

Coupled-Channel Analysis of $D\pi$, $D\eta$ and $D_s\bar{K}$ Scattering using Lattice $\textcolor{red}{Q}\textcolor{green}{C}\textcolor{blue}{D}$

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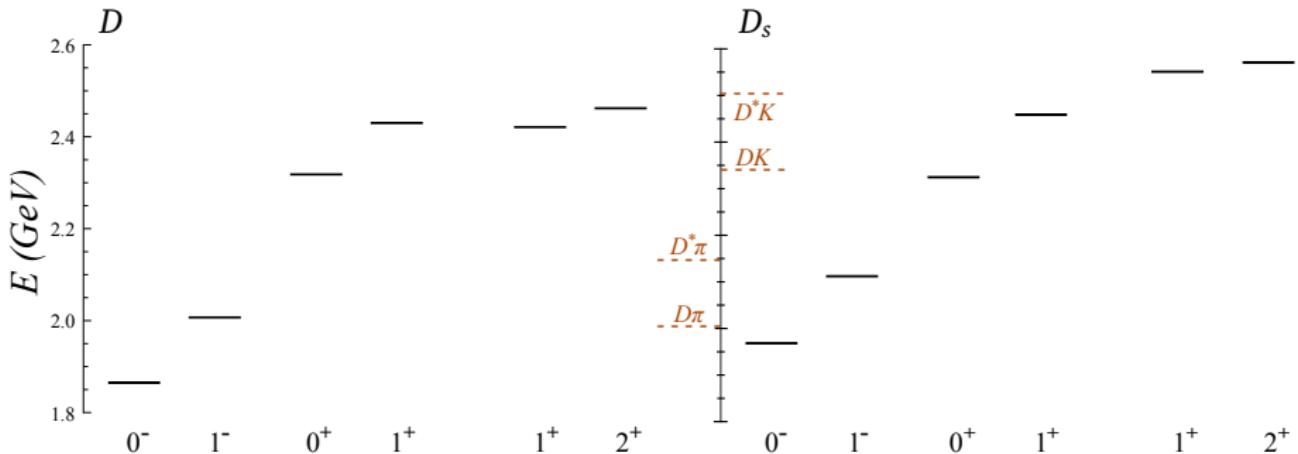


28th July 2016

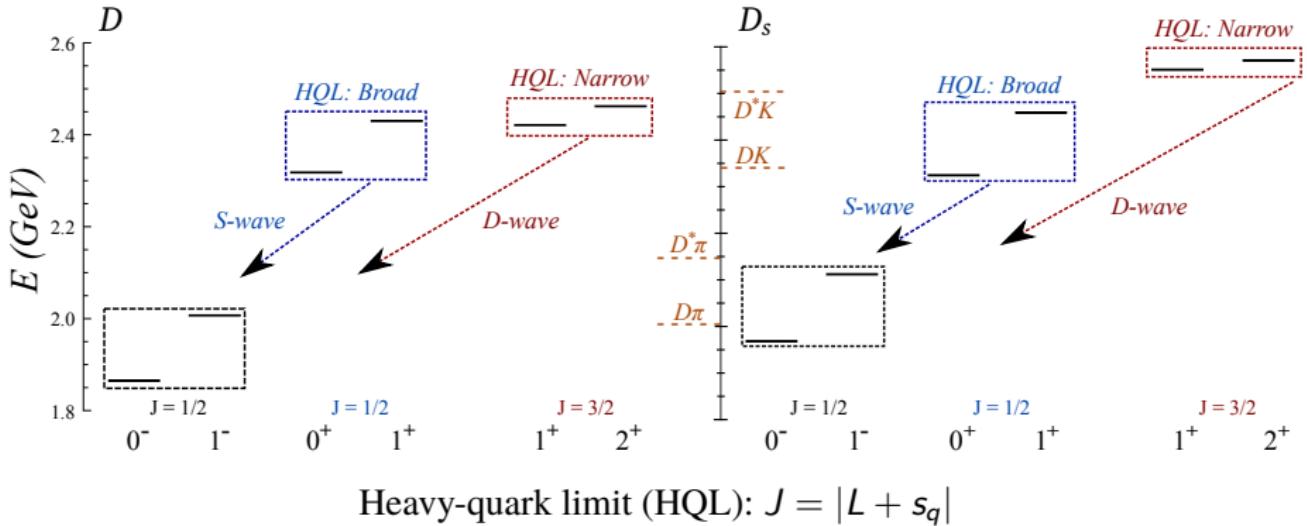
Lattice 2016, Southampton

Based on arXiv:1607.07093

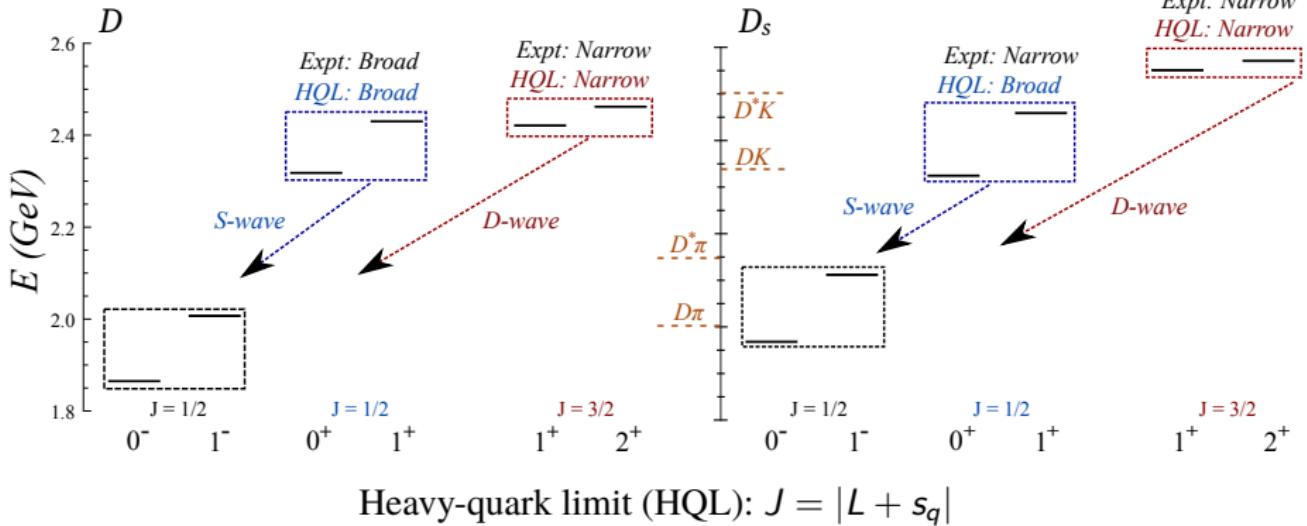
Heavy-Light Mesons



Heavy-Light Mesons



Heavy-Light Mesons



- Influence from thresholds?
- First principles understanding?

Meson	J^P	Mass (MeV)	Width (MeV)
$D_0^+(2400)$	0^+	2318(29)	267(40)
$D_{s0}^+(2317)$	0^+	2317.6(6)	< 3.8

Scattering on the Lattice

- Lüscher method \Rightarrow finite-volume energy levels map out phase shift
- Resonances decay to multiple channels \Rightarrow coupled-channel extensions

Coupled-Channel Lüscher Method

Finite-volume spectrum is given by energies solving

$$\det \left[t_{ij}^{-1}(E) + \mathcal{M}_{ij}(E, L) \right] = 0$$

- i, j label channels ($D\pi, D\eta, D_s\bar{K}$)
- $t_{ij}(E)$ is the **infinite-volume** scattering t -matrix . . . mixes channels
- $\mathcal{M}_{ij}(E, L)$ are known finite-volume functions . . . mixes partial waves

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Idea: Use the lattice QCD spectrum to obtain the infinite volume t -matrix

Scattering on the Lattice - Tackling Coupled-Channels

The bad news

- N coupled-channels $\Rightarrow t$ -matrix has $(N^2 + N)/2$ unknowns (per energy level)
- This is a very under-constrained problem

A work-around

- Parametrise the t -matrix using a “few” free parameters
- Use many more than a “few” lattice QCD energy levels to constrain $t(E)$

We want to

- Preserve S -matrix **unitarity**
- Analytically continue into the complex $s = E_{cm}^2$ plane
- Examine the **pole content** of the parametrised t -matrix (least analysis-dependent way of comparing between different calculations!)

Scattering on the Lattice - Tackling Coupled-Channels

A convenient choice is given by a **K-matrix** description

$$t_{ij}^{-1}(s) = \mathcal{K}_{ij}^{-1}(s) + \mathcal{I}_{ij}(s)$$

Preservation of unitarity

- Above kinematic threshold: $\text{Im}[\mathcal{I}_{ij}(s)] = -\delta_{ij}\rho_i(s)$
- Below kinematic threshold: $\text{Im}[\mathcal{I}_{ij}(s)] = 0$

Flexibility in $\text{Re}[\mathcal{I}_{ij}(s)]$

- Zero above kinematic threshold
- Chew-Mandelstam prescription

Most freedom comes from parametrising $\mathcal{K}_{ij}(s)$, for example,

$$\mathcal{K}_{ij}(s) = \sum_p \frac{g_i^{(p)} g_j^{(p)}}{m_p^s - s} + \sum_n \gamma_{ij}^{(n)} s^n$$

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Also use other parametrisations such as **Breit-Wigner** and **Effective Range Expansion**

Extracting Finite-Volume Energy Levels

Solve a GEVP for a matrix of 2-pt functions constructed using

① Distillation

- Improves overlap with low-lying states
- **Efficient** re-use of perambulators

② Large bases of interpolating operators for each irrep $[\vec{P}]\Lambda$

- “Single-meson” $\sim \bar{\psi}\Gamma D \dots \psi$ with up to 3 derivatives for $\vec{P} = 0$ (2 for $\vec{P} \neq 0$)
- “Two-meson” $\sim \Omega_{\vec{p}_1}^{(1)\dagger} \Omega_{\vec{p}_2}^{(2)\dagger}$ for **variationally optimised** “single-mesons” $\Omega_{\vec{p}}^{(i)\dagger}$

$[000]A_1^+$	$[001]A_1$	$[011]A_1$	$[111]A_1$
$D_{000}\pi_{000}$	$D_{001}\pi_{000}$	$D_{011}\pi_{000}$	$D_{111}\pi_{000}$
$D_{001}\pi_{00-1}$	$D_{000}\pi_{00-1}$	$D_{000}\pi_{-0-1-1}$	$D_{000}\pi_{-1-1-1}$
$D_{011}\pi_{0-1-1}$	$D_{011}\pi_{00-1}$	$D_{001}\pi_{-0-10}$	$D_{011}\pi_{-100}$
* $D_{111}\pi_{-1-1-1}$	$D_{001}\pi_{-0-1-1}$	$D_{011}\pi_{-1-10}$	$D_{001}\pi_{-1-10}$
$D_{000}\eta_{000}$	$D_{111}\pi_{0-1-1}$	$D_{111}\pi_{-00-1}$	$D_{002}\pi_{-1-1-1}$
$D_{001}\eta_{00-1}$	$D_{011}\pi_{-1-1-1}$	$D_{001}\pi_{-1-1-1}$	$D_{111}\pi_{00-2}$
* $D_{011}\eta_{0-1-1}$	$D_{002}\pi_{00-1}$	$D_{002}\pi_{-0-1-1}$	$D_{111}\eta_{000}$
$D_s 000 K_{000}$	$D_{001}\eta_{000}$	$D_{011}\eta_{000}$	$D_{000}\eta_{-1-1-1}$
$D_s 001 K_{00-1}$	$D_{000}\eta_{00-1}$	$D_{000}\eta_{0-1-1}$	$D_{011}\eta_{-100}$
* $D_s 011 K_{0-1-1}$	$D_{011}\eta_{00-1}$	$D_{001}\eta_{0-10}$	$D_{001}\eta_{-1-10}$
	$D_{001}\eta_{0-1-1}$	$D_{111}\eta_{00-1}$	$D_s 111 K_{000}$
	$D_{002}\eta_{00-1}$	$D_s 011 K_{000}$	$D_s 000 K_{-1-1-1}$
	$D_s 001 K_{000}$	$D_s 000 K_{0-1-1}$	$D_s 011 K_{-100}$
	$D_s 000 K_{00-1}$	$D_s 001 K_{0-10}$	$D_s 001 K_{-1-10}$
	$D_s 011 K_{00-1}$	$D_s 111 K_{00-1}$	
	$D_s 001 K_{0-1-1}$		
$(\bar{\psi}\Gamma\psi) \times 11$	$(\bar{\psi}\Gamma\psi) \times 32$	$(\bar{\psi}\Gamma\psi) \times 52$	$(\bar{\psi}\Gamma\psi) \times 37$

Ensemble Details

Calculations are performed on anisotropic $N_f = 2 + 1$ ensembles generated by the **Hadron Spectrum Collaboration**

Anisotropic actions

- Gauge: Symanzik-improved (tree-level tadpole improvement)
- Fermionic: Sheikholeslami-Wohlert (tree-level clover coefficients)

Anisotropy: $\xi = a_s/a_t \approx 3.5$, where $a_s \approx 0.12$ fm

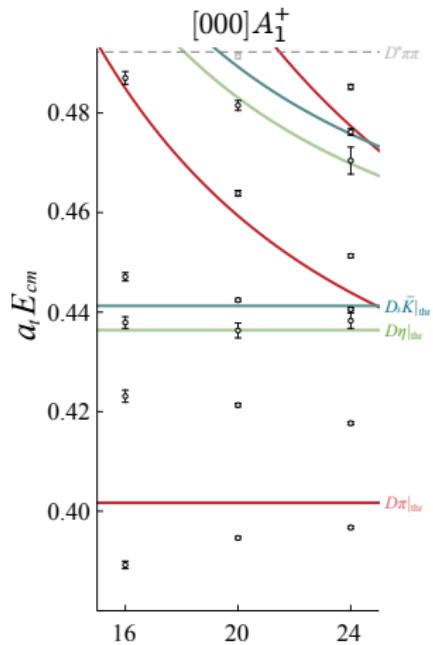
Pion mass: $M_\pi \approx 391$ MeV

Volumes:

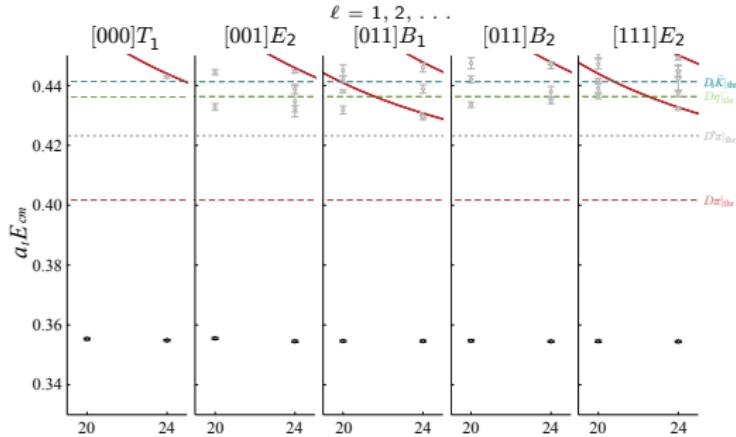
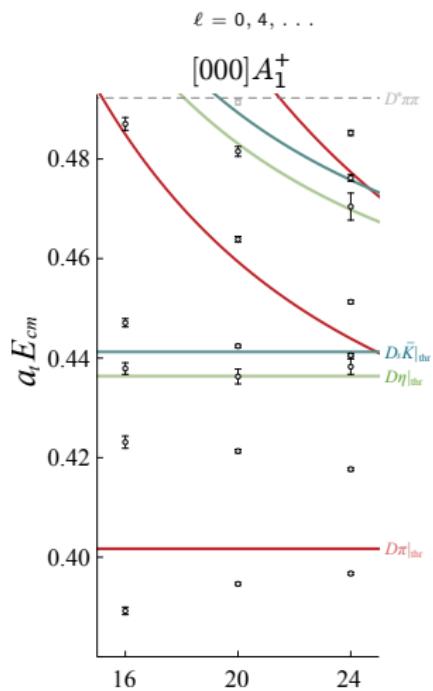
- $L = 16^3 \times 128 \Rightarrow M_\pi L \approx 3.8$ (479 cfgs)
- $L = 20^3 \times 128 \Rightarrow M_\pi L \approx 4.8$ (603 cfgs)
- $L = 24^3 \times 128 \Rightarrow M_\pi L \approx 5.7$ (553 cfgs)

Finite-Volume Energy Levels

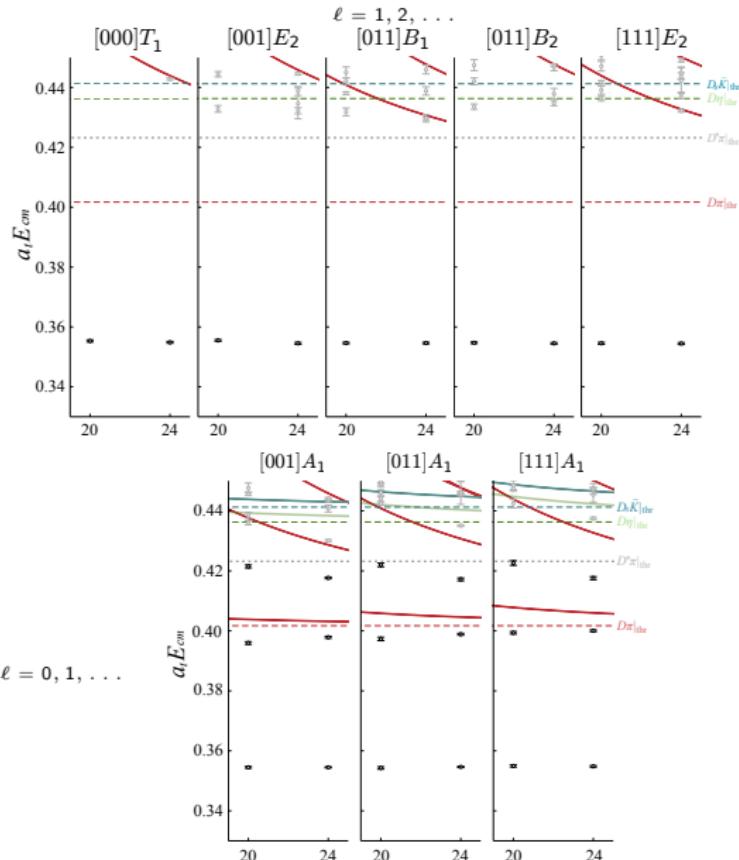
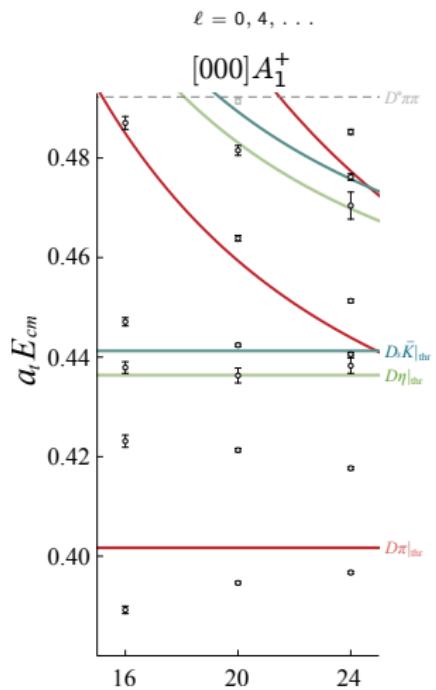
$\ell = 0, 4, \dots$



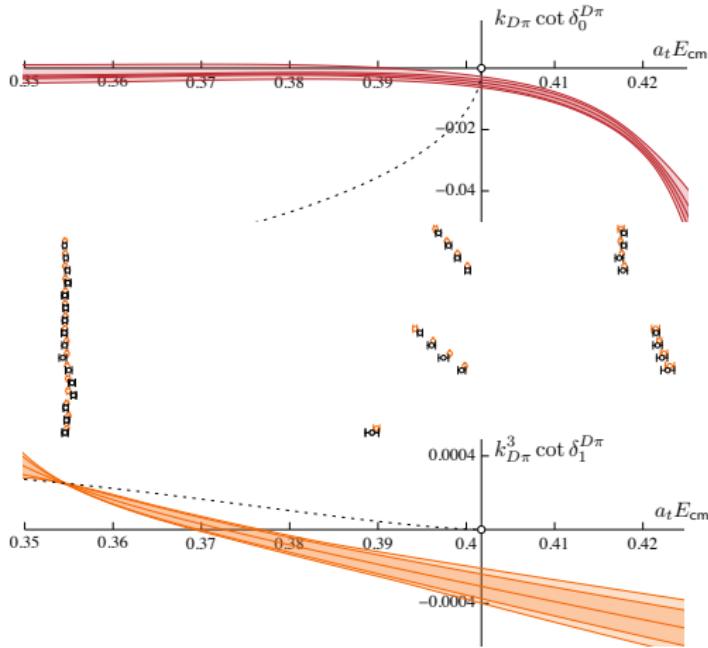
Finite-Volume Energy Levels



Finite-Volume Energy Levels



Elastic $D\pi$ Scattering



- All irreps with $\ell = 0, 1$
- Minimise χ^2 describing difference between *parametrised* and lattice QCD spectra
- “Pole + constant” \mathcal{K} -matrix in S and P -wave:
$$\chi^2/N_{\text{dof}} = \frac{44.2}{33-6} = 1.64$$
- $ik^{2\ell+1}$ (dashed curves) intersects at **location of sub-threshold poles**

Elastic $D\pi$ Scattering

Parametrisation

χ^2/N_{dof}

K-matrix with Chew-Mandelstam $I(s)$ & $K_1 = \frac{g_1^2}{m_1^2 - s} + \gamma_1$

$K = \frac{g^2}{m^2 - s} + \gamma^{(0)}$	1.64
$K = \frac{g^2}{m^2 - s} + \gamma^{(1)}s$	1.63
$K = \frac{(g^{(1)})^2 s}{m^2 - s} + \gamma^{(0)}$	1.64
$K = \frac{(g+g^{(1)})^2 s}{m^2 - s}$	1.66

K-matrix with Chew-Mandelstam $I(s)$ & $K_1 = \frac{g_1^2}{m_1^2 - s}$

$K = \frac{g^2}{m^2 - s} + \gamma^{(0)}$	1.82
$K = \frac{g^2}{m^2 - s} + \gamma^{(1)}s$	1.82

K-matrix with $I(s) = -i\rho(s)$ & $K_1 = \frac{g_1^2}{m_1^2 - s} + \gamma_1$

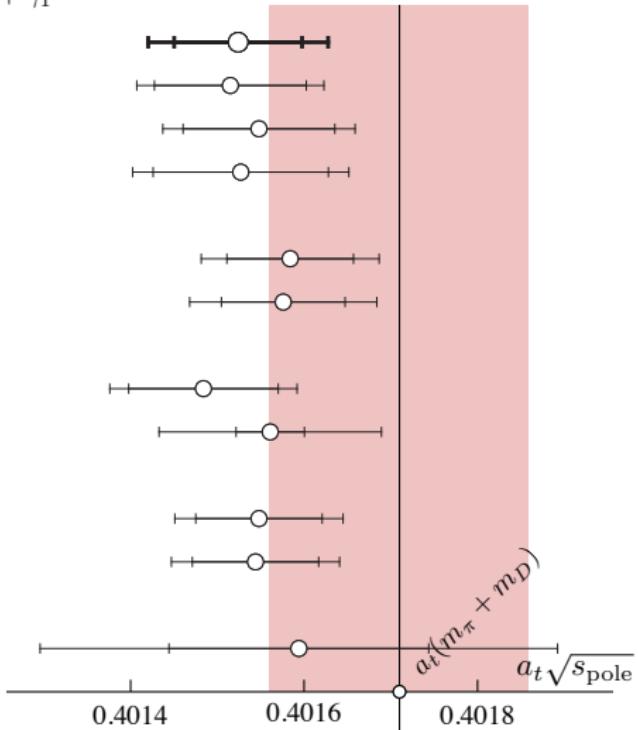
$K = \frac{g^2}{m^2 - s} + \gamma^{(0)}$	1.61
$K = \frac{g^2}{m^2 - s} + \gamma^{(1)}s$	1.64

K-matrix with $I(s) = -i\rho(s)$ & $K_1 = \frac{g_1^2}{m_1^2 - s}$

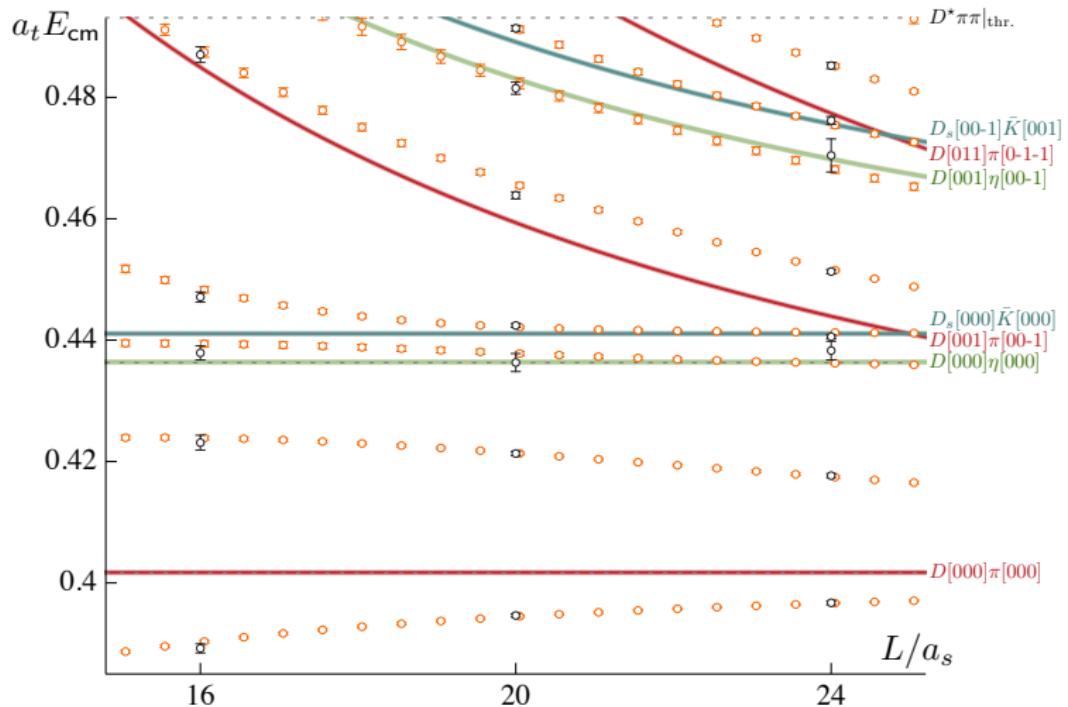
$K = \frac{g^2}{m^2 - s} + \gamma^{(0)}$	1.81
$K = \frac{g^2}{m^2 - s} + \gamma^{(1)}s$	1.80

Effective range expansion in $\ell = 0$ & $K_1 = \frac{g_1^2}{m_1^2 - s} + \gamma_1$

$$k_{D\pi} \cot \delta_0^{D\pi} = \frac{1}{a} + \frac{1}{2} r^2 k_{D\pi}^2 + P_2 k_{D\pi}^4 \quad 1.91$$

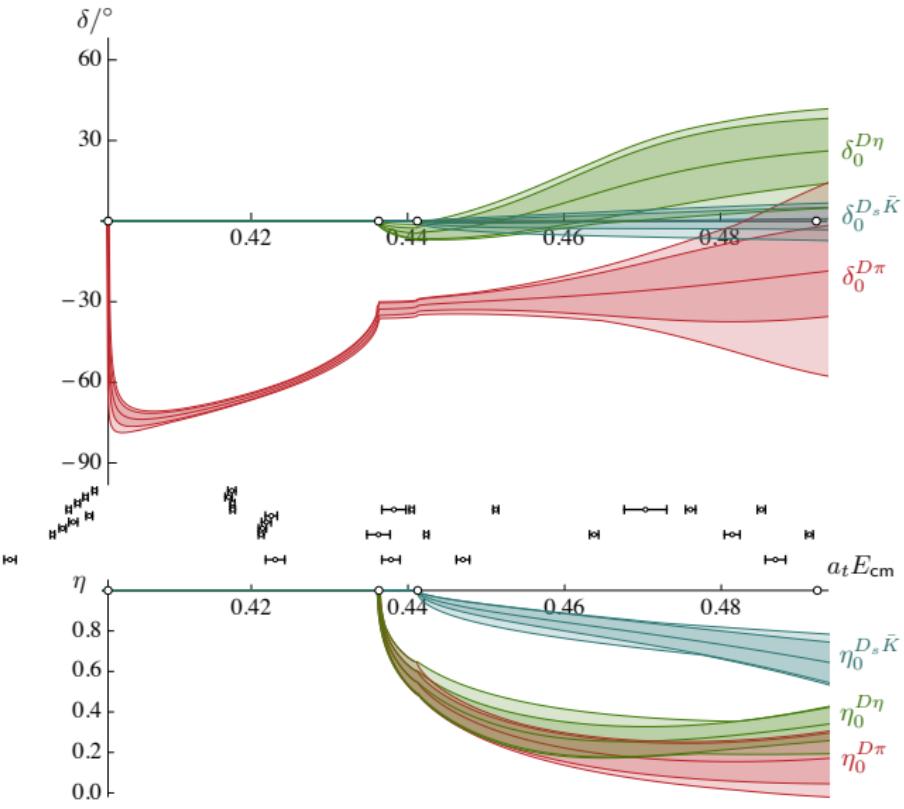


Coupled-Channel $D\pi$, $D\eta$ and $D_s\bar{K}$ Scattering



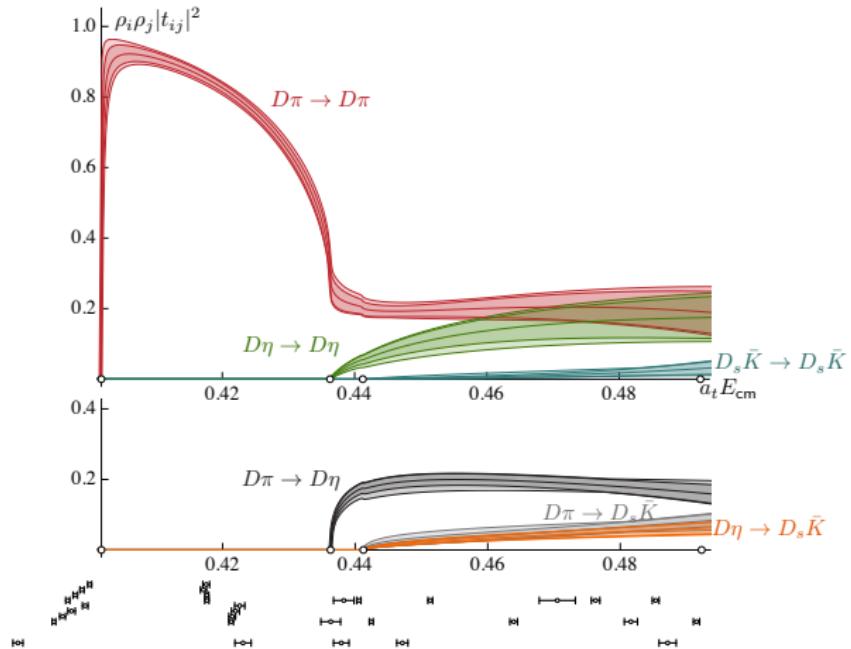
\mathcal{K} -matrix in S and P -wave including all $\ell = 0, 1$ irreps: $\chi^2/N_{\text{dof}} = \frac{61.6}{47-11} = 1.71$

Coupled-Channel $D\pi$, $D\eta$ and $D_s\bar{K}$ Scattering



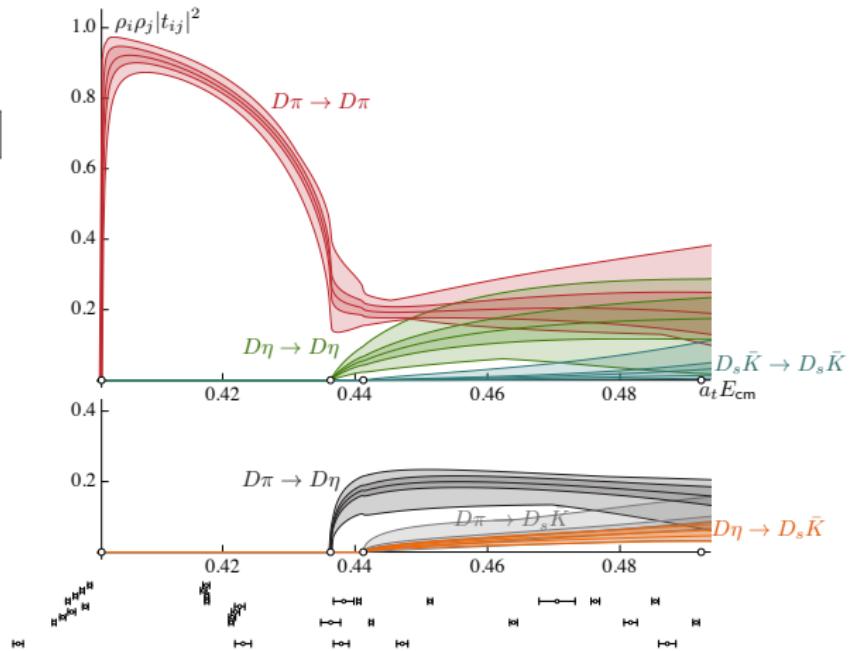
$$t_{jj} = \frac{\eta_j e^{2i\delta_j} - 1}{2i\rho_j}$$

Coupled-Channel $D\pi$, $D\eta$ and $D_s\bar{K}$ Scattering



Coupled-Channel $D\pi$, $D\eta$ and $D_s\bar{K}$ Scattering

All variations



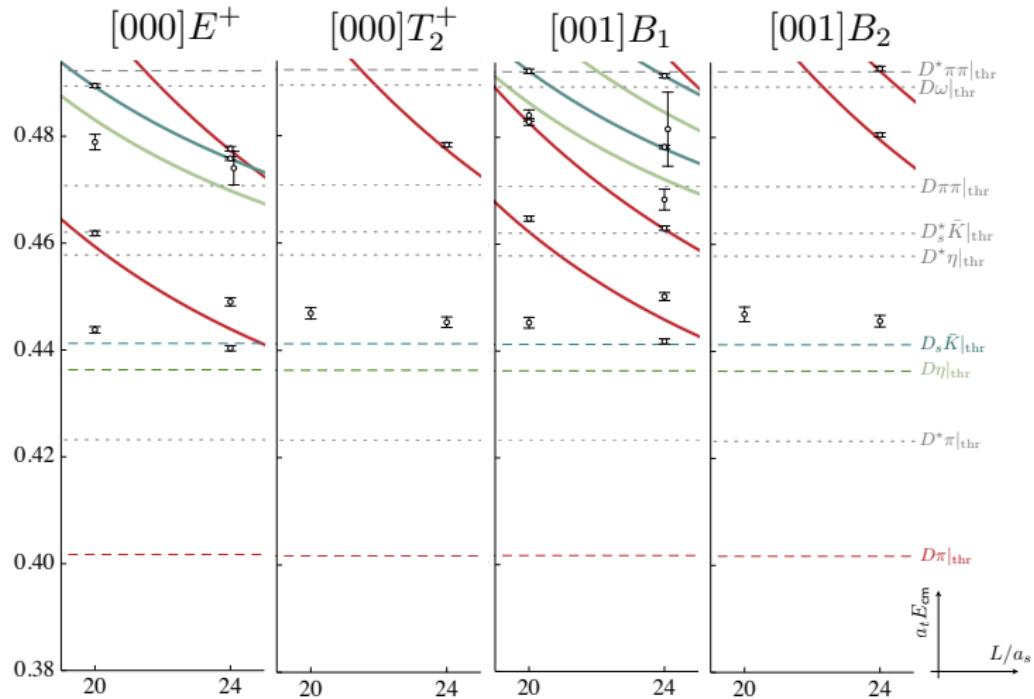
S -wave pole on **real-axis**: $\sqrt{s_{pole}} = 0.40161(15)$

$$t_{ij} \sim c_i c_j / (s_{pole} - s)$$

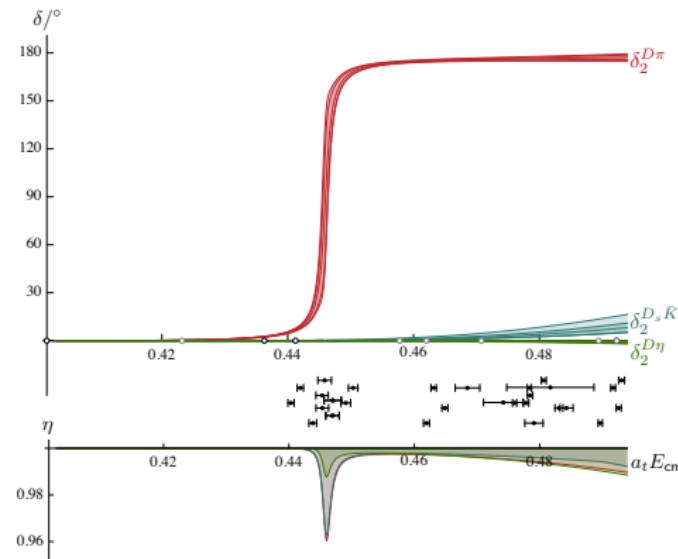
Couplings: $a_t c_{D\pi} = 0.097(28)$; $a_t c_{D\eta} = 0.077(23)$; $a_t c_{D_s\bar{K}} = 0.039(15)$.

Finite-Volume Energies

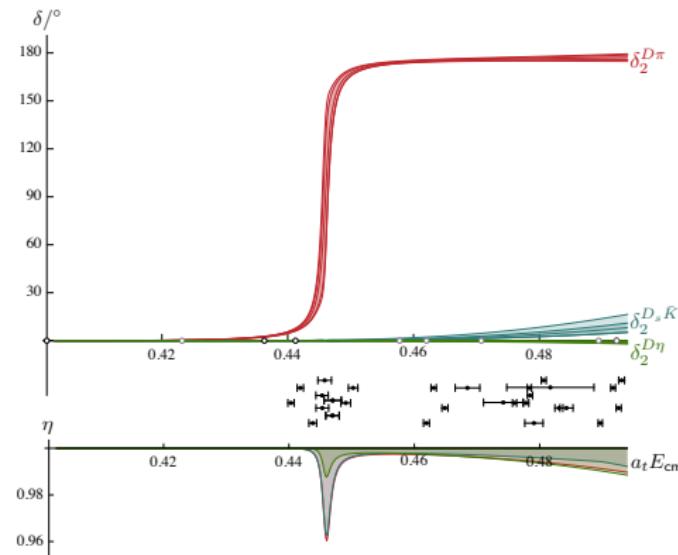
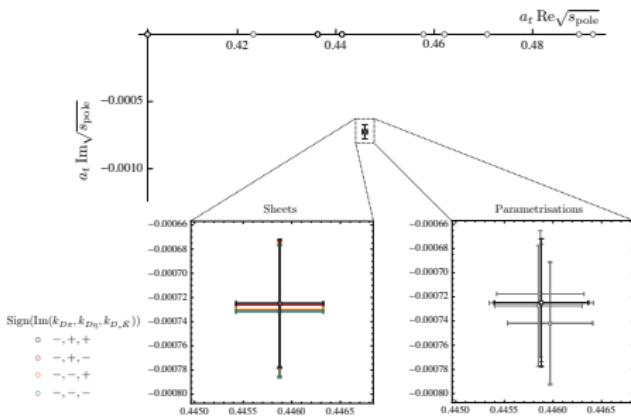
$$\ell = 2, 3, \dots$$



Coupled-Channel $D\pi$, $D\eta$ and $D_s\bar{K}$ Scattering

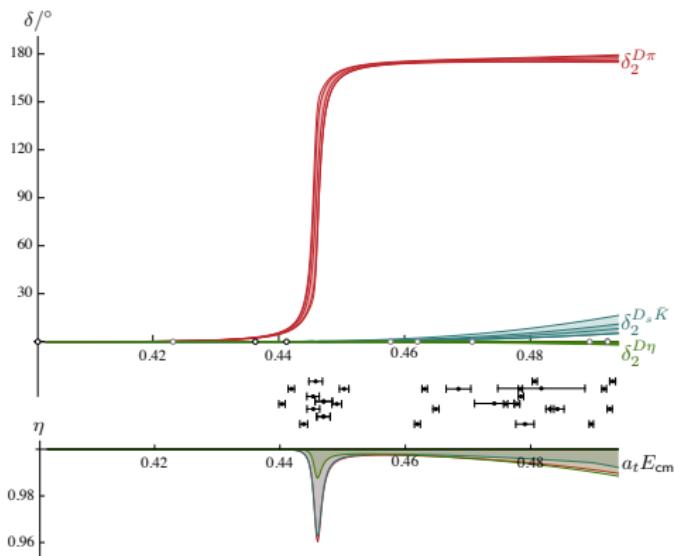
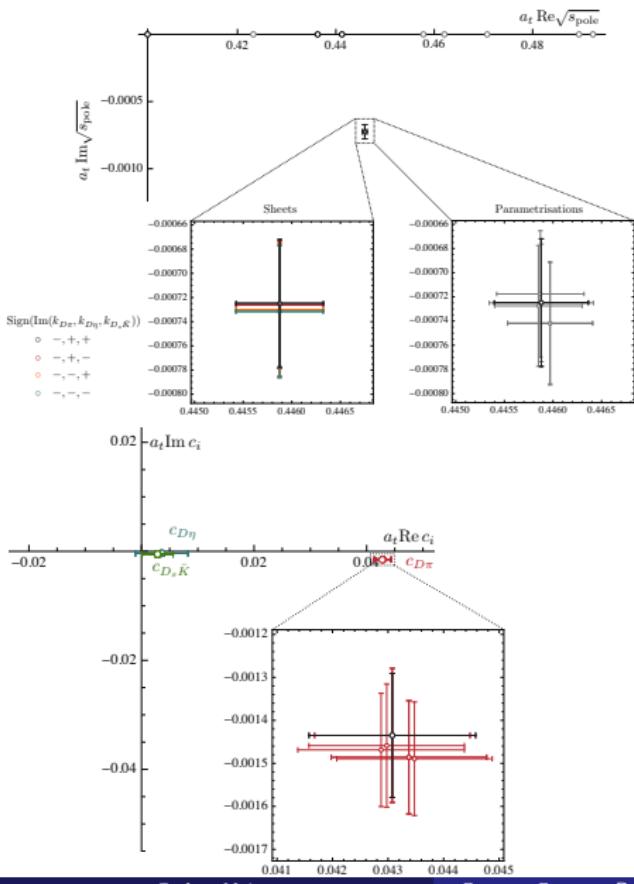


Coupled-Channel $D\pi$, $D\eta$ and $D_s\bar{K}$ Scattering



- D -wave resonance pole on all sheets with $\text{Im}[k_{D\pi}] < 0$

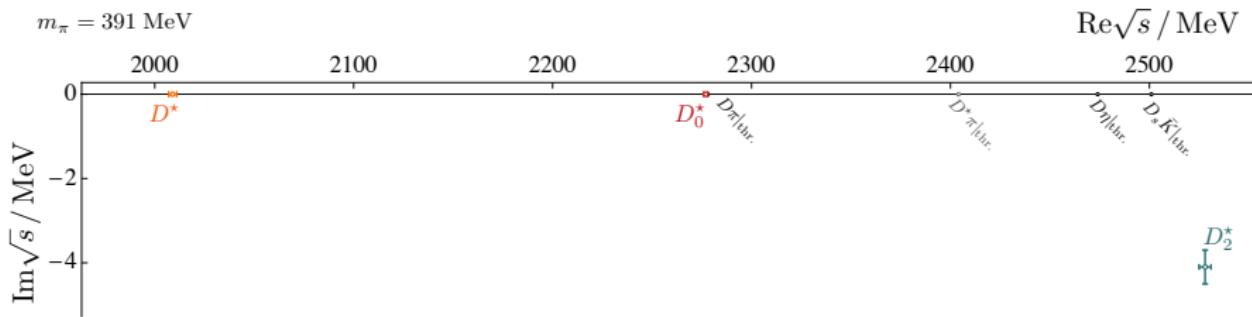
Coupled-Channel $D\pi$, $D\eta$ and $D_s\bar{K}$ Scattering



- D -wave resonance pole on all sheets with $\text{Im}[k_{D\pi}] < 0$
- Negligible coupling to $D\eta$ and $D_s\bar{K}$

Summary and Outlook

- First *ab-initio* coupled-channel scattering calculation including heavy quarks
- Summary of isospin-1/2 channel:

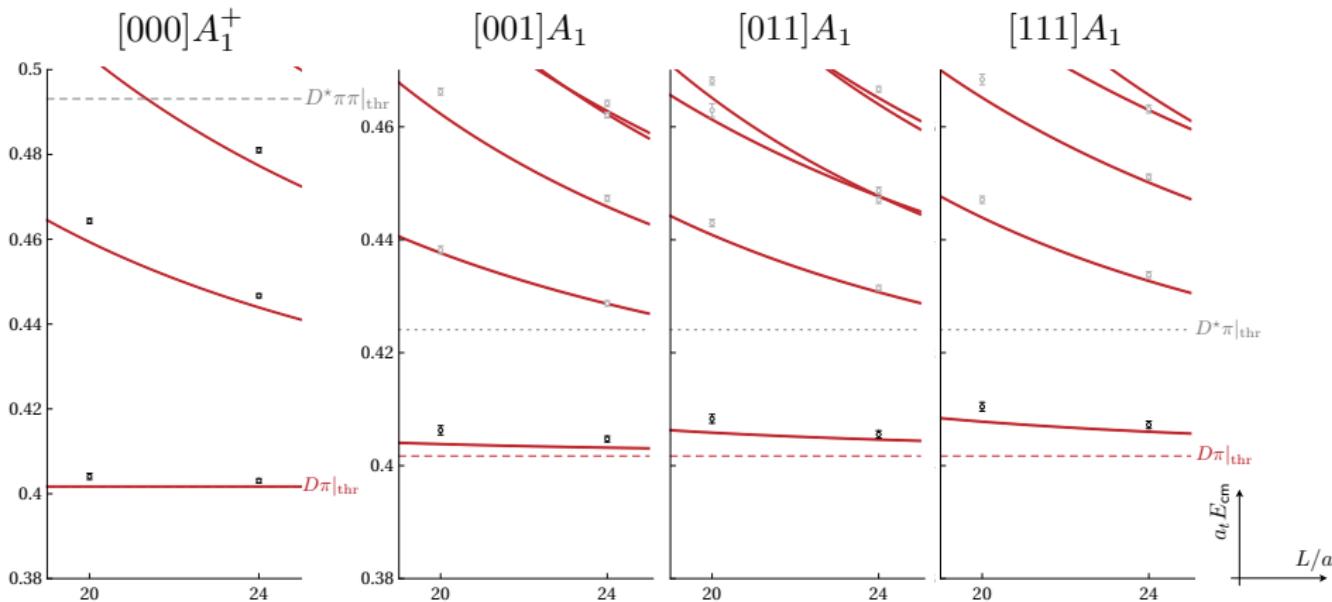


- Also explored isospin-3/2 elastic $D\pi$ scattering (ask if interested)

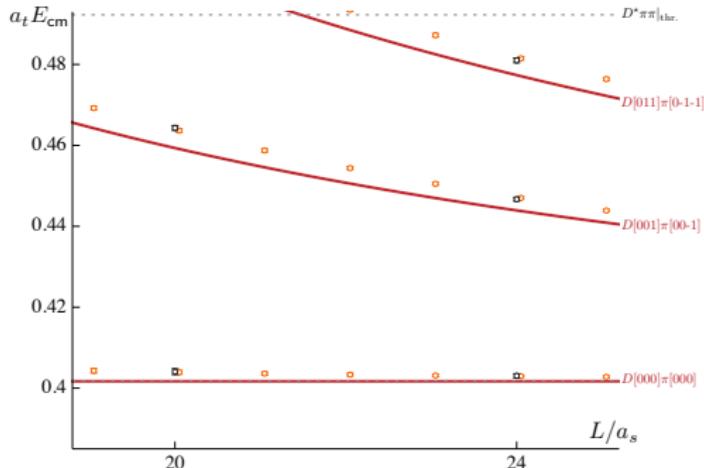
Outlook:

- DK scattering (C. Thomas next)
- Pole dependence on the light quark mass
- Begin to address the “XYZ’s” ?

Isospin-3/2: Finite-Volume Energy Levels



Isospin-3/2: $D\pi$ Phase-Shift



- $\chi^2/N_{\text{dof}} = 1.29$:
- $a_0 = -0.19(5)$ fm
- $r_0 = -0.9(4)$ fm.

