

# 5D Maximally Supersymmetric Yang-Mills on the Lattice

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[JHEP 1606 \(2016\) 030 \(arXiv:1604.02707 \[hep-lat\]\)](#)

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# Motivation

A recent claim

M.R. Douglas, [JHEP 02 (2011) 011]  
N. Lambert, C. Papageorgakis, M. Schmidt-Sommerfeld, [JHEP 01 (2011) 083]

$6D (2, 0)$  superconformal theory =  $5D$  MSYM

$D = 5$  MSYM captures all the degrees of freedom of the parent  
 $6D$  theory  $\implies$  these two theories are the same

# Background

## $6D (2, 0)$ theory

- Lagrangian not known!
- Existence can be probed through dual gravitational theory on  $AdS_7 \times S^4$
- UV finite
- Free energy  $\sim N^3$

## $5D \mathcal{N} = 2$ theory

- 16 supercharges
- Perturbatively non-renormalizable
- Also takes part in AdS/CFT
- Free energy  $\sim N^3$  (at large 't Hooft coupling)

## 5D $\mathcal{N} = 2$ Yang-Mills

Obtained by dimensional reduction of 6D (2, 0) theory on a circle

Action

$$S_5 = \frac{1}{g_5^2} \int d^5x \operatorname{Tr} \left( \frac{1}{4} F_{mn} F_{mn} + \frac{1}{2} D_m \phi_j D_m \phi_j + \frac{1}{4} [\phi_j, \phi_k]^2 - i \bar{\lambda}^{aX} (\gamma^m)_a{}^b D_m \lambda_{bX} - \bar{\lambda}^{aX} (\gamma^j)_X{}^Y [\phi_j, \lambda_{aY}] \right)$$

Euclidean Lorentz rotation group:  $SO_E(5)$

Internal rotation group:  $SO_R(5)$

# Twist - A Lattice Compatible Form

We are interested in formulating  $5D$  theory on the lattice

Our approach: Twist the theory

E. Witten, [[Commun. Math. Phys. 117, 353 \(1988\)](#)]

Define a new rotation group

$$SO'(5) = \text{diag subgroup} \left( SO_E(5) \times SO_R(5) \right)$$

On flat space - just a change of variables

(On twisted  $4D$  MSYM: See talks by [Joel Giedt](#) and [David Schaich](#) [this session].)

# Twisted 5D $\mathcal{N} = 2$ Yang-Mills

Twisted action

$$S_5 = \frac{1}{g_5^2} \int d^5x \operatorname{Tr} \left( \frac{1}{4} \overline{\mathcal{F}}_{mn} \mathcal{F}_{mn} + \frac{1}{2} D_m \phi_m D_n \phi_n \right. \\ \left. - i \bar{\lambda}^{AC} (\gamma^m)_A{}^B D_m \lambda_{BC} - \bar{\lambda}^{CA} (\gamma^m)_A{}^B [\phi_m, \lambda_{CB}] \right)$$

5 scalars  $\rightarrow$  vector field  $\phi_m$

Can combine with gauge field  $A_m$  to form a complexified gauge field  $\mathcal{A}_m$

Theory has  $\mathcal{A}_m$  and  $\overline{\mathcal{A}}_m$

Complexified field strengths:  $\mathcal{F}_{mn}$  and  $\overline{\mathcal{F}}_{mn}$

# Twisted 5D $\mathcal{N} = 2$ Yang-Mills

Fermions

$$\lambda_{AB} = \frac{1}{2\sqrt{2}} \left( \frac{1}{2} (\sigma^{mn})_{AB} \chi_{mn} + (\gamma^m)_{AB} \psi_m - \epsilon_{AB} \eta \right)$$

$\eta, \psi_m$  and  $\chi_{mn}$ : p-forms (Grassmann odd)

Similar decomposition for twisted fields

$Q, Q_m, Q_{mn}$

Important property

$$Q^2 \cdot = 0$$

on the lattice

# Twisted Action

In terms of p-form fields

$$S_5 = \frac{1}{g_5^2} \int d^5x \operatorname{Tr} \left( \frac{1}{4} \overline{\mathcal{F}}_{mn} \mathcal{F}_{mn} - \frac{1}{8} [\mathcal{D}_m, \overline{\mathcal{D}}_m]^2 \right. \\ \left. - i \chi_{mn} \mathcal{D}_m \psi_n - i \psi_m \overline{\mathcal{D}}_m \eta - \frac{i}{8} \epsilon_{mncde} \chi_{de} \overline{\mathcal{D}}_c \chi_{mn} \right)$$

Can be written as a sum of  $\mathcal{Q}$ -exact and  $\mathcal{Q}$ -closed terms

$$S_5 = \mathcal{Q}\Lambda - \frac{1}{g_5^2} \int d^5x \operatorname{Tr} \frac{i}{8} \epsilon_{mncde} \chi_{de} \overline{\mathcal{D}}_c \chi_{mn},$$

$$\Lambda = \frac{1}{g_5^2} \int d^5x \operatorname{Tr} \left( \frac{i}{4} \chi_{mn} \mathcal{F}_{mn} - \eta D_m \phi_n - \frac{1}{2} \eta d \right)$$



# Twisted Action

$Q$  transformations on twisted fields

$$Q\mathcal{A}_m = \psi_m$$

$$Q\psi_m = 0$$

$$Q\bar{\mathcal{A}}_m = 0$$

$$Q\chi_{mn} = -i\bar{\mathcal{F}}_{mn}$$

$$Q\eta = d$$

$$Qd = 0$$

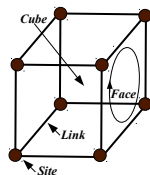
$d$ : auxiliary field

Exactly nilpotent:  $Q^2 \cdot = 0$ .

# On the Lattice

Straightforward. Use geometric discretization

- Discretize theory on a **hypercube**.
- Place p-form fields on p-cells of hypercube



$$\hat{\mu}_1 = (1, 0, 0, 0, 0)$$

$$\hat{\mu}_2 = (0, 1, 0, 0, 0)$$

$$\hat{\mu}_3 = (0, 0, 1, 0, 0)$$

$$\hat{\mu}_4 = (0, 0, 0, 1, 0)$$

$$\hat{\mu}_5 = (0, 0, 0, 0, 1)$$

Note: We have only 16 fermions not 32, to fill the hypercube!

## Lattice covariant difference operators appropriate for p-forms

P. H. Damgaard and S. Matsuura, [JHEP 0707, 051 (2007)]  
S. Catterall, [JHEP 0801, 048 (2008)]

$$\mathcal{D}_m^{(+)} f(\mathbf{n}) = \mathcal{U}_m(\mathbf{n}) f(\mathbf{n} + \hat{\mu}_m) - f(\mathbf{n}) \mathcal{U}_m(\mathbf{n})$$

$$\mathcal{D}_m^{(+)} f_n(\mathbf{n}) = \mathcal{U}_m(\mathbf{n}) f_n(\mathbf{n} + \hat{\mu}_m) - f_n(\mathbf{n}) \mathcal{U}_m(\mathbf{n} + \hat{\mu}_n)$$

$$\overline{\mathcal{D}}_m^{(-)} f_m(\mathbf{n}) = f_m(\mathbf{n}) \overline{\mathcal{U}}_m(\mathbf{n}) - \overline{\mathcal{U}}_m(\mathbf{n} - \hat{\mu}_m) f_m(\mathbf{n} - \hat{\mu}_m)$$

$$\overline{\mathcal{D}}_c^{(+)} f_{mn}(\mathbf{n}) = f_{mn}(\mathbf{n} + \hat{\mu}_c) \overline{\mathcal{U}}_c(\mathbf{n}) - \overline{\mathcal{U}}_c(\mathbf{n} + \hat{\mu}_m + \hat{\mu}_n) f_{mn}(\mathbf{n})$$

$\mathcal{Q}$  transformations on the lattice

$$\begin{aligned}\mathcal{Q}\mathcal{U}_m(\mathbf{n}) &= \psi_m(\mathbf{n}) \\ \mathcal{Q}\psi_m(\mathbf{n}) &= 0 \\ \mathcal{Q}\bar{\mathcal{U}}_m(\mathbf{n}) &= 0 \\ \mathcal{Q}\chi_{mn}(\mathbf{n}) &= -i\left(\bar{\mathcal{D}}_m^{(+)}\bar{\mathcal{U}}_n(\mathbf{n})\right) = -i\bar{\mathcal{F}}_{mn}^L \\ \mathcal{Q}\eta(\mathbf{n}) &= d(\mathbf{n}) \\ \mathcal{Q}d(\mathbf{n}) &= 0\end{aligned}$$

Similar to the continuum transformations.

# Lattice Theory

$\mathcal{Q}$ -exact piece

$$S_{\mathcal{Q}\text{-exact}} = \beta \sum_{\mathbf{n}, m, n} \text{Tr} \left[ -\frac{1}{4} \overline{\mathcal{F}}_{mn}^L(\mathbf{n}) \mathcal{F}_{mn}(\mathbf{n}) - \frac{1}{8} \left( \overline{\mathcal{D}}_m^{(-)} \mathcal{U}_m(\mathbf{n}) \right)^2 \right. \\ \left. - i \chi_{mn}(\mathbf{n}) \mathcal{D}_m^{(+)} \psi_n(\mathbf{n}) - i \eta(\mathbf{n}) \overline{\mathcal{D}}_m^{(-)} \psi_m(\mathbf{n}) \right]$$

Gauge-invariant and local

$\mathcal{Q}$ -closed piece has to be modified

Open loop on the lattice. Need to close it.

# Lattice Theory

Introduce a path ordered link.

Option 1. Use  $\bar{\mathcal{U}}$ -fields

$$\mathcal{P}_{\text{POL}} \equiv \prod_{l \in C_L} \bar{\mathcal{U}}_l$$

Action is  $\mathcal{Q}$ -invariant, gauge-invariant but non-local.

Option 2. Use  $\mathcal{U}$ -fields

$$\mathcal{P}_{\text{POL}} \equiv \prod_{l \in C_L} \mathcal{U}_l$$

Action is  $\mathcal{Q}$  non-invariant, gauge-invariant and local.

Need to introduce additional operator

$$\left( \mathcal{Q} \mathcal{P}_{\text{POL}} \right) \times (\dots)$$

to make it  $\mathcal{Q}$  invariant

# Lattice Theory

Lattice action is:  $S = S_{\mathcal{Q}\text{-exact}} + S_{\mathcal{Q}\text{-closed}}$

$$S_{\mathcal{Q}\text{-exact}} = \beta \sum_{\mathbf{n}, m, n} \text{Tr} \left[ -\frac{1}{4} \overline{\mathcal{F}}_{mn}^L(\mathbf{n}) \mathcal{F}_{mn}(\mathbf{n}) - \frac{1}{8} \left( \overline{\mathcal{D}}_m^{(-)} \mathcal{U}_m(\mathbf{n}) \right)^2 \right. \\ \left. - i \chi_{mn}(\mathbf{n}) \mathcal{D}_m^{(+)} \psi_n(\mathbf{n}) - i \eta(\mathbf{n}) \overline{\mathcal{D}}_m^{(-)} \psi_m(\mathbf{n}) \right]$$

$$S_{\mathcal{Q}\text{-closed}} = -\frac{i\beta}{8} \sum_{\mathbf{n}, m, n, c, d, e} \text{Tr} \epsilon_{mncde} \mathcal{P}_{\text{POL}} \\ \times \chi_{de}(\mathbf{n} + \hat{\mu}_m + \hat{\mu}_n + \hat{\mu}_c) \overline{\mathcal{D}}_c^{(+)} \chi_{mn}(\mathbf{n})$$

$\mathcal{Q}S = 0$ , if we use Option 1. Additional operator needed for Option 2.

# Lattice Theory

- A well defined theory at finite lattice spacing
- Can simulate at **finite coupling** and **finite  $N$**
- No doublers
- Theory has naive continuum limit
- How about quantum continuum limit?
  - Look for 2nd (or higher) order phase transitions



## Future Directions

- Most pressing question:

Does it have a non-trivial UV fixed point on the lattice?

If so, it can give UV completion of the  $5D$  theory and thus a non-perturbative definition

- One can check the scaling of free-energy on the lattice:  
 $N^3$ ? (at all couplings?)
- Validating gauge-gravity duality

Lots of interesting explorations to do...!

THANK YOU!