

Computing the muon anomalous magnetic moment using the hybrid method with physical quark masses

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RBC/UKQCD

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RBC/UKQCD Collaboration

BNL and RBRC

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Motivation

Magnetic moment:

$$\boldsymbol{\mu} = g \frac{e}{2m} \mathbf{S}; \quad U = -\boldsymbol{\mu} \cdot \mathbf{B}; \quad a_\mu = \frac{g_\mu - 2}{2}$$

Contributions to a_μ [PDG, 2015]

Contribution	$a_\mu \times 10^{11}$	Uncertainty
QED (5-loop)	116584718.95	0.08
Electroweak (2-loop)	153.6	1.0
LO hadronic (HVP)	6923	42
NLO hadronic	7	26
Theory Total	116591803	49.4
Experiment Total	116592091	63.3

Hadronic Light-by-Light (HLbL): [Blum et al., 2015; Jin, Tuesday, 14:20]

- 3.6σ discrepancy between theory and experiment.
- Fermilab $g - 2$ and JPARC set to reduce experimental error.

Motivation

Magnetic moment:

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Contributions to a_μ [PDG, 2015]

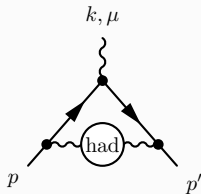
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Hadronic Light-by-Light (HLbL): [Blum et al., 2015; Jin, Tuesday, 14:20]

- 3.6σ discrepancy between theory and experiment.
- Fermilab $g - 2$ and JPARC set to reduce experimental error.
- Greatest uncertainty comes from (non *ab initio*) HVP.

Hadronic Vacuum Polarization

HVP in Euclidean Space [T. Blum, 2002]



- $a_{\mu}^{(2)\text{had}} = 4\alpha^2 \int_0^{\infty} d\hat{q}^2 f(\hat{q}^2) \hat{\Pi}(\hat{q}^2)$
- $\hat{\Pi}(\hat{q}^2) = \Pi(\hat{q}^2) - \Pi(0)$
- $\Pi_{\mu\nu}(\hat{q}) = (\delta_{\mu\nu} \hat{q}^2 - \hat{q}_{\mu} \hat{q}_{\nu}) \Pi(\hat{q}^2)$

Computation:

$$\Pi_{\mu\nu}(\hat{q}) = Z_V \sum_{f,x} Q_f^2 e^{iq \cdot x} \langle \mathcal{V}_{\mu}^f(x) \mathcal{V}_{\nu}^f(0) \rangle; \quad \hat{q} = \frac{2}{a} \sin\left(\frac{aq}{2}\right)$$

Challenges:

- Integrand highly peaked near $\hat{q}^2 = m_{\mu}^2/4$.
- Lattice imposes lower bound on non-zero momenta ($q_{\mu} = \frac{2\pi n_{\mu}}{N_{\mu}}$).
- HVP cannot be directly computed at $\hat{q}^2 = 0$.

HVP Components

Wick contractions give various HVP components: $\Pi_{\mu\nu}^{\ell}(\hat{q}) + \Pi_{\mu\nu}^s(\hat{q}) + \dots$
Separate continuum limits exist for each contraction

[Jüttner and Della Morte, 2010].

Complete

- Strange contribution [Blum et al., 2016].
- Disconnected contribution [Blum et al., 2015].

Ongoing

- Light contribution [C. Lehner, Tuesday, 14:00].
- Charm contribution.
- Isospin breaking effects [Guelpers, Tuesday, 15:20; Harrison, Tuesday, 15:00].

HVP Computation

Strange and charm correlator computations

- Conserved current at sink.
- \mathbb{Z}_2 wall source - Ward identity $q^\mu \Pi_{\mu\nu} = 0$ in large hit limit.

Zero-mode Subtraction

Reduce statistical noise at low- \hat{q}^2 by subtracting $q_t = 0$ component:

$$\Pi_{\mu\nu}(\hat{q}) = \sum_t e^{iq_t \cdot t} C_{\mu\nu}(t) - \sum_t C_{\mu\nu}(t)$$

[Bernecker and Meyer, 2011; C. Lehner and T. Izubuchi, 2014]

Restriction to diagonal of HVP tensor (remove longitudinal part and reduce cut-off effects):

$$\Pi(\hat{q}^2) = \frac{1}{3} \sum_i \frac{\Pi_{ii}(\hat{q})}{\hat{q}^2}; \hat{q}_i = 0$$

Light correlator computation [C. Lehner, Tuesday, 14:00]

- Local current at sink.
- Compute $\Pi_{\mu\nu}(\hat{q})$ using Bernecker-Meyer kernel.

Simulations

Ensembles

RBC/UKQCD 2+1f domain wall fermion ensembles with physical pion masses [RBC/UKQCD, 2014]:

Parameter	48l	64l
$L^3 \times T \times L_s$	$48^3 \times 96 \times 24$	$64^3 \times 128 \times 12$
m_π	139 MeV	139 MeV
m_K	499 MeV	507 MeV
a^{-1}	1.73 GeV	2.36 GeV
$m_\pi L$	3.86	3.78

Target $m_K = 495.7$ MeV

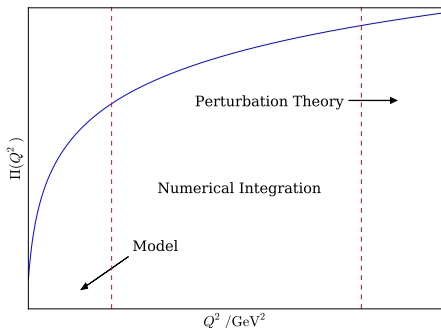
Strange Measurements

Unitary and physical/partially quenched strange masses to account for m_s mistuning.

The Hybrid Method

Motivation

- Systematic error of the parametrisation at low- \hat{q}^2 grows with cut.
- Perturbation theory only valid at high- \hat{q}^2 .
- How are these reconciled? [\[Golterman, Maltman and Peris, 2014\]](#)



Variations

- Parametrisations
- \hat{q}^2 cuts
- Techniques to constrain parametrisation (fits, moments)
- Numerical integration method

Low- q^2 Parametrisations

Padé Approximants

Motivated by the spectral decomposition of the HVP

[Aubin, Blum, Golterman and Peris, 2012]:

$$R_{mn}(\hat{q}^2) = \Pi_0 + \hat{q}^2 \left(\sum_{i=0}^{m-1} \frac{a_i}{b_i + \hat{q}^2} + \delta_{mn} c \right); \quad n = m, m + 1.$$

Conformal Polynomials

Map domain of analyticity onto region within unit disc. Better convergence properties [Golterman, Maltman and Peris, 2014]:

$$P_n^E(\hat{q}^2) = \Pi_0 + \sum_{k=1}^n p_k w^k,$$
$$w = \frac{1 - \sqrt{1+z}}{1 + \sqrt{1+z}}, \quad z = \frac{\hat{q}^2}{E^2}.$$

Time Moments [HPQCD, 2014]

- 1 Tensor decomposition:

$$\sum_t e^{iq_t t} C_{\mu\mu}(t) = \hat{q}_t^2 \Pi(\hat{q}_t^2)$$

- 2 Differentiate w.r.t. q_t :

$$\Delta_{q_t}^{2n} \left(\sum_t e^{iq_t t} C_{\mu\mu}(t) \right) \Big|_{q_t=0} = \Delta_{q_t}^{2n} (\hat{q}_t^2 \Pi(\hat{q}_t^2)) \Big|_{q_t=0}$$

- 3 Plug in a parametrisation for $\Pi(\hat{q}^2)$ and solve for parameters.

Notes

- Account for discrete momenta using finite differences.
- Expansion around $\hat{q}^2 = 0$, data here carry more weight.
- Resulting parameters independent of cut.

Continuous Momenta

- Fourier transform to arbitrary momenta

[Bernecker and Meyer, 2011; Feng et al., 2013]:

$$\Pi_{\mu\nu}(\hat{q}) = \sum_t e^{iq_t \cdot t} C_{\mu\nu}(t) - \sum_t C_{\mu\nu}(t)$$
$$q_t = \frac{2\pi n_t}{T}, \quad n_t \in [-T/2, T/2)$$

- Compute HVP directly at arbitrary \hat{q} .
- No hybrid method, parametrisation independent.
- Can show interpolation effects are $\mathcal{O}(\exp(-m_\pi T))$

[Portelli, Lattice 2015; Del Debbio and Portelli, to appear].

Overview of Strange Analyses

Hybrid Method

- Matching method: diagonal fit, time moments.
- Parametrisations:
 - Padés: $R_{1,1}$, $R_{1,2}$.
 - Conformal polynomials: P_3^E , P_4^E ; $E = 500 \text{ MeV}, 600 \text{ MeV}$.
- Low cuts: $0.5, 0.7, 0.9 \text{ GeV}^2$, $(\frac{2\pi}{aT})^2 \approx 0.013 \text{ GeV}^2$.
- High cut: 5.0 GeV^2 .
- Mid- \hat{q}^2 integration technique: trapezium rule, Simpson's rule.

Sine Cardinal Interpolation

- $\Delta n_t = 0.005$
- High cut: 5.0 GeV^2 .

Extrapolations

- Strange quark mistuning: $\sim 1\%$ on 48l, $\sim 5\%$ on 64l.
- Partially quenched measurements.
- Continuum limit.
- Strange mass extrapolation.

Strange Ansatz

$$a_{\mu}^{(s)\text{had},s}(a^2, am_s) = a_{\mu,0}^{(s)\text{had},s} + \alpha a^2 + \beta \frac{am_s - am_s^{\text{phys}}}{am_s^{\text{phys}} + am_{\text{res}}}$$

Systematic Effects

Accounted For

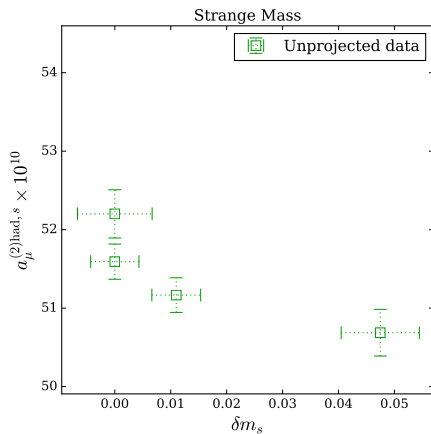
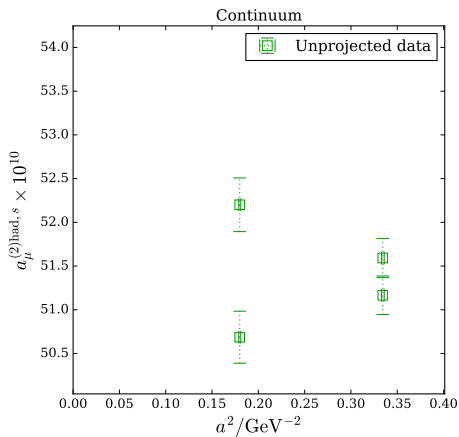
- \hat{q}^2 cuts.
- Low- \hat{q}^2 parametrisation.
- Matching method.
- Mid- \hat{q}^2 integration method.
- Finite volume effects negligible via G -parity.
- Non-unitarity effects negligible.
- Disconnected diagrams [Blum et al., 2015].

Long Term

- Charm quark in the sea.
- Isospin breaking effects, including EM effects
[Guelpers, Tuesday, 15:20; Harrison, Tuesday, 15:00].

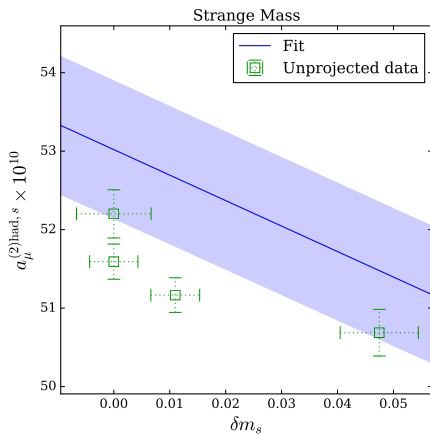
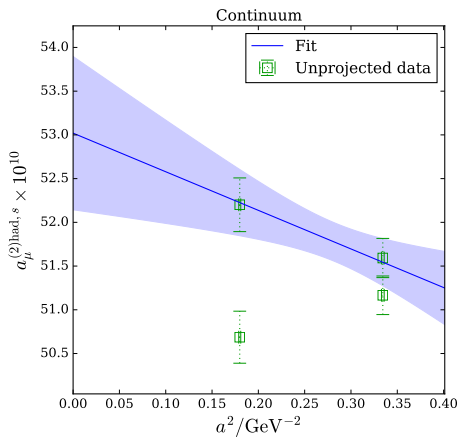
Strange Result: Extrapolations

Moments, $P_3^{0.6\text{GeV}}$, low cut = 0.7 GeV^2 , high cut = 5.0 GeV^2



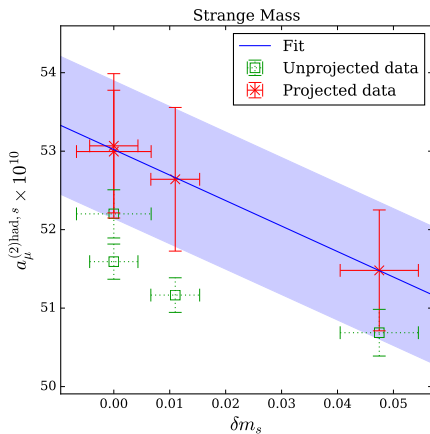
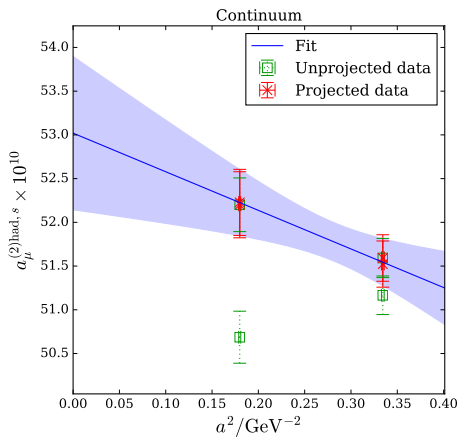
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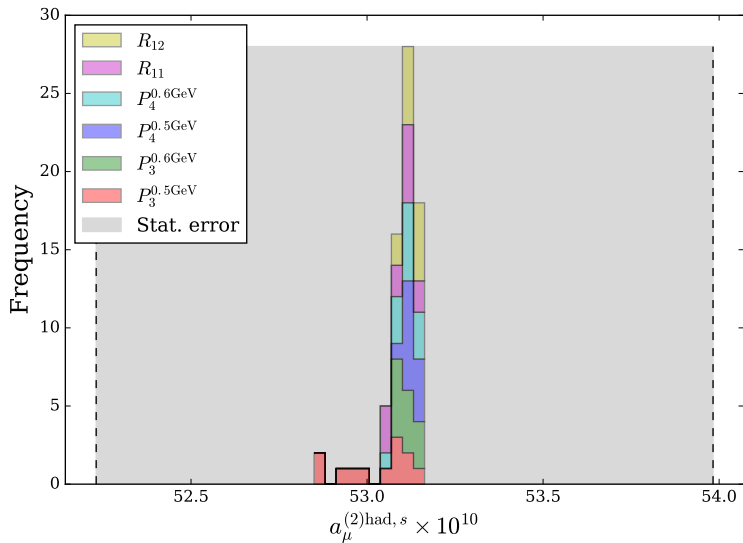


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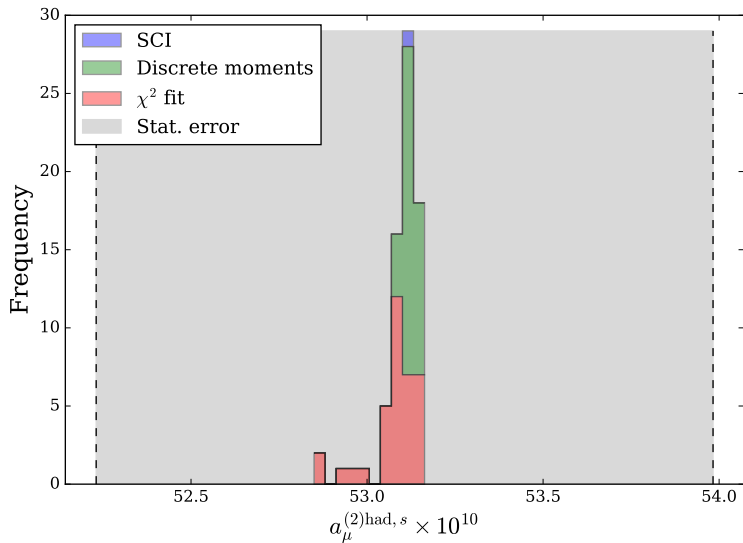
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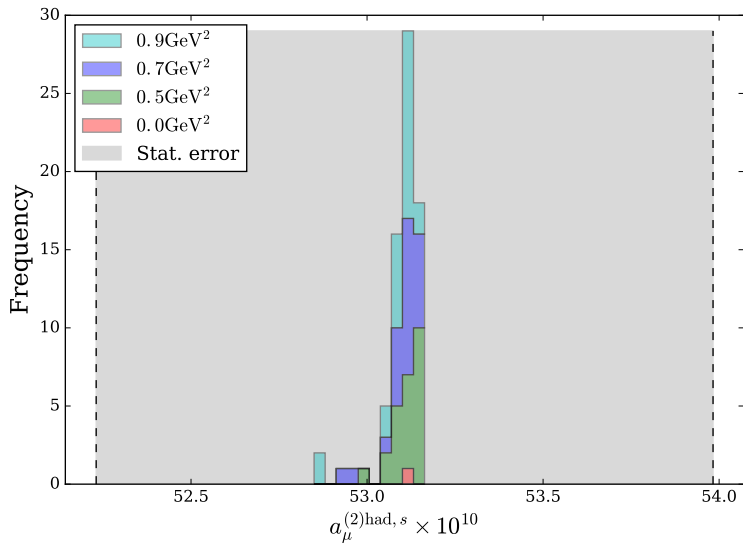
Strange Result: Parametrisations



Strange Result: Matching Method



Strange Result: Low Cuts



Overview of Light Analyses

Hybrid Method

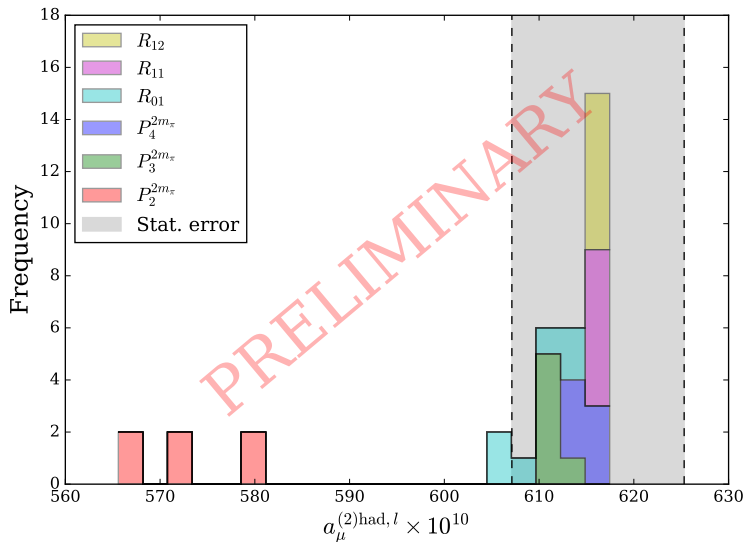
- Matching method: diagonal fit.
- Parametrisations:
 - Padés: $R_{0,1}$, $R_{1,1}$, $R_{1,2}$.
 - Conformal polynomials: $P_2^{2m_\pi}$, $P_3^{2m_\pi}$, $P_4^{2m_\pi}$.
- Low cuts: 0.5, 0.7, 0.9 GeV², $(\frac{2\pi}{aT})^2 \approx 0.013$ GeV².
- High cut: 7.0 GeV².
- Mid- \hat{q}^2 integration technique: trapezium rule, Simpson's rule.

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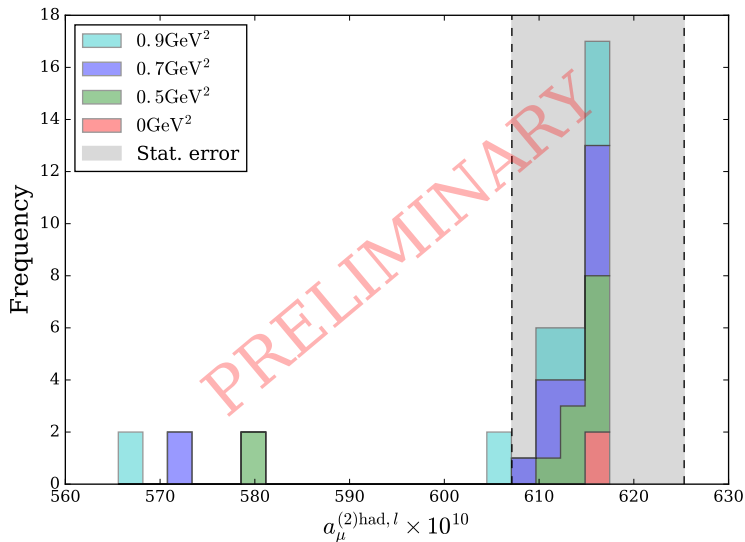
- $\Delta n_t = 0.005$
- High cut: 7.0 GeV².

See also: [\[C. Lehner, Tuesday, 14:00\]](#)

Light Results: Parametrisations



Light Results: Low Cuts



Charm Contribution

HVP measurements computed on seven dynamical ensembles.

- Same computation of $\Pi_{\mu\nu}$ as strange contribution.
- Select single analysis technique motivated by strange analysis.

$L^3 \times T$	a^{-1}/GeV	m_π/MeV	Number of am_h used
$48^3 \times 96$	1.73	139	3
$24^3 \times 64$	1.78	340	3
$24^3 \times 64$	1.78	430	3
$64^3 \times 128$	2.36	139	4
$32^3 \times 64$	2.38	300	4
$32^3 \times 64$	2.38	360	4
$48^3 \times 96$	2.77	230	5

Global fit of charm data ongoing [Tsang, Friday, 14:20].

Summary

Strange Contribution [Blum et al., 2016]

- $a_\mu^{(2)\text{had},s}$ computed using Möbius domain wall fermions with 2+1f.
- Extensive systematic study of analysis techniques.
- Final value of $a_\mu^{(2)\text{had},s}$ largely insensitive to analysis method.
- Results consistent with other studies (HPQCD).

$$a_\mu^{(2)\text{had},s} = 53.1(9) \binom{+1}{-3} \times 10^{-10}$$

Light Contribution

- Investigation into analysis techniques ongoing.

Outlook

- Cut-off dependence of light contribution.
- Charm contribution.

Ensembles

RBC/UKQCD ensembles using 2+1f domain wall fermions with a physical pion mass.

Parameter	48l	64l
$L^3 \times T \times L_s$	$48^3 \times 96 \times 24$	$64^3 \times 128 \times 12$
am_l	0.00078	0.000678
am_s	0.0362	0.02661
am_s^{phys}	0.03580(16)	0.02539(17)
a^{-1} / GeV	1.730(4)	2.359(7)
L / fm	5.476(12)	5.354(16)
$m_\pi L$	3.863(6)	3.778(8)
am_{res}	0.0006012	0.0003116

Strange quark mistuning: partially quenched runs

Analysis Strategy: χ^2 Fit

χ^2 Fit

- Standard χ^2 minimization - covariance approximated by diagonal.

$$\chi^2 = \sum_{\hat{q}^2} \left(\frac{\Pi(\hat{q}^2) - f(\hat{q}^2)}{\delta\Pi(\hat{q}^2)} \right)^2$$

- Fits can be unstable - Z2 wall reduces d.o.f.
- Fit biased towards large \hat{q}^2 .
- Parameters dependent on cut.