

# Moments of The Hadron Vacuum Polarization at Physical Point

The Connected HVP Contribution to Muon  $g - 2$

Kohtaro Miura (CPT, Aix-Marseille Univ.)  
The Budapest-Marseille-Wuppertal Collaboration

34th Lattice Conference, at Southampton UK, 26 July, 2016

# Table of Contents

## 1 Introduction

## 2 Result

- Moments of HVP: Connected Light-Component
- Moments of HVP: Connected Strange/Charm Component

## 3 Discussion: Finite Volume (FV) Effects

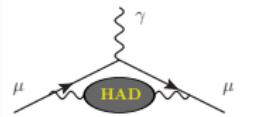
## 4 Summary and Perspective

# This Talk:

- We investigate **Moments** of the Hadronic Vacuum Polarization (HVP):

$$\hat{\Pi}(q^2) = \sum_n q^{2n} \Pi_n ,$$

$$\Pi_n[\mu\nu] = (-)^{n+1} \sum_x \frac{\hat{x}_\nu^{2(n+1)}}{(2n+2)!} \langle j_\mu(x) j_\mu(0) \rangle , j_\mu(x) = \sum_f Q_f (\bar{\psi} \gamma_\mu \psi)(x) .$$



- We have three different correlators (ss, ts, st) because of the box asymmetry ( $T \sim 1.5L$ ):

$$\Pi_{n,ss} = \frac{1}{6} \sum_{i \neq j} \Pi_n[ij] , \Pi_{n,ts} = \frac{1}{3} \sum_j \Pi_n[4j] , \Pi_{n,st} = \frac{1}{3} \sum_i \Pi_n[i4] . \quad (1)$$

- This talk focuses on the **Connected Contributions** to the moments:

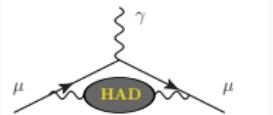
$$\Pi_n = \left( \frac{5}{9} \Pi_n^I + \frac{1}{9} \Pi_n^S + \frac{4}{9} \Pi_n^C \right) + \frac{1}{9} \Pi_n^{\text{disc}} .$$

# This Talk:

- We investigate **Moments** of the Hadronic Vacuum Polarization (HVP):

$$\hat{\Pi}(q^2) = \sum_n q^{2n} \Pi_n ,$$

$$\Pi_n[\mu\nu] = (-)^{n+1} \sum_x \frac{\hat{x}_\nu^{2(n+1)}}{(2n+2)!} \langle j_\mu(x) j_\mu(0) \rangle , j_\mu(x) = \sum_f Q_f (\bar{\psi} \gamma_\mu \psi)(x) .$$



- We have three different correlators (ss, ts, st) because of the box asymmetry ( $T \sim 1.5L$ ):

$$\Pi_{n,ss} = \frac{1}{6} \sum_{i \neq j} \Pi_n[ij] , \quad \Pi_{n,ts} = \frac{1}{3} \sum_j \Pi_n[4j] , \quad \Pi_{n,st} = \frac{1}{3} \sum_i \Pi_n[i4] . \quad (1)$$

- This talk focuses on the **Connected Contributions** to the moments:

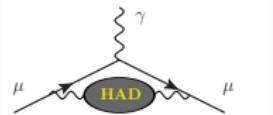
$$\Pi_n = \left( \frac{5}{9} \Pi_n^I + \frac{1}{9} \Pi_n^S + \frac{4}{9} \Pi_n^C \right) + \frac{1}{9} \Pi_n^{\text{disc}} .$$

# This Talk:

- We investigate **Moments** of the Hadronic Vacuum Polarization (HVP):

$$\hat{\Pi}(q^2) = \sum_n q^{2n} \Pi_n ,$$

$$\Pi_n[\mu\nu] = (-)^{n+1} \sum_x \frac{\hat{x}_\nu^{2(n+1)}}{(2n+2)!} \langle j_\mu(x) j_\mu(0) \rangle , j_\mu(x) = \sum_f Q_f (\bar{\psi} \gamma_\mu \psi)(x) .$$



- We have three different correlators (ss, ts, st) because of the box asymmetry ( $T \sim 1.5L$ ):

$$\Pi_{n,ss} = \frac{1}{6} \sum_{i \neq j} \Pi_n[ij] , \quad \Pi_{n,ts} = \frac{1}{3} \sum_j \Pi_n[4j] , \quad \Pi_{n,st} = \frac{1}{3} \sum_i \Pi_n[i4] . \quad (1)$$

- This talk focuses on the **Connected Contributions** to the moments:

$$\Pi_n = \left( \frac{5}{9} \Pi_n^I + \frac{1}{9} \Pi_n^S + \frac{4}{9} \Pi_n^C \right) + \frac{1}{9} \Pi_n^{\text{disc}} .$$

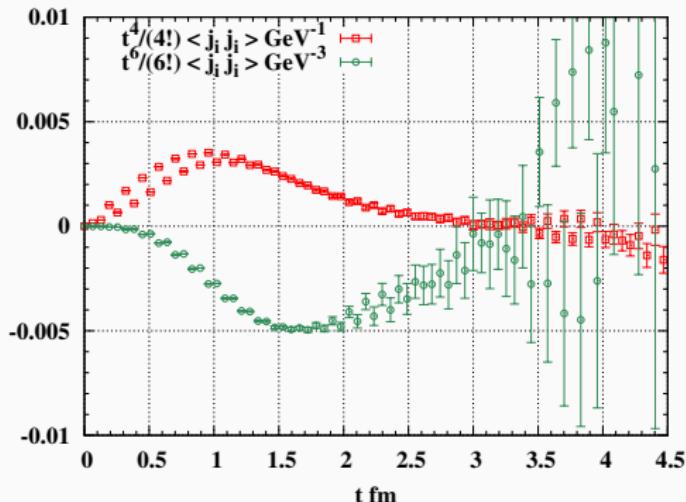
# Simulation Setup

$a[\text{fm}]$	$N_t$	$N_s$	#traj.	$M_\pi[\text{MeV}]$	$M_K[\text{MeV}]$	#SRC (ud,s,c)
0.134	64	48	17000	$\sim 131$	$\sim 479$	(768, 128, 64)
0.118	96	56	15000	$\sim 132$	$\sim 483$	(768, 64, 64)
0.111	84	56	15000	$\sim 133$	$\sim 483$	(768, 64, 64)
0.095	96	64	15000	$\sim 133$	$\sim 488$	(768, 64, 64)
0.078	128	80	10000	$\sim 133$	$\sim 488$	(768, 64, 64)
0.064	144	96	04500	$\sim 133$	$\sim 490$	(768, 64, 64)

## State Of The Art

- Simulations at Physical Pion/Kaon Mass with 6 lattice spacings.
- $(L, T) \sim (6, 9 - 12) \text{ fm}$ .
- $N_f = (2 + 1 + 1)$ : Two degenerate light quarks with strange and charm.

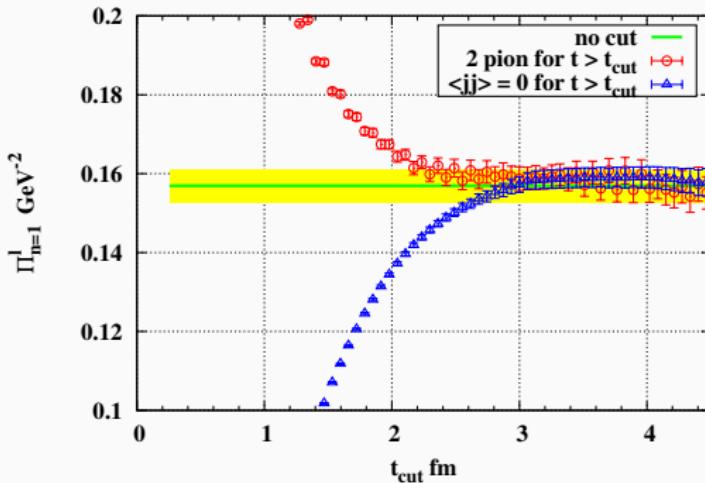
# Moment Density



$$\Pi'_{n=1,st} = \sum_t \left[ \frac{t^4}{4!} \sum_{x,y,z} \frac{\sum_i}{3} \langle j_i(x)j_i(0) \rangle \right], \quad (2)$$

$$\Pi'_{n=2,st} = \sum_t \left[ \frac{t^6}{6!} \sum_{x,y,z} \frac{\sum_i}{3} \langle j_i(x)j_i(0) \rangle \right]. \quad (3)$$

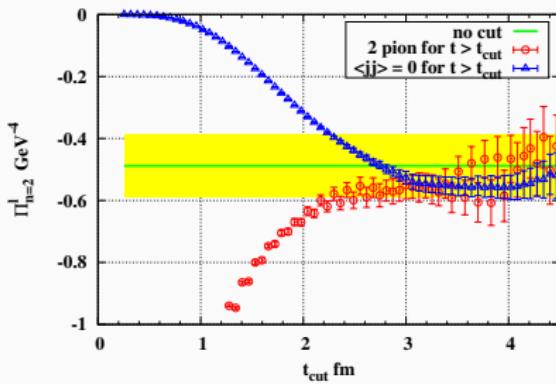
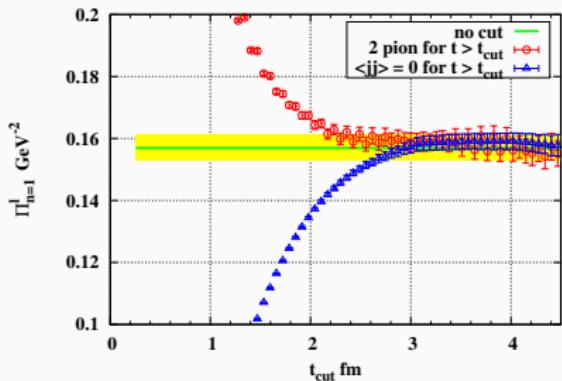
# IR-CUT Dependence of Moment I



- Upper bound: For  $t > t_{\text{cut}}$  (C.Lehner '15),  

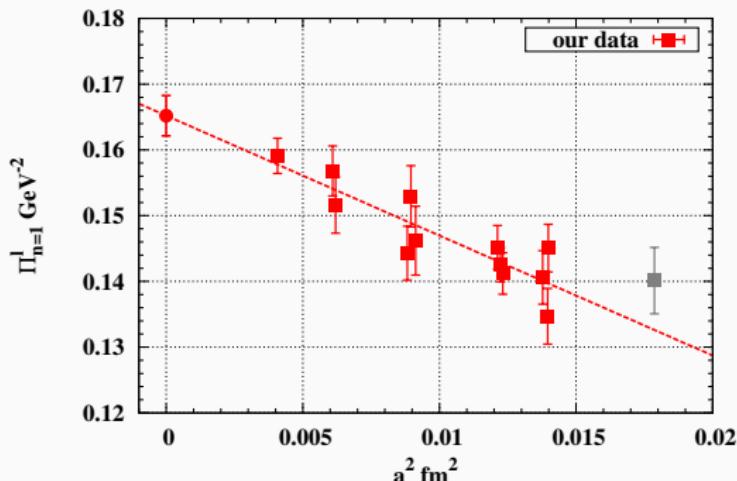
$$\langle jj \rangle_{st}(t > t_{\text{cut}}) = \langle jj \rangle_{st}(t_{\text{cut}}) \frac{\cosh[2E_{2\pi}(T/2-t)]}{\cosh[2E_{2\pi}(T/2-t_{\text{cut}})]}, E_{2\pi} = \sqrt{M_\pi^2 + (\frac{2\pi}{L})^2}.$$
- Lower bound: For  $t > t_{\text{cut}}$ ,  $\langle jj \rangle_{st}(t \geq t_{\text{cut}}) = 0.0.$

# IR-CUT Dependence of Moment II



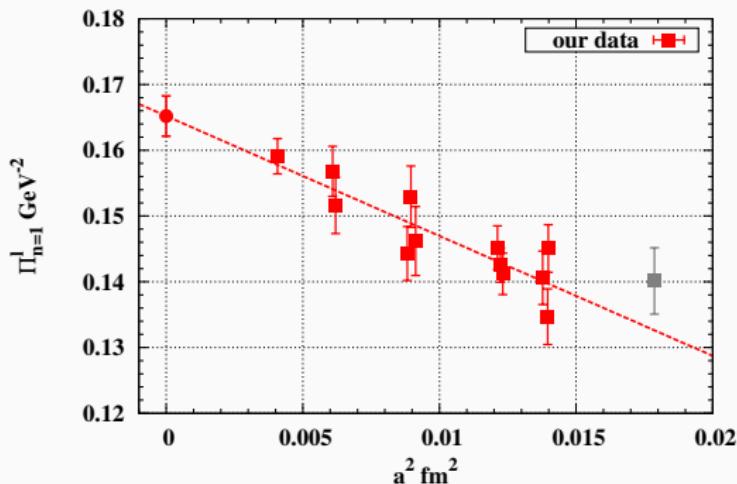
We adopt  $t_{\text{cut}} = 3.1 \text{ fm}$ ,  
 for which the **upper** and lower **bounds** meet within the statistical uncertainty.

# Continuum Extrapolation of $\Pi_{n=1}^I$



$$\Pi_{n=1}^I = \frac{1}{3} (\Pi_{1,ss}^I + \Pi_{1,ts}^I + \Pi_{1,st}^I) . \quad (4)$$

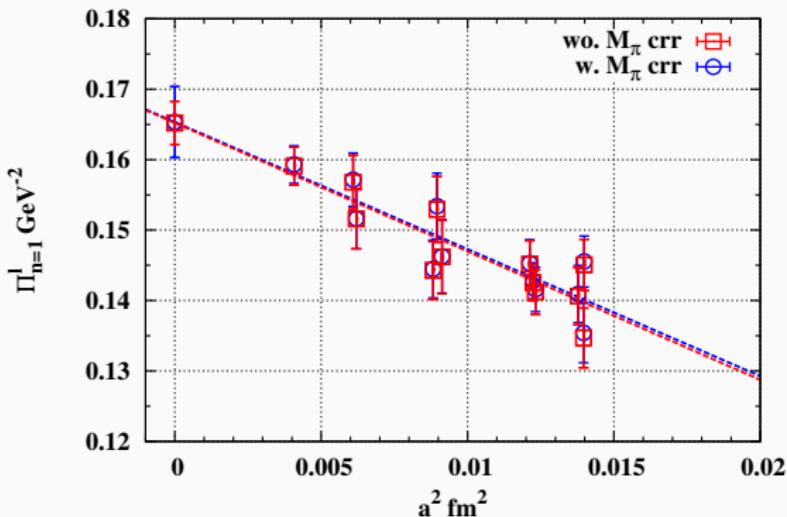
# Continuum Extrapolation of $\Pi_{n=1}^l |_{a^2 \rightarrow 0}$



$$F(C_\Pi^{(2)}, A, B, C_{M_\pi}, C_{M_K}) = \frac{C_\Pi^{(2)}}{a^2} \frac{1 + Aa^2 + \dots}{1 + Ba^2 + \dots} (1 + C_{M_\pi} \Delta M_\pi + C_{M_K} \Delta M_K).$$

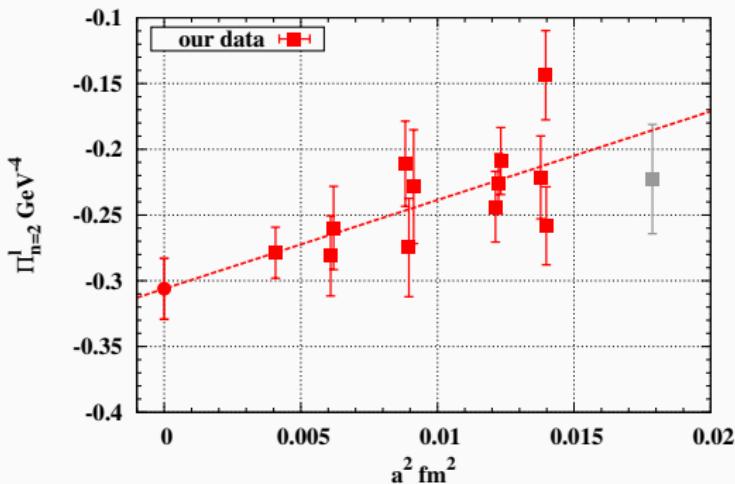
$$\Pi_{n=1}^l |_{a^2 \rightarrow 0} = 0.1652(31), \quad \chi^2/\text{d.o.f.} = 24.3/20$$

# Continuum Extrapolation of $\Pi_{n=1}^l \parallel$



$$F(C_\Pi^{(2)}, A, B, C_{M_\pi}, C_{M_K}) = \frac{C_\Pi^{(2)}}{a^2} \frac{1 + Aa^2 + \dots}{1 + Ba^2 + \dots} (1 + C_{M_\pi} \Delta M_\pi + C_{M_K} \Delta M_K) .$$

# Continuum Extrapolation of $\Pi_{n=2}^I$



$$F(C_\Pi^{(4)}, A, B, C_{M_\pi}, C_{M_K}) = \frac{C_\Pi^{(4)}}{a^4} \frac{1 + Aa^2 + \dots}{1 + Ba^2 + \dots} (1 + C_{M_\pi} \Delta M_\pi + C_{M_K} \Delta M_K) .$$

$$\Pi_{n=2}^I|_{a^2 \rightarrow 0} = -0.306(23) .$$

# Continuum Extrapolation of $\Pi_{n=1,2}^s$

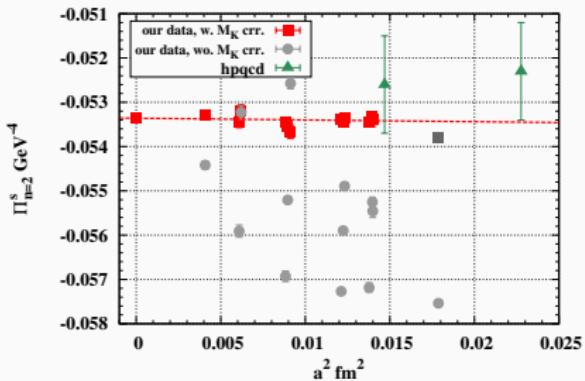
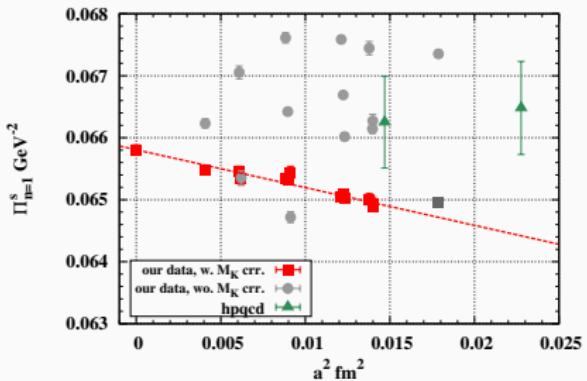
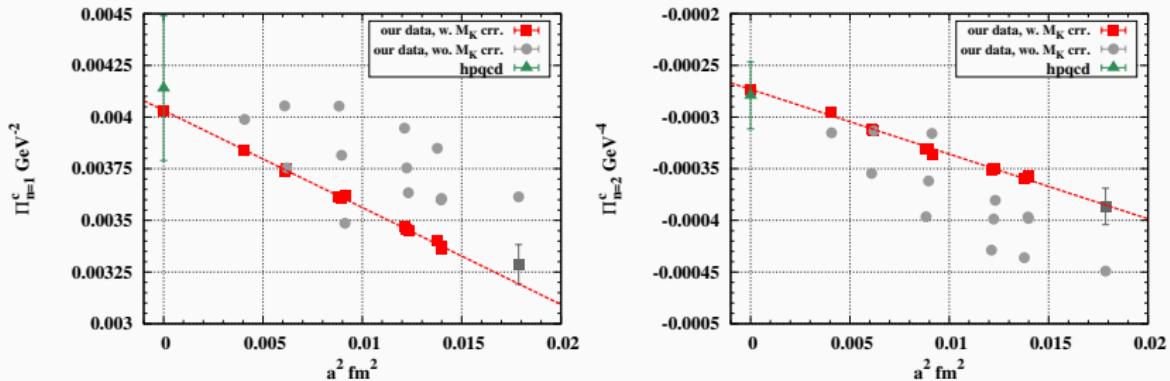


Figure: Red-squares = Our data. Green-triangles = HPQCD, 1403.1778.

$$F(C_\Pi^{(2,4)}, A, B, C_{M_\pi}, C_{M_K}) = \frac{C_\Pi}{a^{2,4}} \frac{1 + Aa^2 + \dots}{1 + Ba^2 + \dots} (1 + C_{M_\pi} \Delta M_\pi + C_{M_K} \Delta M_K).$$

$$\Pi_{n=1}^s|_{a^2 \rightarrow 0} = 0.0658(1), \quad \Pi_{n=2}^s|_{a^2 \rightarrow 0} = -0.0534(2), \quad \chi^2/\text{d.o.f.} = 20.9/18$$

# Continuum Extrapolation of $\Pi_{n=1,2}^c$



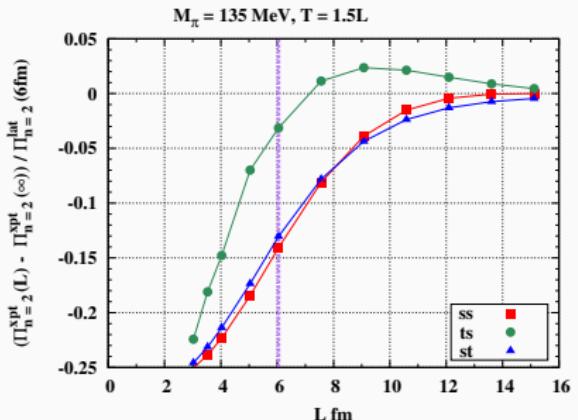
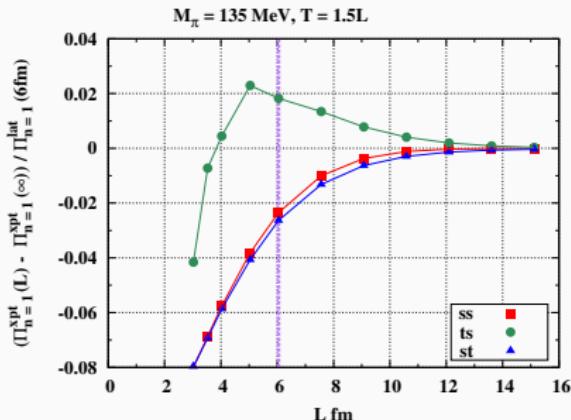
**Figure:** Red-squares = Our data. Green-triangles = HPQCD, 1208.2855.

$$F(C_\Pi^{(2,4)}, A, B, C_{M_\pi}, C_{M_K}) = \frac{C_\Pi}{a^{2,4}} \frac{1 + Aa^2 + \dots}{1 + Ba^2 + \dots} (1 + C_{M_\pi} \Delta M_\pi + C_{M_K} \Delta M_K).$$

$$\Pi_{n=1}^c|_{a^2 \rightarrow 0} = 0.00403(2), \quad \Pi_{n=2}^c|_{a^2 \rightarrow 0} = -2.73(2) \times 10^{-4}.$$

# FV via Box Asymmetry, XPT Estimate for Various $L$ I

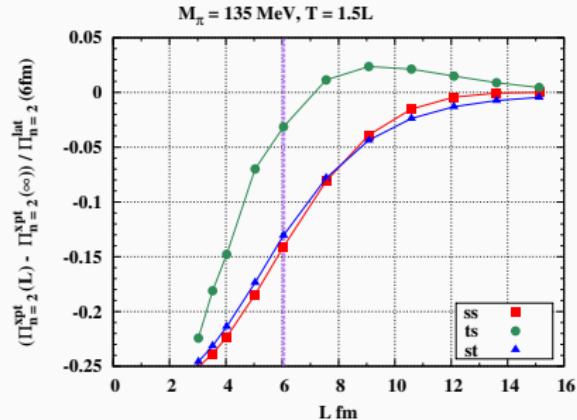
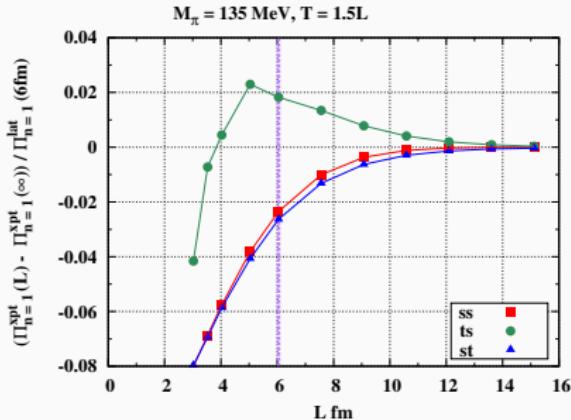
c.f. Aubin et.al., PRD (2016).



$$\begin{aligned} \Delta_{n=1,2}^i(L) &= [\Pi_{n=1,2}^{xpt,i}(L) - \Pi_{n=1,2}^{xpt}(\infty)] , \quad i = ss, ts, st , \\ \frac{\Delta_n^i(L=6\text{fm})}{\Pi_n^{\text{lat},i}} &\sim \begin{cases} 2\% & (\text{for the 1st moment, } n=1) , \\ 10\% & (\text{for the 2nd moment, } n=2) . \end{cases} \end{aligned} \quad (5)$$

## FV via Box Asymmetry, XPT Estimate for Various $L$ II

c.f. Aubin et.al., PRD (2016).



$$\Delta_{n=1,2}^i(L) = [\Pi_{n=1,2}^{xpt,i}(L) - \Pi_{n=1,2}^{xpt}(\infty)] , \quad i = ss, ts, st ,$$

$$FV.(L) \pm dFV.(L) = [\max\{\Delta_i\} + \min\{\Delta_i\}] / 2 \pm [\max\{\Delta_i\} - \min\{\Delta_i\}] / 2$$

$$\xrightarrow[L \rightarrow 6 \text{ fm}]{} \begin{cases} 0.0006(22) & (\text{for the 1st moment, } n = 1) , \\ -0.015(19) & (\text{for the 2nd moment, } n = 2) . \end{cases} \quad (6)$$

# Summary Table of Moments

	$\Pi_1[\text{GeV}^{-2}]$	$\Pi_2[\text{GeV}^{-4}]$
light	0.1657(16)(18)	-0.297(10)(05)
strange	$6.57(1)(2) \times 10^{-2}$	$-5.32(1)(3) \times 10^{-2}$
charm	$4.04(1)(1) \times 10^{-3}$	$-2.68(1)(4) \times 10^{-4}$
disconnected	$-1.5(2)(1) \times 10^{-2}$	$4.6(1.0)(0.4) \times 10^{-2}$
$I = 0$	0.0166(2)(2)	-0.017(1)(1)
$I = 1$	0.0828(8)(9)	-0.148(5)(2)
<b>total</b>	<b>0.0995(9)(10)</b>	<b>-0.166(6)(3)</b>

**Table:** Preliminary results on the first two moments of the HVP function.

TOTAL ERROR: 1.4% for  $\Pi_1$ , and 4.0% for  $\Pi_2$ .

## Summary Table of Moments

	$\Pi_1[\text{GeV}^{-2}]$	$\Pi_2[\text{GeV}^{-4}]$
light	0.1657(16)(18)	-0.297(10)(05)
strange	$6.57(1)(2) \times 10^{-2}$	$-5.32(1)(3) \times 10^{-2}$
charm	$4.04(1)(1) \times 10^{-3}$	$-2.68(1)(4) \times 10^{-4}$
disconnected	$-1.5(2)(1) \times 10^{-2}$	$4.6(1.0)(0.4) \times 10^{-2}$
$I = 0$	0.0166(2)(2)	-0.017(1)(1)
$I = 1$	0.0828(8)(9)	-0.148(5)(2)
<b>total</b>	<b>0.0995(9)(10)</b>	<b>-0.166(6)(3)</b>
$I = 1$ FV corr.	0.0006(23)	-0.015(10)
$I = 1 +$ FV corr.	0.0834(8)(9)(23)	-0.164(5)(2)(10)
<b>total + FV corr.</b>	<b>0.1001(9)(10)(23)</b>	<b>-0.182(6)(3)(10)</b>

**Table:** Preliminary results on the first two moments of the HVP function.

c.f. Phenomenology(Benayoun et.al.1605.04474):

$$\Pi_1 = 0.990(7) \text{ GeV}^{-2}, \quad \Pi_2 = -0.206(2) \text{ GeV}^{-4}.$$

# Summary and Perspective

	$\Pi_1[\text{GeV}^{-2}]$	$\Pi_2[\text{GeV}^{-4}]$
total + FV corr.	0.1001(9)(10)(23)	-0.182(6)(3)(10)

**Table:** Preliminary results on the first two moments of the HVP function.

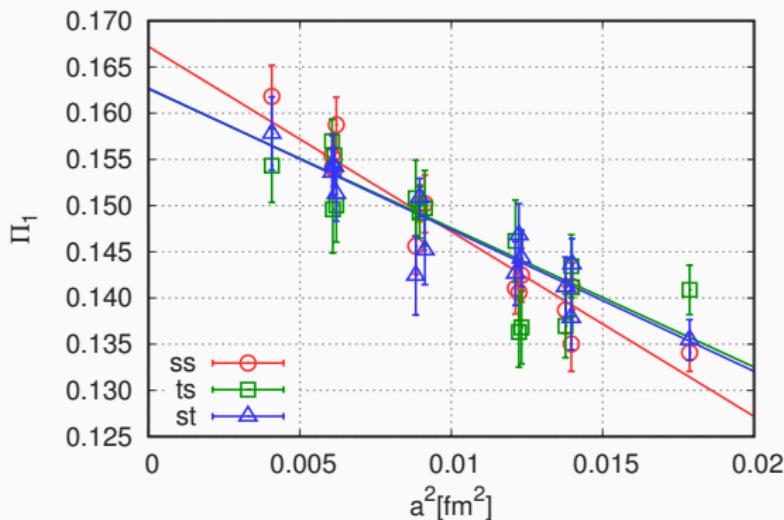
$$\begin{aligned}
 a_\mu^{\text{hvp-lo}} &= \left(\frac{\alpha}{\pi}\right)^2 \int_0^\infty dq^2 w(q^2, m_\mu^2) \frac{q^2 \Pi_1}{1 - q^2 \Pi_2 / \Pi_1} \\
 &\sim [691 \pm \mathcal{O}(10)|_{\text{stat.}} \pm \mathcal{O}(10)|_{\text{sys.}} \pm \mathcal{O}(10)|_{\text{FV}}] \times 10^{-10}. \quad (7)
 \end{aligned}$$

## Future Perspective

- To investigate FV from the lattice data.
- To compute  $a_\mu^{\text{hvp,lo}}|_{\text{conn+disc}}$  with all statistical/systematic uncertainties.
- To take account of the isospin breaking effects and electromagnetism.

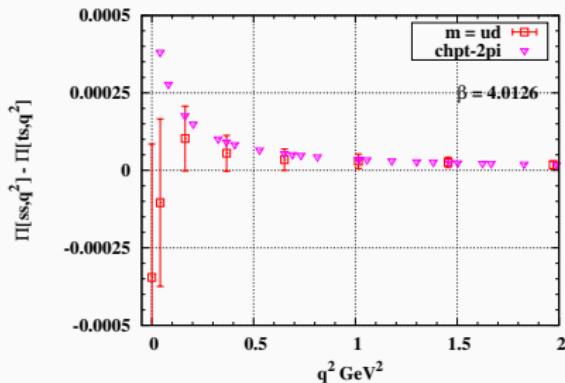
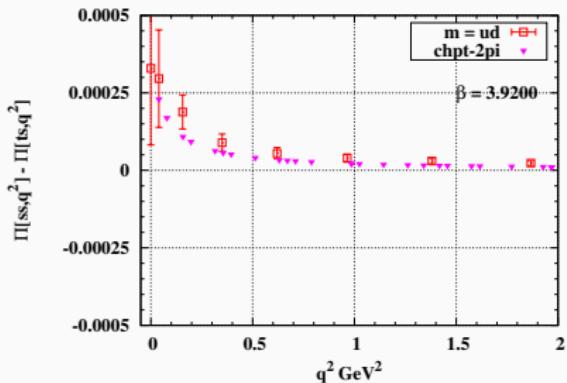
# Backup

# FV via Box Asymmetry



**Figure:** Continuum extrap. of  $\Pi_{n=1}^I$  for each of ss, ts, and st channel.

# FV via Box Asymmetry, Lattice vs XPT



c.f. Aubin et.al., PRD (2016).