

QCD with isospin density: pion condensation

Gergely Endrődi, Bastian Brandt

Goethe University of Frankfurt



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Outline

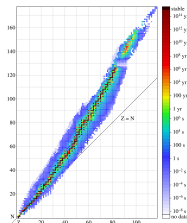
- introduction: QCD with isospin
- relevant phenomena
 - ▶ deconfinement/chiral symmetry breaking at low μ_I
↙ next talk
 - ▶ pion condensation at high μ_I
↙ this talk
- “ λ -extrapolation”
 - ▶ naive method
 - ▶ new method
- outlook and summary

Introduction

- ▶ isospin density $n_I = n_u - n_d$
- ▶ $n_I < 0 \rightarrow$ excess of neutrons over protons
 \rightarrow excess of π^- over π^+

- ▶ applications

- ▶ neutron stars
- ▶ heavy-ion collisions



- ▶ chemical potentials (3-flavor)

$$\mu_B = 3(\mu_u + \mu_d)/2$$

$$\mu_I = (\mu_u - \mu_d)/2$$

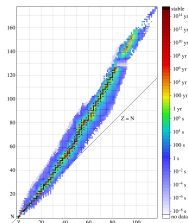
$$\mu_S = 0$$

Introduction

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$$\mu_I = (\mu_u - \mu_d)/2$$

$$\mu_S = 0$$

- ▶ here: zero baryon number but nonzero isospin

$$\mu_u = \mu_I$$

$$\mu_d = -\mu_I$$

Introduction



- ▶ QCD at low energies \approx pions
- ▶ on the level of charged pions: $\mu_\pi = 2\mu_l$
at zero temperature

$$\begin{array}{ll} \mu_\pi < m_\pi & \text{vacuum state} \\ \mu_\pi = m_\pi & \text{Bose-Einstein condensation} \\ \mu_\pi > m_\pi & \text{undefined} \end{array}$$

- ▶ on the level of quarks: lattice simulations
 - ▶ no sign problem
 - ▶ conceptual analogies to baryon density
(Silver Blaze, hadron creation, saturation)
 - ▶ technical similarities
(proliferation of low eigenvalues)

Setup

Symmetry breaking

- ▶ QCD with light quark matrix

$$M = \not{D} + m_{ud}\mathbb{1}$$

- ▶ chiral symmetry (flavor-nontrivial)

$$SU(2)_V$$

Symmetry breaking

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$$SU(2)_V \rightarrow U(1)_{\tau_3}$$

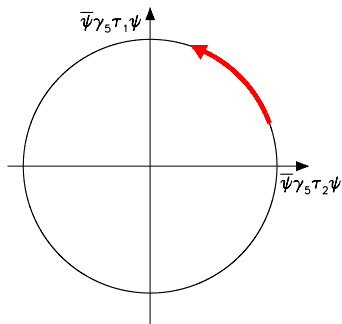
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- ▶ spontaneously broken by a pion condensate

$$\langle \bar{\psi}\gamma_5\tau_{1,2}\psi \rangle$$

- ▶ a Goldstone mode appears

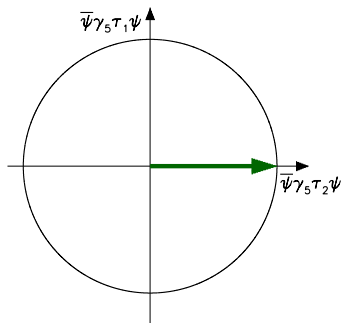
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- ▶ add small **explicit breaking**

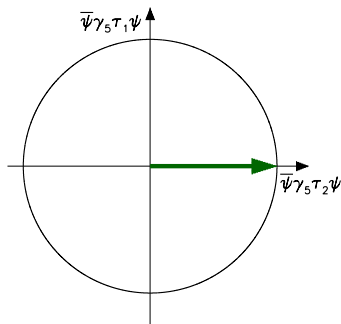
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- ▶ extrapolate results $\lambda \rightarrow 0$

Simulation details

- ▶ staggered light quark matrix with $\eta_5 = (-1)^{n_x+n_y+n_z+n_t}$

$$M = \begin{pmatrix} \not{D}_\mu + m_{ud} & \lambda\eta_5 \\ -\lambda\eta_5 & \not{D}_{-\mu} + m_{ud} \end{pmatrix}$$

- ▶ $\gamma_5\tau_1$ -hermiticity

$$\eta_5\tau_1\not{D}_\mu\tau_1\eta_5 = \not{D}_\mu^\dagger$$

→ determinant is real and positive

- ▶ first done by [Kogut, Sinclair '02]
- ▶ here: $N_f = 2 + 1$ rooted stout-smear staggered quarks + tree-level Symanzik improved gluons

Pion condensate: definition and renormalization

- ▶ condensate

$$\langle \pi \rangle = \frac{T}{V} \frac{\partial \log \mathcal{Z}}{\partial \lambda}$$

- ▶ additive divergences cancel in

$$\lim_{\lambda \rightarrow 0} \langle \pi \rangle$$

- ▶ multiplicative renormalization

$$Z_\pi = Z_\lambda^{-1} = Z_{m_{ud}}^{-1}$$

- ▶ renormalization + convenient normalization

$$\Sigma_\pi \equiv m_{ud} \cdot \langle \pi \rangle \cdot \frac{1}{m_\pi^2 f_\pi^2}$$

- ▶ so that in leading-order chiral PT [Son, Stephanov '00]

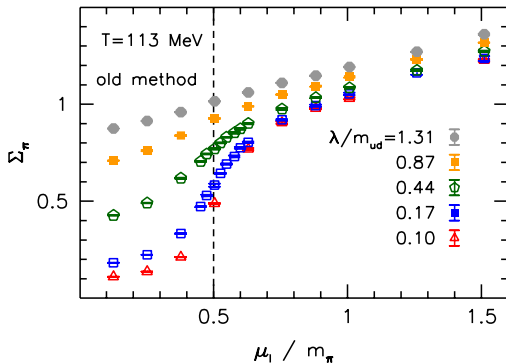
$$\Sigma_{\bar{\psi}\psi}^2(\mu_I) + \Sigma_\pi^2(\mu_I) = 1$$

Pion condensate: old method

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- ▶ traditional method [Kogut, Sinclair '02]
measure full operator at nonzero λ (via noisy estimators)

$$\Sigma_\pi \propto \left\langle \text{Tr} M^{-1} \eta_5 \tau_2 \right\rangle$$

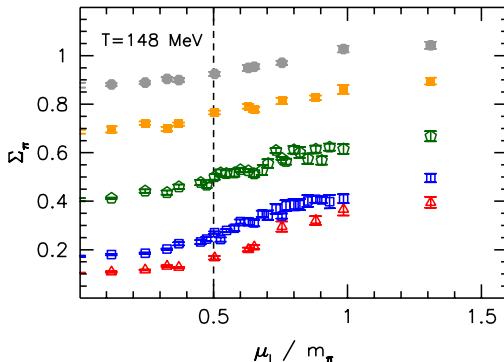


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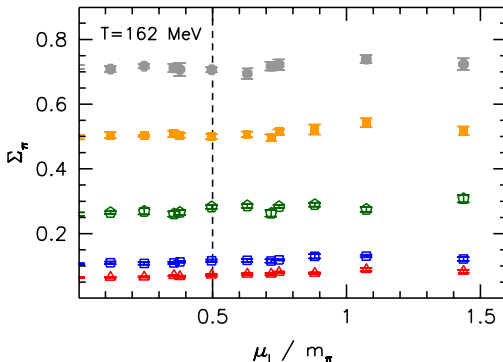


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Pion condensate: new method

Singular value representation

- ▶ pion condensate

$$\pi = i\text{Tr}(M^{-1}\eta_5\tau_2) = \text{Tr} \frac{2\lambda}{(\not{D}_\mu + m)^\dagger(\not{D}_\mu + m) + \lambda^2}$$

- ▶ singular values

$$(\not{D}_\mu + m)^\dagger(\not{D}_\mu + m)\psi_i = \xi_i^2\psi_i$$

- ▶ spectral representation

$$\pi = \frac{T}{V} \sum_i \frac{2\lambda}{\xi_i^2 + \lambda^2} = \int d\xi \rho(\xi) \frac{2\lambda}{\xi^2 + \lambda^2} \xrightarrow{\lambda \rightarrow 0} \pi\rho(0)$$

first derived in [Kanazawa, Wettig, Yamamoto '11]

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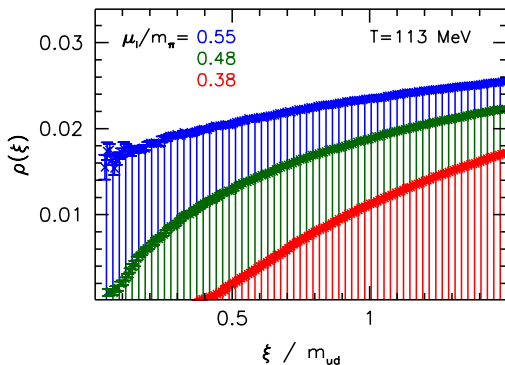
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- ▶ compare to Banks-Casher-relation at $\mu_l = 0$

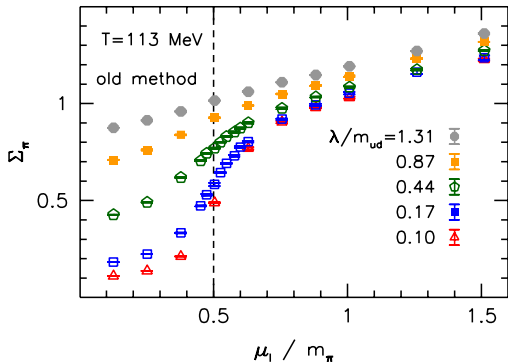
Singular value density

- ▶ spectral densities at $\lambda/m_{ud} = 0.17$



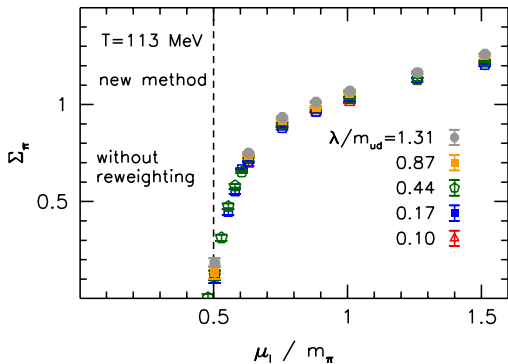
Density at zero

- ▶ scaling with λ is improved drastically



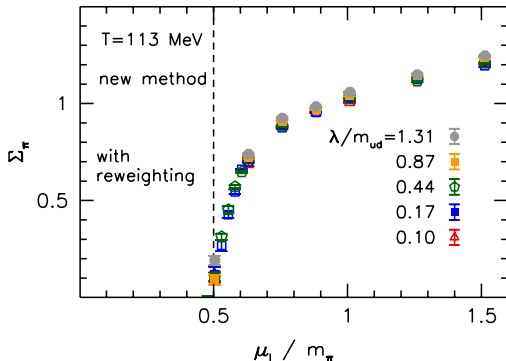
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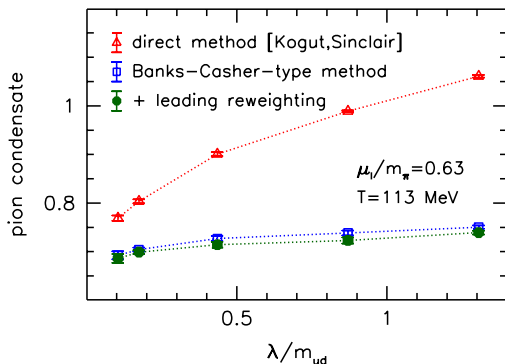
- ▶ leading-order reweighting

$$\langle \pi \rangle_{\text{rew}} = \langle \pi W_\lambda \rangle / \langle W_\lambda \rangle$$

$$W_\lambda = \exp[-\lambda V_4 \pi + \mathcal{O}(\lambda^2)]$$

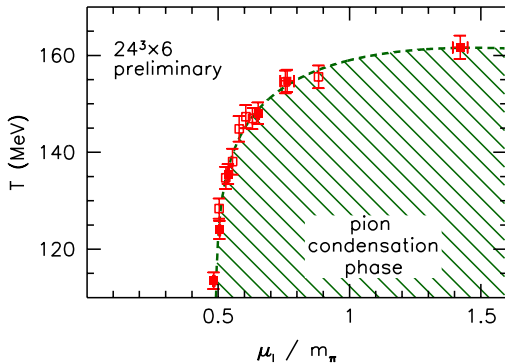
Comparison between old and new methods

- ▶ extrapolation in λ gets almost completely flat



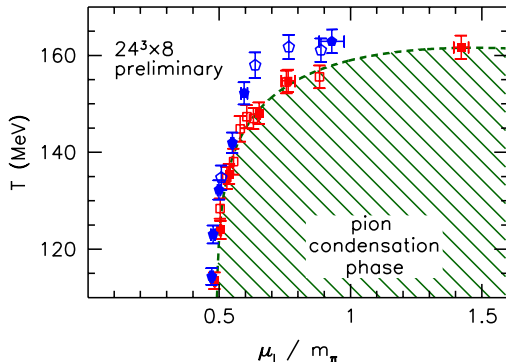
Phase boundary

- ▶ interpolate $\rho(0)$ as function of μ_I to find phase boundary



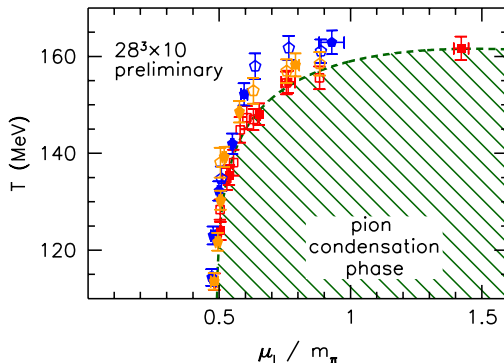
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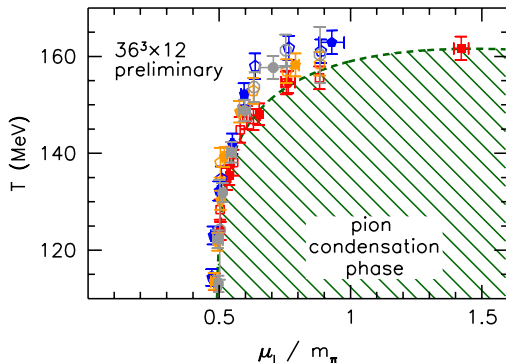
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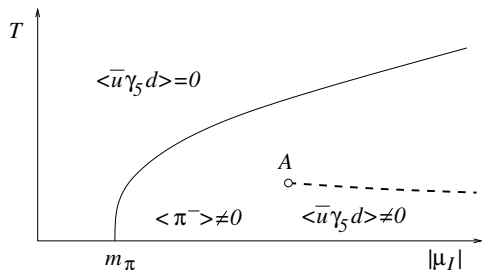
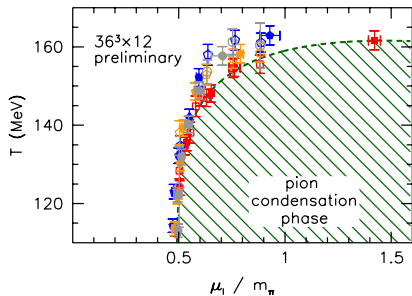
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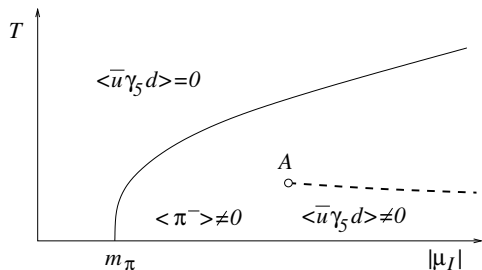
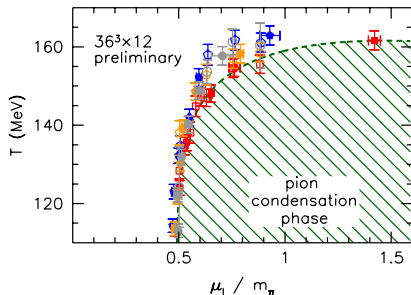
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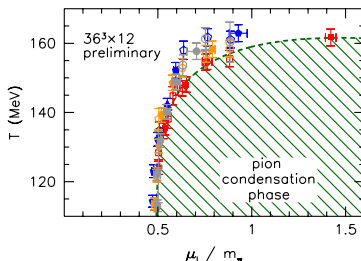
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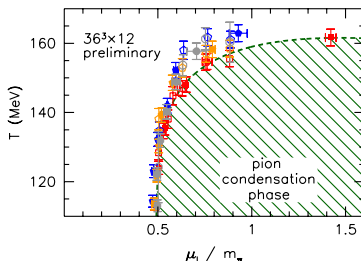
- ▶ compare to expectations from χ PT [Son, Stephanov '00]
- ▶ no pion condensate above $T \approx 160$ MeV

Outlook



- ▶ order of transition?
- ▶ deconfinement/chiral symmetry breaking transition?
- ▶ asymptotic- μ_I limit?
- ▶ BCS phase at large μ_I ?
- ▶ analogies to two-color QCD [Holicki, Thu]
- ▶ test Taylor-expansion in μ_I

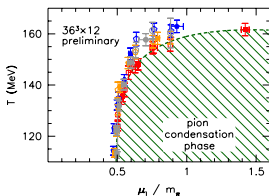
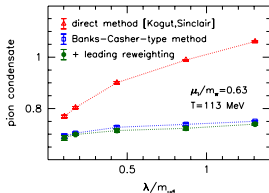
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- ▶ test Taylor-expansion in μ_I
- ▶ stay for next talk [Brandt, Thu]

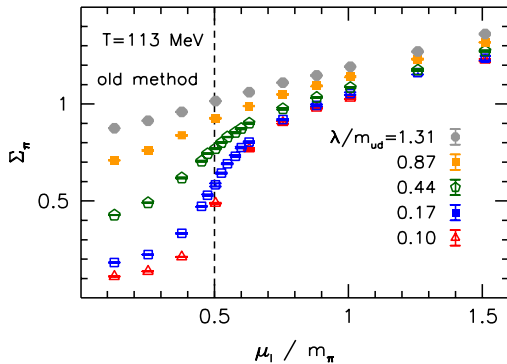
Summary

- ▶ QCD with isospin chemical potentials via lattice simulations at the physical point
- ▶ determine pion condensate via Banks-Casher-type relation
 \rightsquigarrow flat extrapolation in pion source
- ▶ phase boundary surprisingly flat for intermediate μ_I
- ▶ chance to test effective theories and low-energy models



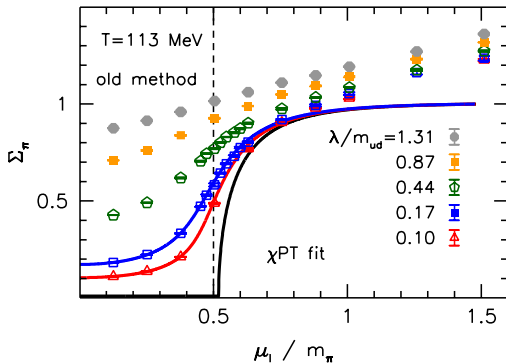
Backup

Order of the transition – fits



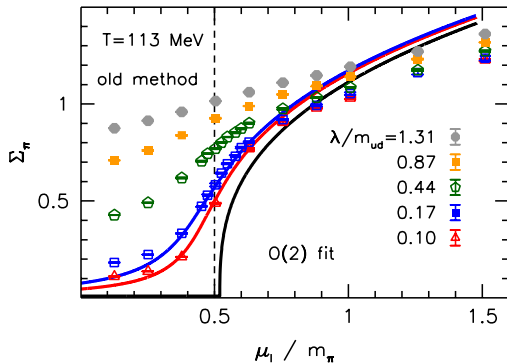
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 - ▶ chiral perturbation theory [Splittorff et al '02, Endrődi '14]
 - ▶ $O(2)$ scaling [Ejiri et al '09]