

Update of $N_f=3$ finite temperature QCD phase structure with Wilson-Clover fermions

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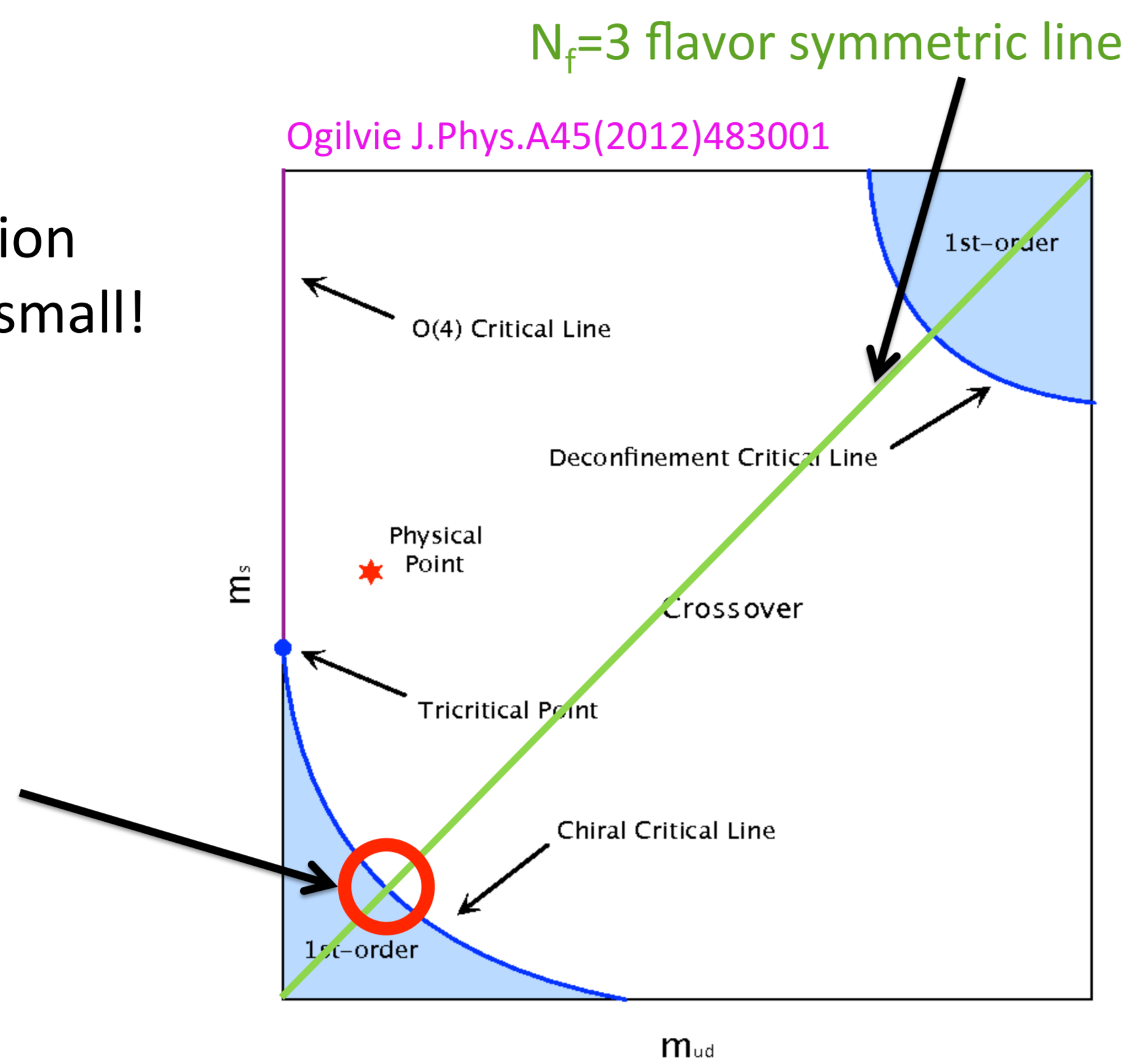
Phase structure of 3-flavor QCD at finite temperature

Columbia plot

1st order region in chiral region has been reported to be so small! in KS fermion study

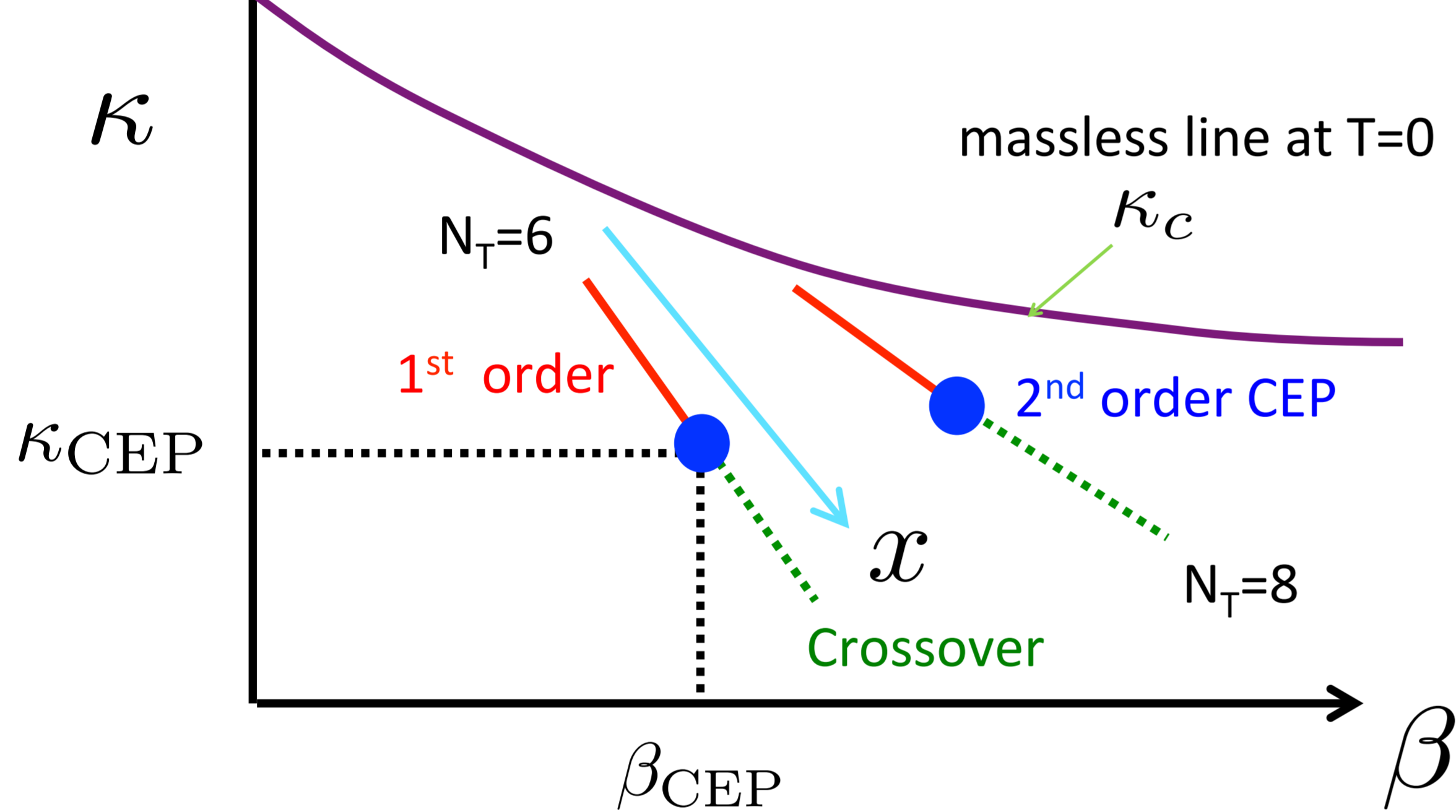
Endrodi PoS LAT2007, 182
de Forcrand & Philipsen 2007
Ding PoS LAT2011, 191

We study the location of critical point with $N_f=3$ by Wilson-type fermions



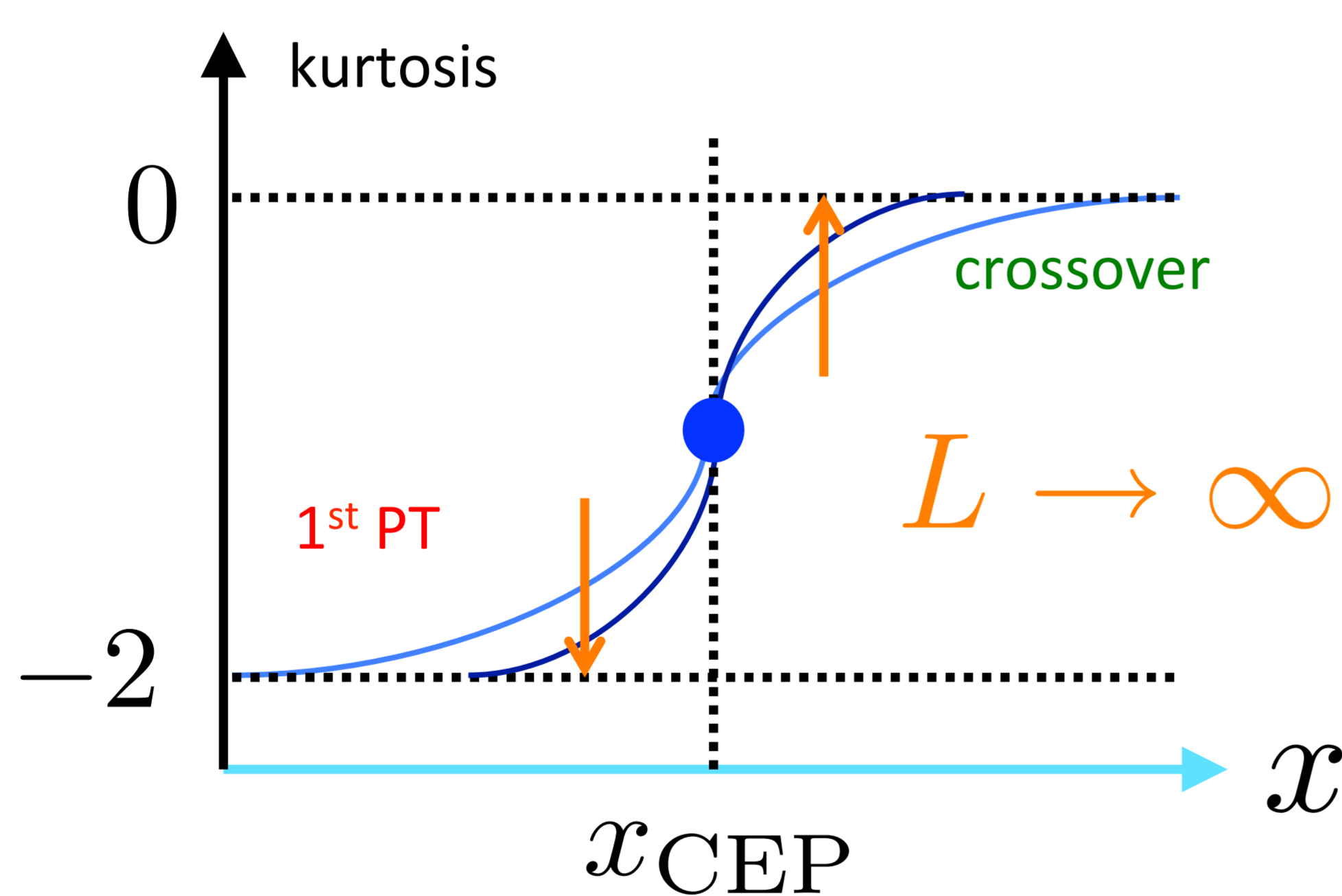
Strategy

bare parameter phase structure for Wilson-type fermions



Kurtosis intersection

kurtosis
 $K = \frac{\kappa_4}{\kappa_2^2}$
cumulants of order parameter



Set up & Simulation parameters

- NP O(a) improved Wilson fermions & Iwasaki gauge
- BQCD code
- RHMC, accept. rate is around 0.8
- measured naïve chiral condensate (order parameter)
- 10-20 noises (Z_2 noise) are used for estimating $\text{tr}[D^{-1,-2,-3,-4}]$ needed for calculating cumulants

N_T	N_L	β	κ
6	10-24	1.715, 1.73, 1.75	0.139620 - 0.141020
8	16-24	1.74995, 1.76019	0.139950 - 0.140240
10	16-28	1.77, 1.78, 1.79	0.140000 - 0.139250

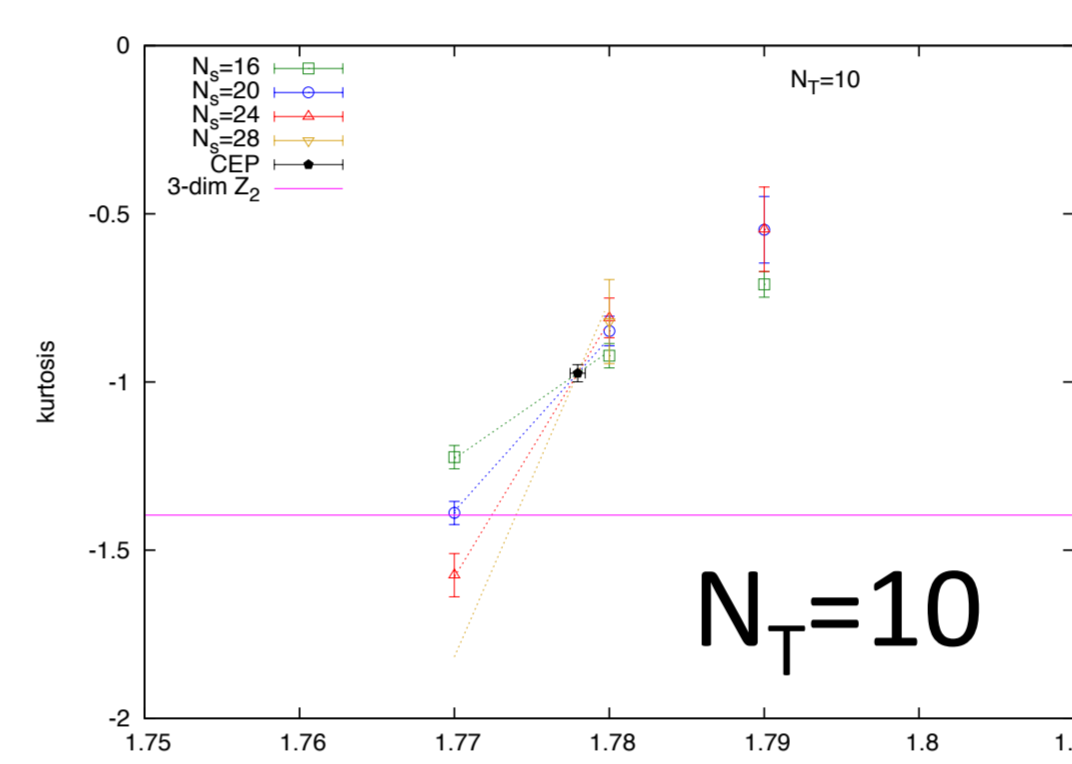
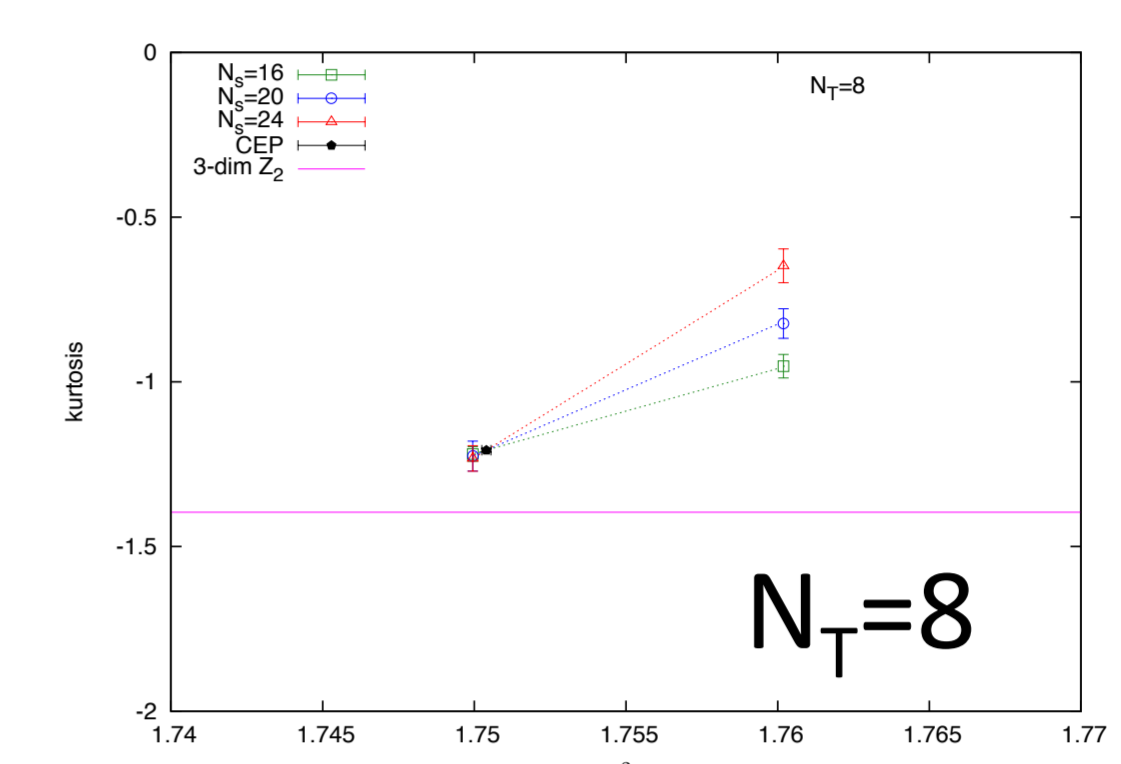
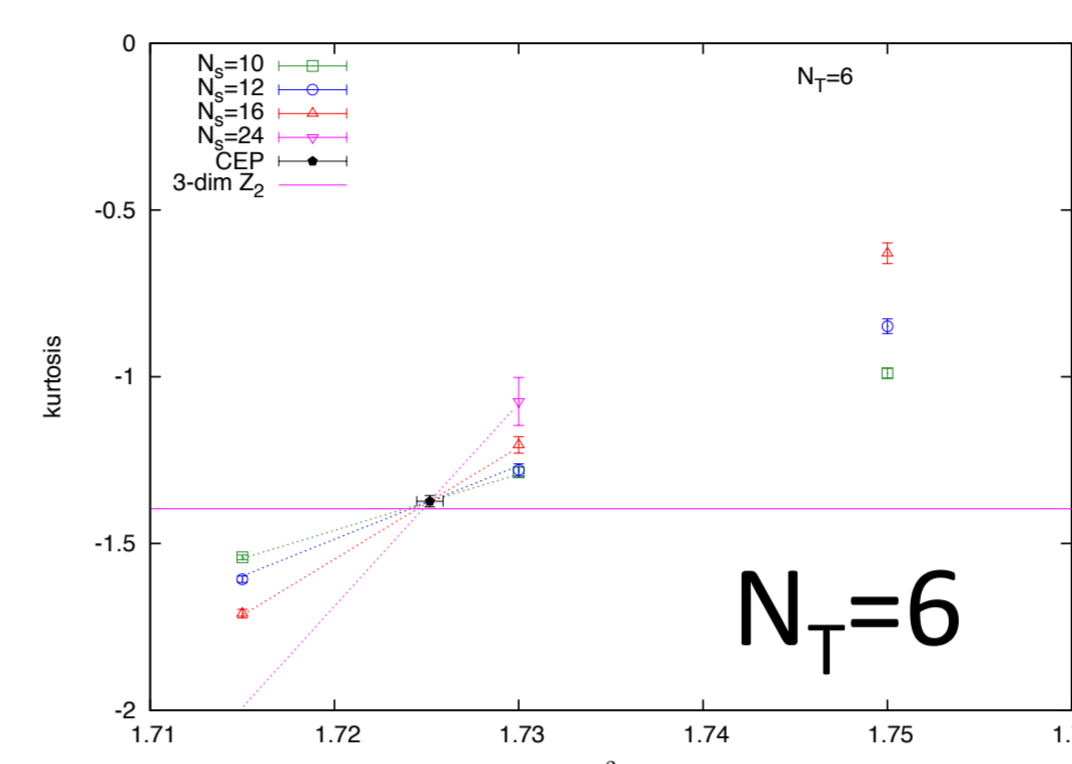
new data

Results

$$K = K_E + AN_L^{1/\nu}(\beta - \beta_{\text{CEP}})$$

Kurtosis intersection

Consistent with Z_2 universality class



Inconsistent with Z_2

Chiral condensate tends to largely overlap with energy-like obs. for $N_T=8$ & 10?

new data

Scaling of kurtosis for mixed observable

Any observable (say chiral condensate) can be considered as a mixture of magnetization (purely order parameter) and energy

$$\mathcal{O} = c_M \mathcal{M} + c_E \mathcal{E}$$

Karsch et al., 2001

For such an observable, the associated kurtosis shows the scaling behavior around the critical point

Jin et al., 2015

$$K_{\mathcal{O}} = \left(K_E + AN_L^{1/\nu}(\beta - \beta_{\text{CEP}}) \right) \left(1 + \frac{BN_L^{y_t - y_h}}{\kappa} \right)$$

new term

$$y_t - y_h \approx -1 \quad \text{for } Z_2, O(2), O(4) \text{ in 3D}$$

$$B \propto c_E$$

fit form	N_t	β_E	K_E	ν	A	B	$y_t - y_h$	χ^2/dof
Old-0	6	1.72518(71)	-1.373(17)	0.683(54)	0.58(17)	x	x	0.68
Old-1	6	1.72429(26)	-1.396	0.643(44)	0.45(12)	x	x	0.85
Old-2	6	1.72480(56)	-1.384(12)	0.63	0.420(11)	x	x	0.71
Old-3	6	1.72431(24)	-1.396	0.63	0.418(11)	x	x	0.70
New-0	6	1.7211(22)	-1.63(13)	0.777(70)	1.00(33)	-1.19(50)	-1	0.32
New-1	6	1.72489(44)	-1.396	0.696(55)	0.63(19)	-0.136(85)	-1	0.59
New-2	6	1.7245(20)	-1.41(12)	0.63	0.423(23)	-0.12(67)	-1	0.93
New-3	6	1.72461(39)	-1.396	0.63	0.422(12)	-0.065(65)	-1	0.70
New-4	6	1.72462(40)	-1.396	0.63	0.422(12)	-0.052(52)	-0.894	0.70
Old-0	8	1.75040(15)	-1.2075(55)	0.513(13)	0.116(18)	x	x	0.01
Old-1	8	1.7453(10)	-1.396	0.96(29)	1.7(1.7)	x	x	1.78
Old-2	8	1.74978(67)	-1.231(24)	0.63	0.345(12)	x	x	0.21
Old-3	8	1.7454(12)	-1.396	0.63	0.338(43)	x	x	2.60
New-0	8	1.7495(11)	-1.31(12)	0.529(25)	0.145(47)	-1.0(1.1)	-1	0.01
New-1	8	1.74876(11)	-1.396	0.543(13)	0.175(24)	-1.681(54)	-1	0.01
New-2	8	1.7460(17)	-1.69(21)	0.63	0.419(33)	-3.4(1.2)	-1	0.09
New-3	8	1.74836(36)	-1.396	0.63	0.374(10)	-1.47(15)	-1	0.12
New-4	8	1.74846(38)	-1.396	0.63	0.375(11)	-1.12(12)	-0.894	0.13
Old-0	10	1.77796(48)	-0.974(25)	0.466(45)	0.084(52)	x	x	0.22
Old-1	10	1.7690(14)	-1.396	2.8(5.4)	16(34)	x	x	7.08
Old-2	10	1.7785(10)	-0.940(51)	0.63	0.448(23)	x	x	0.61
Old-3	10	1.7694(16)	-1.396	0.63	0.421(95)	x	x	10.03
New-0	10	1.7796(18)	-0.70(30)	0.419(63)	0.032(43)	5.2(8.0)	-1	0.25
New-1	10	1.77529(52)	-1.396	0.554(92)	0.29(26)	-3.91(33)	-1	0.50
New-2	10	1.7757(29)	-1.33(41)	0.63	0.539(88)	-3.5(2.7)	-1	0.58
New-3	10	1.77524(52)	-1.396	0.63	0.552(28)	-3.90(34)	-1	0.44
New-4	10	1.77545(53)	-1.396	0.63	0.559(29)	-2.97(25)	-0.894	0.43

One additional parameter in the fitting form reduces the tension in large $N_T=8, 10$

Consistent with Z_2

Summary

- We derive a new fitting form of kurtosis intersection method including the effect of the energy-like observable
- New fitting result is consistent with Z_2 for larger $N_T=8, 10$
- We shall try the mixed observable analysis
- Continuum limit ($N_f=6, 8, 10$)
- $N_f=2+1$ and study the shape of critical line