An a_0 resonance in strongly coupled $\pi\eta, K\overline{K}$ scattering from lattice QCD

PRD93 094506 (2016) (with David Wilson and Robert Edwards)

Jozef Dudek

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the light scalar mesons - empirically

conventional to put them in an 'inverted' mass nonet

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let's study their appearance within QCD ...



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the light scalar mesons - empirically

conventional to put them in an 'inverted' mass nonet





'start' with the a_0 resonance ...



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the *a*₀(980) as it really is - a *resonance*

• sharp experimental enhancement at $K\overline{K}$ threshold decaying to $\pi\eta$



• usually observed in 'less-simple' production processes

• amplitude models typically give

$$\frac{g^2(K\overline{K})}{g^2(\pi\eta)} \sim 1$$

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e.g. $p\overline{p} \rightarrow \pi \pi \eta$

 $\phi \rightarrow \gamma \pi \eta$

Baru et. al. EPJA23 523 (2005) עת → תח "KLOE" $-1\overline{00}$ -5050 100 0 $E-2m_K$ "Bugg" -5050 -100100 0 "Achasov" -100 -50 50 100 0 "E852" -50 50 -1000 100 Jefferson Lab



the $a_0(980)$ as it really is - a resonance



• matrices of correlation functions with a large operator basis

"
$$q\bar{q}$$
"-like $\bar{\psi}\Gamma D...D\psi$
 $\pi\eta$ -like $\sum_{\hat{p}_1,\hat{p}_2} C(\vec{p}_1,\vec{p}_2;\vec{P}) \pi(\vec{p}_1)\eta(\vec{p}_2)$
 $K\bar{K}$ -like $\sum_{\hat{p}_1,\hat{p}_2} C(\vec{p}_1,\vec{p}_2;\vec{P}) K(\vec{p}_1)\overline{K}(\vec{p}_2)$
 $\pi\eta$ '-like $\sum_{\hat{p}_1,\hat{p}_2} C(\vec{p}_1,\vec{p}_2;\vec{P}) \pi(\vec{p}_1)\eta'(\vec{p}_2)$

(with optimized pseudoscalar operators)



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(with optimized pseudoscalar operators)

many Wick contractions required, including annihilations ...



... distillation

totally straightforward, massive reuse of propagators





I=1, G=- spectra

spectra obtained from variational analysis

 $m_{\pi} \sim 391 \,\mathrm{MeV}$

 16^3 , 20^3 , 24^3

- *a*_s ~ 0.12 fm
- $a_t \sim a_s/3.5$

$[000] A_1^+$ $[100] A_1$ $[110] A_1$ $[111] A_1$ $[200] A_1$ $a_t E_{\mathsf{cm}}$ $\omega \pi \pi |_{\text{thr.}}$ 0.29 0.29 0.29 0.29 0.29 Ī $\pi\eta\eta\big|_{\rm thr.}$ 0.27 0.27 0.27 0.27 0.27 $\overline{KK}|_{\text{thr.}}$ 互 0.25 0.25 0.25 0.25 0.25 $\left.\pi\eta'\right|_{\mathrm{thr.}}$ 0.23 0.23 0.23 0.23 0.23 Ξ Þ 互 Ξ 互 0.21 0.21 0.21 0.21 0.21 $\pi\pi\pi\pi|_{\rm thr}$ $\frac{1}{\pi} K\overline{K}\Big|_{\text{thr.}}$ 0.19 0.19 0.19 Ξ 0.19 Π 0.19 互 Ŧ Ę $- = \pi \eta |_{\text{thr.}}$ 互 0.17 - 포 0.17 0.17 0.17 - Ţ 0.17 20 24 16 20 16 16 24 20 24 16 20 24 20 24 16 (13) (14) (20) (18) (11) (21)(12) (16) (19) (12) (11) (12) (13)(26) (17)

(*conservatively*) **47** levels in the relevant energy region

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coupled channel scattering

• finite-volume formalism established

$$\det \begin{bmatrix} \mathbf{t}^{-1}(E) + i\boldsymbol{\rho}(E) - \mathbf{M}(E,L) \end{bmatrix} = 0$$
scattering
matrix
space
tin the two-chapped case
t =
$$\begin{bmatrix} t_{\pi\eta\to\pi\eta} & t_{\pi\eta\to K\overline{K}} \end{bmatrix}$$

HE, JHEP 0507 011 HANSEN, PRD86 016007 BRICENO, PRD88 094507 GUO, PRD88 014051

in the two-channel case
$$\mathbf{t} = \begin{bmatrix} t_{\pi\eta\to\pi\eta} & t_{\pi\eta\to K\overline{K}} \\ t_{K\overline{K}\to\pi\eta} & t_{K\overline{K}\to K\overline{K}} \end{bmatrix}$$







coupled channel scattering



$$\det \begin{bmatrix} \mathbf{t}^{-1}(E) + i\boldsymbol{\rho}(E) - \mathbf{M}(E,L) \end{bmatrix} = 0$$
scattering
matrix
phase
space
functions
 $\begin{bmatrix} t_{\pi\eta\to\pi\eta} & t_{\pi\eta\to K\overline{K}} \end{bmatrix}$

HE, JHEP 0507 011 HANSEN, PRD86 016007 BRICENO, PRD88 094507 GUO, PRD88 014051

in the two-channel case
$$\mathbf{t} = \begin{bmatrix} t_{K\overline{K} \to \pi\eta} & t_{K\overline{K} \to K\overline{K}} \end{bmatrix}$$

• parameterizing the energy-dependence in a unitarity-preserving way

K-matrix approach

$$\mathbf{t}^{-1}(E) = \mathbf{K}^{-1}(E) + \mathbf{I}(E) \qquad \text{Im } (\mathbf{I}(E))_{ij} = -\delta_{ij} \rho_i(E)$$

e.g. "Chew-Mandelstam" phase-space

e.g. poles plus polynomial form

$$(\mathbf{K}(E))_{ij} = \sum_{p} \frac{g_i^{(p)} g_j^{(p)}}{m_p^2 - E^2} + \sum_{n} \gamma_{ij}^{(n)} (E^2)^n$$





$\pi\eta/K\overline{K}$ scattering describing the spectra

 $m_{\pi} \sim 391 \,\mathrm{MeV}$ *K*-matrix parameterization: one pole plus constant matrix (6 free parameters) $[000] A_1^+$ $[110] A_1$ $[100] A_1$ $[200] A_1$ $[111] A_1$ $\pi \eta' |_{\text{thr.}}$ 0.23 0.23 0.23 0.23 0.23 0.21 0.21 0.21 0.21 0.21 $\pi\pi\pi$ KK0.19 0.19 0.19 0.19 0.19 0.17 0.17 0.17 0.17 0.17 16 24 20 24 20 16 20 24 16 20 24 20 16 16 24 $\chi^2 / N_{\rm dof} = \frac{58.0}{47 - 6} = 1.4$



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$\pi \eta / K \overline{K}$ scattering in $J^P = 0^+$

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parameterization variation

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$\pi\eta/KK$ scattering in $J^P = 0^+$







$\pi \eta / K \overline{K}$ scattering in $J^P = 0^+$



how do we determine rigorously if an amplitude is resonant?

look for a **pole singularity** at **complex energy**

$$t_{ij}(s) \sim \frac{g_i g_j}{s_0 - s}$$

Re[$\int s_0$] ~ 'mass' 2·Im[$\int s_0$] ~ 'width'



$\pi\eta/K\overline{K}$ scattering in $J^P = 0^+$

• we find a single dominant (nearby) pole

COMPLEX ENERGY PLANE



 $m_{\pi} \sim 391 \,\mathrm{MeV}$

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COMPLEX MOMENTUM PLANE



Sheet	$\mathrm{Im}k_{\pi\eta}$	$\mathrm{Im}k_{K\bar{K}}$
I	+	+
II		+
III	_	_
IV	+	_



$\pi \eta / K\overline{K}$ scattering in $J^P = 0^+$

• we find a single dominant (nearby) pole

 $m_{\pi} \sim 391 \,\mathrm{MeV}$



$\pi\eta/KK$ scattering in $J^P = 0^+$

18 scale setting using the Ω baryon mass



 $\text{Re}\sqrt{s_0} - 2m_K = 79 \pm 27$

pole coupling ratio $\left|\frac{g_{K\bar{K}}}{g_{\pi\eta}}\right| = 1.3(4)$



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complete study of scalar mesons ?

Raul Briceño (previous session) presented $\pi\pi$ elastic scattering in *I*=0 with the σ

- bound state at m_{π} =391 MeV
- broad resonance at m_{π} =236 MeV ARXIV: 1607.05900

extension to coupled channel $(\pi\pi, K\overline{K}, ...)$ and likely $f_0('980')$ coming up



 π K, ηK scattering already done at m_{π} =391 MeV κ as a virtual bound-state ? ... at m_{π} =236 MeV to come soon ... PRL113 182001 (2014) PRD91 054008 (2015)

 $\pi\eta$, $K\overline{K}$ at m_{π} =236 MeV to come soon ...

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what are the scalar mesons ?

- what tools do we have at our disposal ?
 - quark mass dependence of the pole positions and channel couplings
 - distribution of poles across Riemann sheets

may be relatable to loose meson-meson molecule versus tightly bound object

• coupling to external currents

form-factors from residue at the resonance pole

finite-volume formalism demonstrated with $\gamma \pi \rightarrow \pi \pi$





• and things we haven't thought of yet ...



hadspec

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an a_0 resonance ...

Nilmani Mathur

 $(\dot{\mathbf{I}})$

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MESON SPECTRUM

PRL103 262001 (2009)	I = 1
PRD82 034508 (2010)	$I = 1, K^{\star}$
PRD83 111502 (2011)	I = 0
JHEP07 126 (2011)	CĒ
PRD88 094505 (2013)	I = 0
JHEP05 021 (2013)	D, D_s

HADRON SCATTERING

PRD83 071504 (2011)	$\pi\pi I = 2$
PRD86 034031 (2012)	$\pi\pi I = 2$
PRD87 034505 (2013)	$\pi\pi I = 1, \rho$
PRL113 182001 (2014)	$\pi K, \eta K : K^{\star}$
PRD91 054008 (2015)	$\pi K, \eta K : K^*$
PRD92 094502 (2015)	$\pi\pi, K\bar{K}:\rho$
PRD93 094506 (2016)	$\pi\eta, K\bar{K}: a_0$
ARXIV:1607.05900	$\pi\pi I = 0, \sigma$
ARXIV:1607.07093	$D\pi, D\eta, D_s\overline{K}$

BARYON SPECTRUM

PRD84 074508 (2011) $(N, \Delta)^{\star}$ $(N, \Delta)_{\rm hvb}$ PRD85 054016 (2012) $(N \dots \Xi)^{\star}$ PRD87 054506 (2013) $\begin{array}{c} \Omega_{ccc}^{\star} \\ \Xi_{cc}^{\star} \end{array}$ PRD90 074504 (2014) PRD91 094502 (2015)

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MATRIX ELEMENTS

PRD90 014511 (2014) $f_{\pi^{\star}}$ **PRD91 114501 (2015)** $M' \to \gamma M$ **PRL115 242001 (2015)** $\gamma^* \pi \to \pi \pi$ **PRD93 114508 (2016)** $\gamma^* \pi \to \pi \pi$

LATTICE TECH.

PRD79 034502 (2009) PRD80 054506 (2009) PRD85 014507 (2012)

lattices distillation $\vec{p} > 0$

