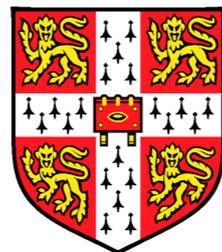


Resonances in Coupled-Channel Scattering

David Wilson

Lattice 2016
University of Southampton
24-30 July 2016



**UNIVERSITY OF
CAMBRIDGE**

Coupled-channel scattering

This talk:

Topical report of recent coupled-channel scattering results from the Hadron Spectrum Collaboration

The method:

- Build large correlation matrices with a diverse range of operators
- Extract many energy levels using the variational method
- Use these energies with extensions of Lüscher's method to obtain infinite volume scattering amps
- Investigate the poles of the scattering amplitudes to obtain resonance information

Topics I won't cover:

The HALQCD method, Finite Volume Hamiltonian, EFTs in a box, etc.

Coupled-channel scattering

Sounds hard... why bother?

$a_0(980)$, $f_0(980)$

$a_1(1260)$

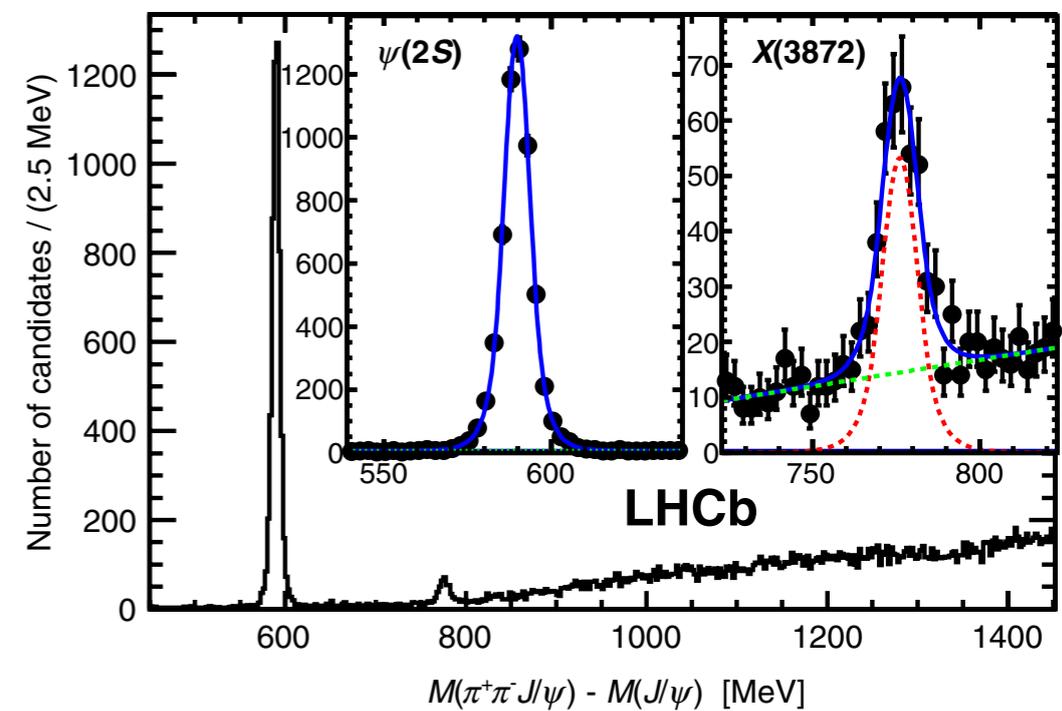
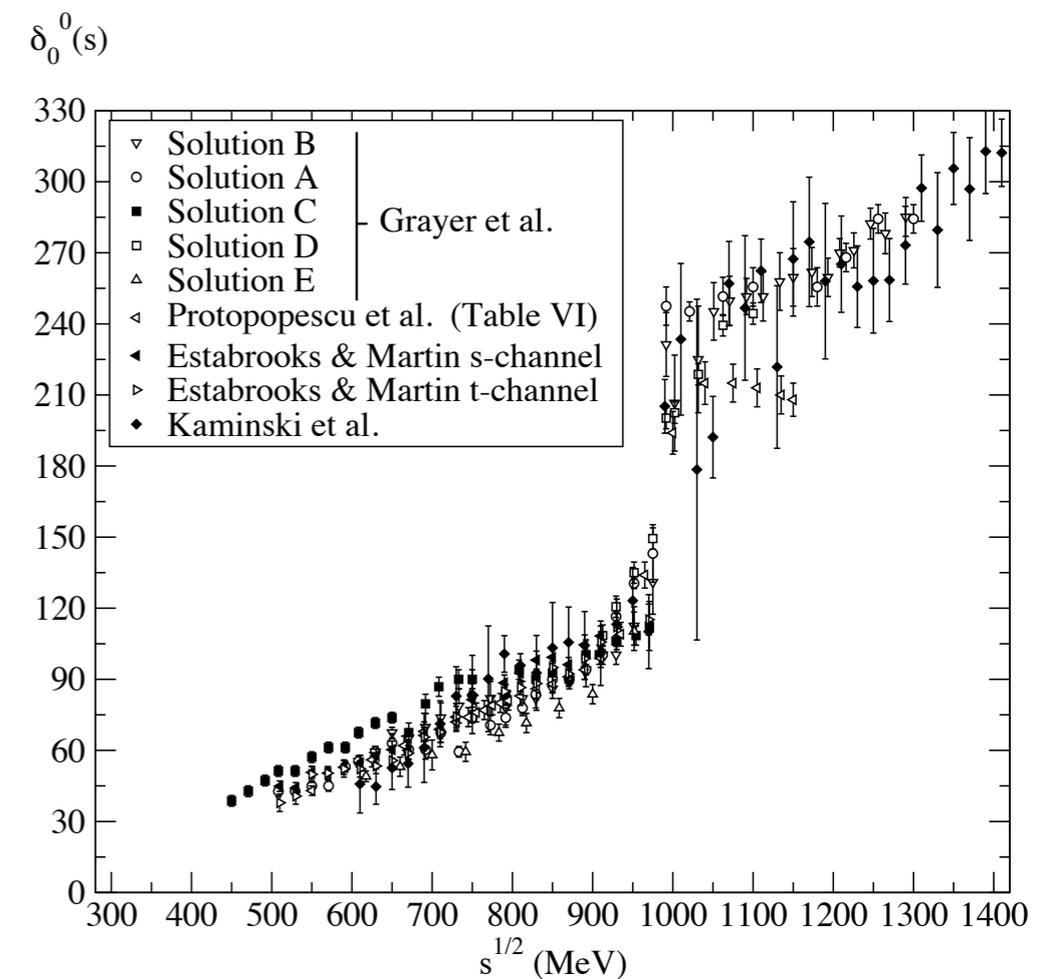
$X(3872)$, and other XYZ states

$N^*(1440)$, $\Lambda(1405)$, ...

all decay into multiple final states

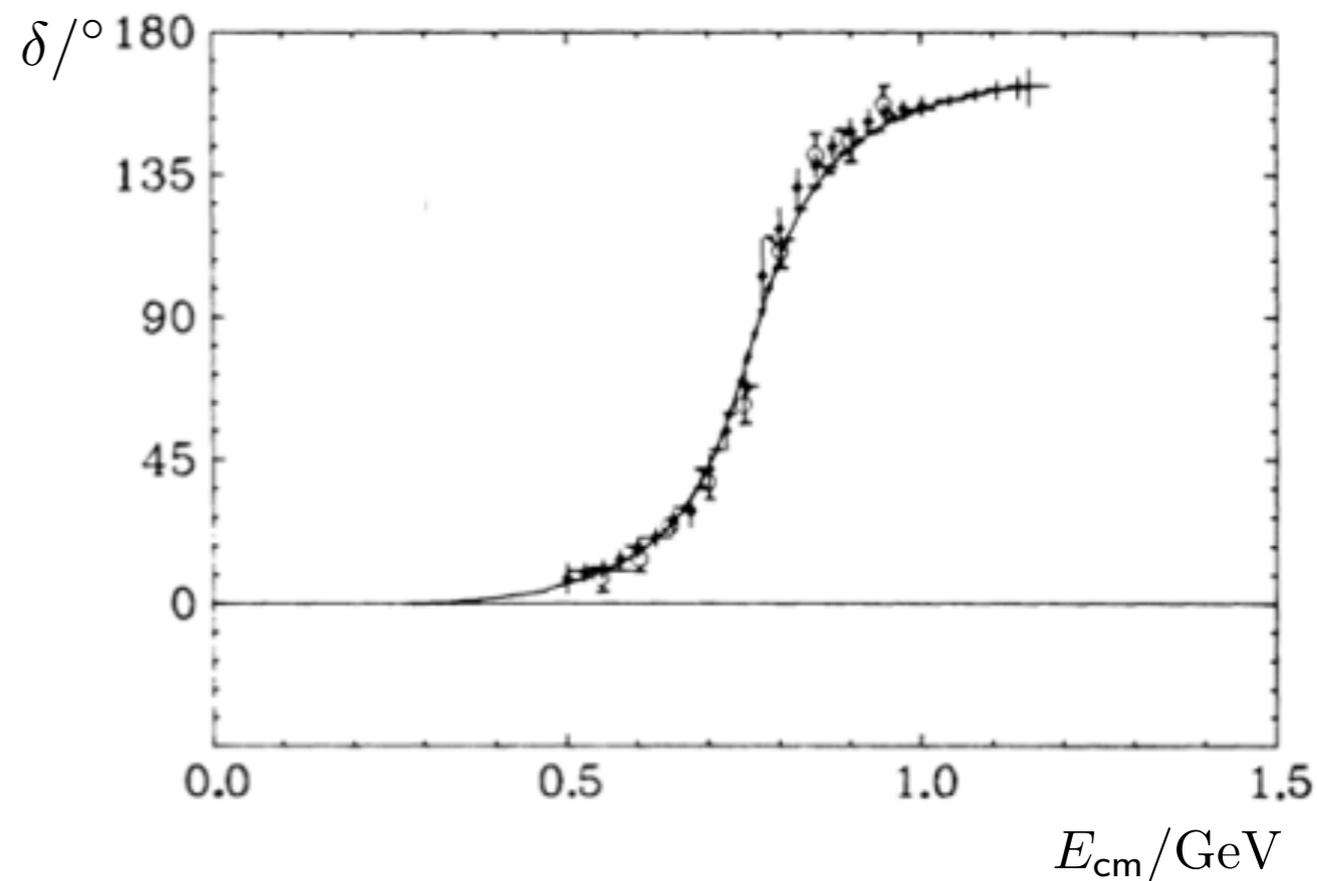
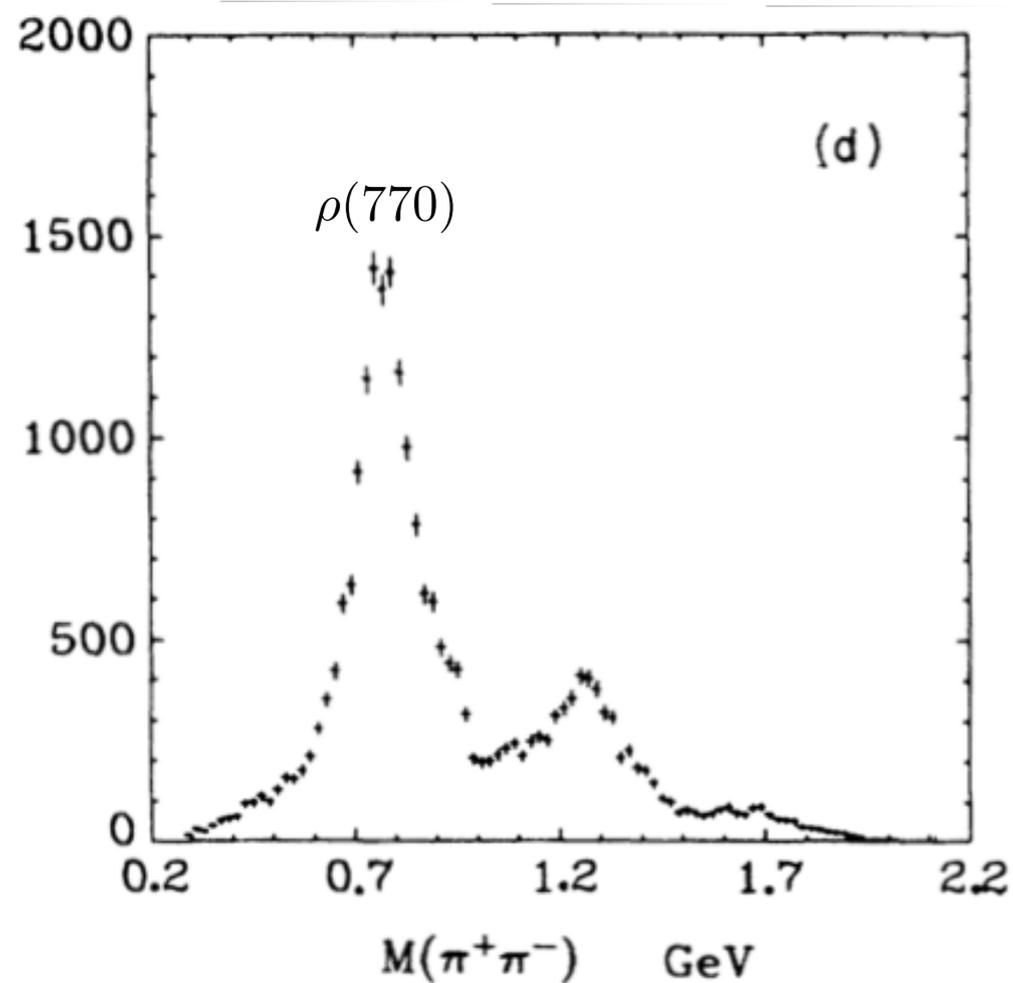
all are resonant enhancements in multiple channels

to understand these rigorously, we need coupled-channel analyses



Extracting resonance properties

excited states seen as resonant enhancements
in the scattering of lighter stable particles

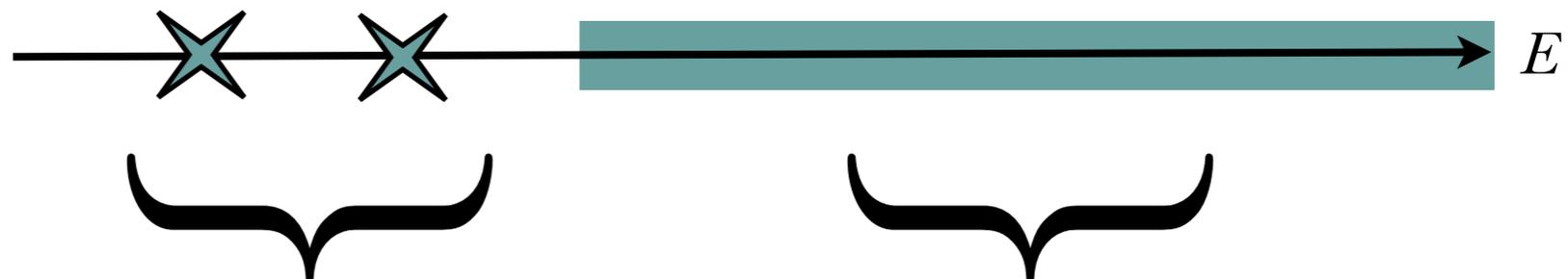


Extracting resonance properties

excited states seen as resonant enhancements
in the scattering of lighter stable particles



Infinite volume



Bound states

Meson-meson continuum

Finite volume

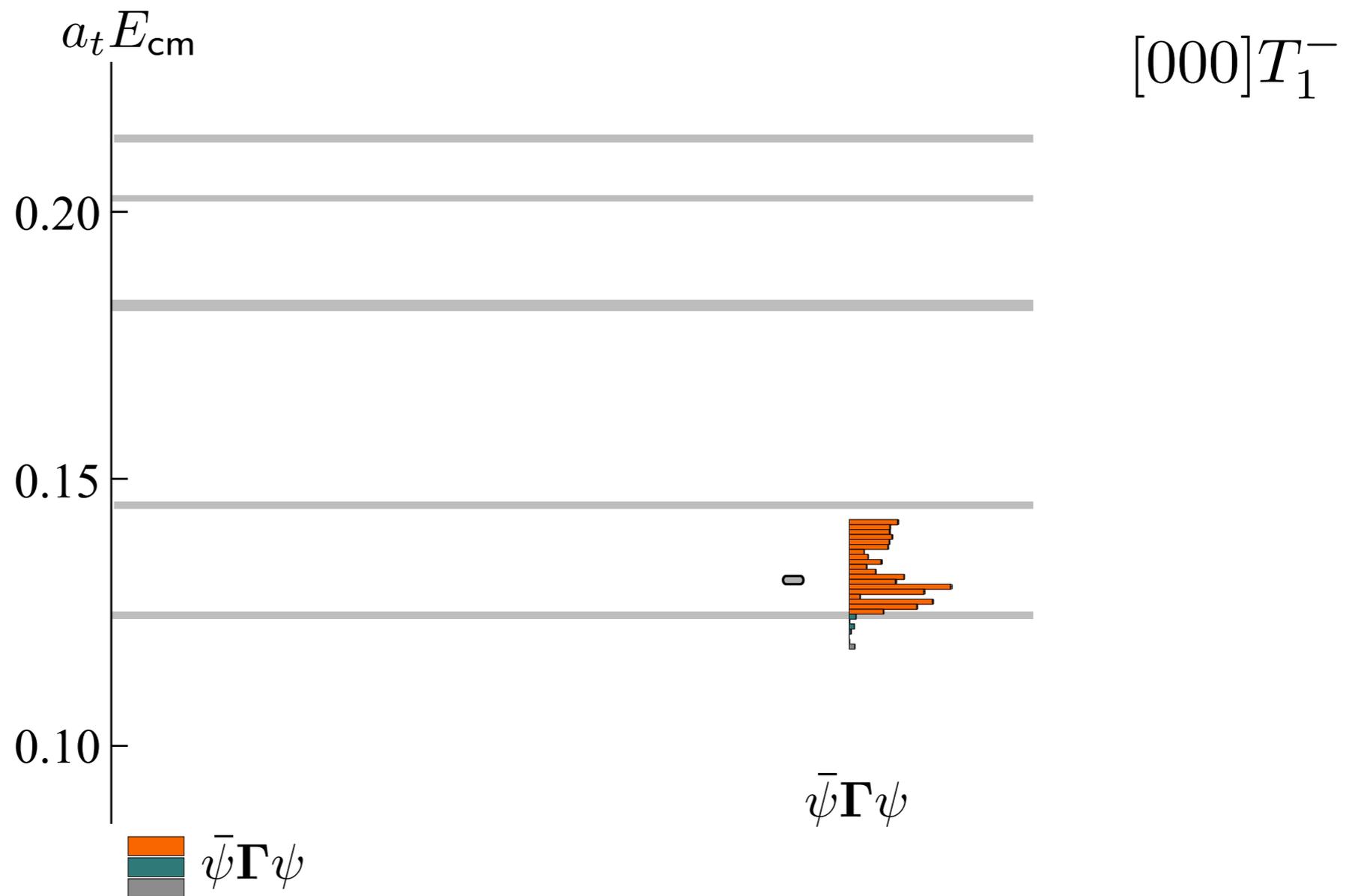


Extracting resonance properties

build a large basis of operators: $\mathcal{O}^\dagger \sim \bar{\psi} \Gamma \overleftrightarrow{D} \dots \overleftrightarrow{D} \psi$

compute large correlation matrices: $C_{ij}(t) = \langle 0 | \mathcal{O}_i(t) \mathcal{O}_j^\dagger(0) | 0 \rangle$

solve GEVP: $C_{ij}(t) v_j^n = \lambda_n(t, t_0) C_{ij}(t_0) v_j^n$

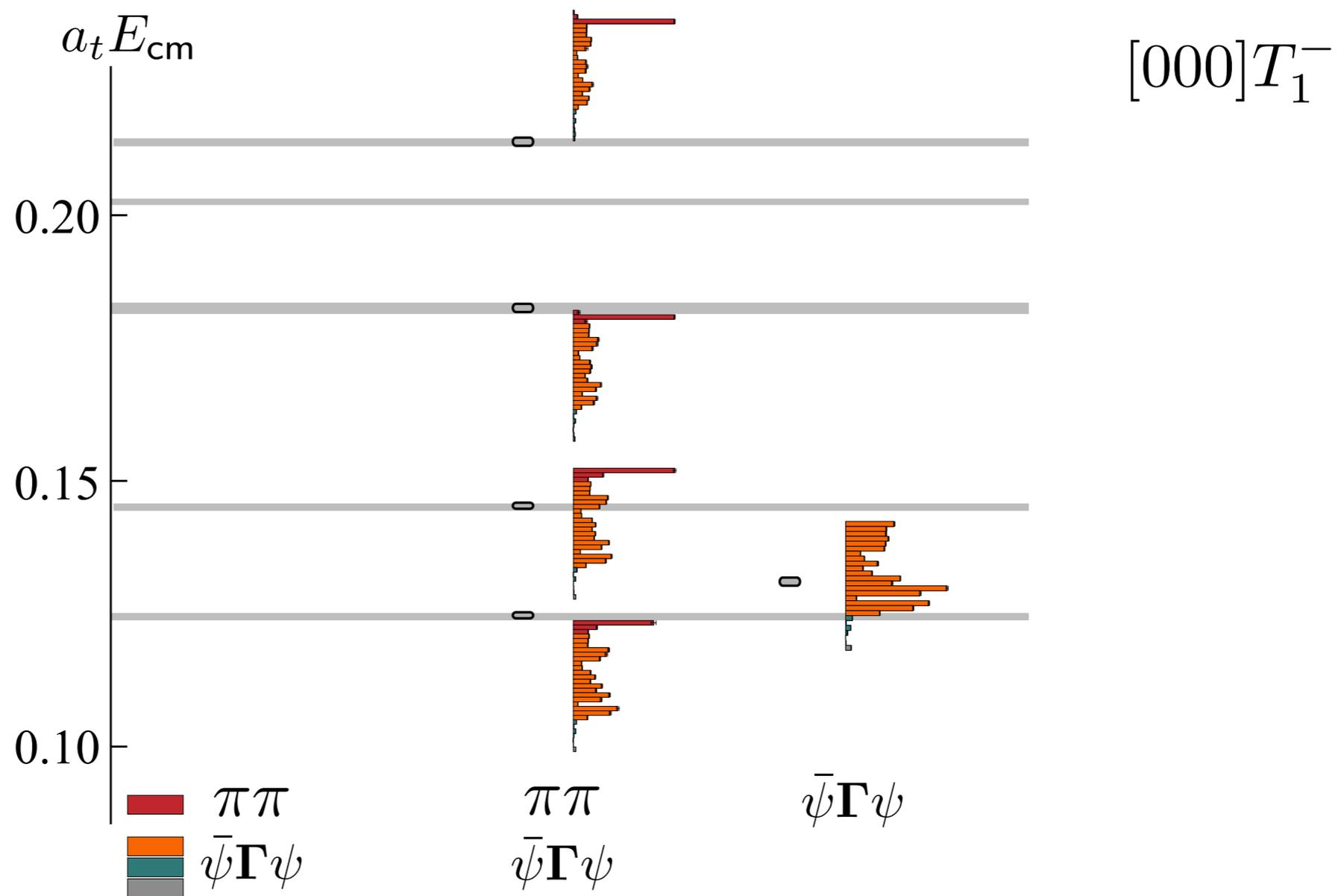


$m_\pi = 236 \text{ MeV}$

Extracting resonance properties

add in $\pi\pi$ operators using a variationally optimal pion $\pi^\dagger = \sum_i v_i^\pi \mathcal{O}_i^\dagger$

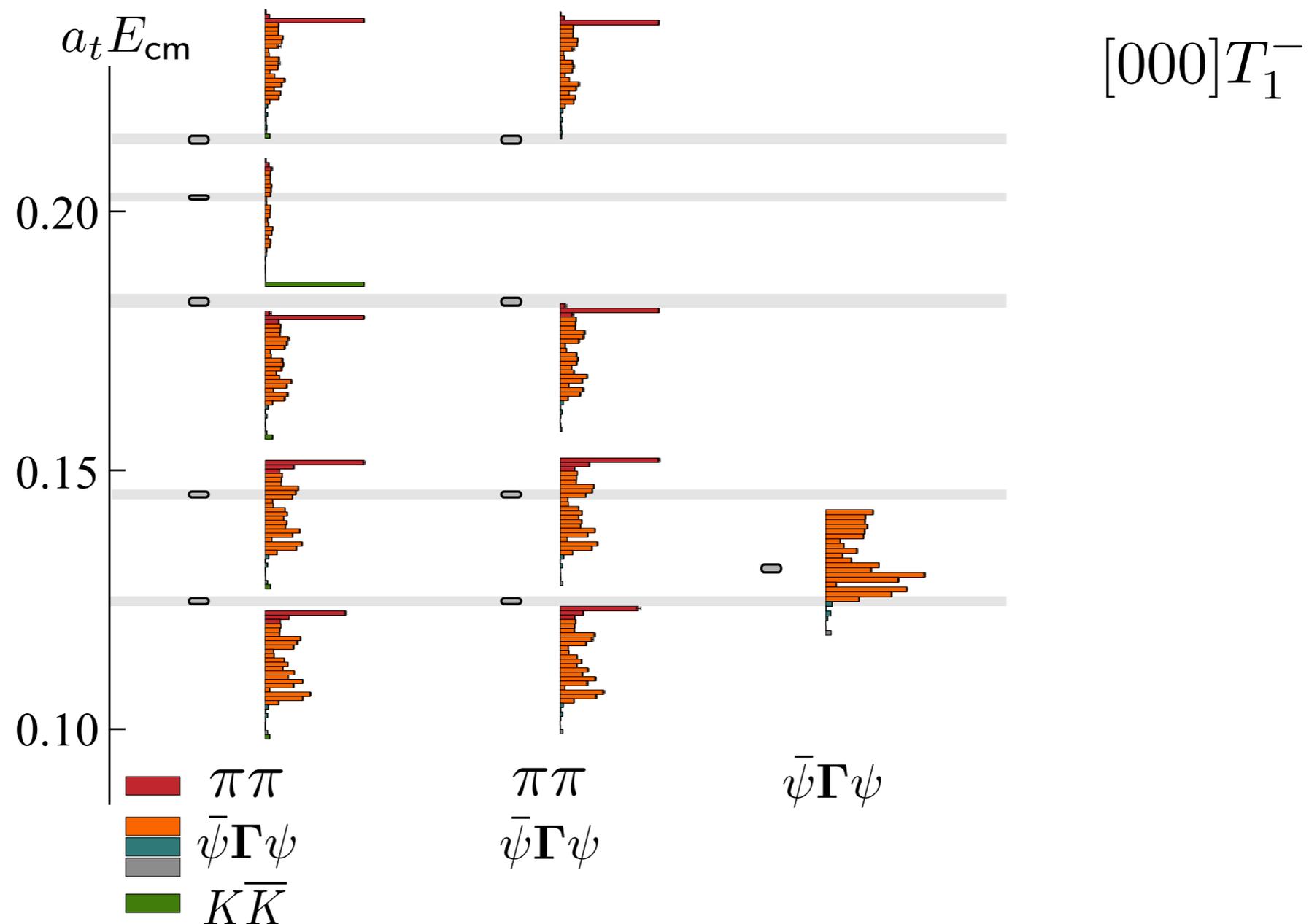
combine in pairs $(\pi\pi)^\dagger = \sum_{\vec{p}_1 + \vec{p}_2 = \vec{P}} \mathcal{C}(\vec{p}_1, \vec{p}_2) \pi^\dagger(\vec{p}_1) \pi^\dagger(\vec{p}_2)$



$m_\pi = 236 \text{ MeV}$

Extracting resonance properties

essential to have operators that overlap onto “meson” and “meson-meson” contributions to the physical spectrum

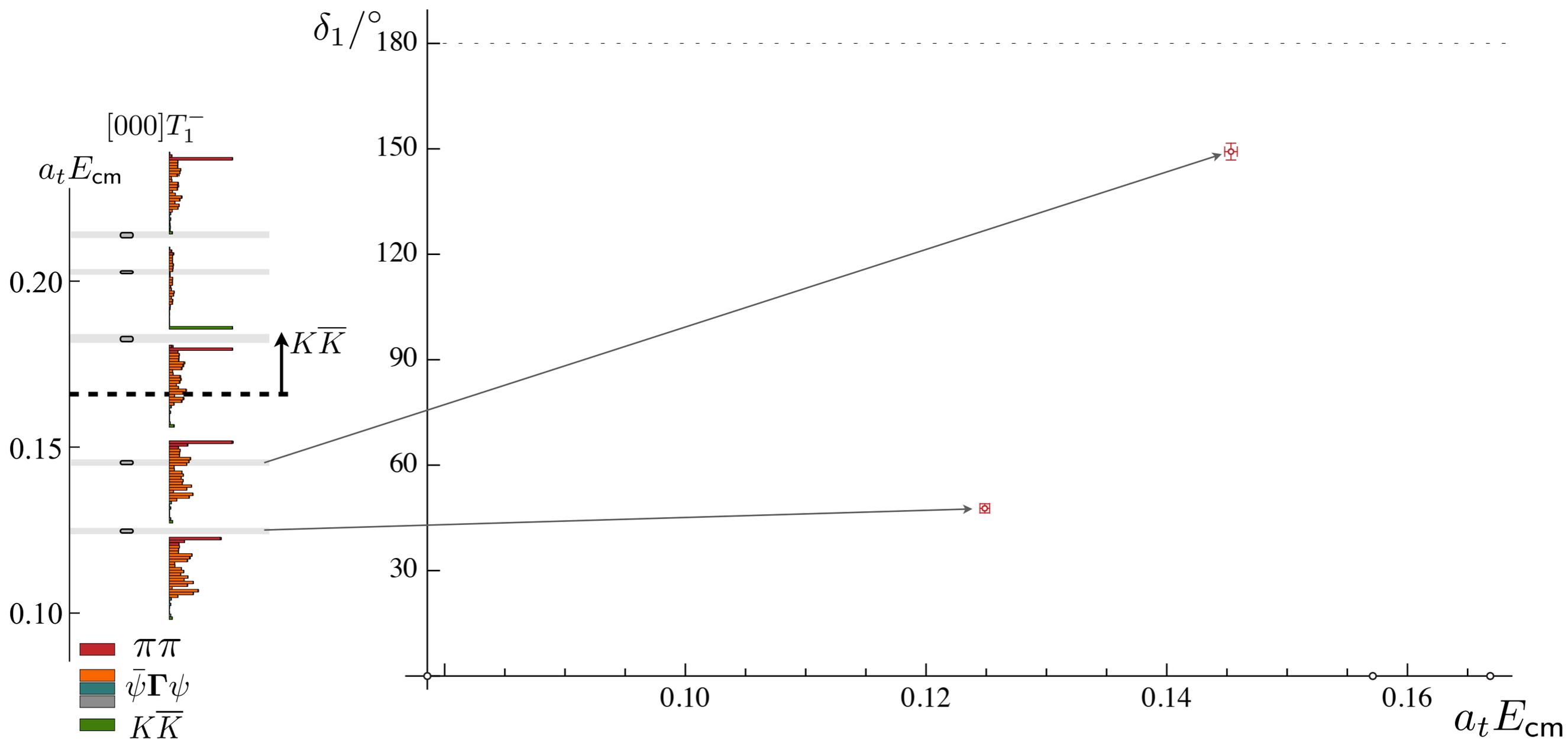


$$m_\pi = 236 \text{ MeV}$$

ρ resonance

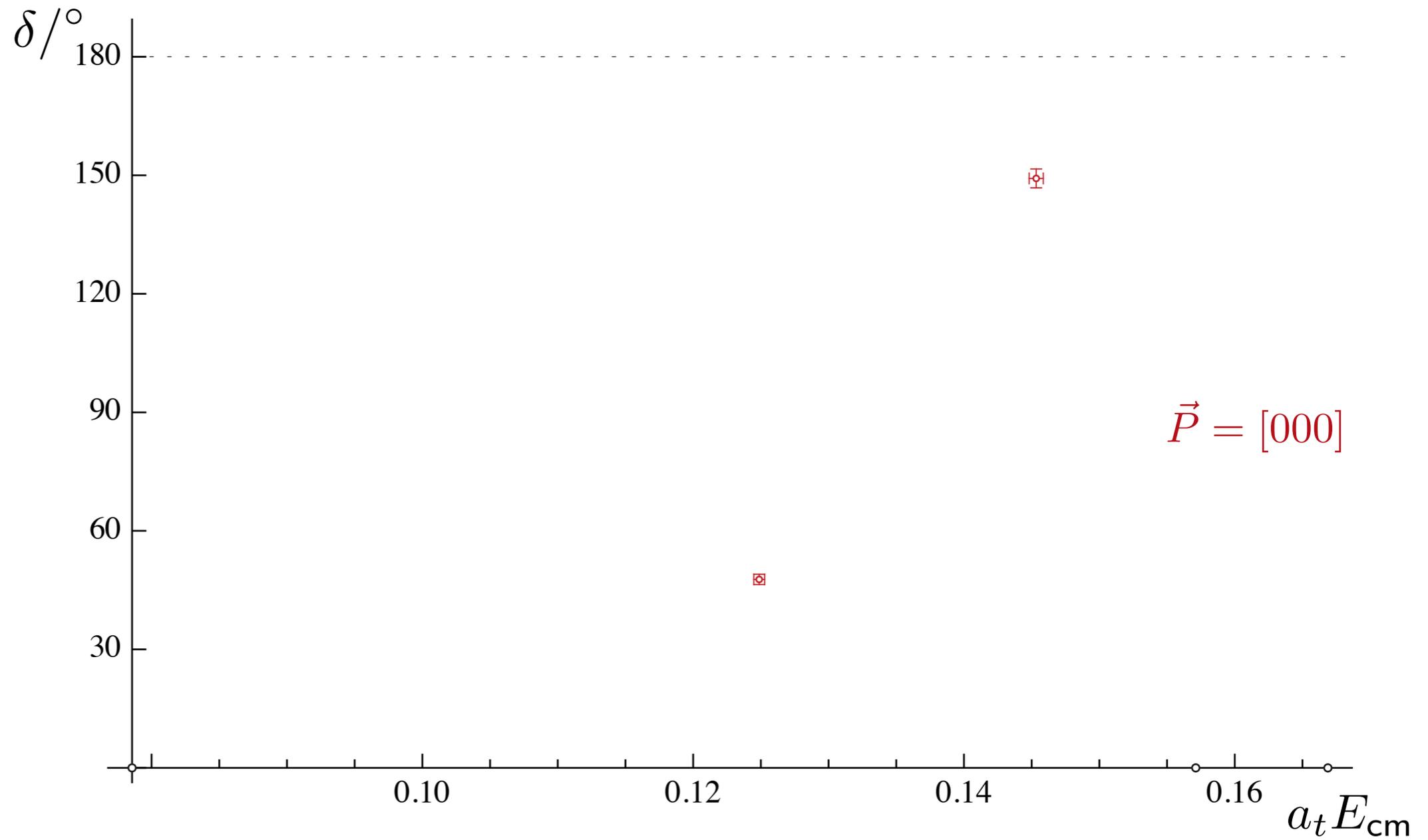
Phase shifts via the Lüscher method: $\tan \delta_1 = \frac{\pi^{3/2} q}{\mathcal{Z}_{00}(1; q^2)}$

$$\mathcal{Z}_{00}(1; q^2) = \sum_{n \in \mathbb{Z}^3} \frac{1}{|\vec{n}|^2 - q^2}$$



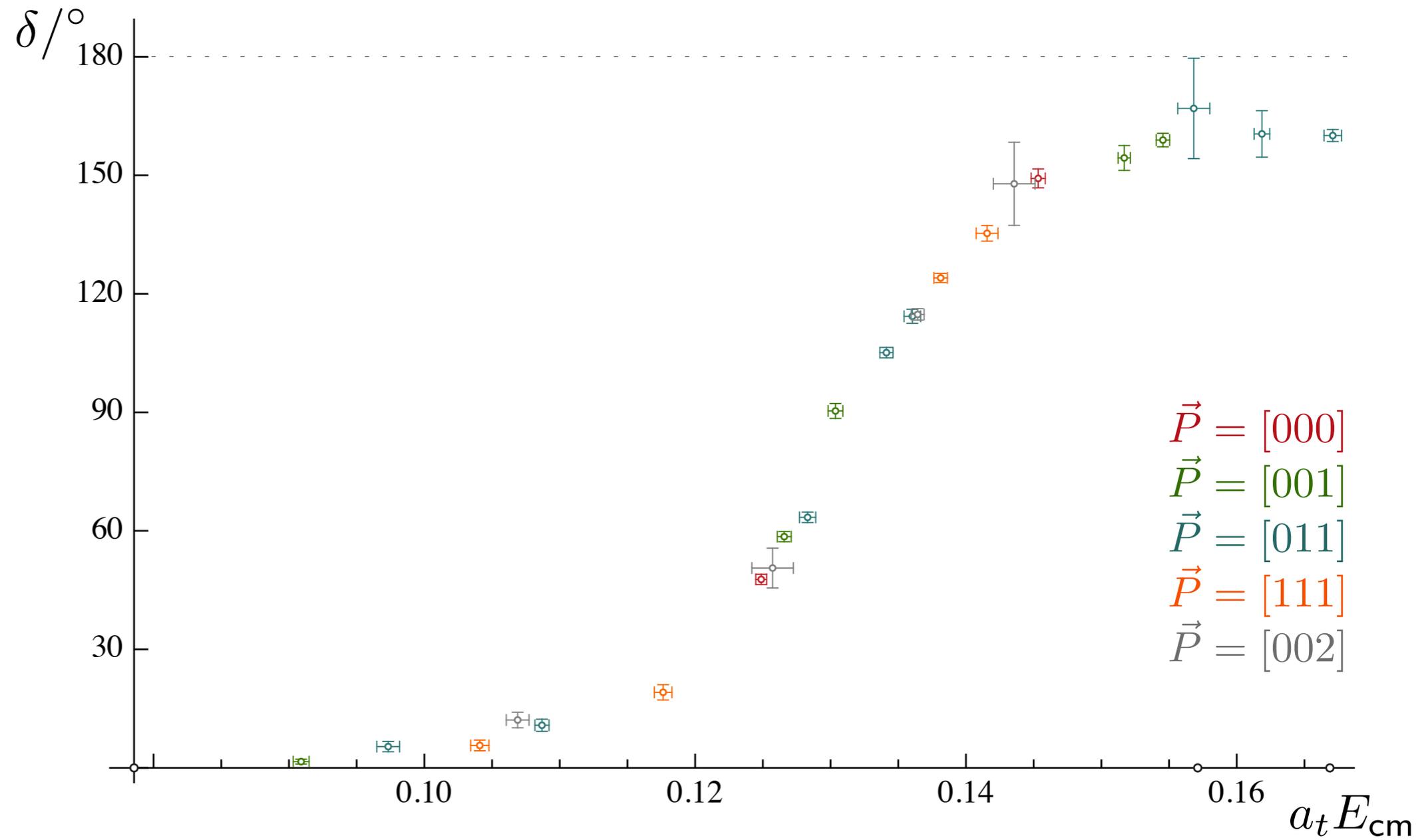
$m_\pi = 236 \text{ MeV}$

ρ resonance with moving frames



$m_\pi = 236 \text{ MeV}$

ρ resonance with moving frames

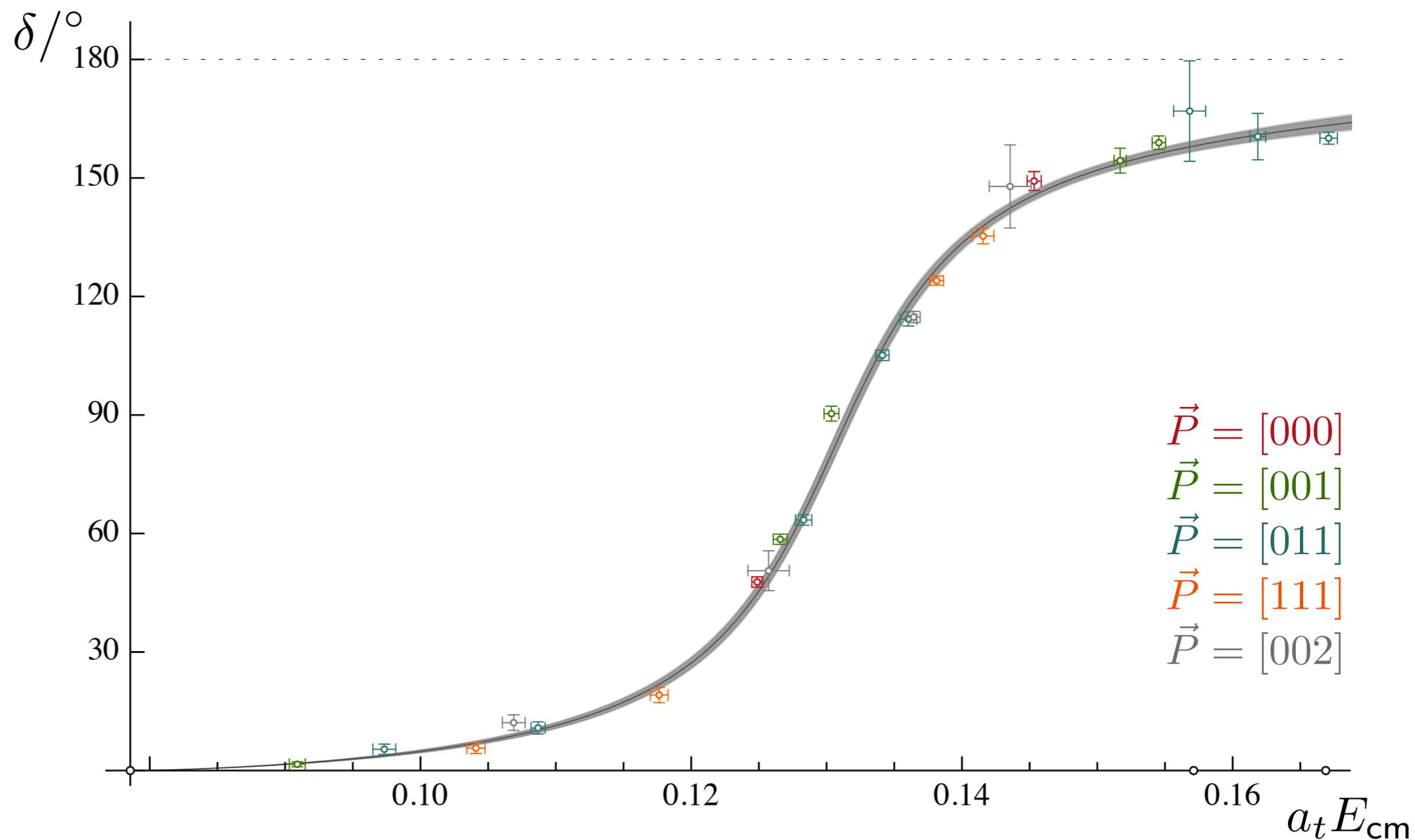


$m_\pi = 236 \text{ MeV}$

ρ resonance with moving frames

PRD 92 094502, arXiv:1507.02599

- for more see Antoni Woss Tuesday 26 Jul 2016 at 14:40



one volume, 22 energy levels... lots of constraint

$m_\pi = 236 \text{ MeV}$

Coupled-channel scattering

Direct extension of the elastic quantization condition derived by Lüscher

$$\det [\mathbf{1} + i\rho(E) \cdot t(E) \cdot (\mathbf{1} + i\mathcal{M}(E, L))] = 0$$

phase space

infinite volume scattering
 t -matrix

known finite-volume
functions

Many derivations, **all in agreement**:

He, Feng, Liu 2005 - two channel QM, strong coupling

Hansen & Sharpe 2012 - field theory, multiple two-body channels

Briceño & Davoudi 2012 - strongly-coupled Bethe-Salpeter amplitudes

Guo et al 2012 - Hamiltonian & Lippmann-Schwinger

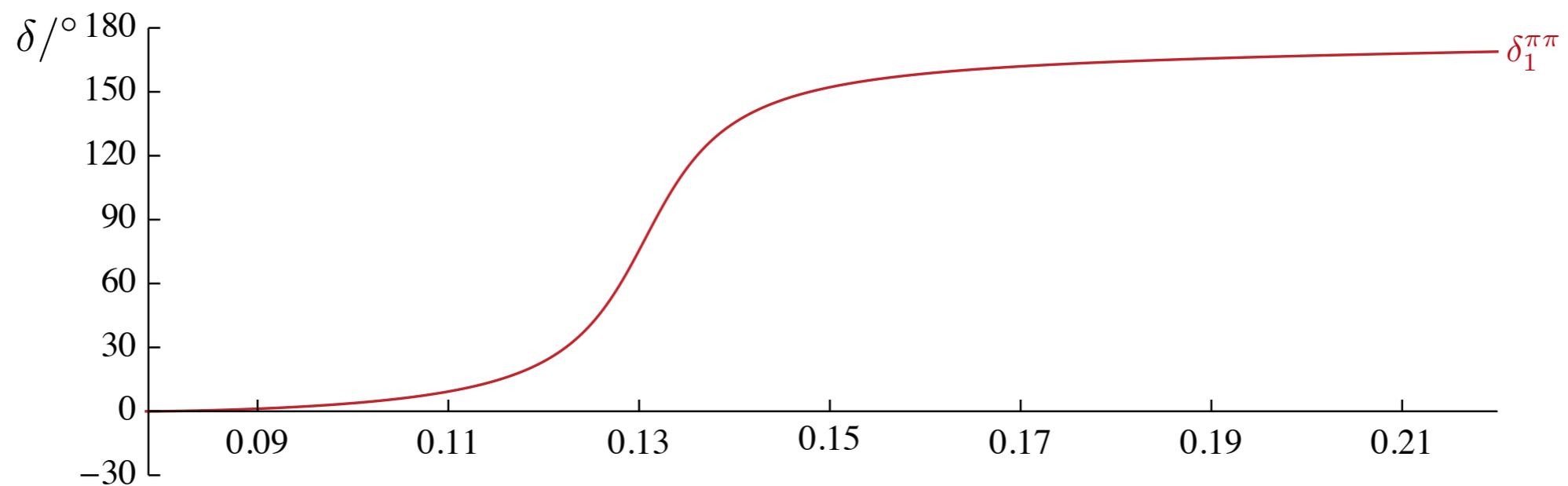
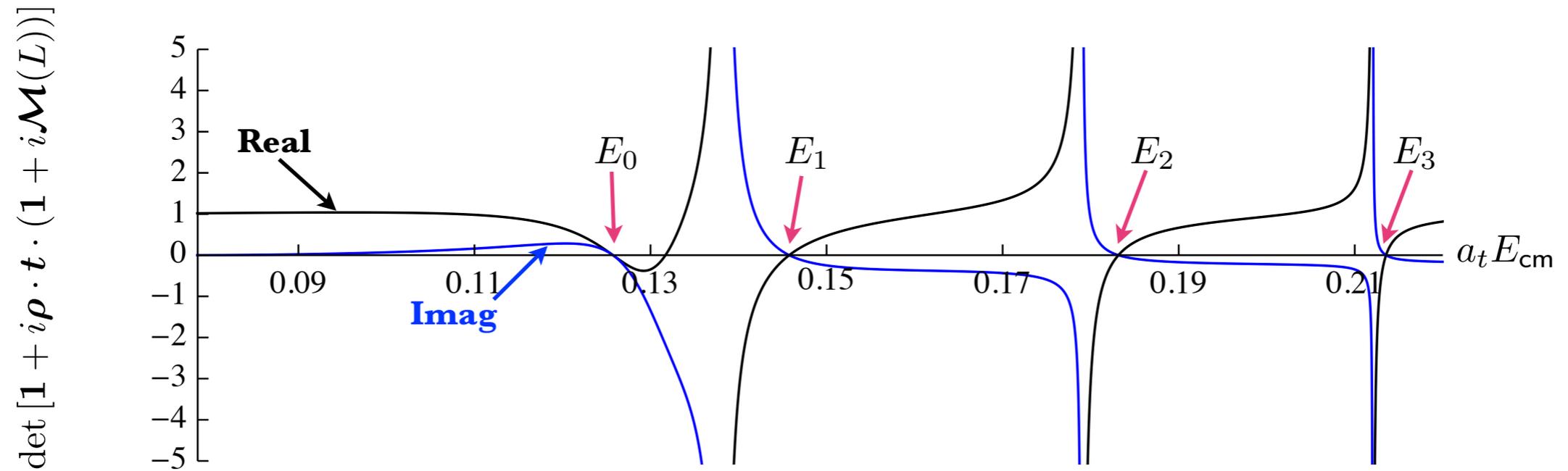
Also derivations in specific channels, or for a specific parameterization of the interactions like NREFT, chiral PT, Finite Volume Hamiltonian, etc.

Briceño 2014 - Generalised to scattering of particles with non-zero spin, and spin- $1/2$.

Significant steps towards a general 3-body quantization condition have been made
- see Stephen Sharpe on Tuesday 26 Jul 2016 at 15:40 for the latest

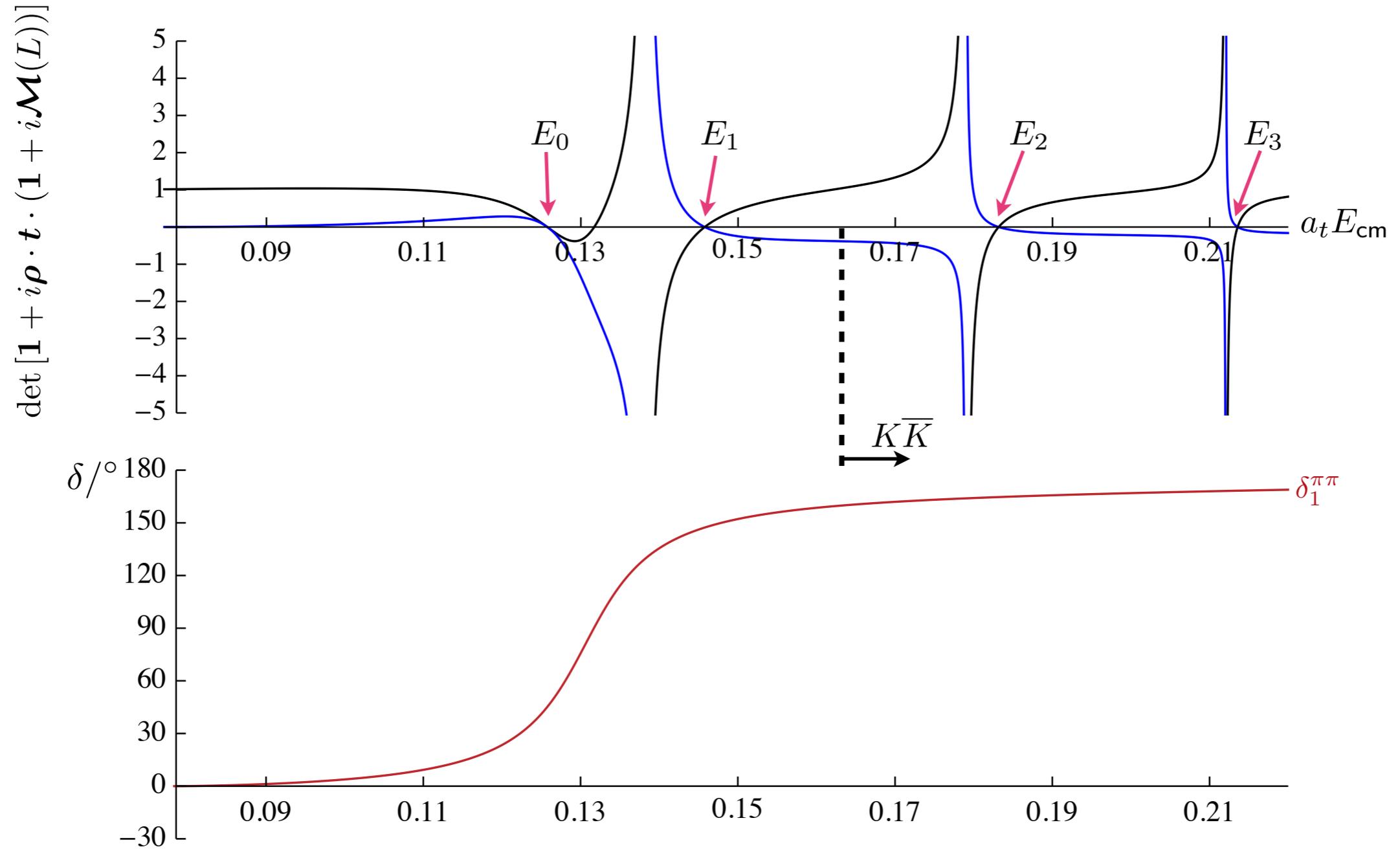
Determinant

$$t = (\pi\pi \rightarrow \pi\pi)$$



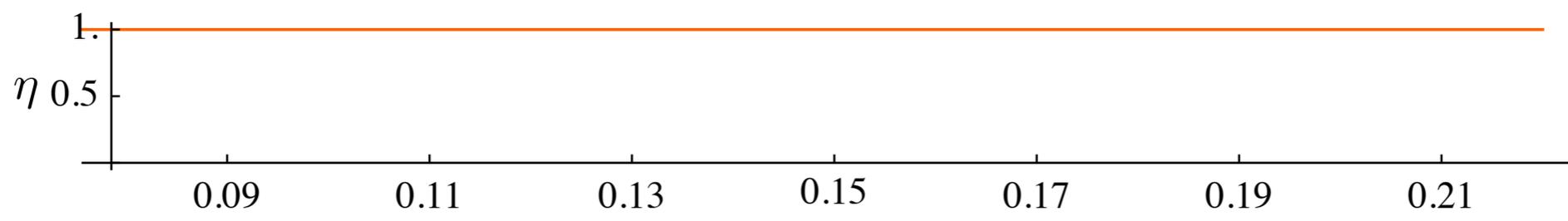
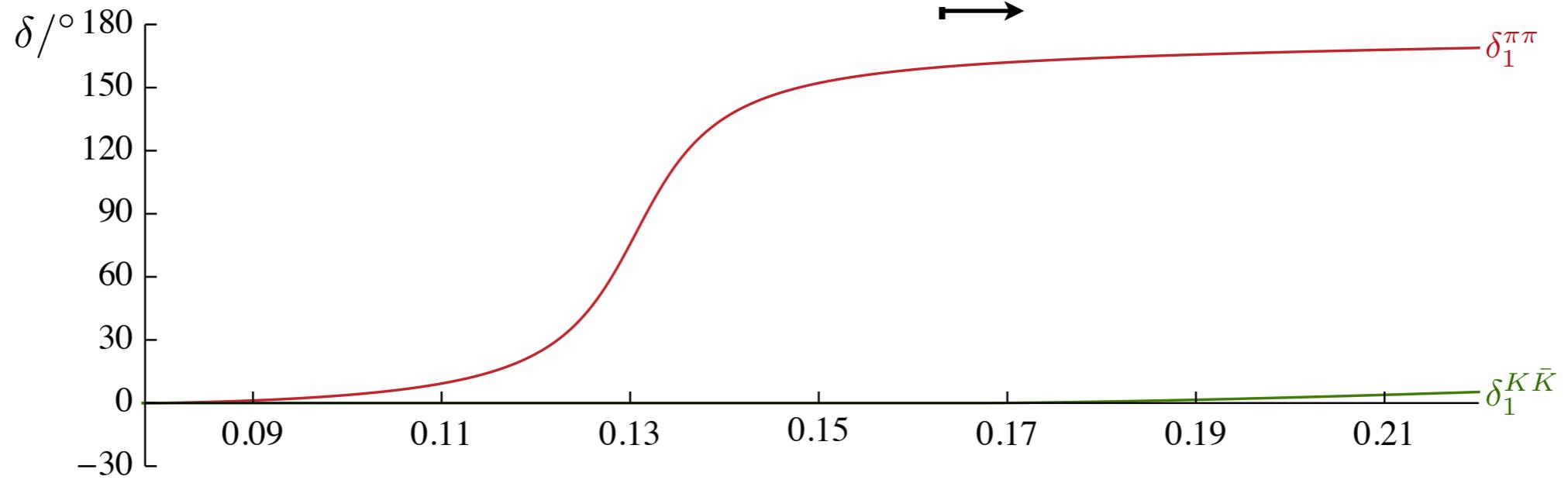
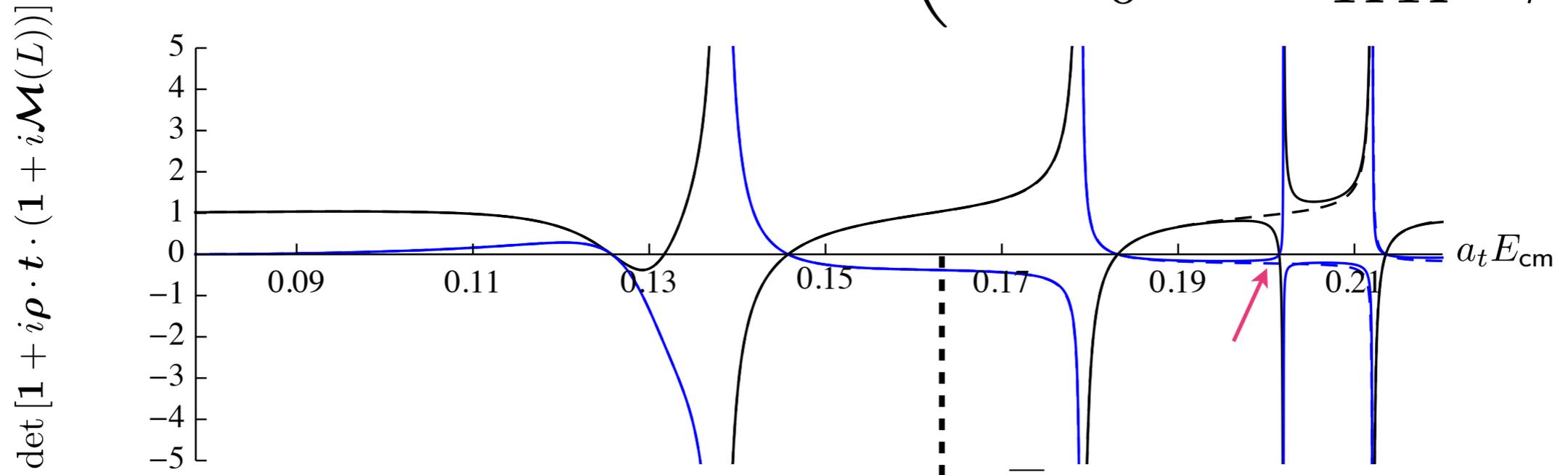
Determinant

$$t = (\pi\pi \rightarrow \pi\pi)$$



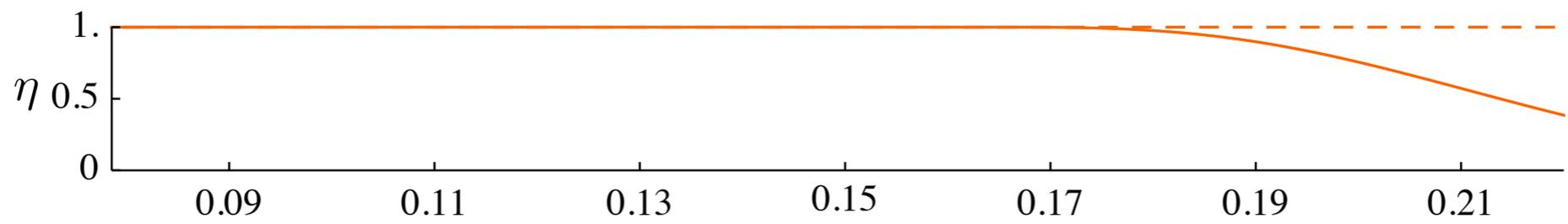
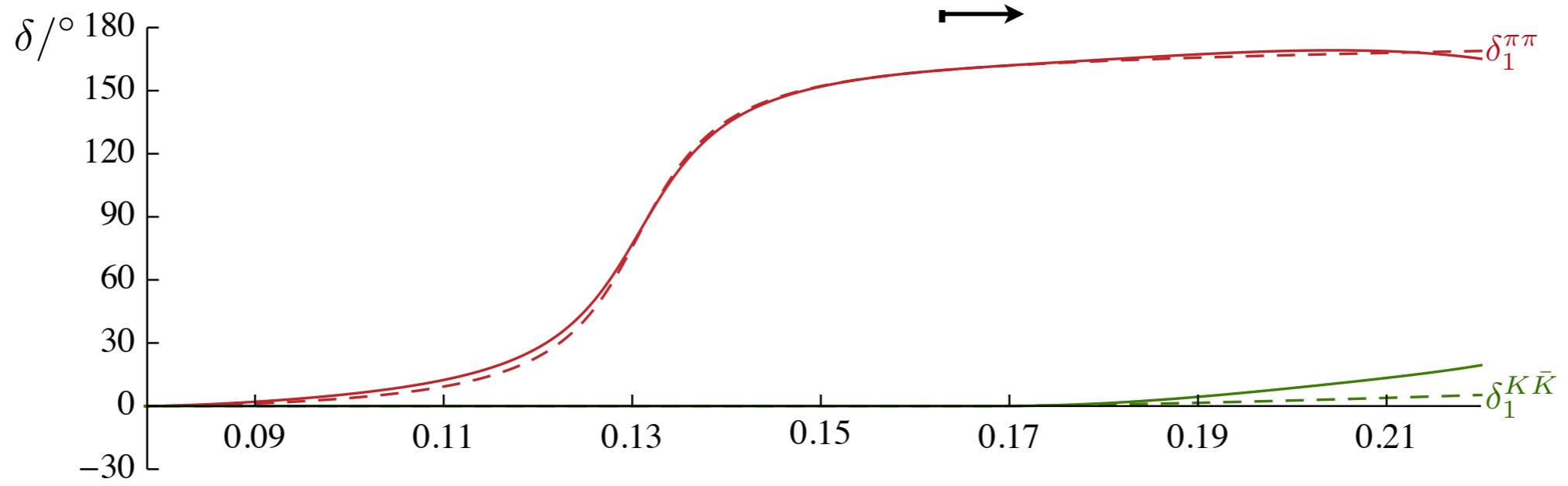
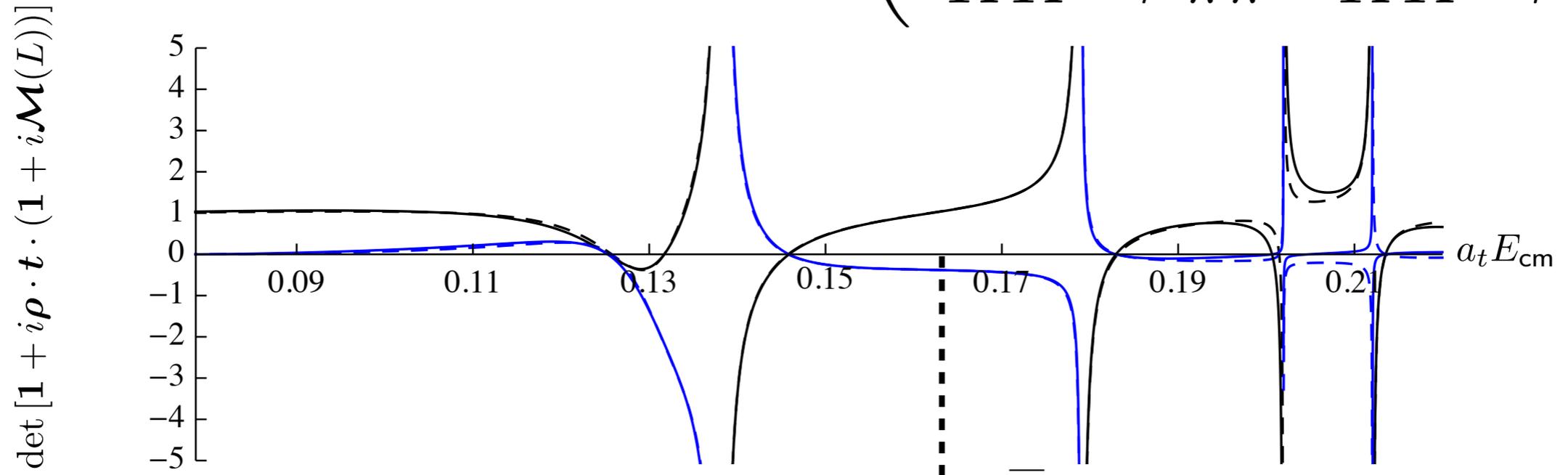
Determinant

$$t = \begin{pmatrix} \pi\pi \rightarrow \pi\pi & 0 \\ 0 & K\bar{K} \rightarrow K\bar{K} \end{pmatrix}$$



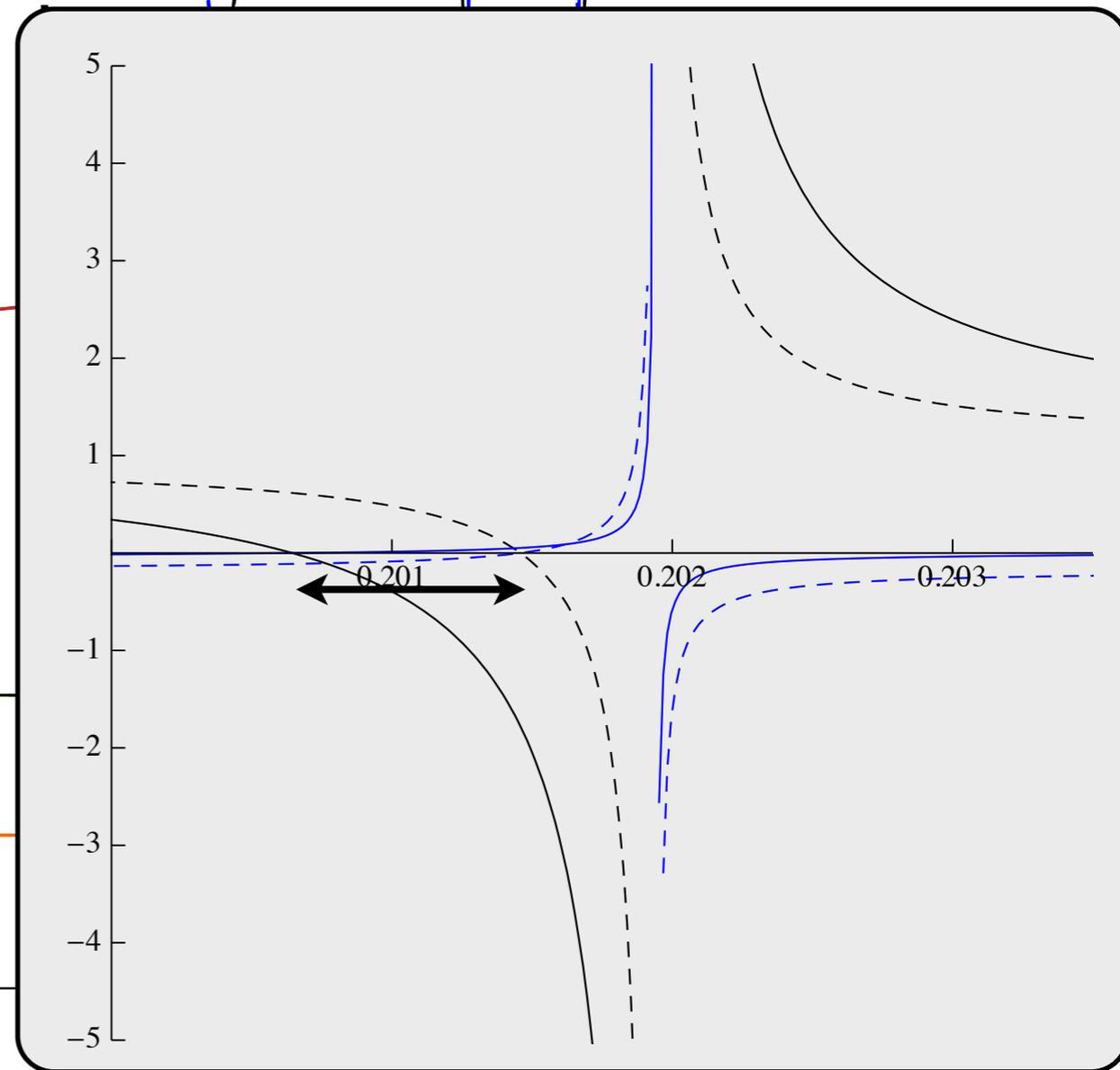
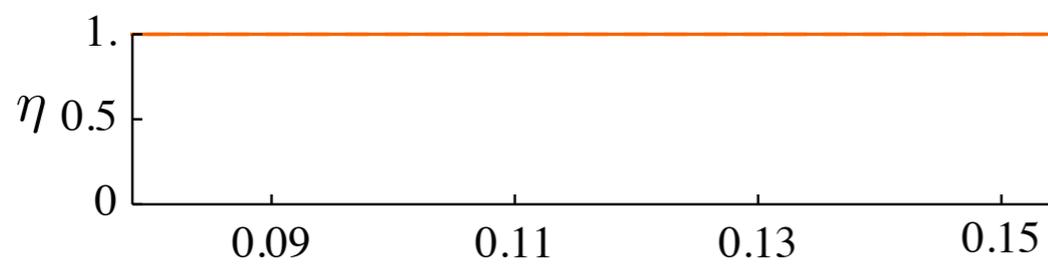
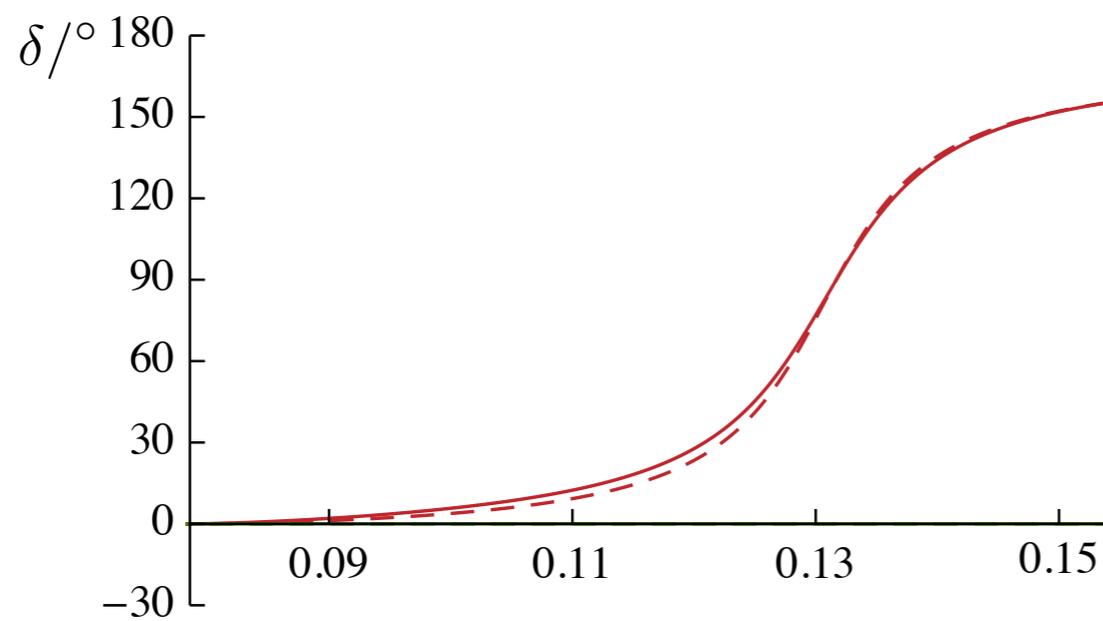
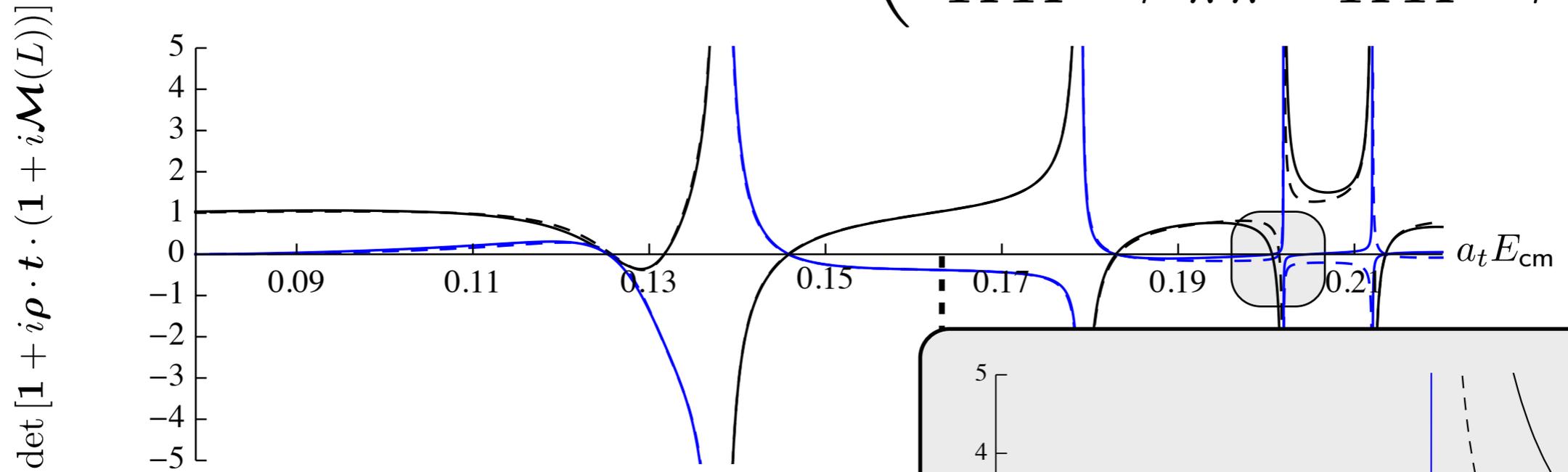
Determinant

$$t = \begin{pmatrix} \pi\pi \rightarrow \pi\pi & \pi\pi \rightarrow K\bar{K} \\ K\bar{K} \rightarrow \pi\pi & K\bar{K} \rightarrow K\bar{K} \end{pmatrix}$$



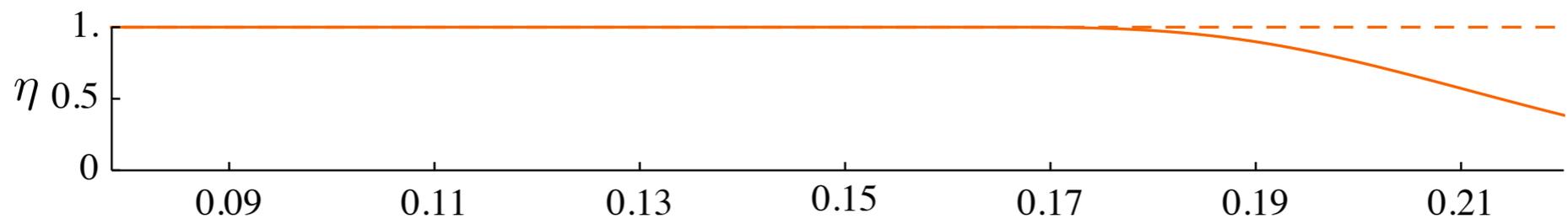
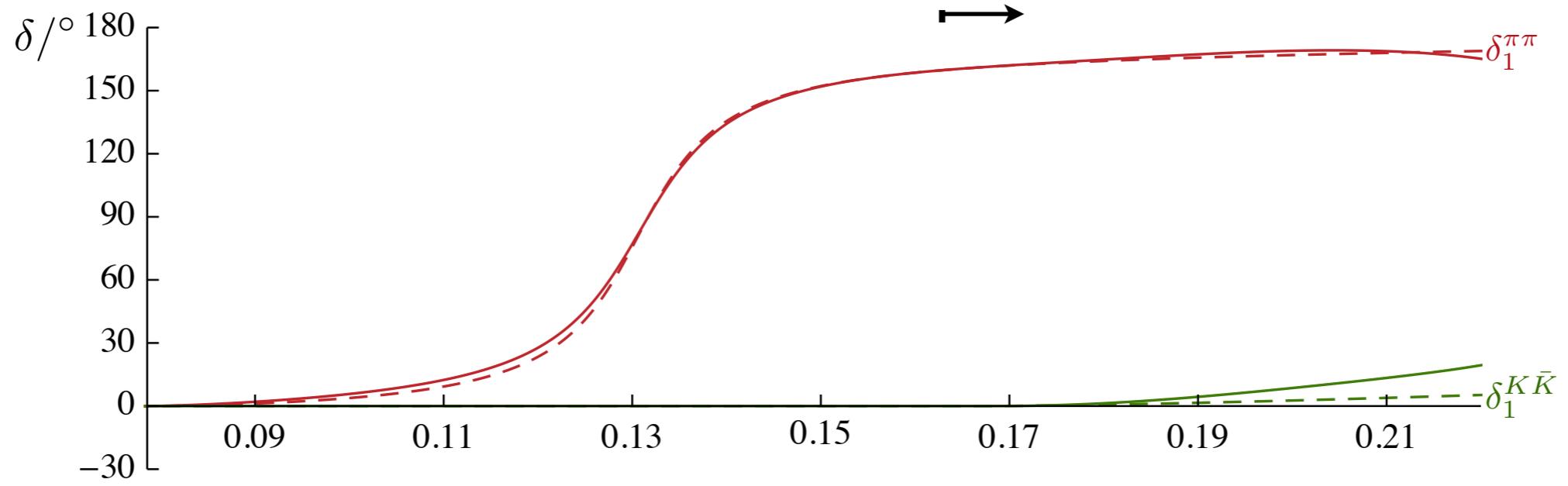
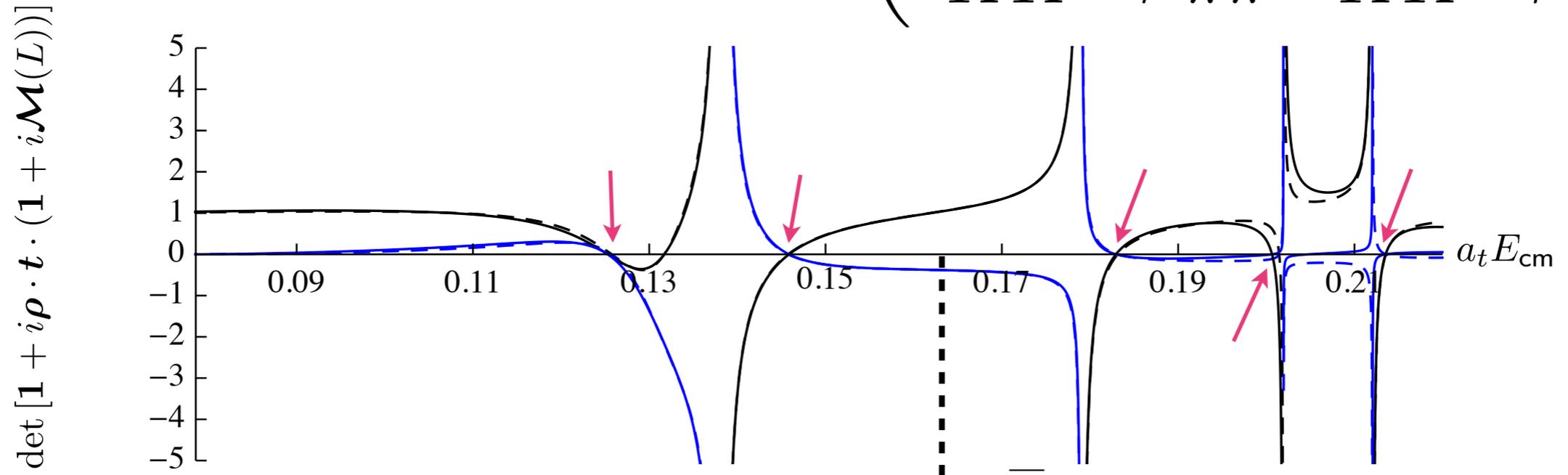
Determinant

$$t = \begin{pmatrix} \pi\pi \rightarrow \pi\pi & \pi\pi \rightarrow K\bar{K} \\ K\bar{K} \rightarrow \pi\pi & K\bar{K} \rightarrow K\bar{K} \end{pmatrix}$$



Determinant

$$t = \begin{pmatrix} \pi\pi \rightarrow \pi\pi & \pi\pi \rightarrow K\bar{K} \\ K\bar{K} \rightarrow \pi\pi & K\bar{K} \rightarrow K\bar{K} \end{pmatrix}$$



Amplitude parameterization

$$\mathbf{t} = \begin{pmatrix} \pi\pi \rightarrow \pi\pi & \pi\pi \rightarrow K\bar{K} \\ K\bar{K} \rightarrow \pi\pi & K\bar{K} \rightarrow K\bar{K} \end{pmatrix}$$

$$\det [\mathbf{1} + i\rho(E) \cdot \mathbf{t}(E) \cdot (\mathbf{1} + i\mathcal{M}(E, L))] = 0$$

determinant condition:

- several unknowns at each value of energy
- energy levels typically do not coincide
- underconstrained problem for a single energy

one solution: use energy dependent parameterizations

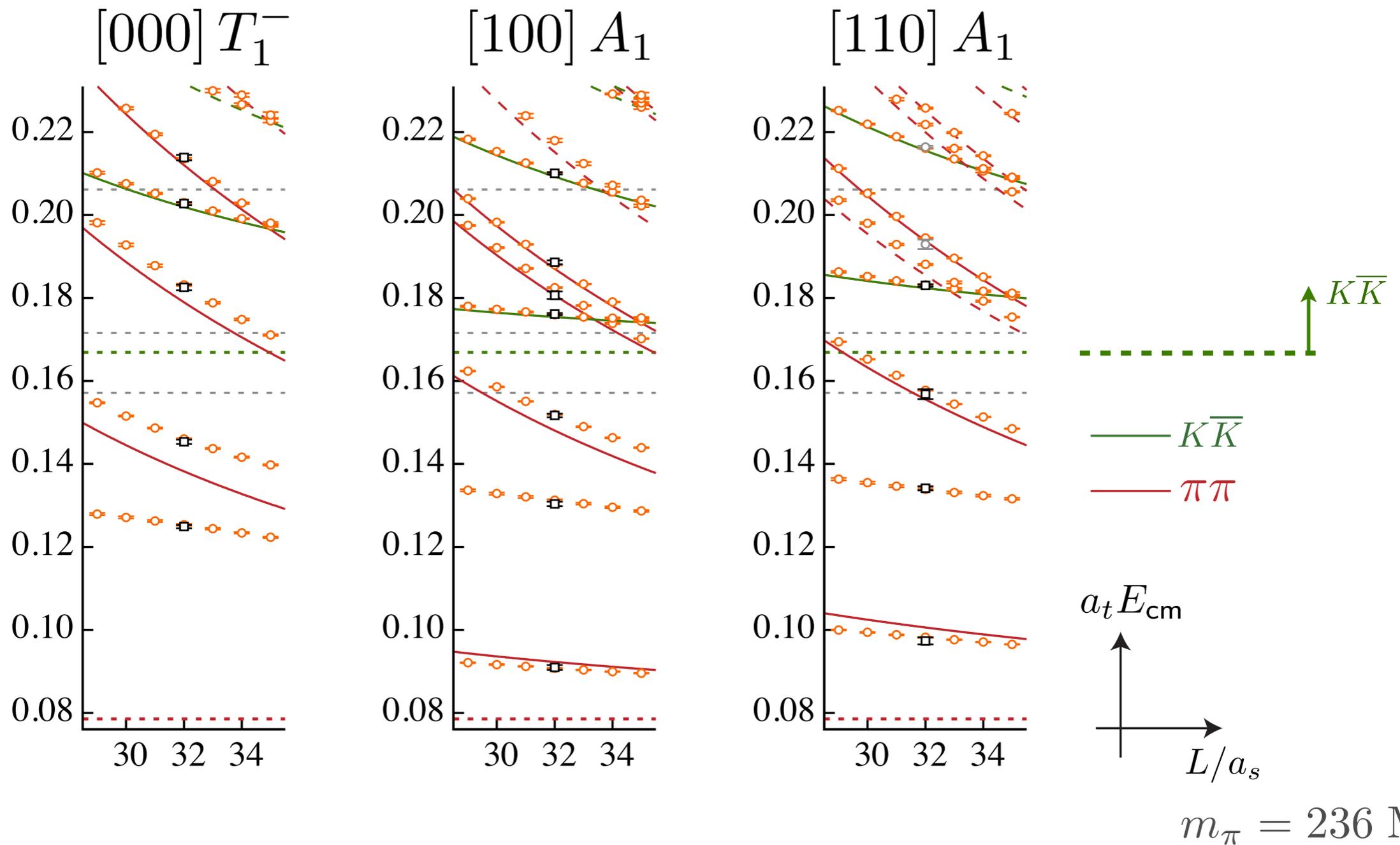
- Constrained problem when #(energy levels) > #(parameters)
- Essential amplitudes respect unitarity of the S-matrix

$$\mathbf{S}^\dagger \mathbf{S} = \mathbf{1} \quad \rightarrow \quad \text{Im } \mathbf{t}^{-1} = -\rho \quad \rho_{ij} = \delta_{ij} \frac{2k_i}{E_{\text{cm}}}$$

K-matrix approach:

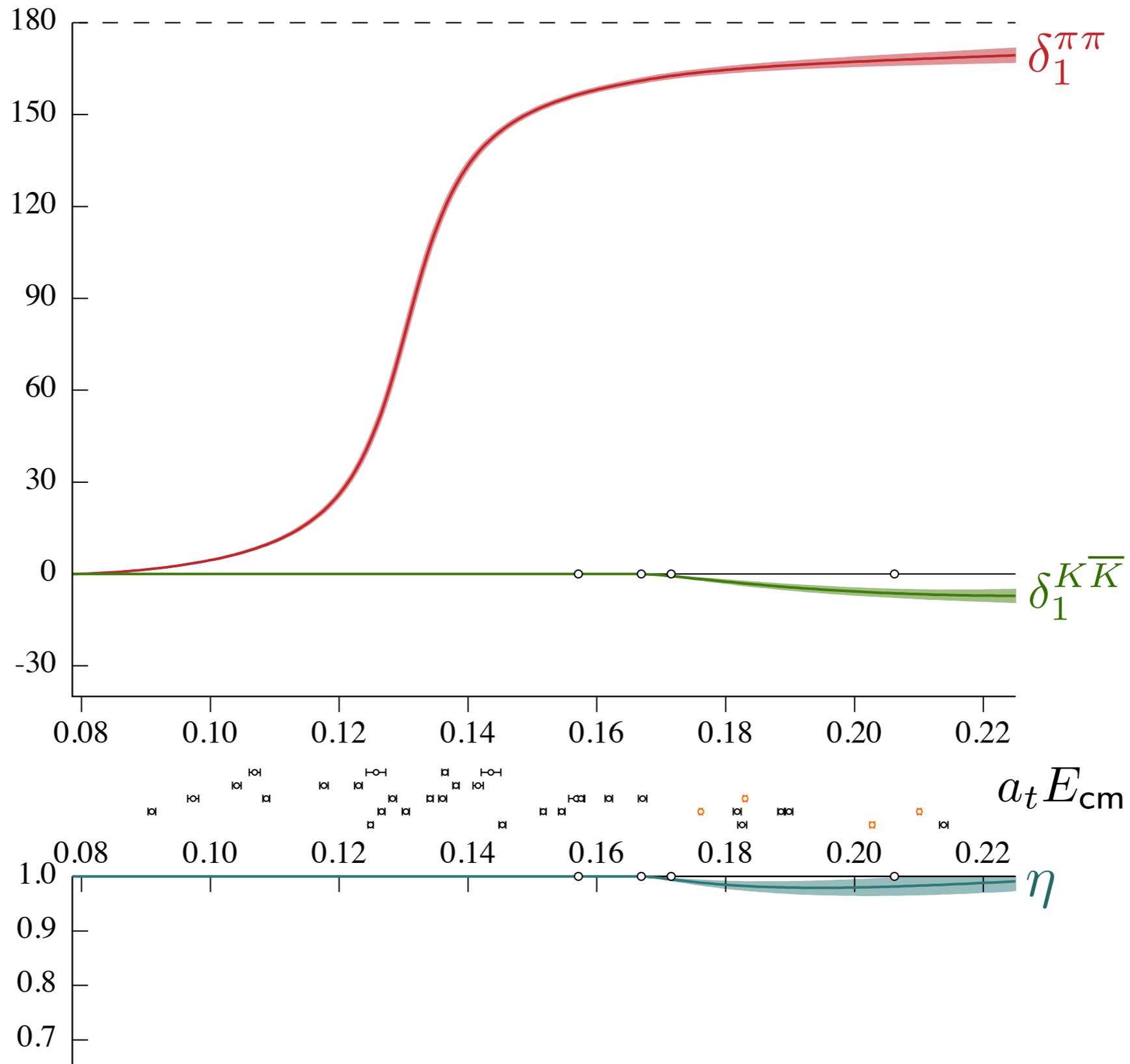
$$\mathbf{t}^{-1} = \mathbf{K}^{-1} - i\rho \quad \text{e.g.: } K_{ij} = \frac{g_i g_j}{m^2 - s} + \gamma_{ij}$$

ρ resonance into the coupled-channel region



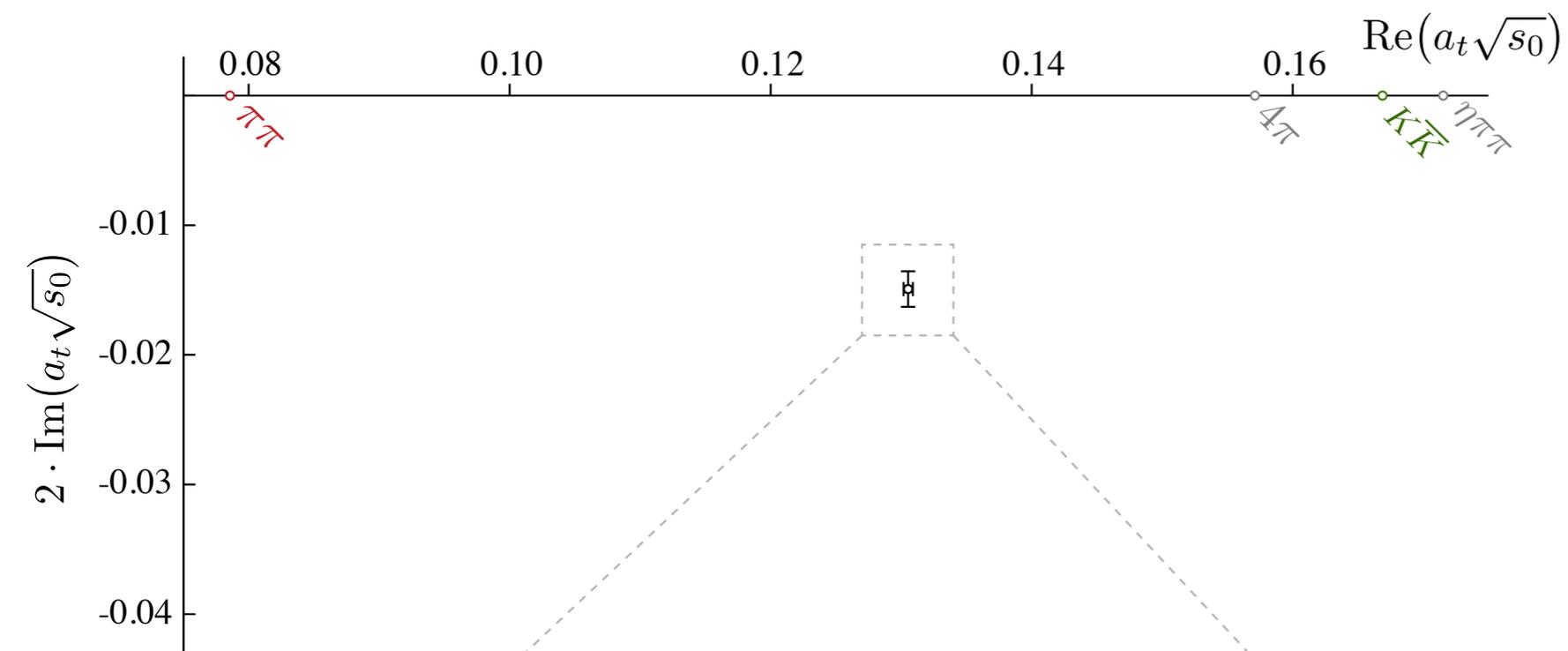
ρ resonance into the coupled-channel region

PRD 92 094502, arXiv:1507.02599



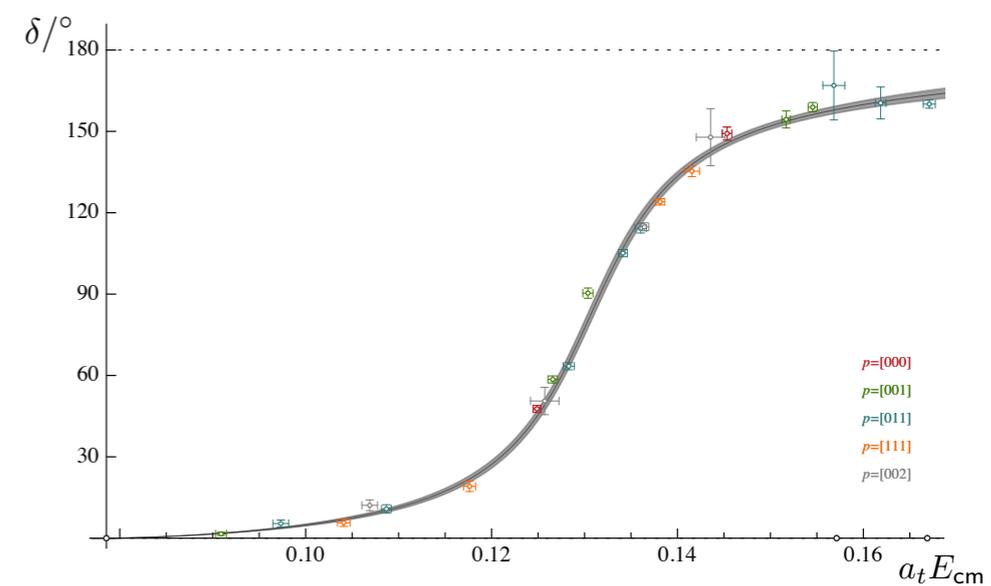
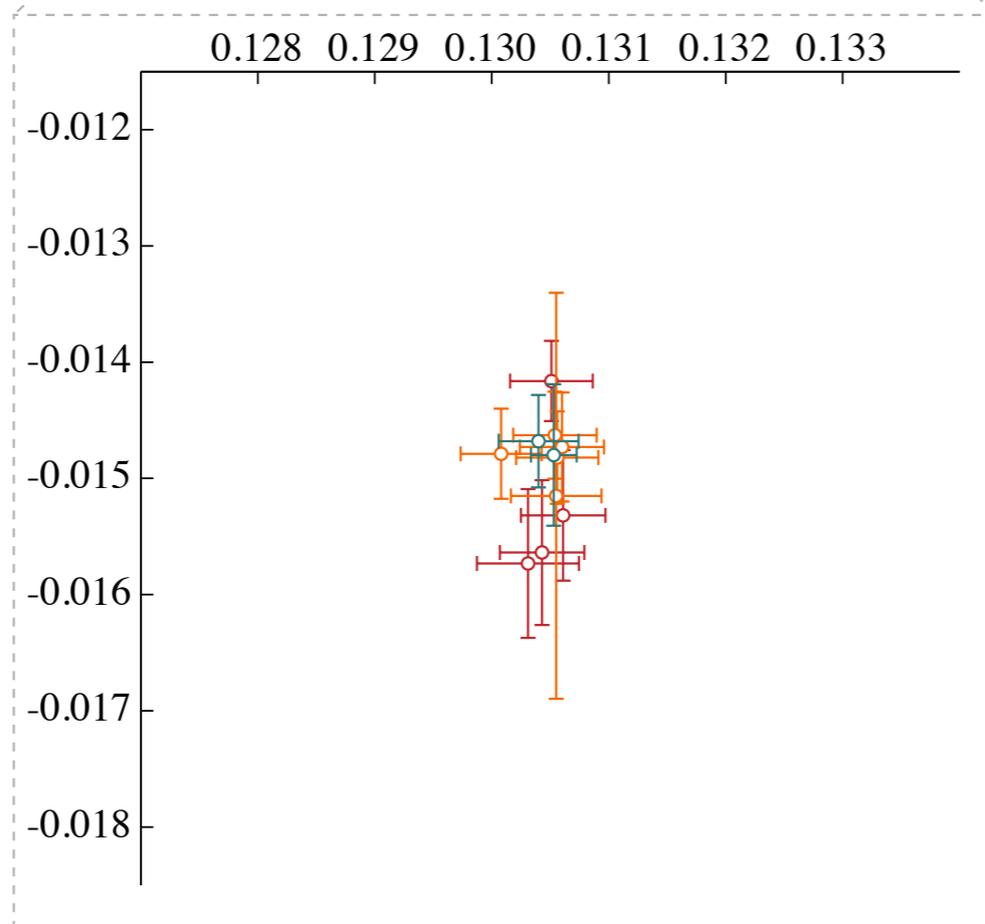
$m_\pi = 236 \text{ MeV}$

ρ resonance pole



near a pole:

$$t_{ij} \sim \frac{C_i C_j}{s_0 - s}$$



$$m_\pi = 236 \text{ MeV}$$

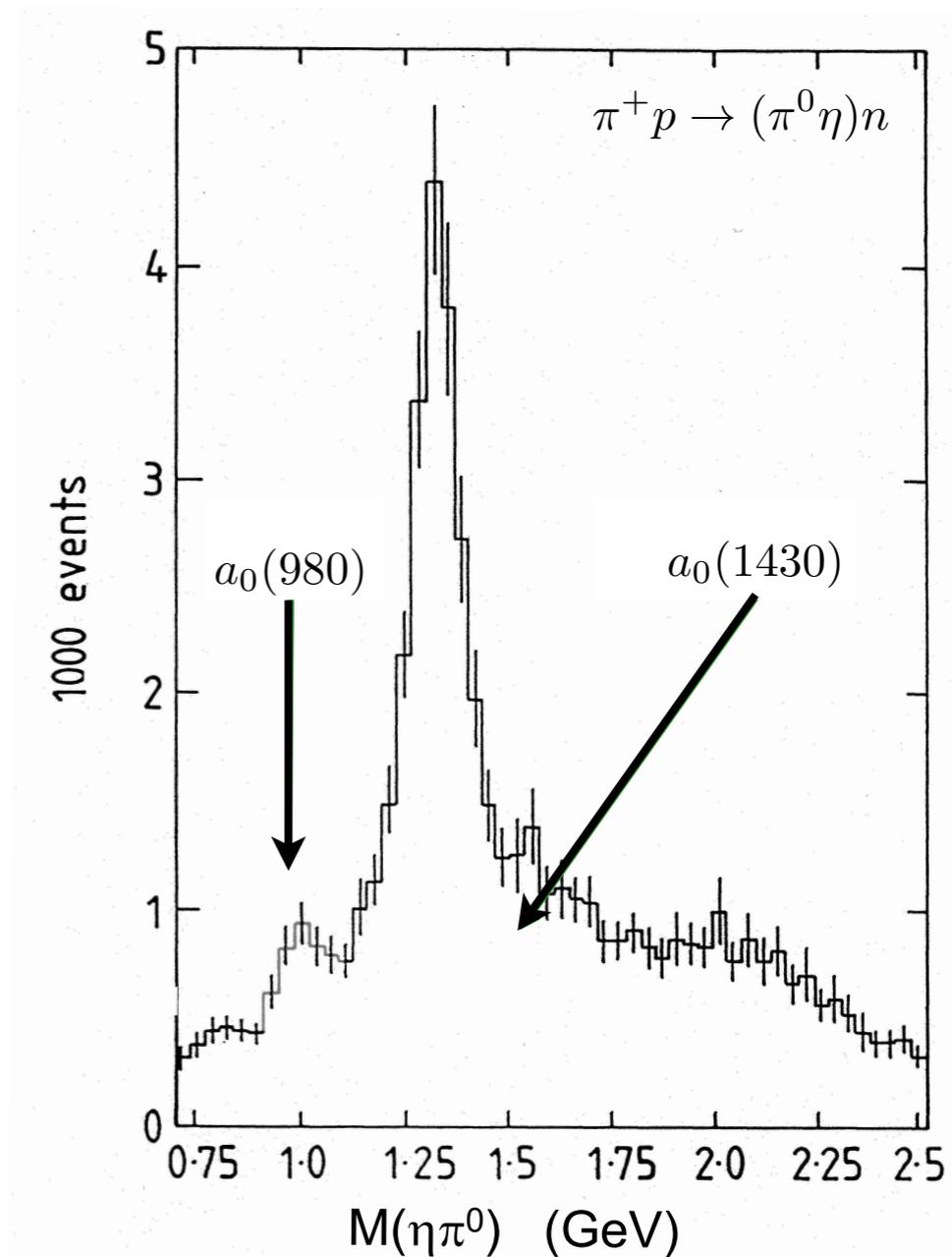
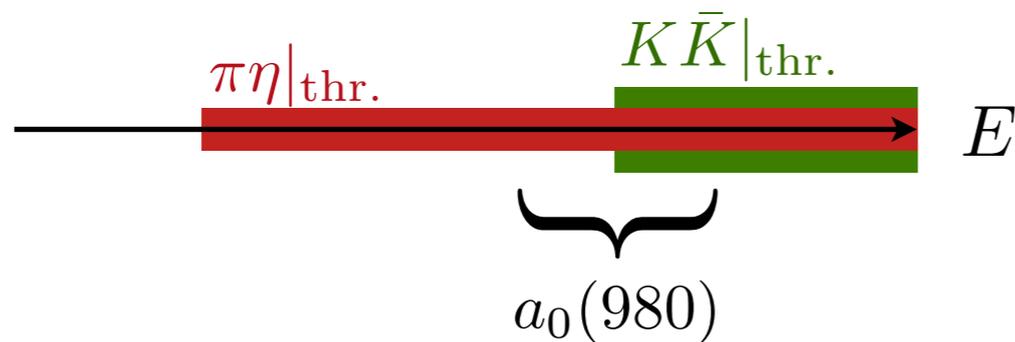
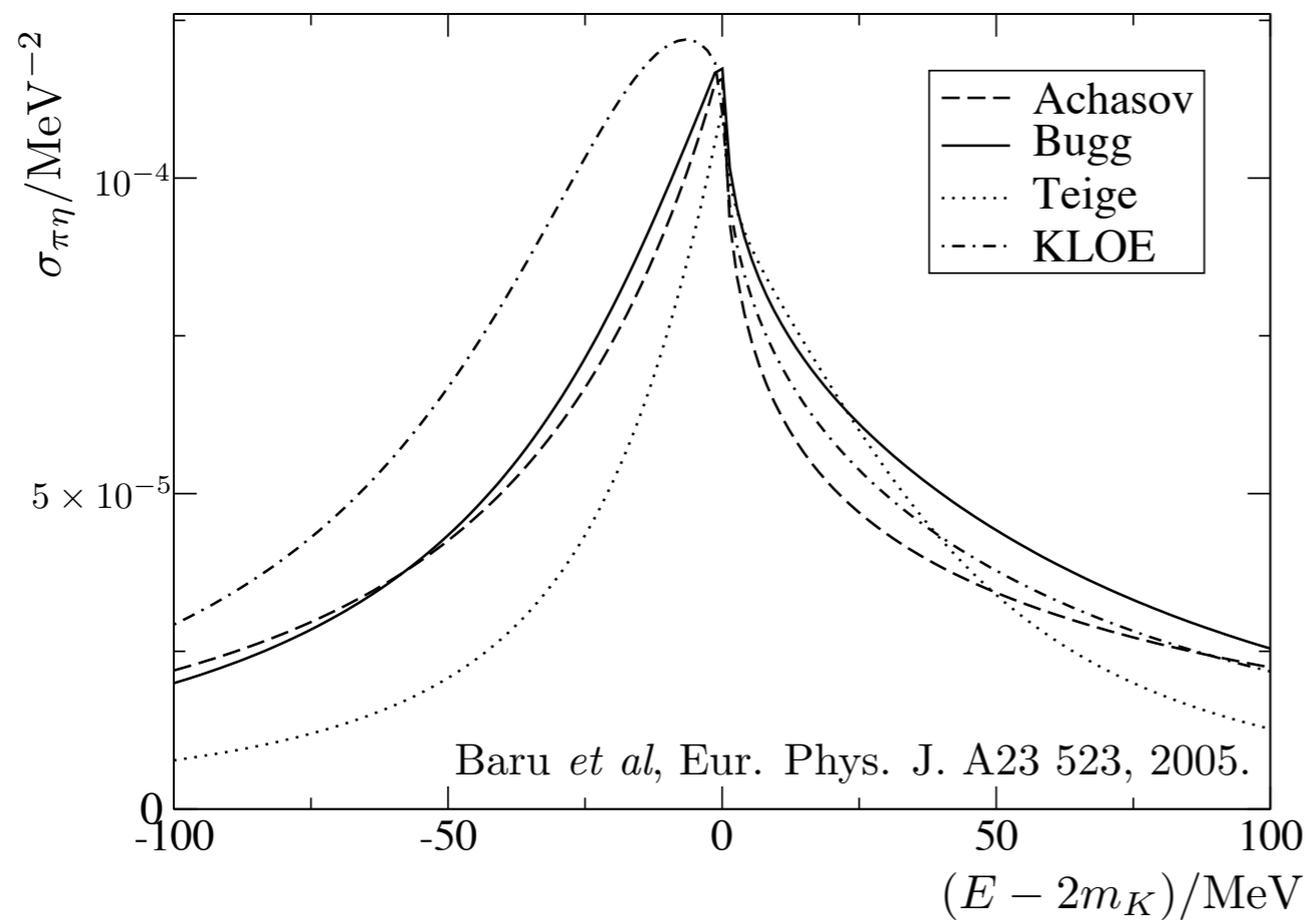
An a_0 resonance

- for more see Jozef Dudek, Tuesday 26 Jul 2016 at 16:50

PRD 93 094506, arXiv:1602.05122

$$\pi\eta - K\bar{K} - \pi\eta'$$

$$I = 1 \quad J = 0$$

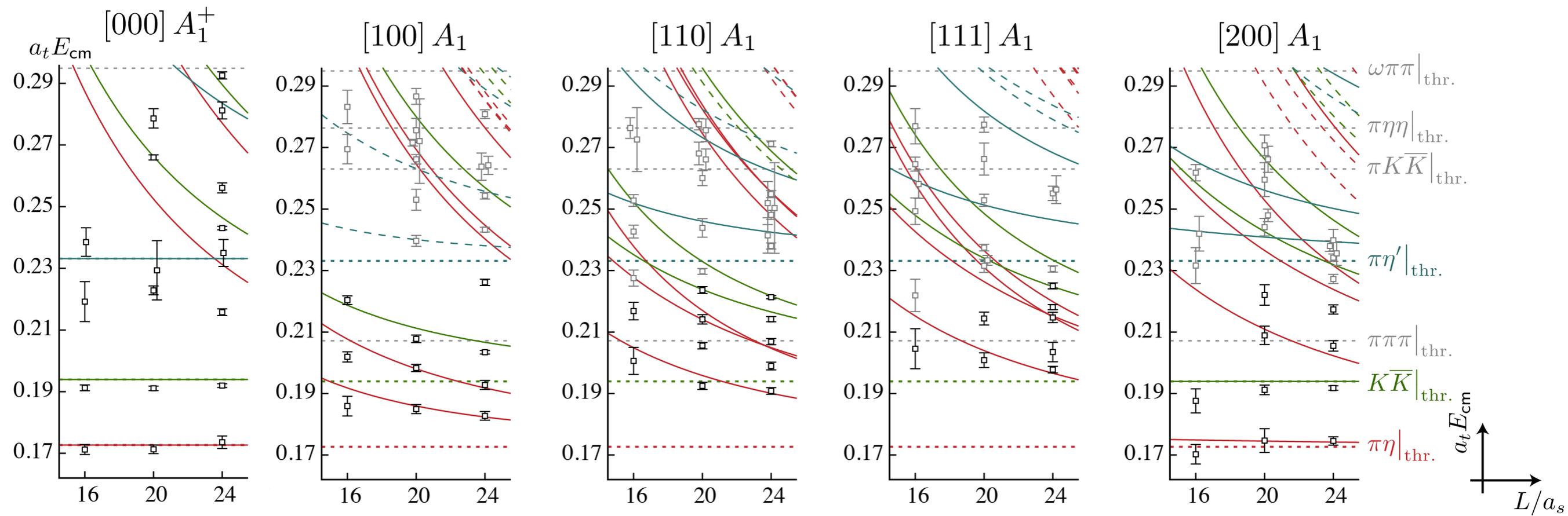


GAMS, Alde *et al* PLB 203 397, 1988.

$$m_\pi = 391 \text{ MeV}$$

An a_0 resonance

$$\pi\eta - K\bar{K} - \pi\eta'$$

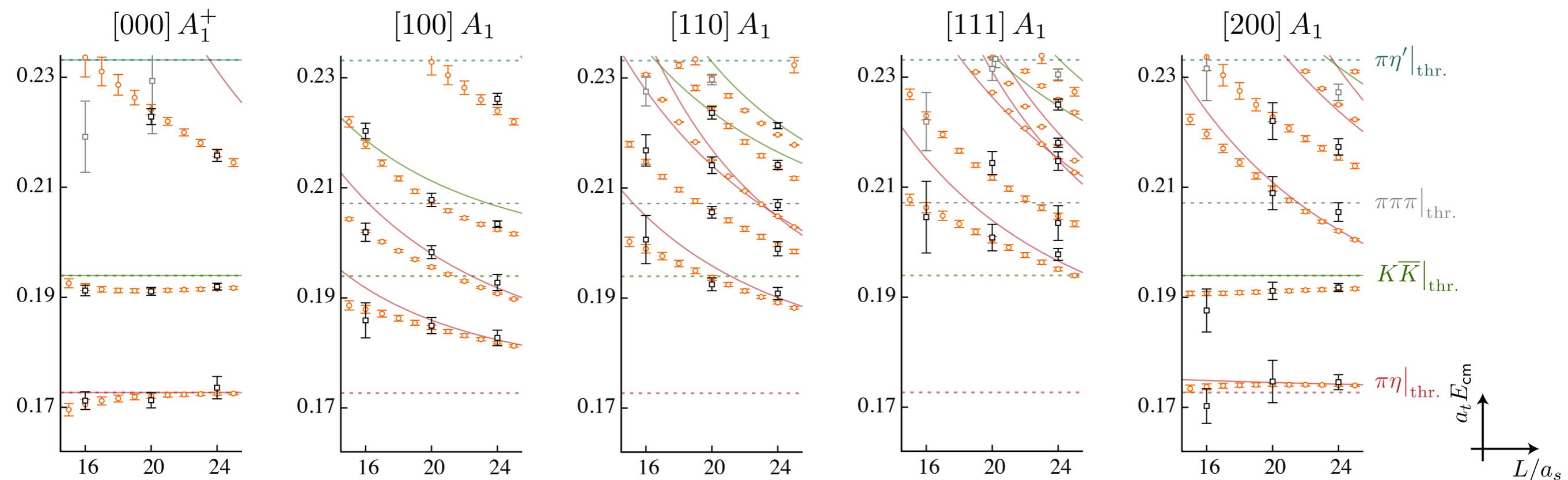


$$m_\pi = 391 \text{ MeV}$$

a_0 resonance - two channel region

$\pi\eta$ - $K\bar{K}$

using 47 energy levels



$$K_{ij} = \frac{g_i g_j}{m^2 - s} + \gamma_{ij}$$

$$\begin{aligned} m &= (0.2214 \pm 0.0029 \pm 0.0004) \cdot a_t^{-1} \\ g_{\pi\eta} &= (0.091 \pm 0.016 \pm 0.009) \cdot a_t^{-1} \\ g_{K\bar{K}} &= (-0.129 \pm 0.015 \pm 0.002) \cdot a_t^{-1} \\ \gamma_{\pi\eta, \pi\eta} &= -0.16 \pm 0.24 \pm 0.03 \\ \gamma_{\pi\eta, K\bar{K}} &= -0.56 \pm 0.29 \pm 0.04 \\ \gamma_{K\bar{K}, K\bar{K}} &= 0.12 \pm 0.38 \pm 0.08 \end{aligned}$$

$$\chi^2/N_{\text{dof}} = \frac{58.0}{47-6} = 1.41$$

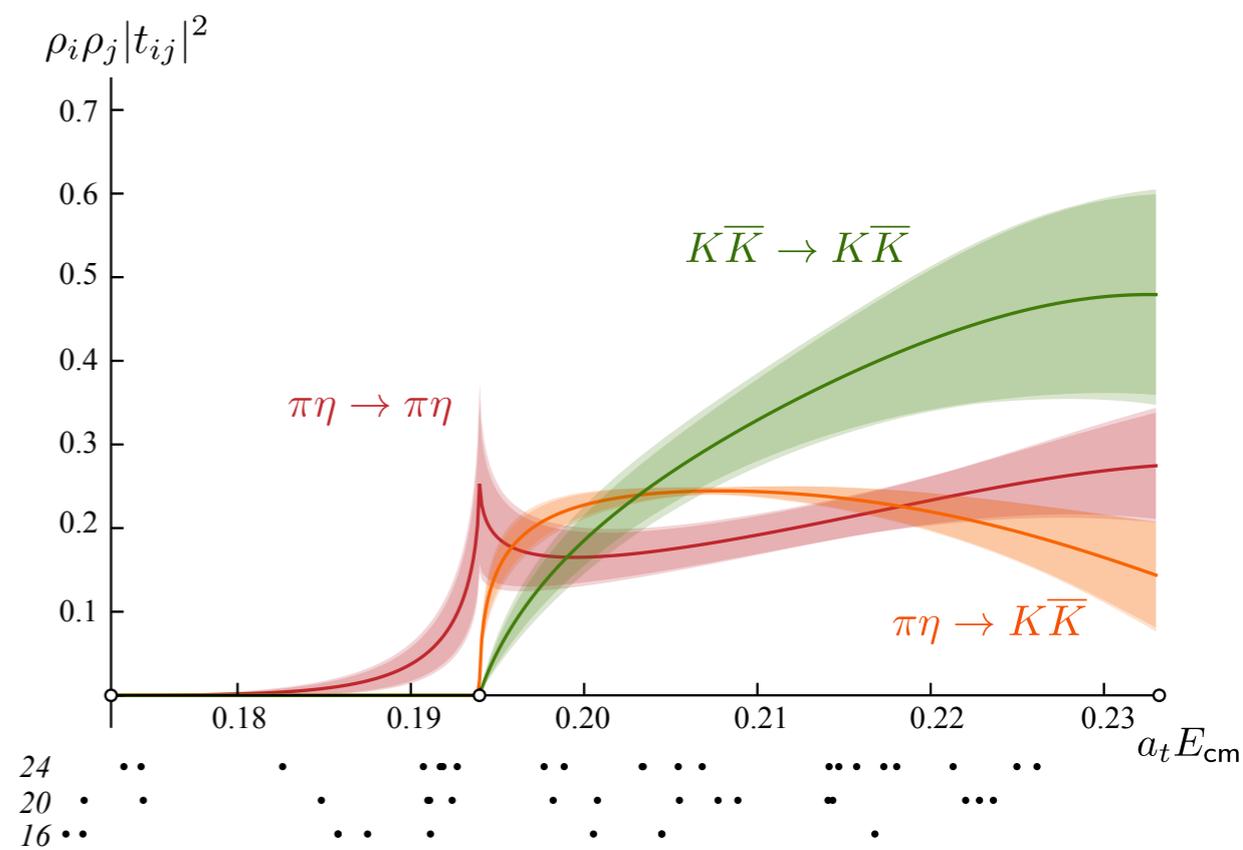
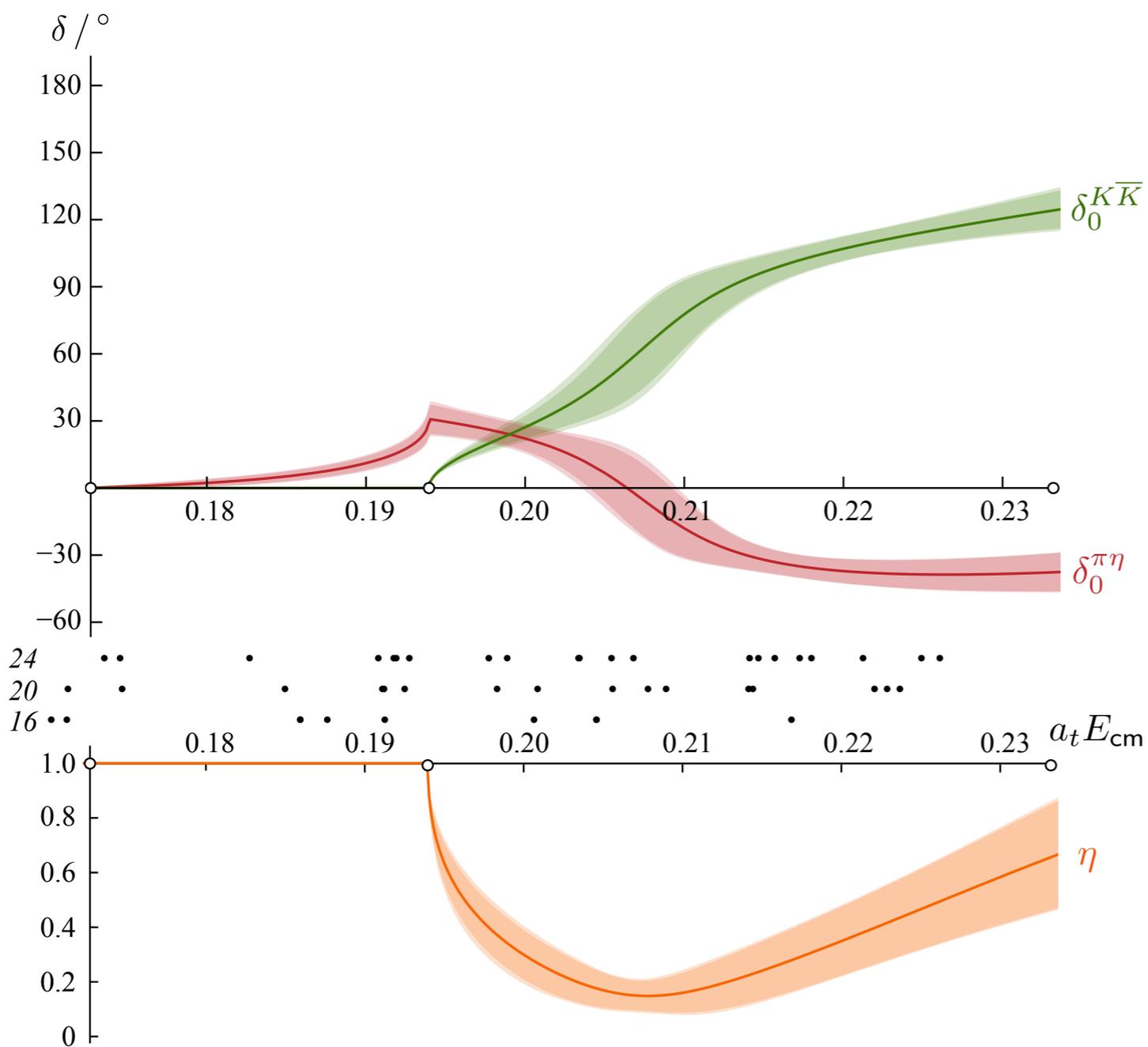
$$\begin{bmatrix} 1 & 0.58 & -0.06 & -0.51 & 0.39 & 0.02 \\ & 1 & -0.63 & -0.87 & 0.84 & -0.49 \\ & & 1 & 0.52 & -0.68 & 0.83 \\ & & & 1 & -0.90 & 0.53 \\ & & & & 1 & -0.78 \\ & & & & & 1 \end{bmatrix}$$

$$m_\pi = 391 \text{ MeV}$$

a_0 resonance - two channel region

S-wave $\pi\eta$ - $K\bar{K}$

from 47 energy levels

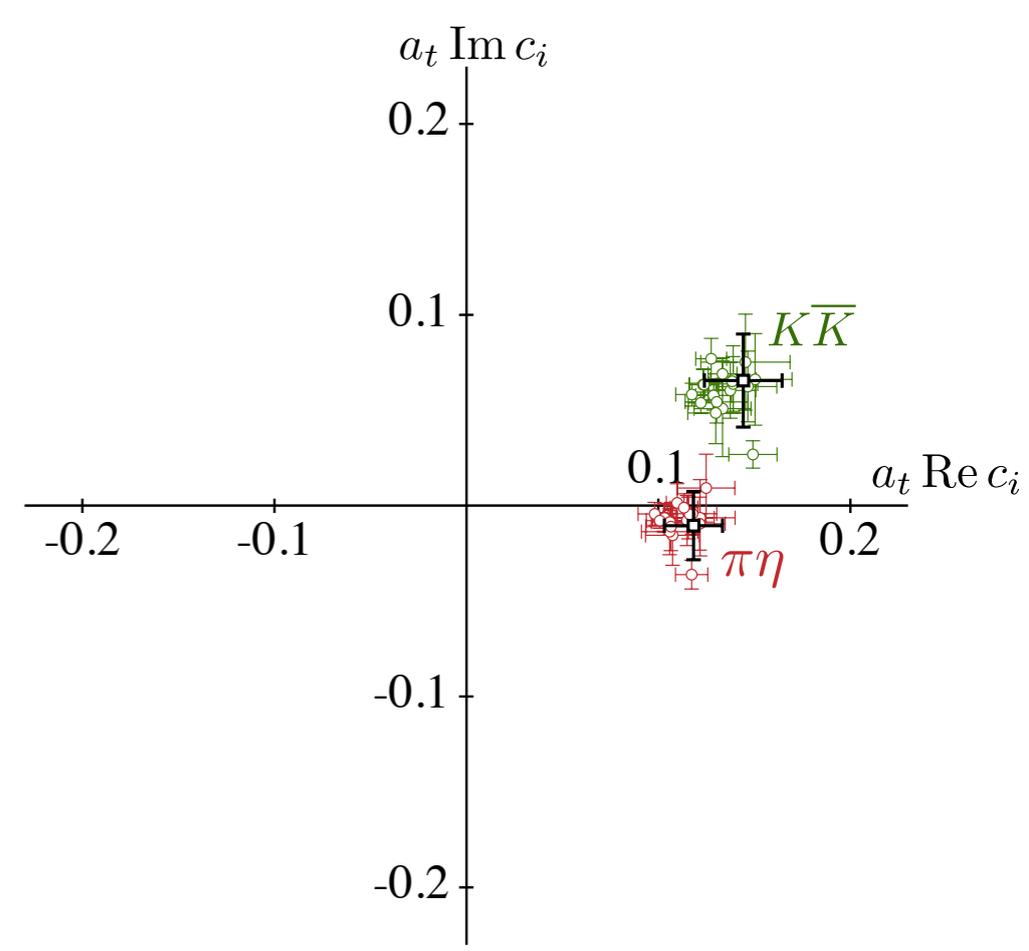
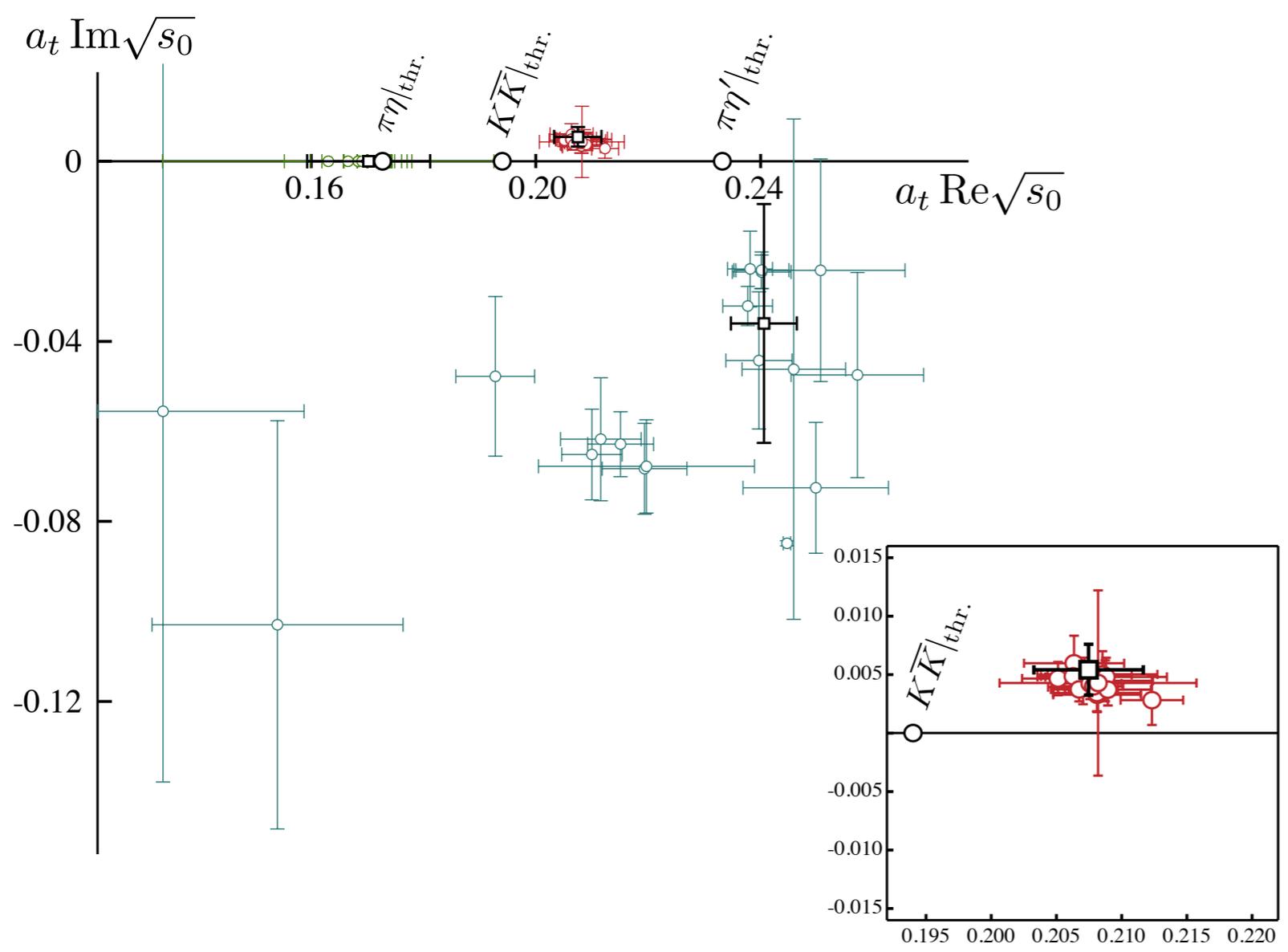


$$m_\pi = 391 \text{ MeV}$$

a_0 resonance pole

- for more see Jozef Dudek, Tuesday 26 July 2016 at 16:50

$$t_{ij} \sim \frac{C_i C_j}{s_0 - s}$$

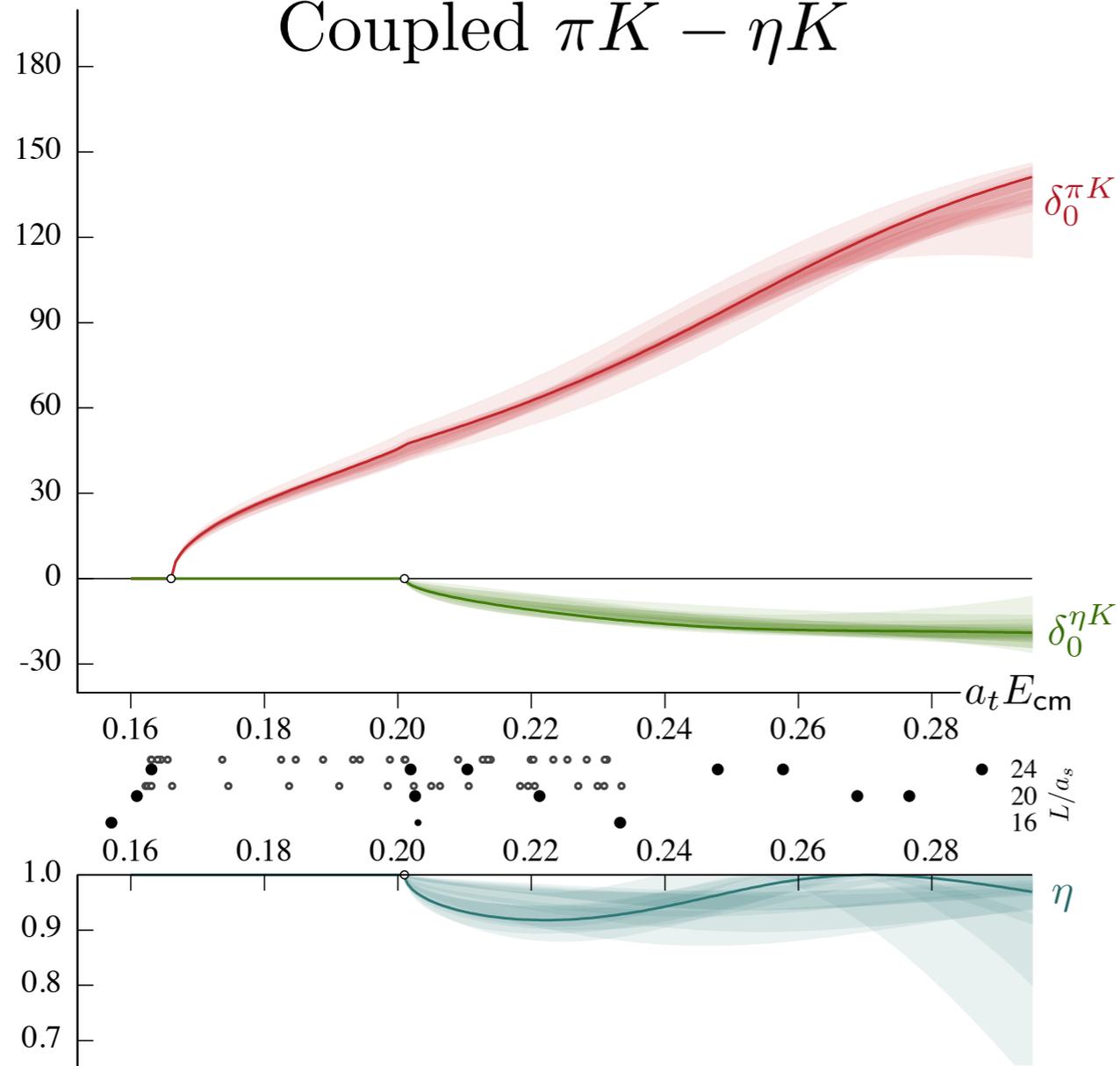


$$m_\pi = 391 \text{ MeV}$$

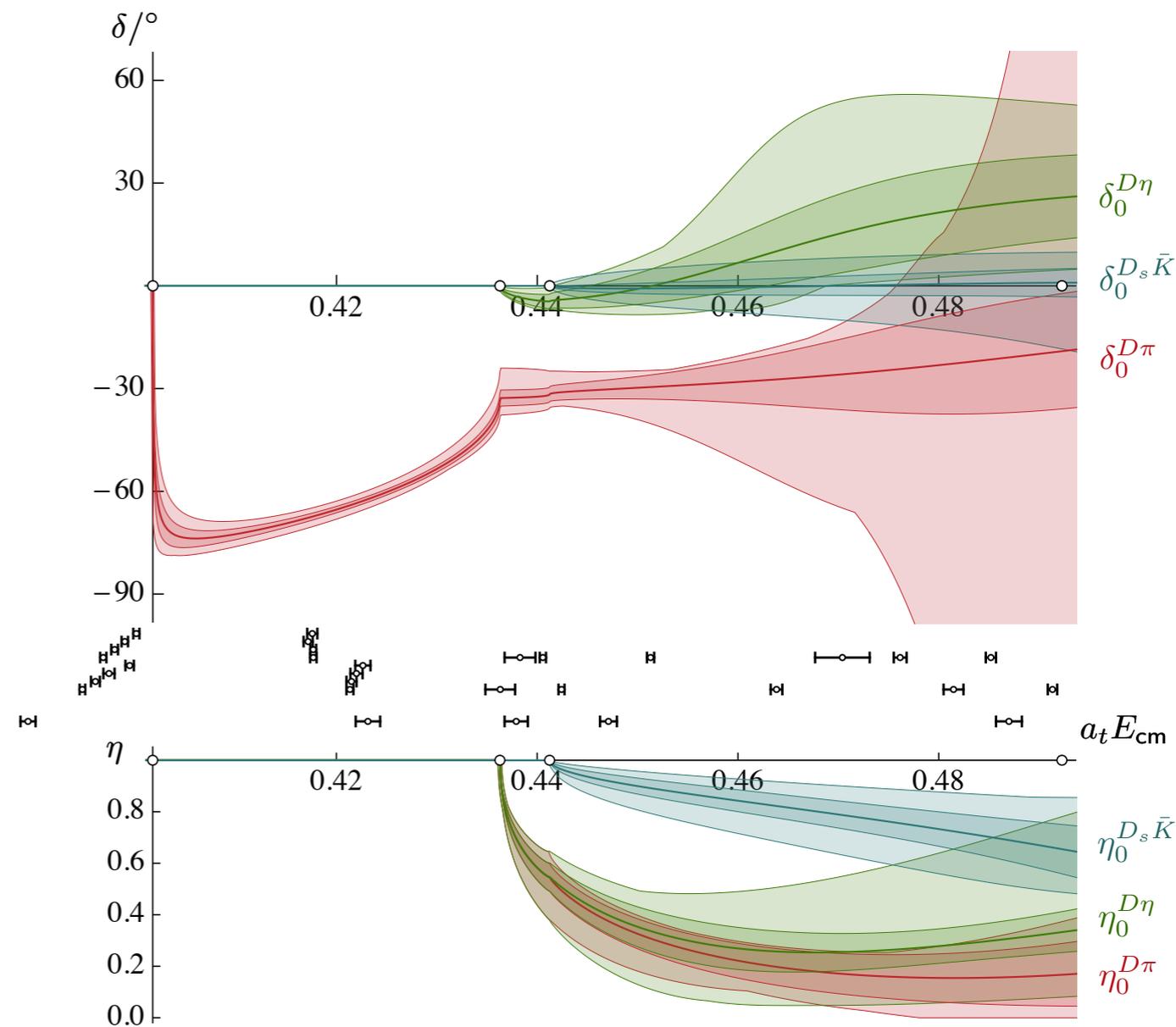
Other calculations

- Graham Moir, Thursday 28 July 2016 at 15:00

Coupled $\pi K - \eta K$



Coupled $D\pi - D\eta - D_s \bar{K}$



Combined S & P-wave analysis
80 energy levels from 3 volumes
arXiv:1406.4158, PRL 113 (2014) no.18, 182001

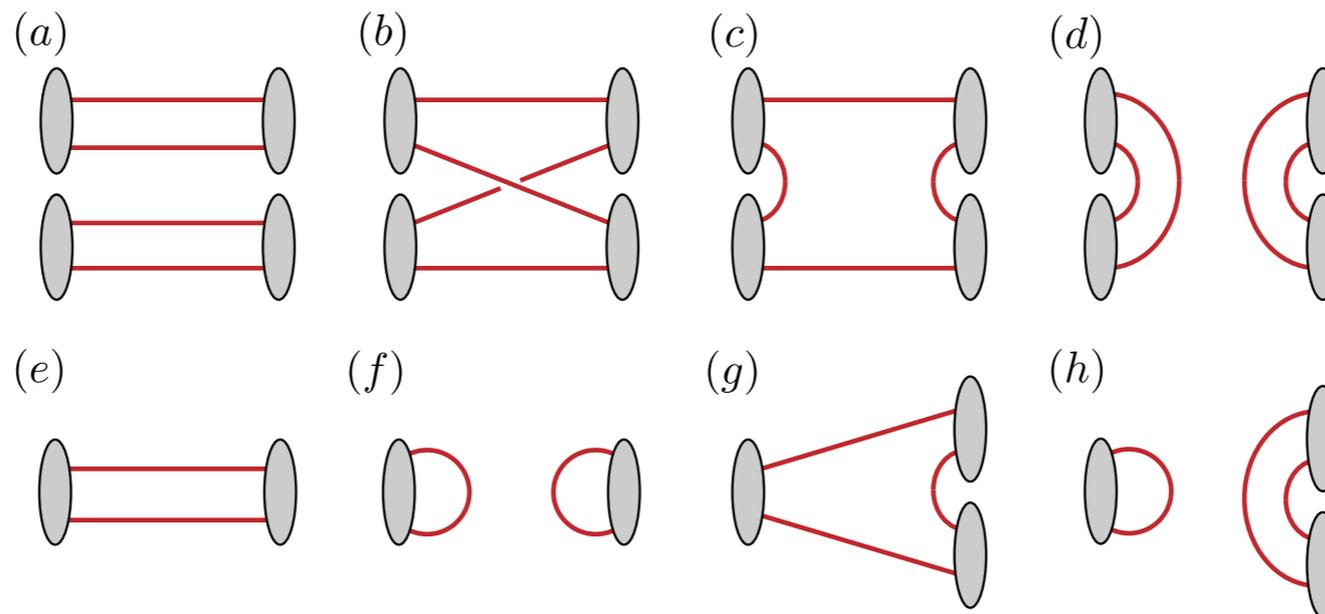
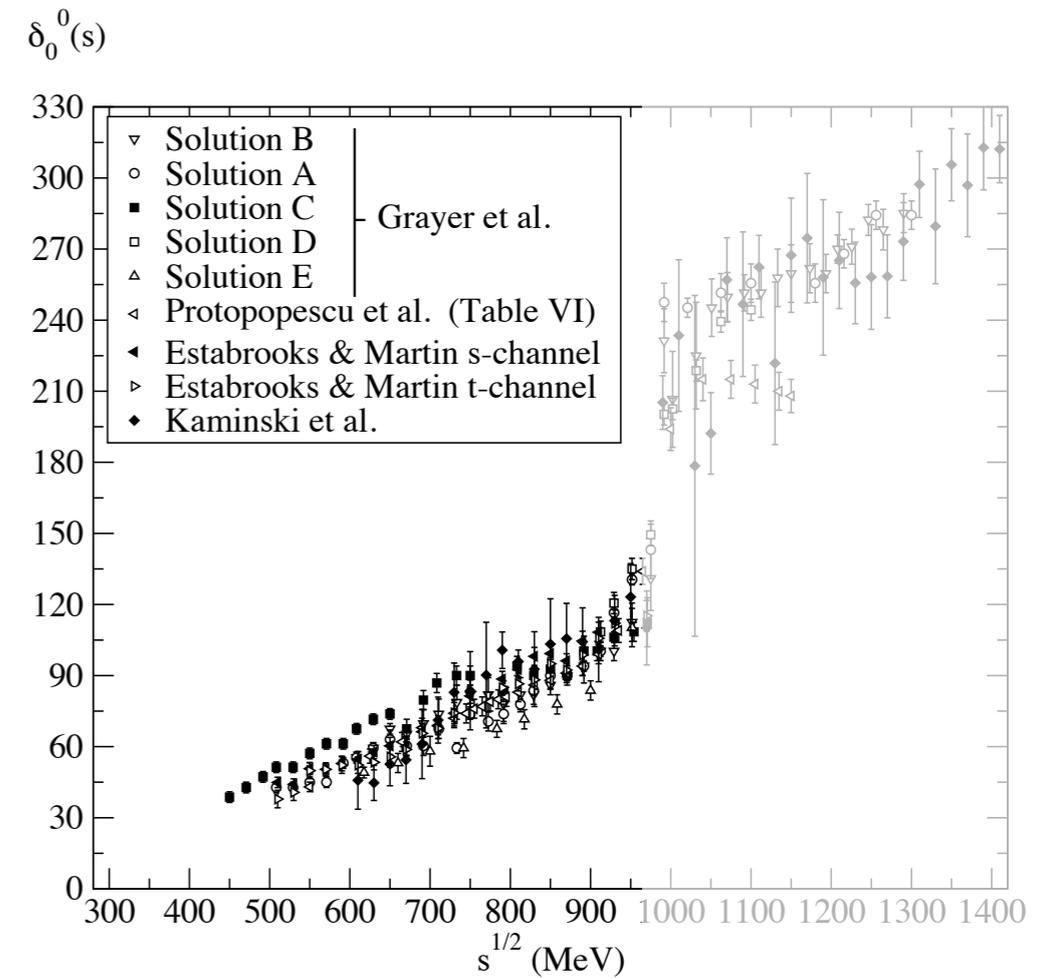
Combined S & P-wave analysis
3 coupled channels in S-wave
47 energy levels from 3 volumes
arXiv:1607.????

$$m_\pi = 391 \text{ MeV}$$

The $f_0(500)/\sigma$ resonance

- see Raul Briceño, Tuesday 26 July 2016 at 15:20
arXiv:1607.05900

elastic scattering with
vacuum quantum numbers
 $\pi\pi$ in $I = 0, J = 0$

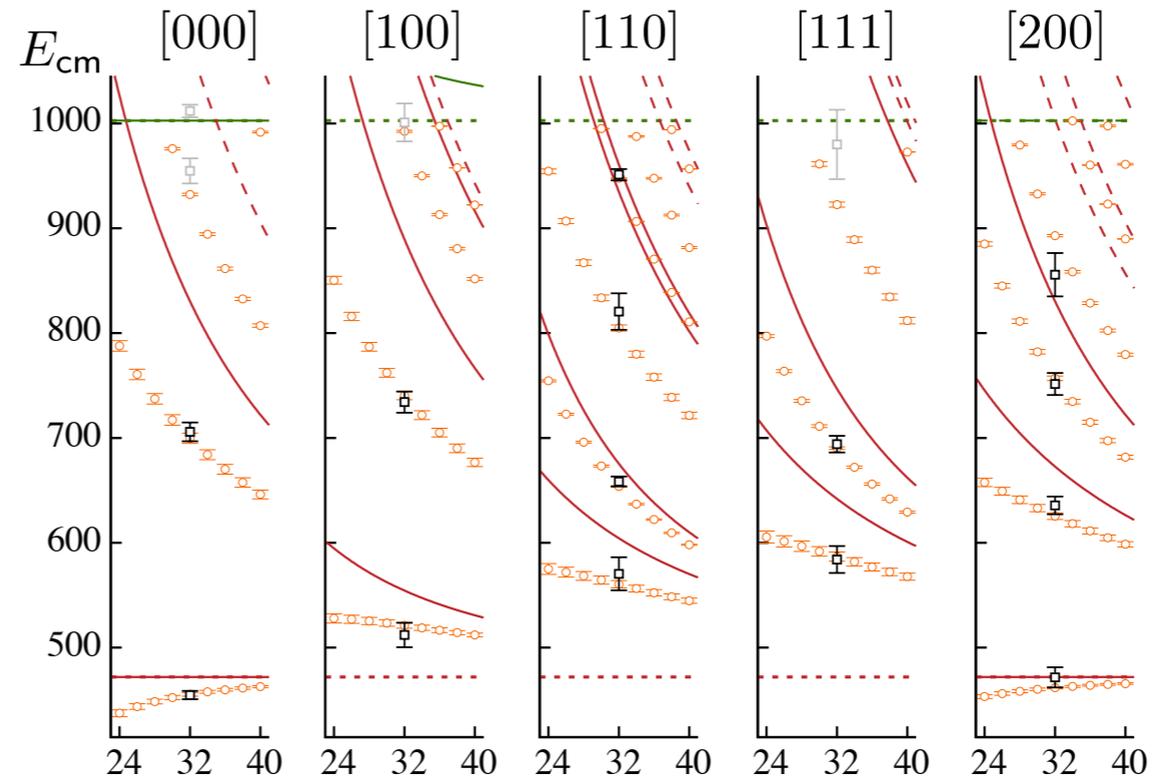


The $f_0(500)/\sigma$ resonance

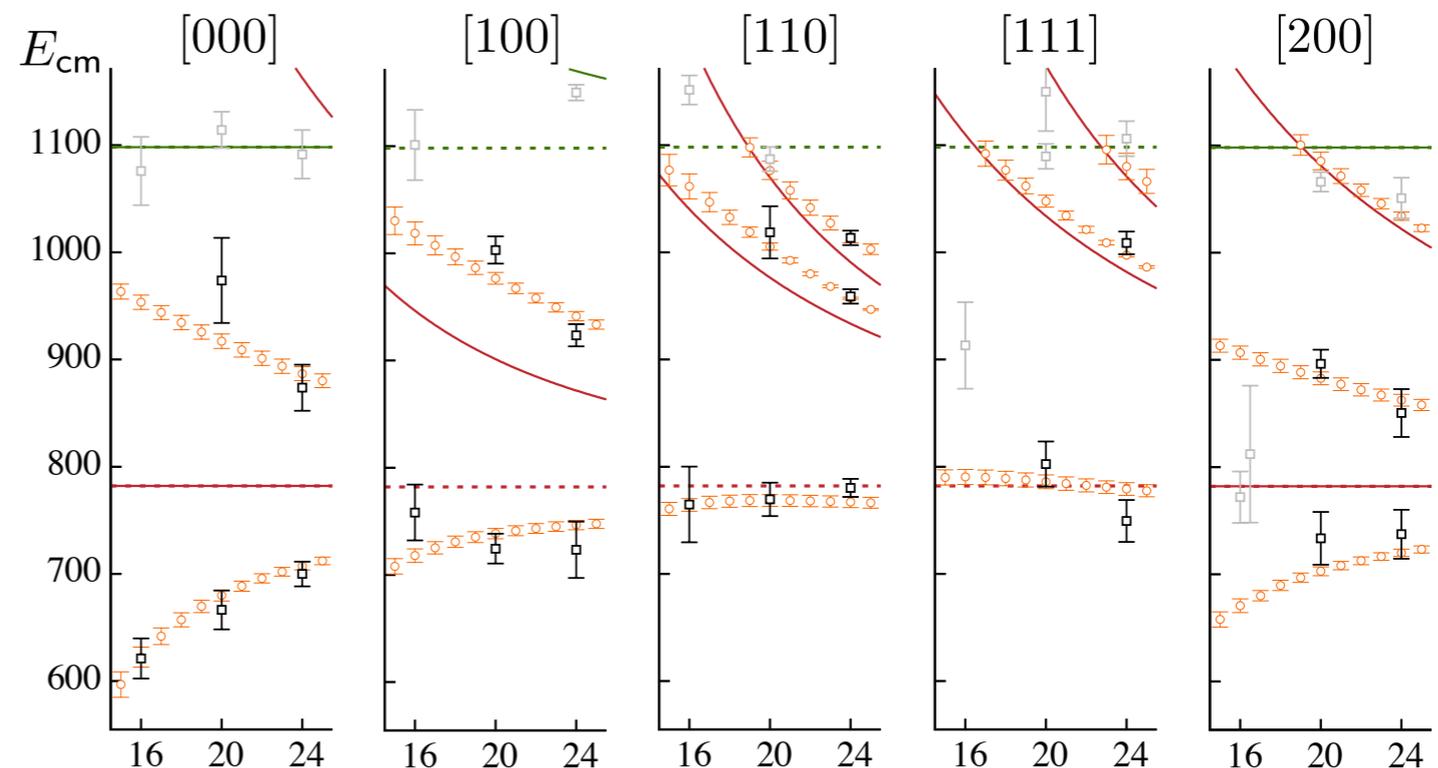
- see Raul Briceño, Tuesday 26 July 2016 at 15:20

elastic scattering with
vacuum quantum numbers
 $\pi\pi$ in $I = 0, J = 0$

$$m_\pi = 236 \text{ MeV}$$

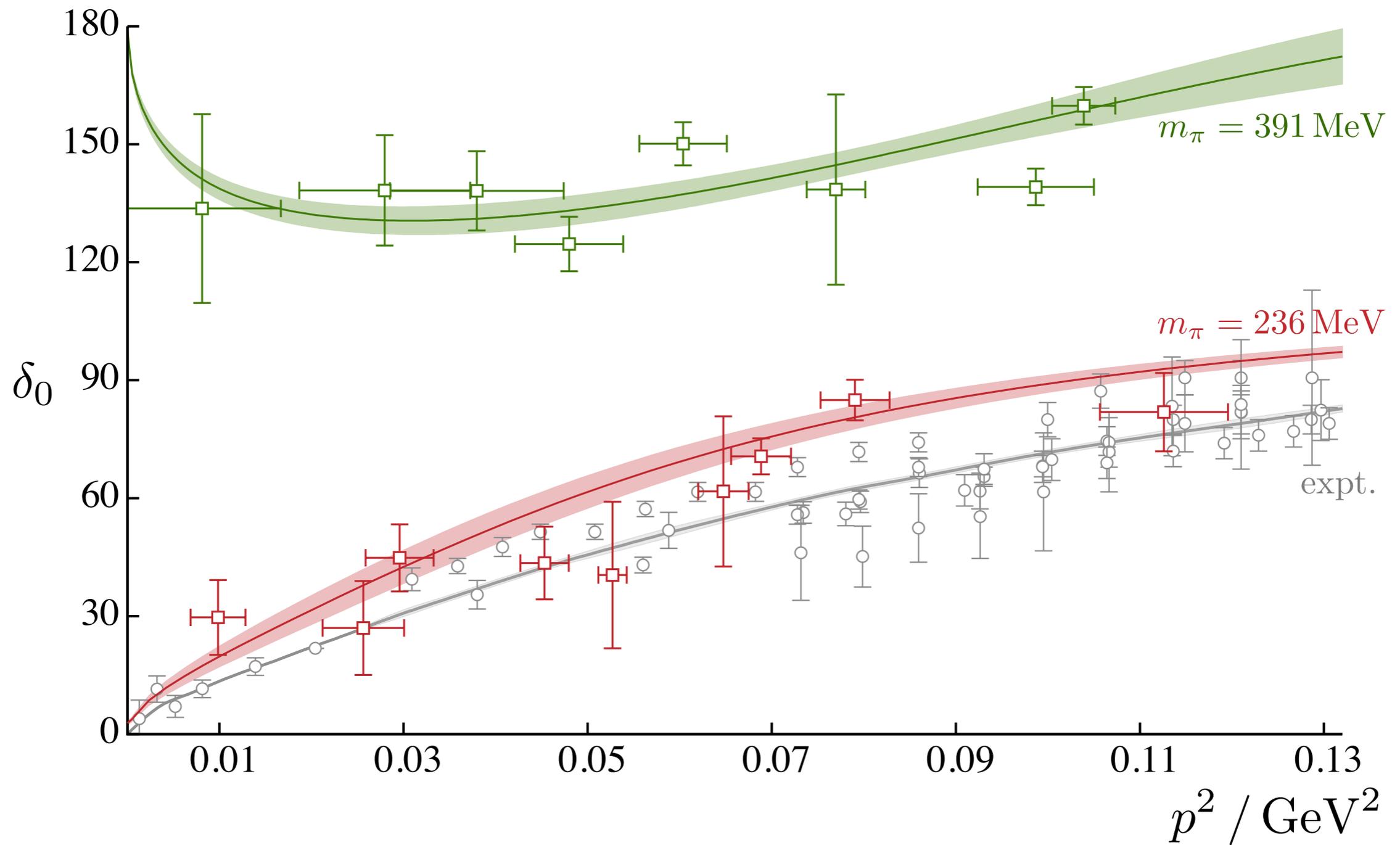


$$m_\pi = 391 \text{ MeV}$$



The $f_0(500)/\sigma$ resonance

- see Raul Briceño, Tuesday 26 July 2016 at 15:20
arXiv:1607.05900



Future directions

two-body coupled-channel

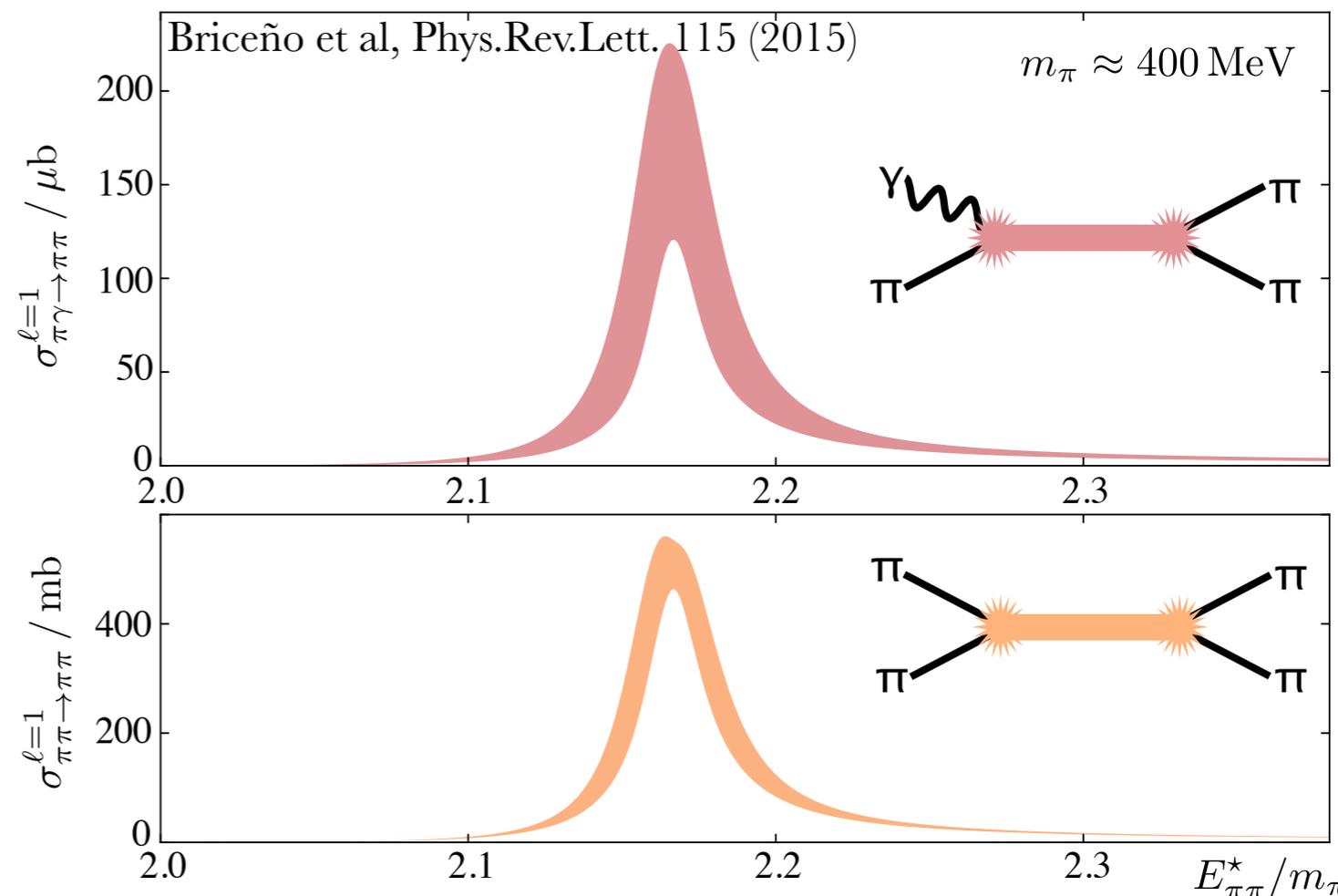
$f_0(980)$

$D\bar{D}$

$D\bar{D}^*$

$N\pi$

$\gamma a \rightarrow bc$



- see also Luka Leskovec on 29 July 2016 at 17:50

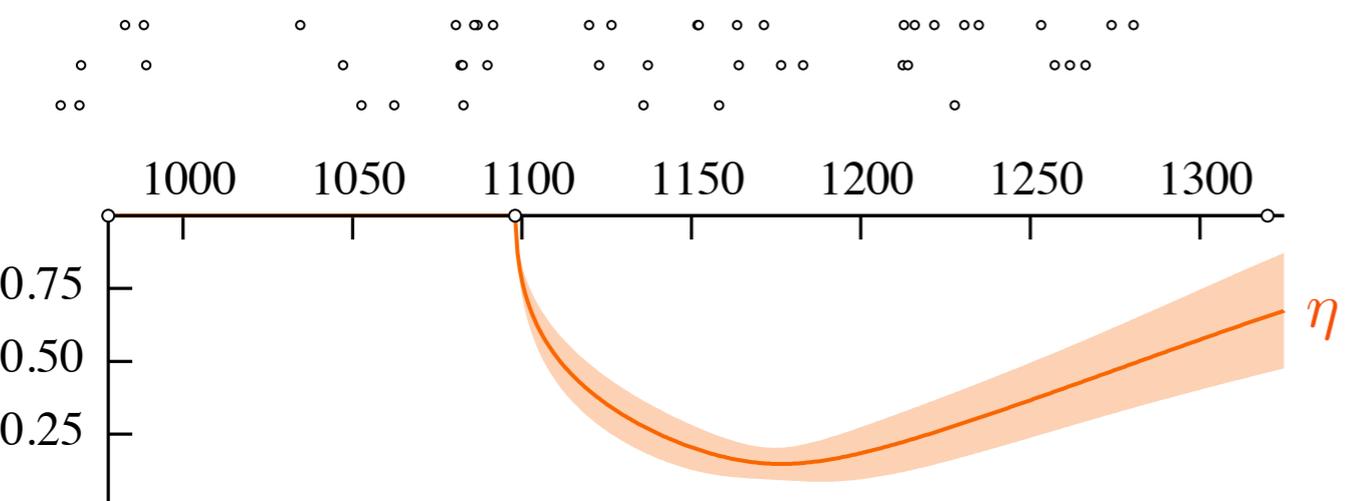
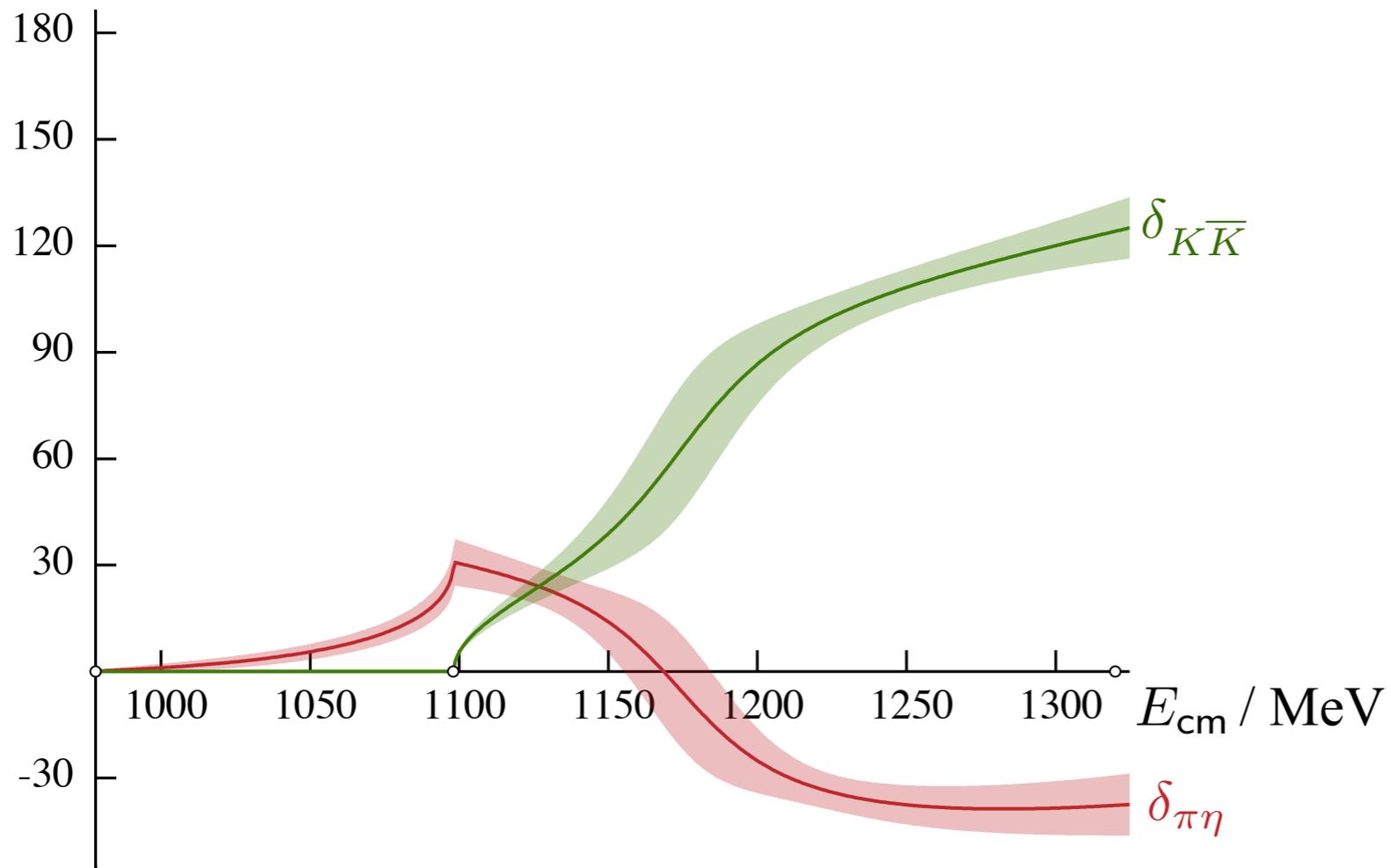
further operator structures - glueball, tetraquark, ... - see Gavin Cheung, Monday 25 July 2016 at 14:55

formalism for three-body and beyond

- needed for higher energies

- needed to get closer to the physical mass

- see Stephen Sharpe, Tuesday 26 July 2016 at 15:40



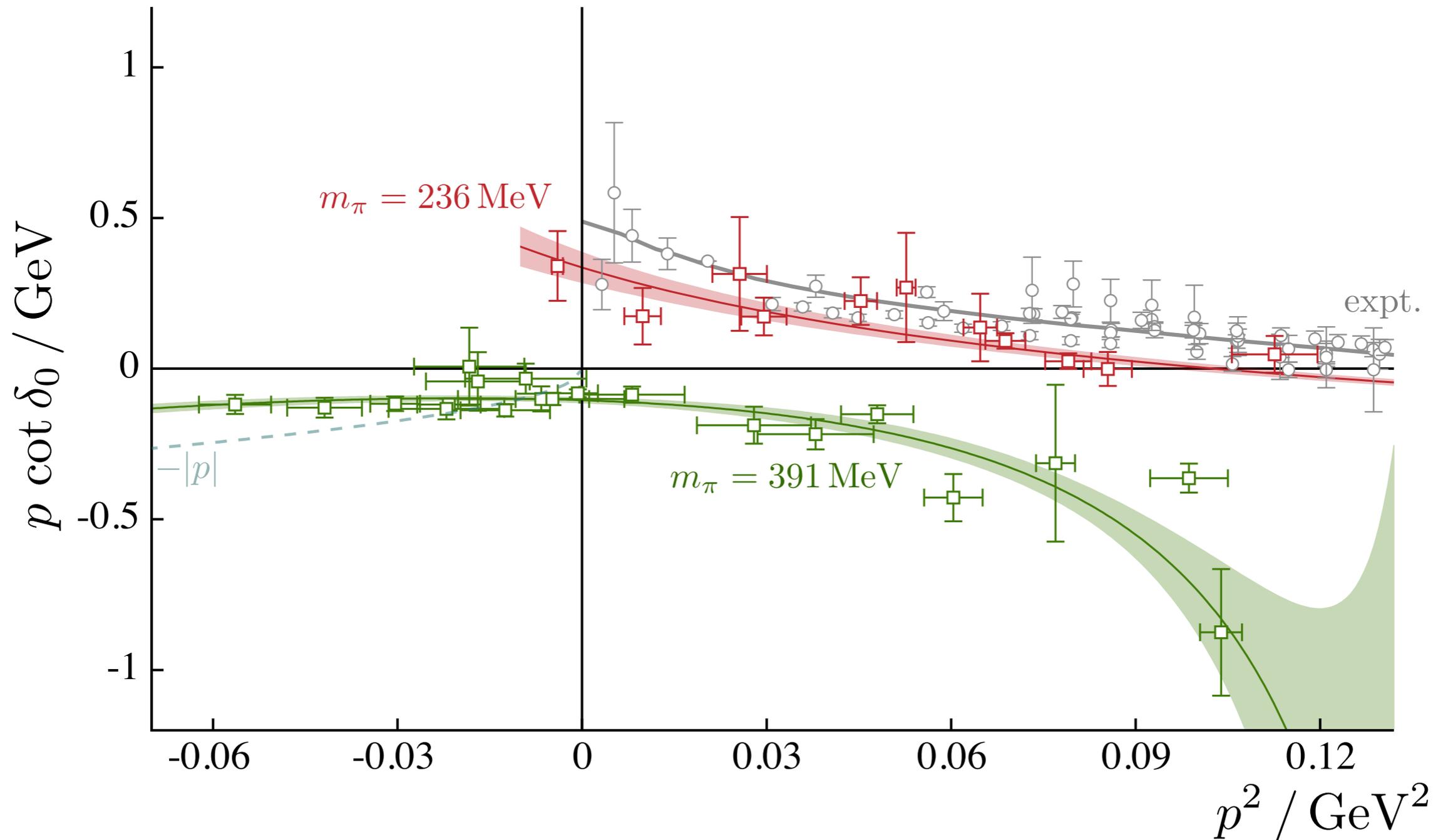
$m_\pi = 391 \text{ MeV}$

Promising results so far ...
 ... but a lot still to do

Backup

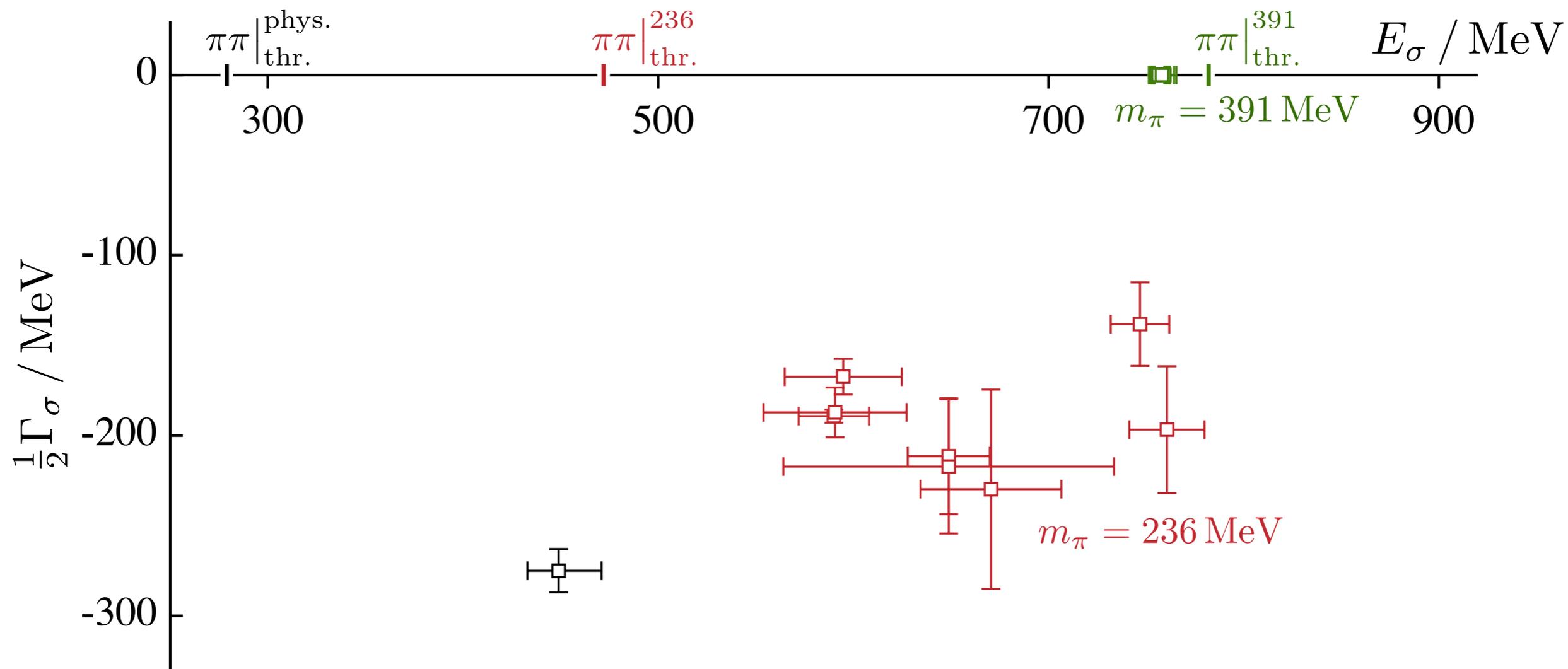
$$m_\pi = 391 \text{ MeV}$$

The $f_0(500)/\sigma$ resonance



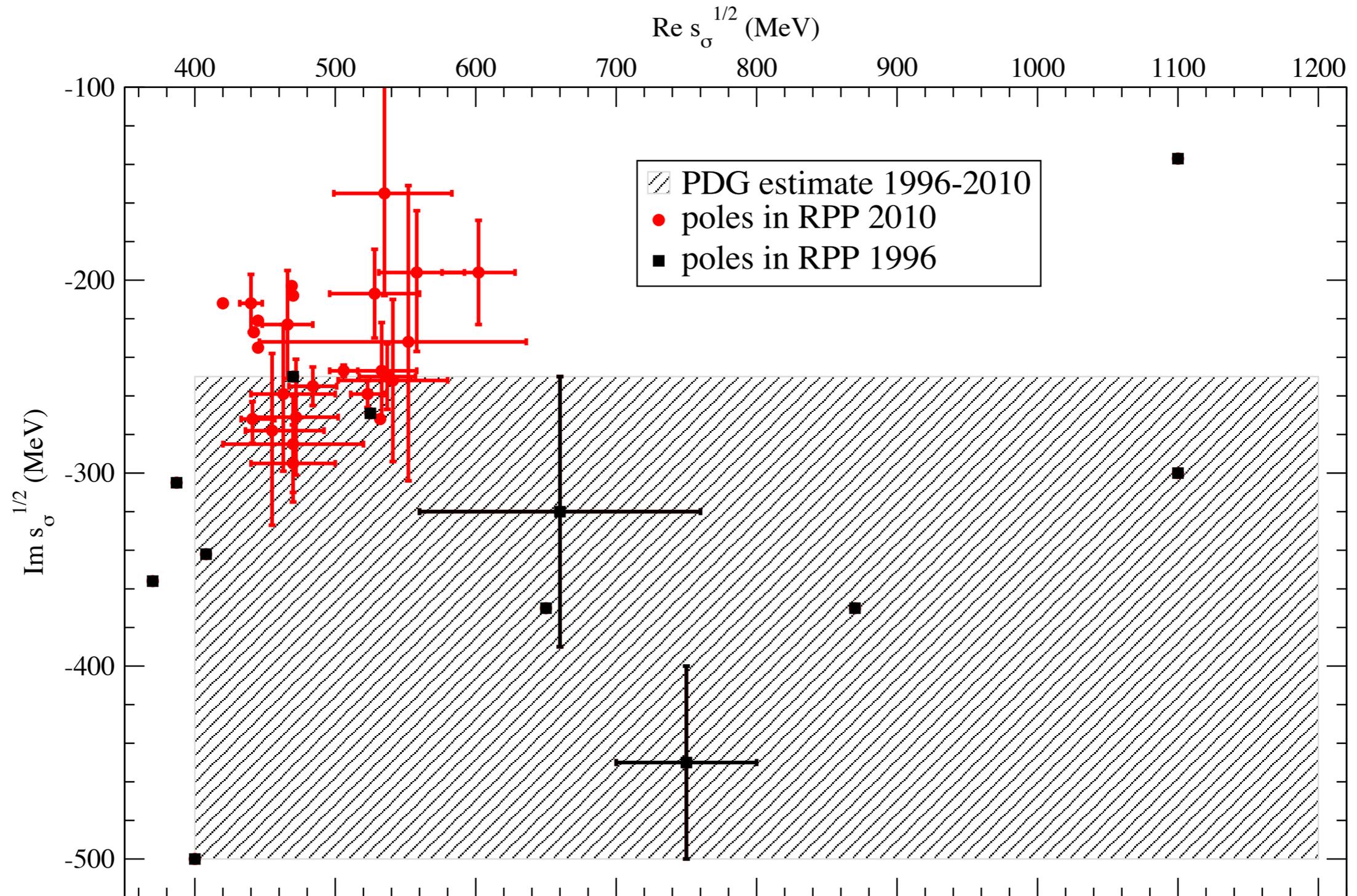
The $f_0(500)/\sigma$ resonance

- see Raul Briceño, Tuesday 26 Jul 2016 at 15:20



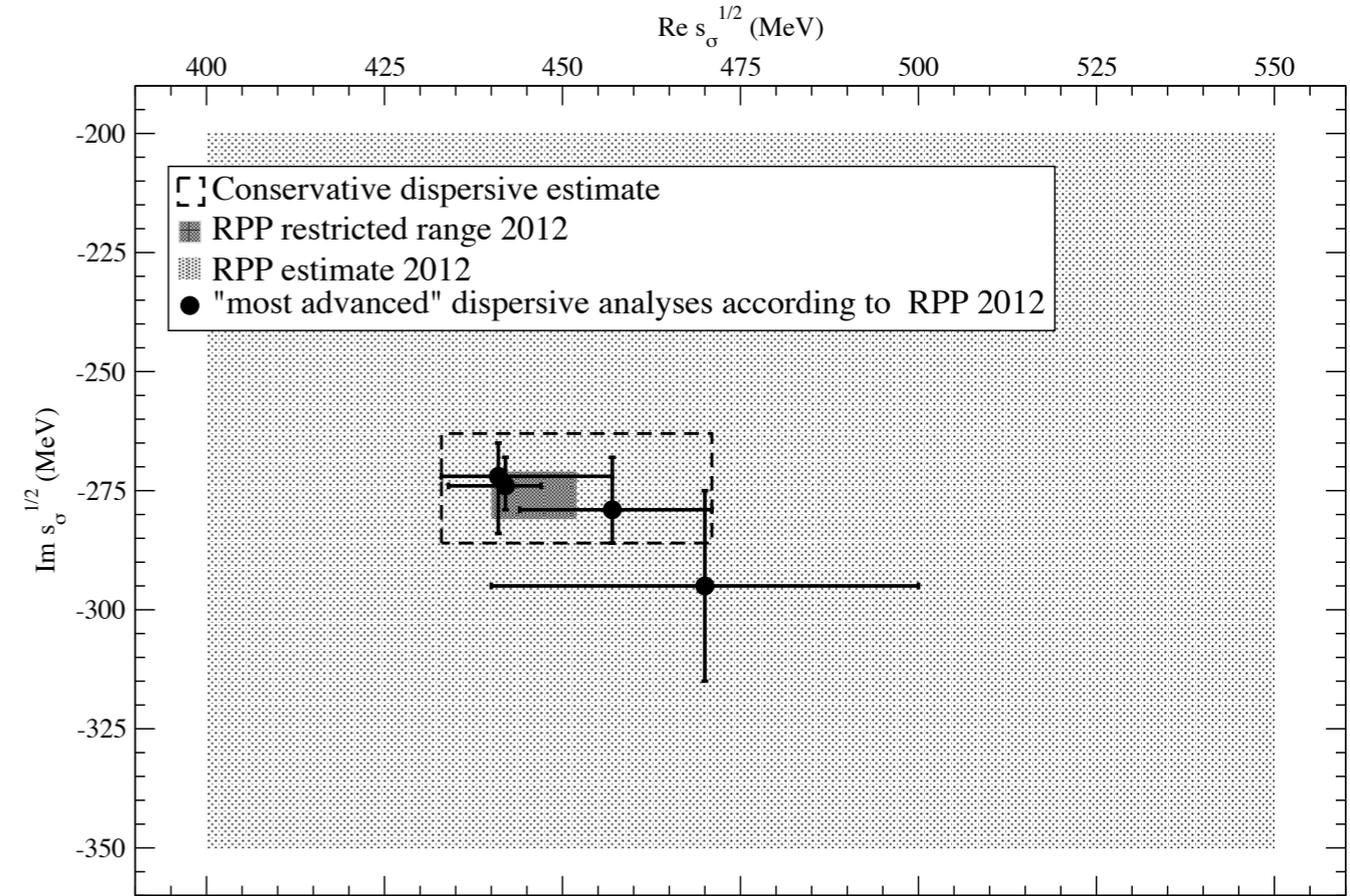
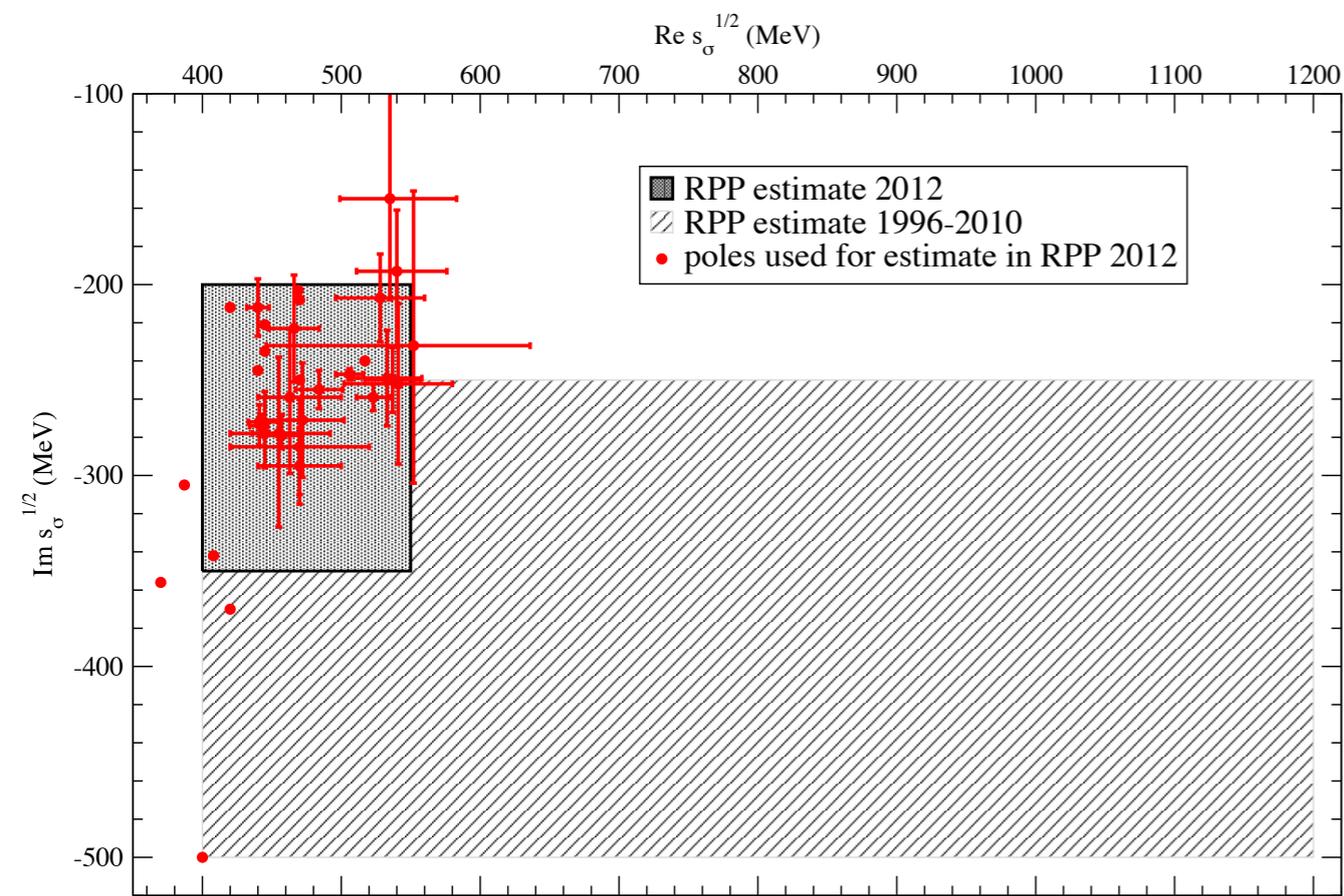
The $f_0(500)/\sigma$ resonance

from J. R. Pelaez, arXiv:1510.00653



The $f_0(500)/\sigma$ resonance

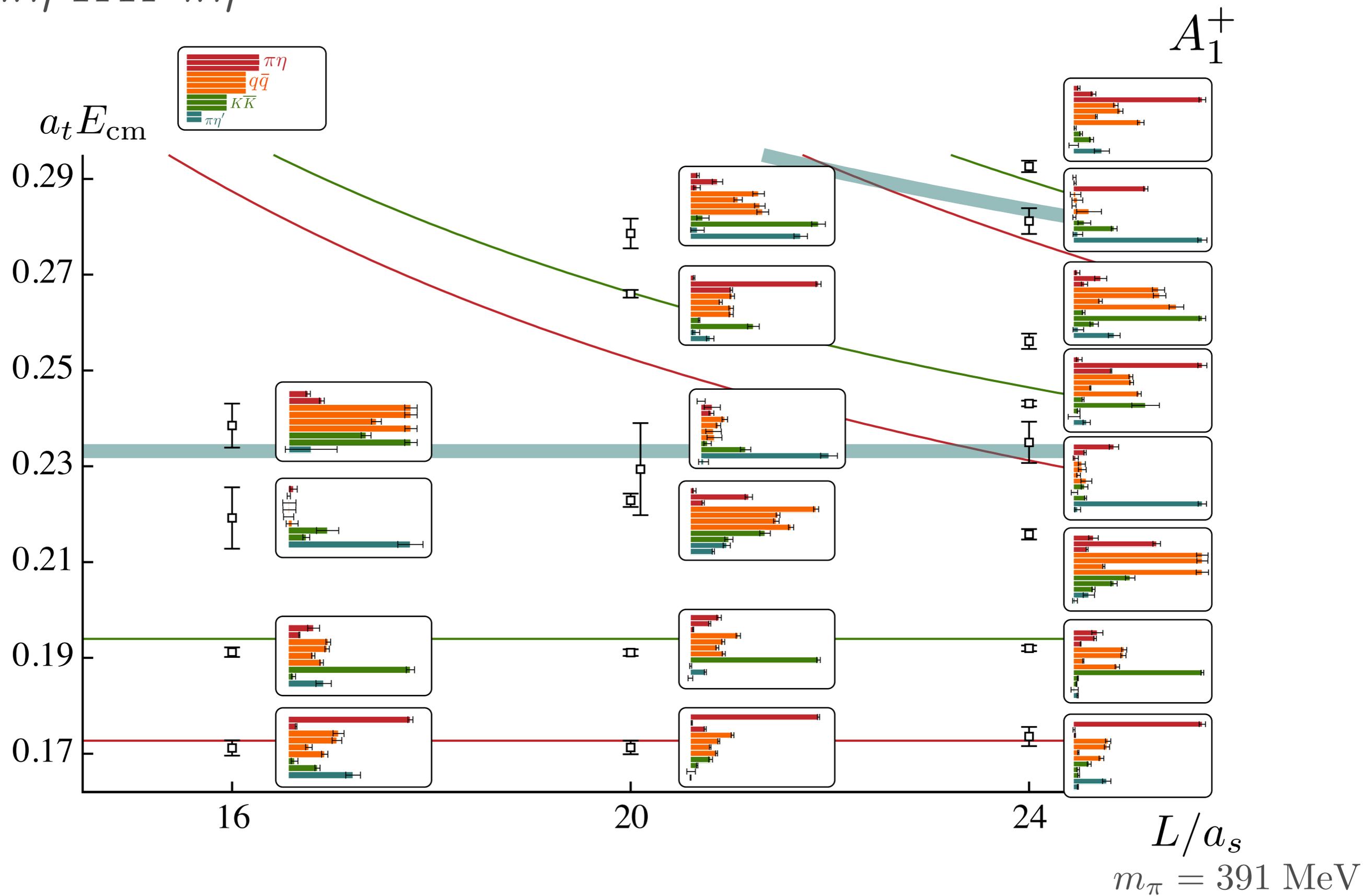
from J. R. Pelaez, arXiv:1510.00653

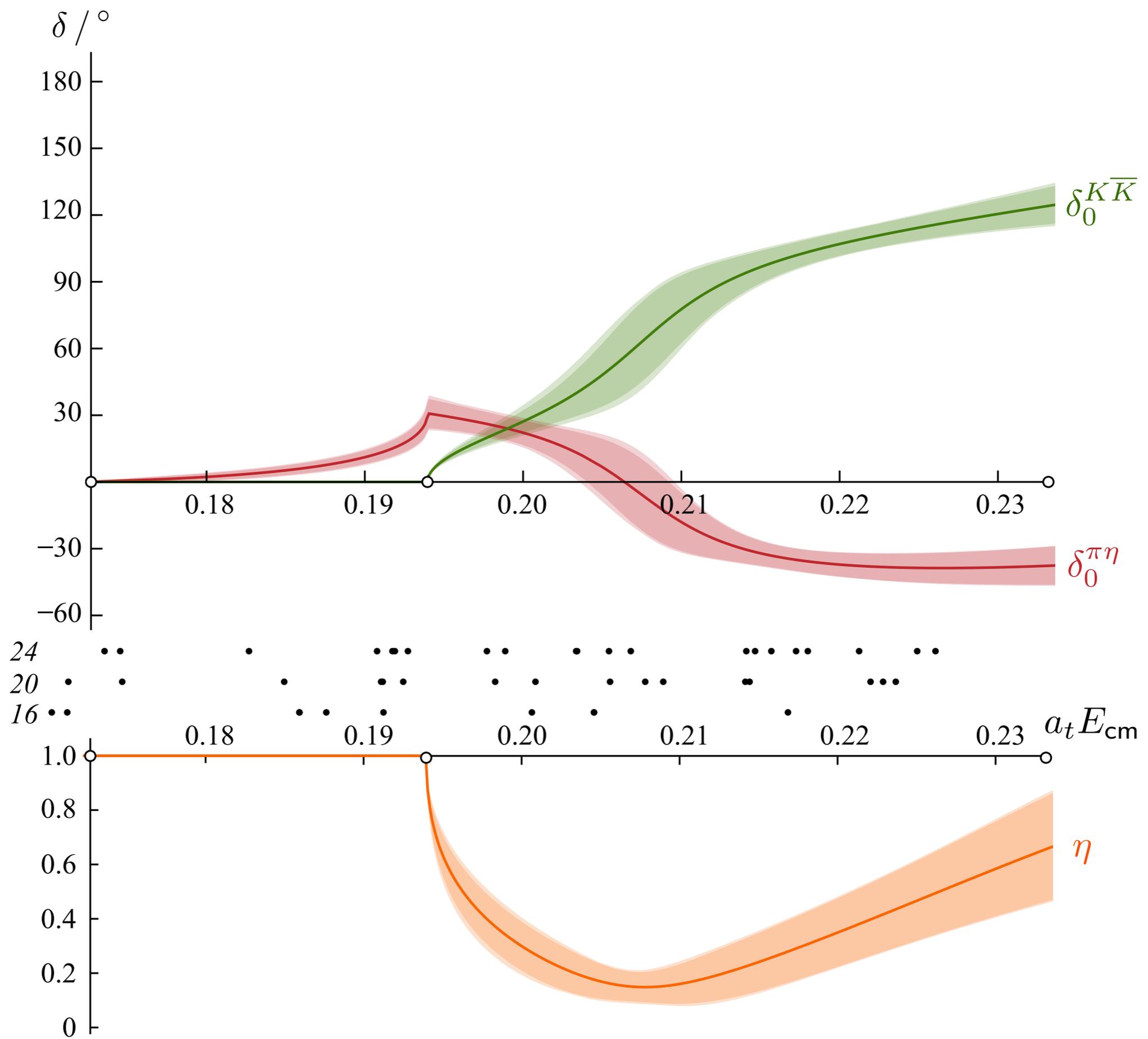


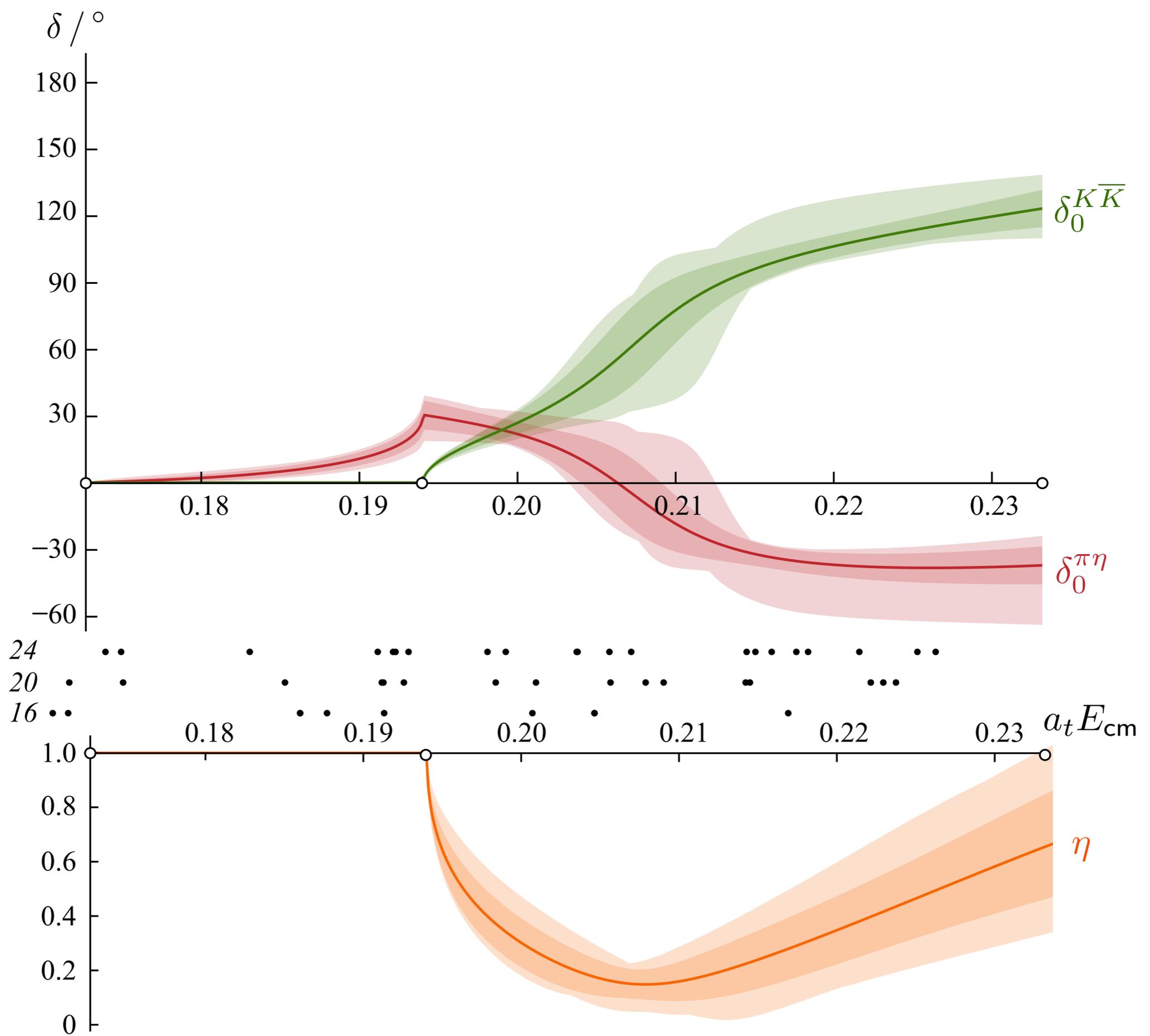
An a_0 resonance

- for more see Jozef Dudek, Tuesday 26 Jul 2016 at 16:50

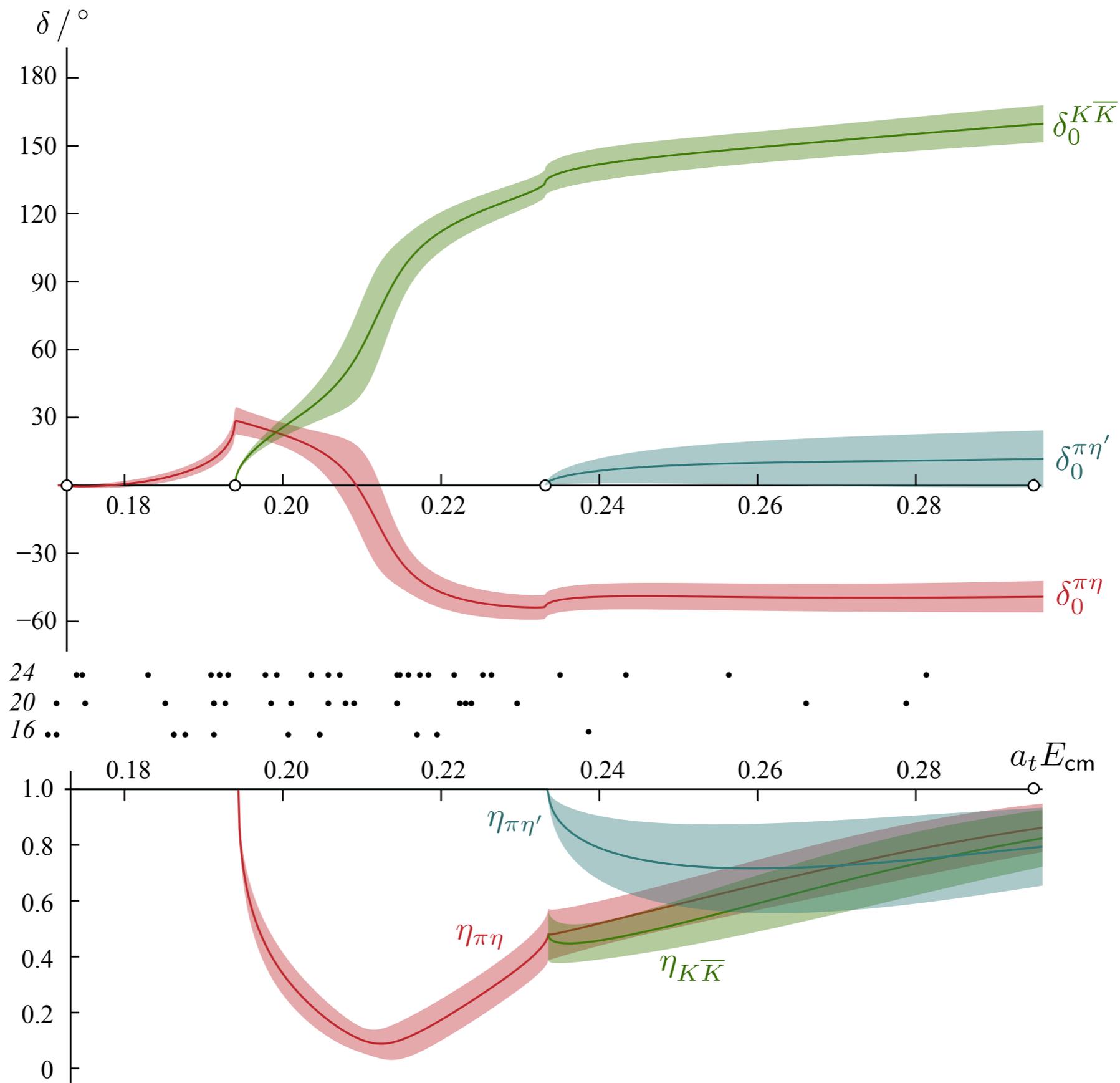
$\pi\eta$ - $K\bar{K}$ - $\pi\eta'$







An a_0 resonance - three channel region



$$m_\pi = 391 \text{ MeV}$$