

# Minimal Flavour Violation Scenarios

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## Origin of Quark and Lepton Flavour in the SM remains a mystery!

- Only **parametrization in terms of Yukawa matrices**, yielding fermion masses and mixing angles after EWSB.
- In case of **new physics at the TeV scale** ( $\rightarrow$  hierarchy problem), and assuming generic new flavour couplings:

*Why NP hasn't been observed in rare flavour decays?*

- Postulate additional constraints on NP flavour sector.
- Popular concept: **Minimal Flavour Violation (MFV)**:

*New flavour couplings related to SM mixing angles and masses.*

[D'Ambrosio et al. '02, Ciuchini et al. '98, Buras et al. '00]

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## MFV does ...

- provide **additional suppression factors** for NP flavour transitions, based on a **symmetry principle**.
- imply (SM-like) **correlations** between different flavour observables.
- give a **systematic book-keeping device** to parametrize flavour couplings in specific new-physics models.  
(**reduction of free parameters** in NP flavour sector)
- allow to formulate flavour violation within **effective theory approach**.

## MFV does not ...

- represent a *theory* of flavour violation.  
(overall couplings in front of MFV structures still undetermined)
- *explain* the size of fermion masses and mixing angles.  
(origin of Yukawa matrices still unclear)

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## Effective-Theory Framework for MFV

- MFV in the quark sector:

- 1-Higgs-Doublet Models.
- [Large  $\tan\beta$  effects in 2-Higgs Models.]

[D'Ambrosio et al., hep-ph/0207036]

- MFV in the lepton sector:

- minimal field content (+ dim-5 LN-violating operator)
- [extended field content (right-handed neutrino)]

[Cirigliano et al., hep-ph/0507001, hep-ph/0601111]

- [from R-parity violating tri-linear coupling]

[Davidson/Palorini, hep-ph/0607329]



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# Quark flavour in the SM

- Flavour symmetry group that commutes with SM gauge group:

$$F = SU(3)_{Q_L} \times SU(3)_{U_R} \times SU(3)_{D_R}$$

- Quark multiplets for three families transform as

$$Q_L = \begin{pmatrix} U_L \\ D_L \end{pmatrix} \sim (3, 1, 1), \quad U_R \sim (1, 3, 1), \quad D_R \sim (1, 1, 3).$$

- SM Yukawa matrices  $Y_U, Y_D$  break flavour symmetry  $F$ :
- Diagonalized by bi-unitary transformations:  $\Rightarrow m_q^{\text{diag}} = y_q v$
- $V_{CKM} = V_{u_L}^\dagger V_{d_L}$  parameterizing the mismatch between the left-handed up- and down-quark rotations.

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# Yukawa matrices as Spurions

- Consider  $Y_U, Y_D$  as **spurion fields**, transforming as

$$Y_U \sim (3, \bar{3}, 1), \quad Y_D \sim (3, 1, \bar{3}).$$

under flavour group  $F = SU(3)_{Q_L} \times SU(3)_{U_R} \times SU(3)_{D_R}$ .

- Yukawa terms **formally** invariant under flavour group.
- Background values for  $Y_U$  and  $Y_D$  finally break  $F$ .
- Concept can be generalized to higher-dimensional operators!

**Analogy:** Treatment of quark mass matrix in ChPT ...

# New Physics and Minimal Flavour Violation

- Consider SM as **Effective Theory** at the electroweak scale.
- Allow for higher-dimensional operators , induced by **NP effects** above the electroweak scale.
- **Assume** that Flavour- *and* CP-symmetry breaking is still governed by SM spurions  $Y_U$  and  $Y_D$ , only.

NP contributions to flavour transitions **suppressed** by same factors of  $V_{CKM}$  and  $m_Q$  as in the SM.

# Example: MFV in the MSSM

- New flavour structures arise from "soft SUSY-breaking" terms.
  - Squark mass matrices:

$$\tilde{m}_{Q_L}^2 = \tilde{m}^2 \left( a_1 + b_1 Y_U Y_U^\dagger + b_2 Y_D Y_D^\dagger + \dots \right),$$

$$\tilde{m}_{U_R}^2 = \tilde{m}^2 \left( a_2 + b_5 Y_U^\dagger Y_U \right),$$

$$\tilde{m}_{D_R}^2 = \tilde{m}^2 \left( a_3 + b_6 Y_D^\dagger Y_D \right)$$

- Squark tri-linear couplings:

$$A_U = A \left( a_4 + b_7 Y_D^\dagger Y_D \right) Y_U,$$

$$A_D = A \left( a_5 + b_8 Y_U^\dagger Y_U \right) Y_D,$$

- Integrating out SUSY particles yields MFV structures at EW scale.

# Wolfenstein expansion

- Off-diagonal CKM elements suppressed by Wolfenstein- $\lambda \sim 0.2$ .
- Quark masses small compared to EW scale. (except for  $m_t$ )

⇒ The more spurion insertions, the more suppression.

- Only need to consider minimal number of spurion insertions:
  - no insertion necessary for charged left-handed decays,
  - 1 insertion necessary for charged right-handed decays,
  - 4 insertions necessary for neutral right-handed decays.

- All FCNCs with external down-type quarks controlled by

$$(\Delta_{LL}^q)_{i \neq j} = (Y_U Y_U^\dagger)_{ij} = y_t^2 (V_{CKM}^*)_{3i} (V_{CKM})_{3j} \quad (y_t = \mathcal{O}(1))$$



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# Example: $b \rightarrow s\gamma$

- Leading operators in weak effective Hamiltonian:

$$O_7 = \bar{s}_L \sigma_{\mu\nu} b_R F^{\mu\nu}, \quad O'_7 = \bar{s}_R \sigma_{\mu\nu} b_L F^{\mu\nu}.$$

- Construct FCNC  $D_L$ - $D_R$  transitions, invariant under  $F$ :

$\bar{Q}'_L D'_R$				not invariant
$\bar{Q}'_L Y_D D'_R$	$\rightarrow$	$y_D^{\text{diag}} \bar{D}_L D_R$		flavour-diagonal
$\bar{Q}'_L Y_U Y_U^\dagger Y_D D'_R$	$\rightarrow$	$\bar{D}_L V_{CKM}^\dagger (y_U^{\text{diag}})^2 V_{CKM} y_D^{\text{diag}} D_R$		✓
$\bar{D}'_R Y_D^\dagger Y_U Y_U^\dagger Q'_L$	$\rightarrow$	$\bar{D}_R y_D^{\text{diag}} V_{CKM}^\dagger (y_U^{\text{diag}})^2 V_{CKM} D_L$		✓

- FCNC require (at least) two CKM elements.
  - right-handed quarks suppressed by quark mass.
- $O_7$  pre-factor:  $\sim \frac{m_b}{v} y_t^2 V_{ts}^* V_{tb}$
  - $O'_7$  pre-factor:  $\sim \frac{m_s}{v} y_t^2 V_{ts}^* V_{tb}$

The same flavour factors as from penguin diagrams in the SM !

- No significant NP impact on **CP asymmetries**:
  - e.g.  $A_{CP}^{\text{dir,mix}}$  for  $B \rightarrow J/\psi K_s$  and  $B \rightarrow \phi K_s$  as in SM
- NP contributions to different **FCNC decays** are related:
  - e.g. ratio of  $b \rightarrow s\gamma$  and  $b \rightarrow d\gamma$  still determined by  $|V_{ts}/V_{td}|$
  - $\Delta M_{B_s}/\Delta M_{B_d}$  still given by  $|V_{ts}/V_{td}|$
- NP contributions to **charged decays** insignificant:
  - e.g.  $|V_{ub}|$  from  $b \rightarrow u\ell\nu$
- Corrections to flavour observables smaller than in generic case  
⇒ NP scales in the TeV range still allowed.

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# Phenomenological Signatures for beyond MFV ?

- Violation of MFV relations in FCNC:

- $b \rightarrow s\gamma$  vs.  $b \rightarrow d\gamma$
- $\Delta M_{B_d}$  vs.  $\Delta M_{B_s}$
- $B \rightarrow J/\psi K_S$  vs.  $B \rightarrow \phi K_S$
- ...

SM, MFV ✓  
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( $\rightarrow$  nMFV)

- Sizeable right-handed contributions to charged decays:

- test of left-handedness in  $b \rightarrow cl\nu$   
from moment analysis of inclusive decay spectra

SM, MFV ✓

[Mannel et al, hep-ph/0701054]

- inconsistencies in  $|V_{ub}/V_{cb}|$  measurements  
(exclusive vs. inclusive vs.  $\sin 2\beta$ )
- ...

(( $\rightarrow$  nMFV))



## Variations on MFV / Beyond MFV

- Extended Scalar Sector:
  - 2-Higgs models:  $y_b \cdot \tan \beta \sim \mathcal{O}(1)$  [D'Ambrosio et al.]
  - MSSM with large  $\mu$ :  $y_b \cdot |\mu| \sim \mathcal{O}(\tilde{m})$  [Buras et al.]
  - ...
- Extended Fermion Sector:
  - Littlest Higgs Models with  $T$ -parity [Buras et al., Hubisz et al.]
  - Froggatt-Nielsen scenarios
  - ...
- Extended Gauge Sector:
  - MFV and  $SU(5)$  [Grinstein et al., hep-ph/0608123]
- Dominant couplings to 3rd generation [Agashe et al., hep-ph/0509117]

General framework would allow for several new spurion fields, with different transformations under flavour group, but still small expansion parameters in off-diagonal elements.

[TF/Mannel 2006]



# Lepton flavour in the (minimally extended) SM

- Flavour symmetry group:  $F_{\text{lept.}} = SU(3)_{L_L} \times SU(3)_{E_R}$
- Lepton multiplets:  $L_L = \begin{pmatrix} \nu_\ell \\ \ell \end{pmatrix}_L \sim (3, 1), \quad E_R \sim (1, 3).$
- Charged-lepton Yukawas  $Y_\ell \sim (3, \bar{3})$  break flavour symmetry.
- No neutrino Yukawas  $\Rightarrow$  No lepton-mixing in the SM

## Including Neutrino Masses

- The combination  $N = H^T \tau_2 L_L$  is a SM gauge singlet.
- ⇒ Gauge-invariant dim-5 operator:

$$\mathcal{L}_{\text{Maj}} = \frac{1}{2\Lambda_{LN}} (N^T g N)$$

- Does not conserve lepton number.
- Generates neutrino Majorana mass terms of order  $v^2/\Lambda_{LN} \lll 1$

- New flavour matrix: (in basis where  $Y_\ell$  is diagonal)

$$g = \frac{\Lambda_{LN}}{v^2} U_{\text{PMNS}}^* m_\nu^{\text{diag}} U_{\text{PMNS}}^\dagger$$

**MLFV:** Treat new flavour matrix as spurion field:  $g \sim (\bar{6}, 1)$

## Uniqueness of MLFV ?

- Mechanism for neutrino mass generation unknown !
- Implementation of MFV in Lepton Sector not unambiguous !

see the discussion in [Davidson/Palorini, hep-ph/0607329]

# Lepton-flavour transitions induced by spurion $g$

- Consider all gauge-invariant operators of dim-6 (and higher)

[Cirigliano et al., hep-ph/0507001]

- Bi-linear building blocks:

$$(\bar{L} g^\dagger g L), \quad (\bar{L} g^\dagger g Y_\ell^\dagger R), \quad (\bar{R} Y_\ell g^\dagger g Y_\ell^\dagger R)$$

⇒ e.g. radiative decays  $l \rightarrow l' \gamma$  controlled by  $y_\ell$  and

$$(g^\dagger g)_\tau^e \simeq \frac{\Lambda_{\text{LN}}^2}{\sqrt{2}v^4} \left( -s_{12}c_{12} \Delta m_{\text{sol}}^2 \pm s_{13}e^{i\delta} \Delta m_{\text{atm}}^2 \right),$$

$$(g^\dagger g)_\tau^\mu \simeq \frac{\Lambda_{\text{LN}}^2}{2v^4} \left( -c_{12}^2 \Delta m_{\text{sol}}^2 \pm \Delta m_{\text{atm}}^2 \right),$$

$$(g^\dagger g)_\mu^e \simeq \frac{\Lambda_{\text{LN}}^2}{\sqrt{2}v^4} \left( s_{12}c_{12} \Delta m_{\text{sol}}^2 \pm s_{13}e^{i\delta} \Delta m_{\text{atm}}^2 \right),$$

(insensitive to absolute neutrino mass scale and Majorana phases)

# Lepton-flavour transitions induced by spurion $g$

(continued)

- 4-lepton operators:

$$(\bar{L} g^\dagger g L)(\bar{L} L), \quad (\bar{L} g^\dagger g L)(\bar{R} R), \quad \text{etc.}$$

include the same flavour structures as radiative decays.

- independent flavour structure from:

$$(g)_{ij}(g^*)^{kl} (\bar{L}_k L^i)(\bar{L}_l L^j)$$

(corresponds to 27plet in  $\bar{6} \otimes 6 = 1 + 8 + 27$ )

- Sensitive to **absolute neutrino mass scale** and **Majorana phases** !

[Cirigliano/Grinstein, hep-ph/0601111]

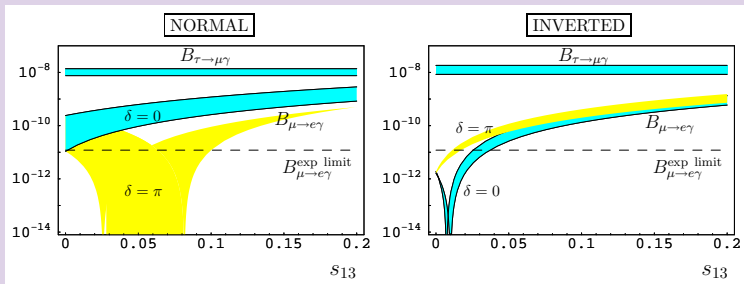
# MLFV phenomenology

- dim-6 operators suppressed by  $1/\Lambda_{\text{NP}}^2$
- dominating flavour coefficients  $|(g^\dagger g)_{ij}| \lesssim 1$
- Sizeable LFV effects possible for  $\Lambda_{\text{NP}}$  in the TeV range  
(i.e.  $\Lambda_{\text{NP}} \ll \Lambda_{\text{LN}}$ )

- radiative decays  $\ell \rightarrow \ell' \gamma$  →
- decays into charged leptons  $\tau \rightarrow 3\ell$  and  $\mu \rightarrow 3\ell$  →
- hadronic decays ( $\tau \rightarrow \mu \pi^0 \dots$ ),  
 $\mu - e$  conversion in nuclei. [hep-ph/0601111]
- leptogenesis [hep-ph/0607068, hep-ph/0609067, hep-ph/0612262]



- $s_{13} \leq 0.25$  and  $\Delta m_{\text{atm}}^2 \gg \Delta m_{\text{sol}}^2$ ,
- fix remaining parameters such that  $B_{\tau \rightarrow \mu \gamma} = \frac{\Gamma(\tau \rightarrow \mu \gamma)}{\Gamma(\tau \rightarrow \mu \nu_{\tau} \bar{\nu}_{\mu})} \sim 10^{-8}$



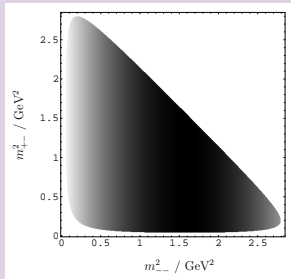
- $\mu \rightarrow e \gamma$  most sensitive probe of MLFV scenario.
- With existing bounds on  $\mu \rightarrow e \gamma$  and for  $s_{13} \gtrsim 0.1$ , observation of  $\tau \rightarrow \mu \gamma$  at near future experiments unlikely within MLFV.

## Leptonic decays $\ell \rightarrow 3\ell'$

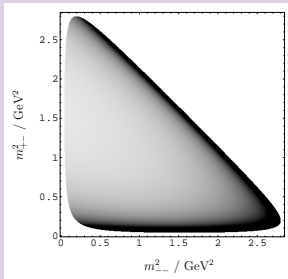
- two competing flavour structures:  $(gg^\dagger)_{\ell\ell'}$  vs.  $g_{\ell\ell'} g_{\ell'\ell}^*$   
 $\Rightarrow$  Study ratios  $\Gamma[\ell \rightarrow 3\ell']/\Gamma[\ell \rightarrow \ell'\gamma]$ . [hep-ph/0601111]
- 4-lepton operators vs. radiative decays  $\Rightarrow$  study decay distributions.

Example: Dalitz plots for  $\tau^- \rightarrow \mu^- \mu^+ \mu^-$  contributions:

$$(\bar{\mu}_L \gamma^\mu \tau_L)(\bar{\mu}_L \gamma_\mu \mu_L)$$



$$\tau \rightarrow \mu \gamma^* \rightarrow 3\mu$$



[Dassinger/TF/Mannel/Turczyk, preprint HEP-2007-06]

## Exploring NP flavour sector:

(1.) Establish NP in **discovery channels**

(2.) Test **flavour precision data** against MFV hypothesis:

**either:** disprove MFV (violation of correlations)

↪ Allow for new flavour structures (spurion fields).

↪ Flavour puzzle even more pronounced.

**or:** determine parameters within your favourite NP model for (1.)

↪ use MFV for book-keeping.