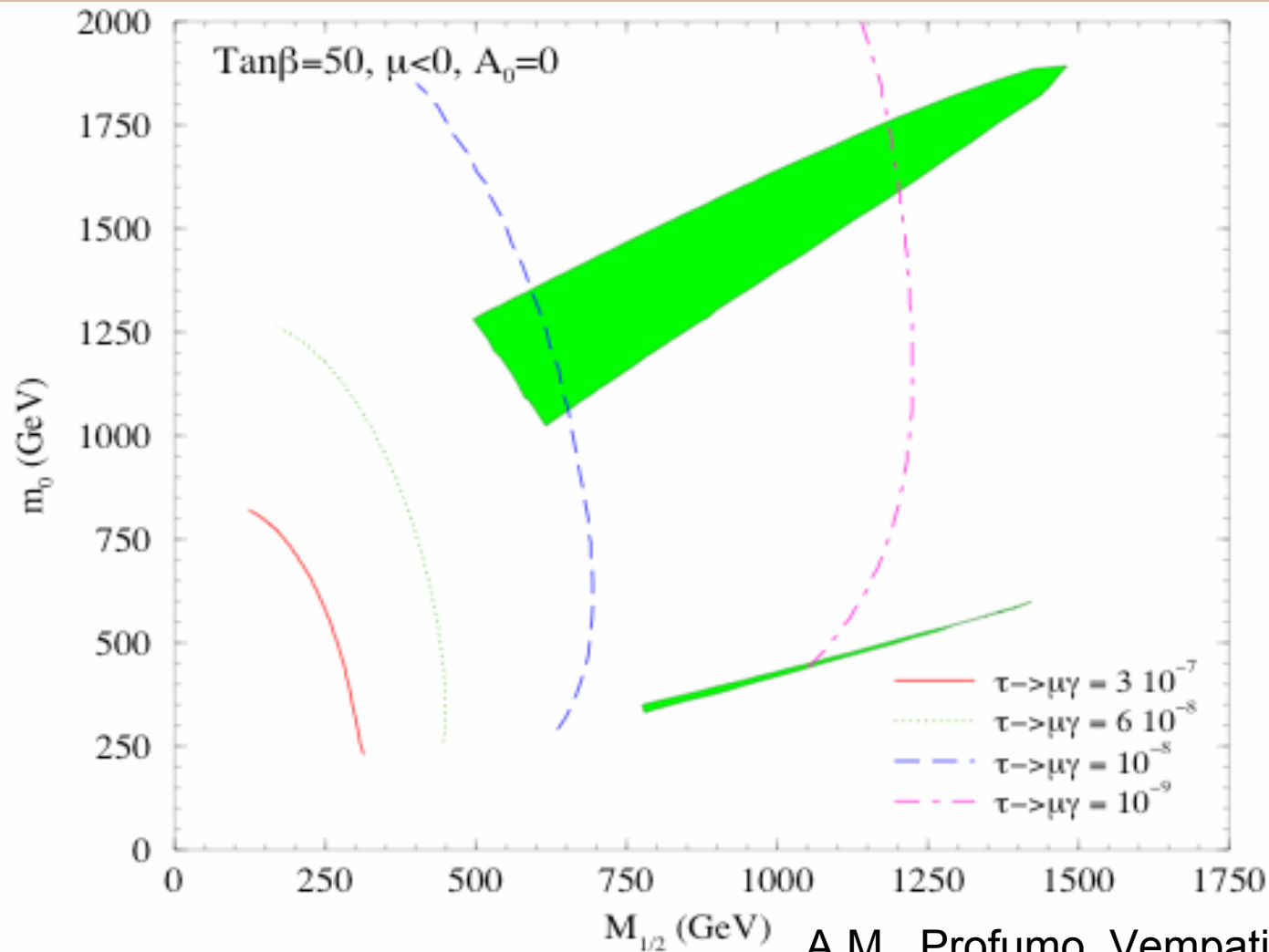
general analysis Ciuchini, A.M., Paradisi, Silvestrini, Vempati, Vives (to appear next week)

Bounds on the hadronic $(\delta_{23})_{RR}$ as modified by the inclusion of the LFV correlated bound

CMPSVV



LFV - DM CONSTRAINTS IN MINIMAL SUPERGRAVITY



SEARCHING FOR WIMPs

WIMPS HYPOTHESIS

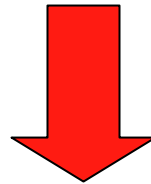
DM made of particles with mass 10Gev - 1Tev

ELW scale

With WEAK INTERACT.

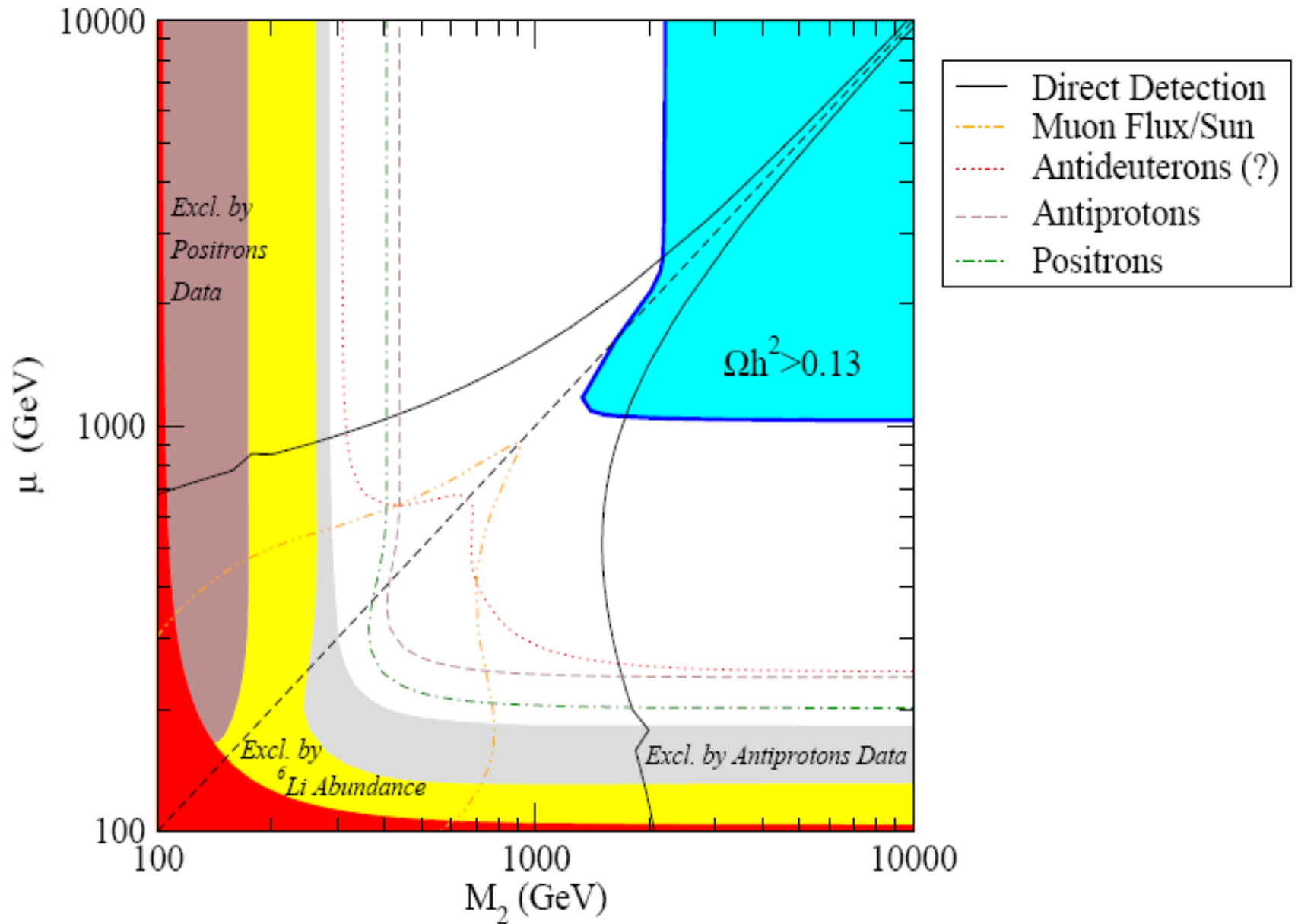
LHC, ILC may PRODUCE WIMPS

WIMPS escape the detector
→ MISSING ENERGY SIGNATURE



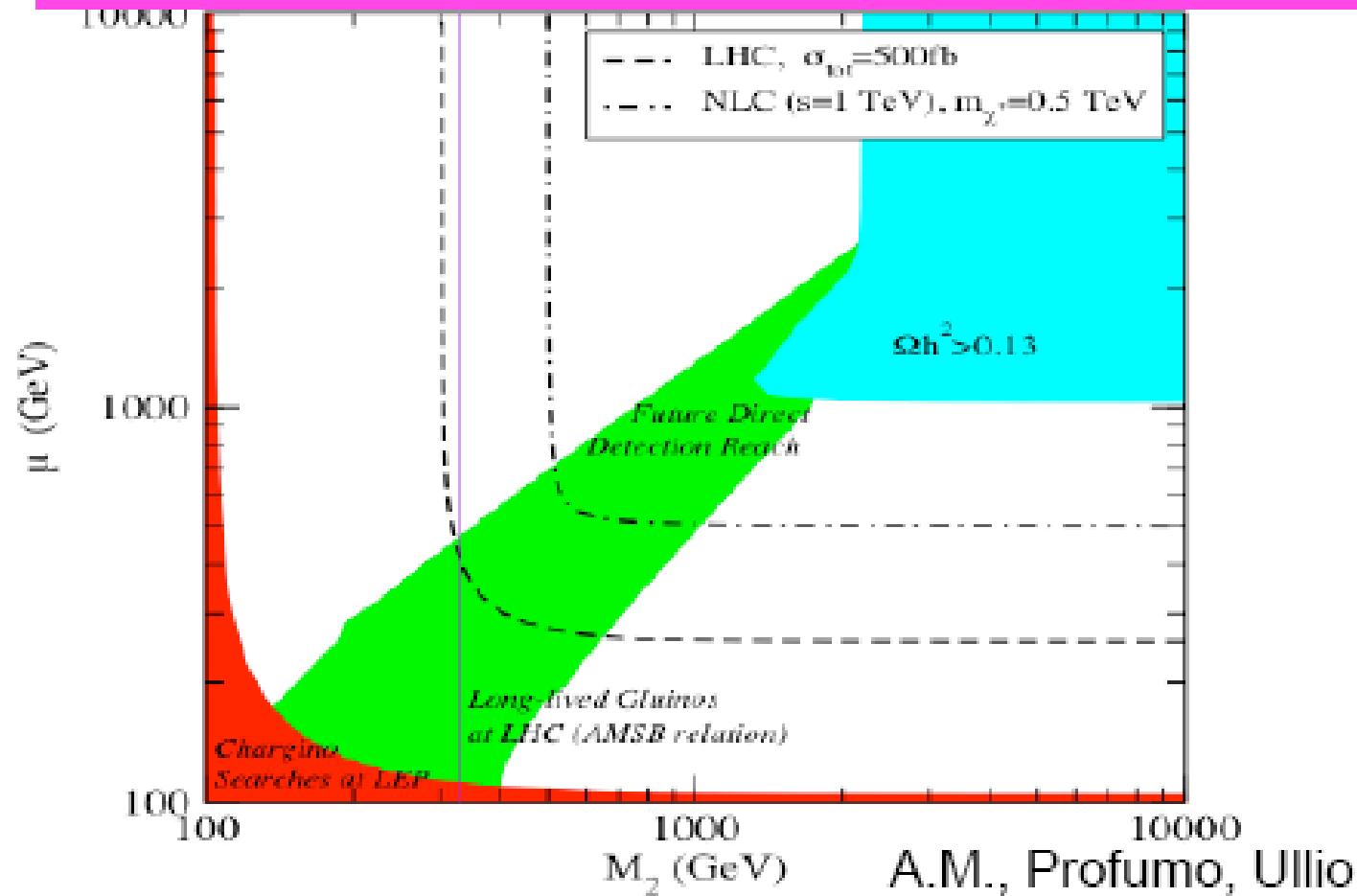
FROM "KNOWN" COSM. ABUNDANCE OF WIMPs → PREDICTION FOR WIMP PRODUCTION AT COLLIDERS WITHOUT SPECIFYING THE PART. PHYSICS MODEL OF WIMPs

BIRKEDAL, MATCHEV, PERELSTEIN, FENG, SU, TAKAYAMA



A.M., PROFUMO, ULLIO

LHC, ILC, DM SEARCHES SENSITIVITIES



Final thoughts on the “complementarity” of flavor physics in our search for NP

- **“Slow” decoupling**: sensitivity to masses of NP larger than what can be explored with LHC (even in strict MFV “exploration power” of flavor physics is in the TeV range)
- At least in SUSY, it is possible, through low-energy FCNC effects induced by the running, to **get access to some large scale** (SUSY SeeSaw scale , Supergravity breaking scale)
- Possible correlation of **hadronic and leptonic FCNC** in SUGRAGUTs
- **“Reconstruction of the fundamental theory”**: ex., once LHC fixes the scale of the NP particles, we can go back to flavor knowledge and try to understand the flavor structure of such NP

TEVATRON → LHC → ILC

DM - FLAVOR
for DISCOVERY
and/or FUND. TH.
RECONSTRUCTION

A MAJOR
LEAP AHEAD
IS NEEDED

NEW
PHYSICS AT
THE ELW
SCALE

DARK MATTER

"LOW ENERGY"

PRECISION PHYSICS

$m_\chi, n_\chi, \sigma_\chi \dots$

LINKED TO COSMOLOGICAL EVOLUTION

→ Possible interplay with dynamical DE

FCNC, CP ≠, (g-2), $(\beta\beta)_{0\nu\nu}$

LFV

BARYO-LEPTO-GENESIS