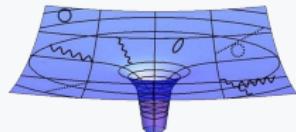


Four-Fermion Theories with Exact Chiral Symmetry in Three Dimensions

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RESEARCH TRAINING GROUP
QUANTUM AND GRAVITATIONAL FIELDS



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Outline

- 1 Introduction to Four-Fermion Theories
- 2 Two Four-Fermion Interactions
- 3 Fierz Identities for the Thirring Model
- 4 Fermion Bag Approach
- 5 Summary



What are Four-Fermion Theories?

QFTs of fermions with 4th power of fermion fields as interaction:

$$\mathcal{L} = \bar{\psi}_j (\partial + m) \psi_j + \sum_{\alpha} \frac{g_{\alpha}^2}{2N_f} (\bar{\psi}_j \Gamma_{\alpha} \psi_j)^2 \quad j = 1, \dots, N_f$$

Main Models

Thirring 1958, soluble fermionic theory in 2D

$$\Gamma_{\alpha} = \gamma_{\mu}$$

Nambu & Jona-Lassino 1961, dynamical mass generation in 4D

$$\Gamma_1 = \mathbb{1}, \Gamma_2 = i\gamma_5$$

Gross & Neveu 1974, asymptotic freedom, chiral symmetry breaking in 2D

$$\Gamma = \mathbb{1}$$



What is the Thirring model?

QFT with N_f flavours of massless fermions with **current** interaction

Euclidean Lagrangian

$$\mathcal{L} = \bar{\psi}_j \partial^\mu \psi_j + \frac{g^2}{2N_f} \sum_{\mu=1}^3 (\bar{\psi}_j \gamma^\mu \psi_j)^2 \quad j = 1, \dots, N_f$$

- **3-dimensional** euclidean spacetime
- representation of Clifford algebra:
 - irreducible 2-dimensional (later in this talk)
 - reducible 4-dimensional (now)

Motivation

Similarity to QED₃ and possible applications in superconductors, graphene, ...

Why is the Thirring model interesting?

Symmetries

- chiral symmetry, generated by $\{\mathbb{1}, \gamma_4, \gamma_5, i\gamma_4\gamma_5\}$
- flavour symmetry

Result: $U(N_f, N_f)$, can be spontaneously broken to $U(N_f) \otimes U(N_f)$
 \Rightarrow chiral condensate $\langle \bar{\psi}\psi \rangle \neq 0$

Chiral Symmetry Breaking

$"N_f = 0.5"$ irreducible representation for $N_{f,\text{irr}} = 1$ corresponds to Gross-Neveu model with chiral symmetry breaking

$N_f \rightarrow \infty$ no chiral symmetry breaking

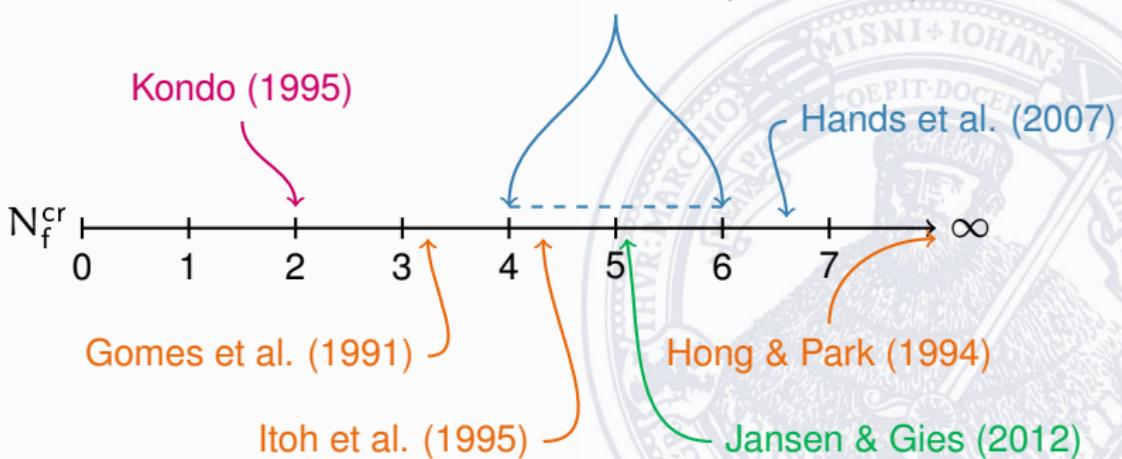
\Rightarrow There is a N_f^{cr} where chiral behaviour changes.

What is the value of N_f^{cr} ?

Results for N_f^{cr} from **Schwinger-Dyson equations**, $\frac{1}{N_f}$ -**expansion**, **functional renormalization group**, **Lattice simulation with staggered fermions**:

Kim & Kim (1996)

Del Debbio, Hands et al. (1996-1999)



Chiral Symmetry on the Lattice

Nielsen-Ninomiya Theorem

It is not possible to have a **chiral**, **local** and **translational invariant** Dirac operator with **correct continuum limit** without doublers.

older results: staggered fermions with mass

- mass breaks symmetry explicitly
- still **doublers** and no full **chiral** symmetry
- symmetry correct in the **continuum limit?**

our approach: SLAC derivative [fine for non-gauge theories:
Bergner et al. arXiv:0705.2212; Wozar, Wipf arXiv:1107.3324]

- in momentum space: multiplication by $i\gamma^\mu p_\mu$
- exact chiral symmetry
- **not local**: need to do Fourier transformation

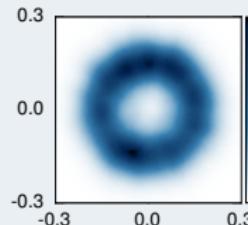
Problems

Technical Problems

- Chiral condensate always zero due to exact chiral symmetry and integration over fermions.
- Peak in susceptibility may indicate lattice artefact phase.

Coupling to Global Model [PoS(LATTICE 2015)050]

- Can obtain nice histograms of $\langle \bar{\psi} \psi \rangle$.
- Hard to recover Thirring model.
⇒ No reliable conclusion regarding N_f^{cr} ,
but likely $N_f^{cr} \leq 2$.



Scanning Larger Theory-Space

Why Coupled Four-Fermion Interactions?

- Easy to study chiral symmetry breaking in Gross-Neveu model.
⇒ Gain new insights into Thirring model and the larger theory space by coupling these two models.
- Test predictions from functional renormalization group.

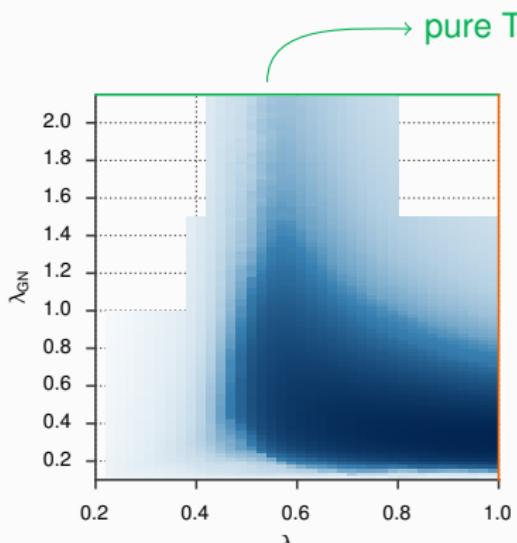
Interaction with $i\gamma_4\gamma_5$

$(\bar{\Psi} i\gamma_4\gamma_5 \Psi)^2$ is interesting because

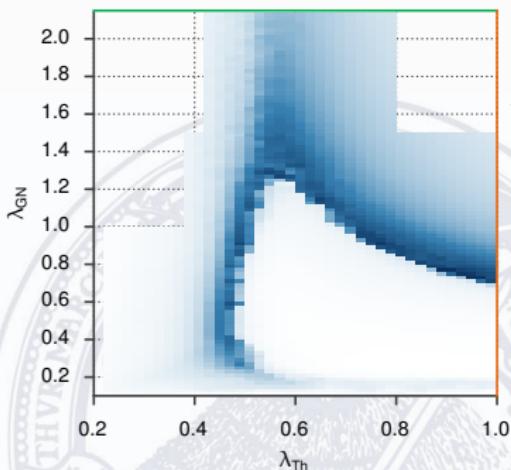
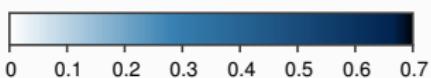
- same symmetry (and problems) as Thirring, while
- expected to be Gross-Neveu-like, corresponds to irreducible Gross-Neveu model.

Gross-Neveu and Thirring

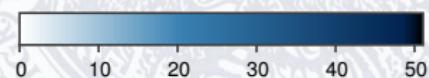
$$\lambda = \frac{1}{2g^2}, \text{ Lattice } 8 \times 7 \times 7, N_f = 1$$



Chiral Condensate

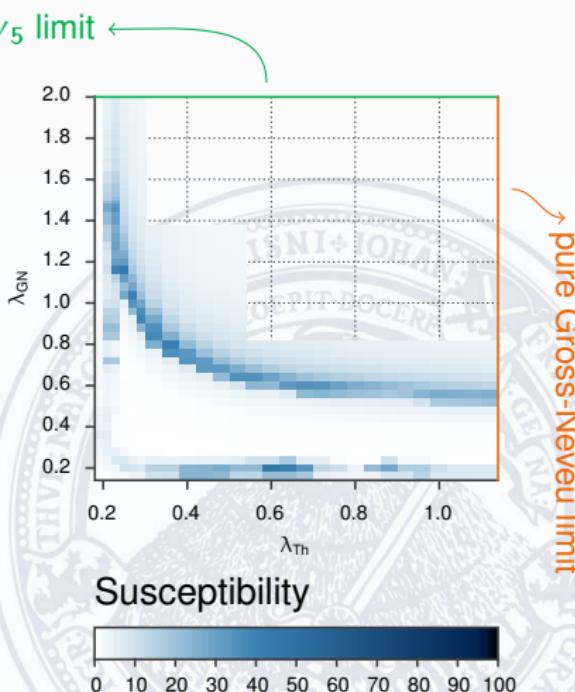
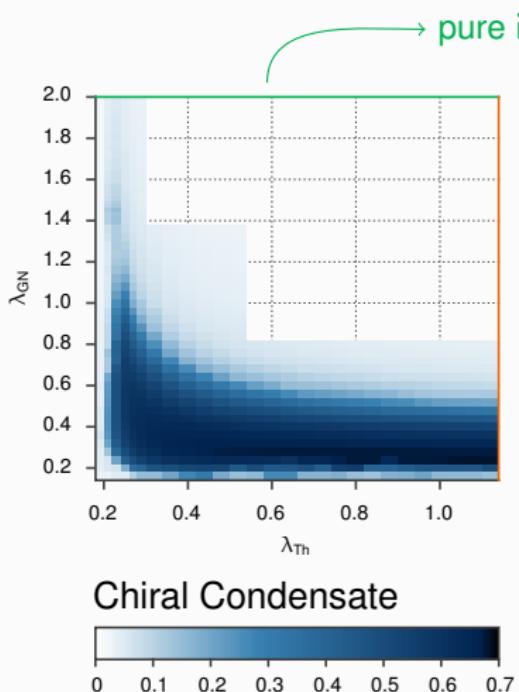


Susceptibility



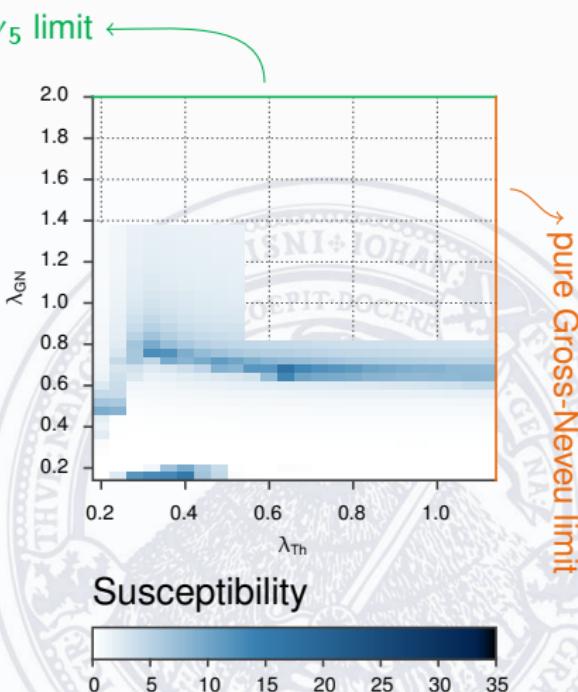
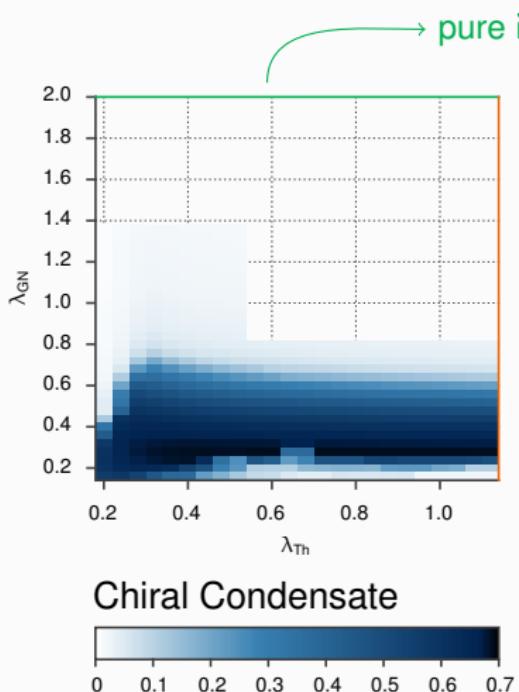
Gross-Neveu and $i\gamma_4\gamma_5$

$\lambda = \frac{1}{2g^2}$, Lattice $12 \times 11 \times 11$, $N_f = 1$



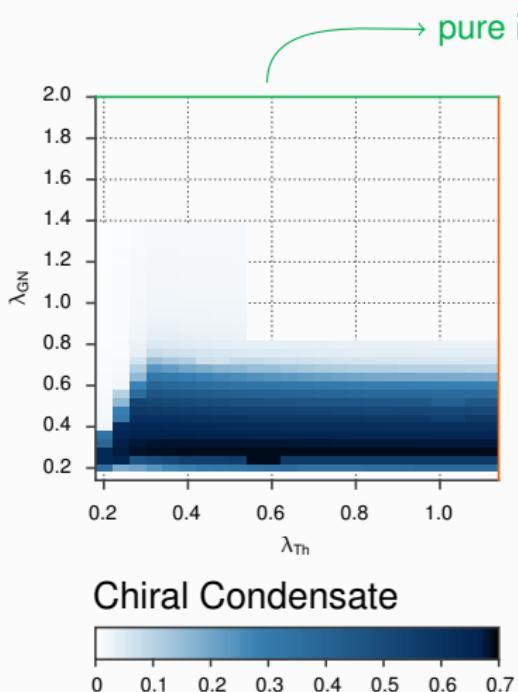
Gross-Neveu and $i\gamma_4\gamma_5$

$\lambda = \frac{1}{2g^2}$, Lattice $12 \times 11 \times 11$, $N_f = 2$

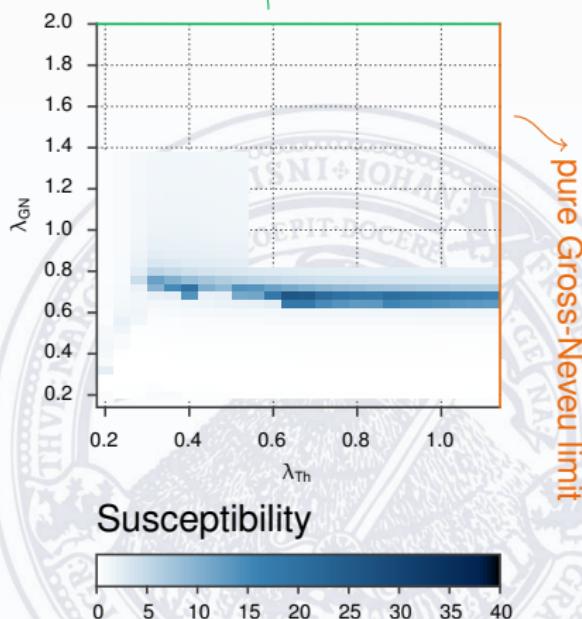


Gross-Neveu and $i\gamma_4\gamma_5$

$\lambda = \frac{1}{2g^2}$, Lattice $12 \times 11 \times 11$, $N_f = 3$



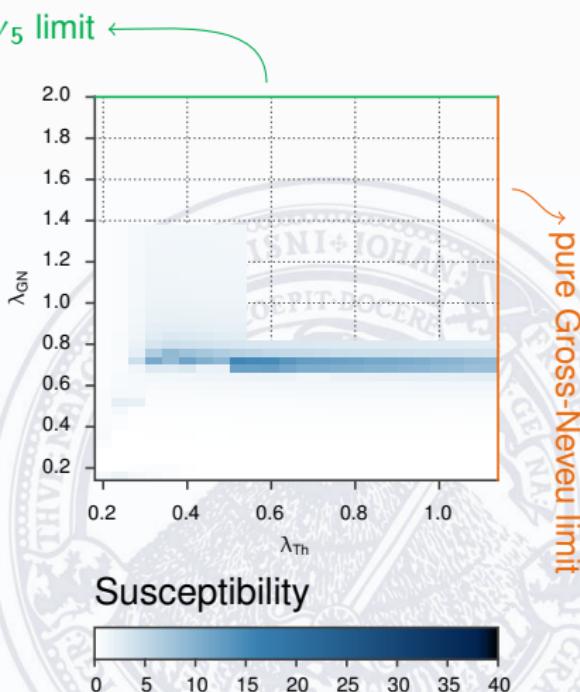
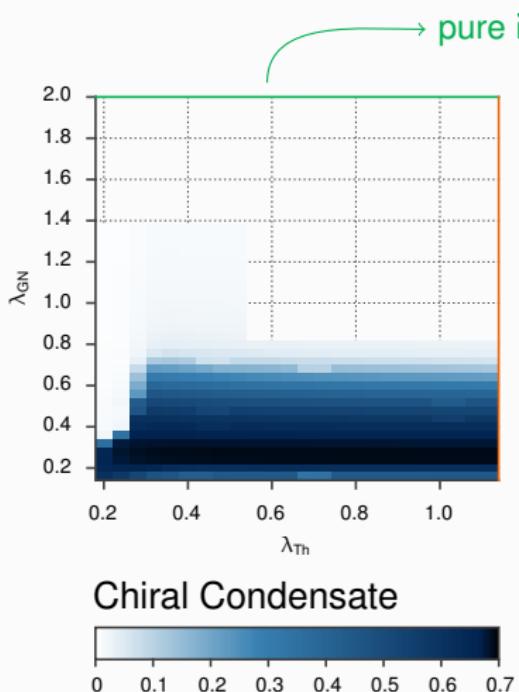
Chiral Condensate



Susceptibility

Gross-Neveu and $i\gamma_4\gamma_5$

$\lambda = \frac{1}{2g^2}$, Lattice $12 \times 11 \times 11$, $N_f = 4$



Fierz Identity (1) for Thirring Interaction

Goal: rewrite interaction term, use irreducible representation:

2-component spinors χ^a , $a = 1, \dots, 2N_f := N_{f,\text{irr}}$

$$(\bar{\chi}^a \sigma_\mu \chi^a) (\bar{\chi}^b \sigma^\mu \chi^b) = -(\bar{\chi}^a \chi^a) (\bar{\chi}^b \chi^b) - 2 (\bar{\chi}^a \chi^b) (\bar{\chi}^b \chi^a) \quad (1)$$



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Lagrangian after Hubbard-Stratonovich Transformations

ϕ real scalar field, T^{ab} hermitian, traceless matrix field

$$\mathcal{L}_1 = \bar{\chi}_a \underbrace{[(\not{d} + \phi) \delta^{ab} + T^{ab}]}_{=D_1^{ab}} \chi_b + \frac{N_{f,\text{irr}}}{4g^2} T_{ab} T^{ba} + \frac{N_{f,\text{irr}}^2}{2g^2(2 + N_{f,\text{irr}})} \phi^2$$

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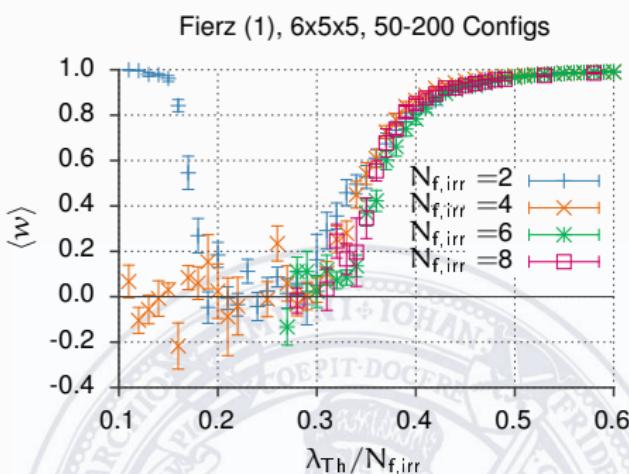
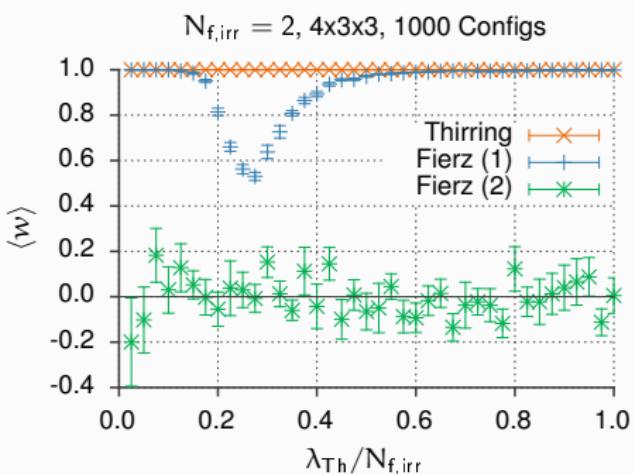
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D_1^{ab} has no special properties \Rightarrow complex eigenvalues

Sign Problem after Fierz Rearrangement



Thirring has no sign problem

Fierz (1) has a complex phase with $\langle w \rangle := \langle e^{-i \operatorname{Im} S} \rangle \approx 0$ for $\lambda_{Th} \in [0.2, 0.3]$

Fierz (2) another identity, imaginary eigenvalues, switches sign on nearly every update $\Rightarrow \langle w \rangle \approx 0 \forall \lambda_{Th}$

Fermion Bag Approach

Current approach to get informations about chiral symmetry breaking and solving the sign problem of [Fierz \(1\)](#):

⇒ Introduce a spin field $k_{xi}^{ab} \in \{0, 1\}$, integrate fermions, T^{ab} and ϕ :

Final Partition Sum

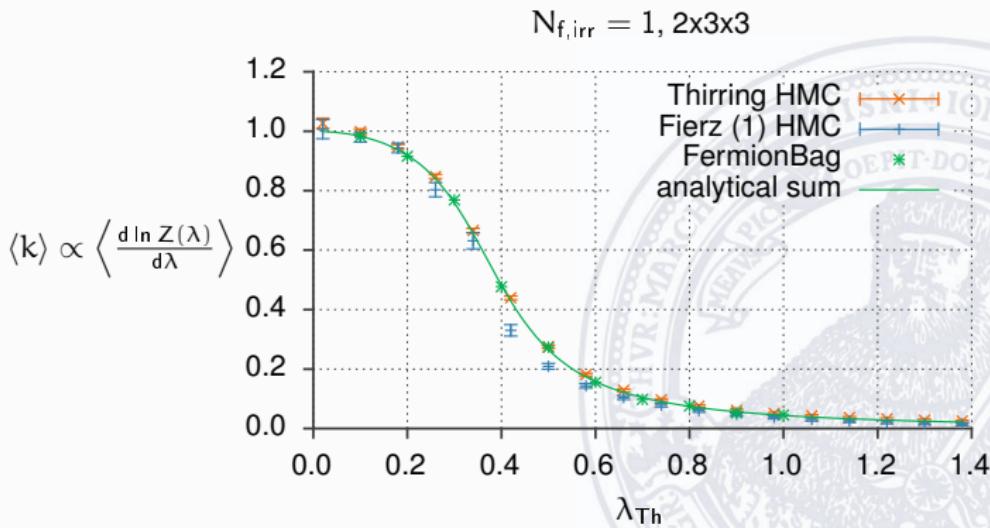
$$Z(\lambda) \propto \sum_k (-\lambda)^{-\frac{k}{2}} \det(i\cancel{\partial}[k]) 2^{\tilde{n}_x^2} \prod_x f(n_x^1, n_x^2)$$

- $\cancel{\partial}[k]$ is the SLAC operator matrix with columns and rows deleted corresponding to k .
- n_x^1 , n_x^2 and \tilde{n}_x^2 count certain entries of k_{xi}^{ab}
- $f(a, b)$ is a product of gamma- and confluent hypergeometric functions

Simulation Results

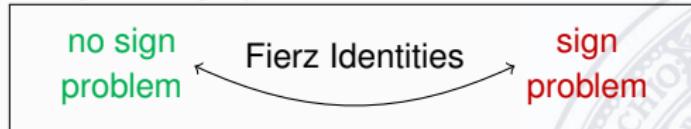
Simple Metropolis simulation for $N_{f,\text{irr}} = 1$:

- Agreement with **analytical** calculation and **original** formulation.
- Deviations for **Fierz(1)** (sign problem!).



The End

- The Thirring model shows chiral symmetry breaking for $N_f < N_f^{cr}$, for which there are many different predictions.
- Simulations with multiple four-fermion couplings may allow new insights.
- We get a sign problem after Fierz transformation.



- Fermion bag approach may solve the sign problem and allow access to the chiral condensate.

Thank you for your attention!