

Phenomenology of a UV complete model

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Criteria for UV completion

how can we build a BSM model given the lack of experimental evidence?

- (1) Higgs boson is a pNGB
- (2) top quark mass due to mixing with a composite state

global symmetry: $G_F \rightarrow H_F$, $H_F \supset G_{\text{cus}} \supset G_{\text{SM}}$

$$G_{\text{cus}} = \text{SU}(3)_c \times \text{SU}(2)_L \times \text{SU}(2)_R \times U(1)_X$$

$$G_F/H_F \ni (1, 2, 2)_0$$

asymptotic freedom + unitary gauge group + several irreps (LH Weyl)

$$\text{fermions} \in mR_1 + 3R_2 + 3R_3$$

$$G_F = \text{SU}(m) \times \text{SU}(3) \times \text{SU}(3)' \times U(1) \times U(1)'$$

[Ferretti & Karateev 13]

Ferretti's model - UV description

	G_{HC}	G_F				
	$SU(4)$	$SU(5)$	$SU(3)$	$SU(3)'$	$U(1)_X$	$U(1)'$
ψ	6	5	1	1	0	-1
χ	4	1	3	1	-1/3	5/3
$\tilde{\chi}$	$\bar{\mathbf{4}}$	1	1	$\bar{\mathbf{3}}$	1/3	5/3

[Ferretti 14]

strong dynamics induces SSB of the global symmetry

$$\begin{aligned}
 SU(5) &\longrightarrow SO(5) \\
 SU(3) \times SU(3)' &\longrightarrow SU(3)_c
 \end{aligned}$$

compute the spectrum

[plenary talk by Claudio]

Ferretti's model - composite Higgs model

$$SU(5) \longrightarrow SO(5) \supset SO(4) \simeq SU(2)_L \times SU(2)_R$$

SM group: $T_L^a \in SU(2)_L$, $T_R^3 \in U(1)$

$$14 \text{ NGB} : 1_0 + 2_{\pm 1/2} + 3_0 + 3_{\pm 1}$$

$$\eta \quad H \quad \Phi_0 \quad \Phi_{\pm}$$

$$\Pi = H + H^\dagger + \Phi_0 + \Phi_+ + \Phi_+^\dagger, \quad \Sigma = \exp\left(\frac{i\Pi}{f}\right)$$

$$U = \Sigma \Sigma^T, \quad U \mapsto g U g^T$$

using the notation $\phi_{m_R}^{m_L}$, exotic Higgses with $Q = m_L + m_R$:

$$\Phi_0 \supset (\phi_0^-, \phi_0^0, \phi_0^+)$$

$$\Phi_+ \supset (\phi_+^-, \phi_+^0, \phi_+^+)$$

Gauging the EW - pNGB potential

$$\mathcal{L}_{\text{eff}} \supset \frac{f^2}{16} \text{tr} \left[(D_\mu U)^\dagger D_\mu U \right]$$
$$D_\mu U = \partial_\mu U - ig W_\mu^a [T_L^a, U] - ig' B_\mu [T_R^3, U]$$

Coleman-Weinberg potential due to the gauge bosons:

$$V_g = -B f^4 \text{tr} \left[g^2 T_L^a U T_L^a U^\dagger + g'^2 T_R^3 U T_R^3 U^\dagger \right]$$

B is a LEC, V_g provides a mass for the extra pNGBs in the spectrum

[Ferretti 16]

$$B f^4 = C_{LR} \propto \int_0^\infty dq^2 q^2 \Pi_{LR}(q^2)$$
$$(q^2 \delta_{\mu\nu} - q_\mu q_\nu) \Pi_{LR}(q^2) = \int d^4x e^{iqx} \langle J_\mu^L(x) J_\nu^R(0) \rangle$$

[Golterman & Shamir 15]

Partial compositeness

generates top quark mass, and EW symmetry breaking

[Kaplan 85]

$$\mathcal{L}_{\text{UV}} \supset \frac{\lambda_L}{\Lambda_{\text{UV}}^{d_L-5/2}} \bar{q}_L \mathcal{O}_L + \frac{\lambda_R}{\Lambda_{\text{UV}}^{d_R-5/2}} \bar{t}_R \mathcal{O}_R + \text{h.c.}$$

[Panico & Wulzer 15]

in the UV complete model:

$$\mathcal{L}_{\text{UV}} \supset \frac{\lambda_L}{\Lambda_{\text{UV}}^{d_L-5/2}} \bar{q}_L \mathcal{B}_R + \frac{\lambda_R}{\Lambda_{\text{UV}}^{d_R-5/2}} \bar{t}_R \mathcal{B}_L + \text{h.c.}$$



q_L, t_R spurion fields, $\mathcal{B}_{L,R}$ hyperbaryon fields ($\chi\psi\chi$), in the $(\mathbf{5}, \mathbf{3})$ of H_F

[Golterman & Shamir 15]

Fermionic mass term

$$\mathcal{L}_{\text{eff}} \supset \frac{M}{2} \bar{\Psi} \Psi + \lambda_q f \bar{q}_L \Sigma \Psi_R + \lambda_t f \bar{t}_R \Sigma^* \Psi_L + \dots$$

$$\lambda_q = \lambda_L \left(\frac{\Lambda}{\Lambda_{\text{UV}}} \right)^{\gamma_L} \sqrt{Z_B}, \quad \lambda_t = \lambda_R \left(\frac{\Lambda}{\Lambda_{\text{UV}}} \right)^{\gamma_R} \sqrt{Z_B}$$

$$\Psi = \frac{1}{\sqrt{2}} \begin{pmatrix} iB - iX \\ B + X \\ iT + iY \\ -T + Y \\ \sqrt{2}iR \end{pmatrix}$$

$$(T, B) \in (\mathbf{3}, \mathbf{2})_{1/6}, \quad R \in (\mathbf{3}, \mathbf{1})_{2/3}, \quad (X, Y) \in (\mathbf{3}, \mathbf{2})_{7/6}$$

Fermionic mass term

X exotic state, with charge 5/3

$$\mathcal{L}_{\text{eff}} \supset (\bar{t}_L, \bar{T}_L, \bar{Y}_L, \bar{R}_L) \cdot \mathcal{M}_T \cdot \begin{pmatrix} t_R \\ T_R \\ Y_R \\ R_R \end{pmatrix} + \dots$$

top quark mass:

$$m_t/v = \sqrt{2} \frac{f}{M} \frac{1}{\sqrt{1 + \lambda_q^2 \frac{f^2}{M^2}} \sqrt{1 + \lambda_t^2 \frac{f^2}{M^2}}}$$

Higgs potential

$$V(h) = -\alpha f^2 \sin^2 \frac{h}{f} + \beta f^2 \sin^4 \frac{h}{f}$$

[Panico & Wulzer 15]

where

$$\alpha = C_{LR}(3g^2 + g'^2) + 2y^2 C_{\text{top}} < 0$$

$$\beta = -2y^2 C_{\text{top}}$$

$$C_{\text{top}} \propto \int \frac{d^4 p}{(2\pi)^4} \frac{d^4 q}{(2\pi)^4} \frac{p_\mu q_\nu}{p^2 q^2} \langle \bar{\mathcal{B}}_{Rk}(p) \gamma_\mu P_R \mathcal{B}_{Ri}(p) \bar{\mathcal{B}}_{Li}(q) \gamma_\nu P_L \mathcal{B}_{lk}(p) \rangle$$

[Golterman & Shamir 15]

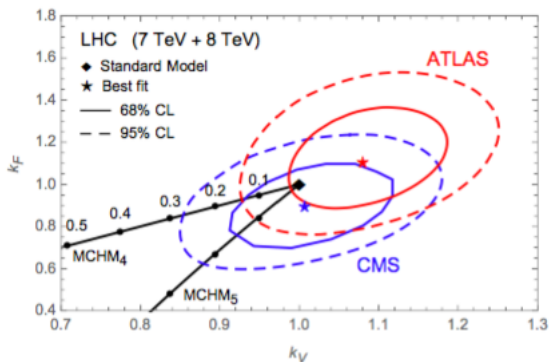
$$\xi = v^2 / f^2 = \alpha / (2\beta)$$

$$\alpha = \frac{1}{4} \frac{m_h^2}{v^2} \frac{\xi}{1 - \xi}$$

Phenomenology I - modified coupling

Ferretti's model similar to generic MCHM5, $k \sim \sqrt{1 - \xi}$

[Contino et al 06]



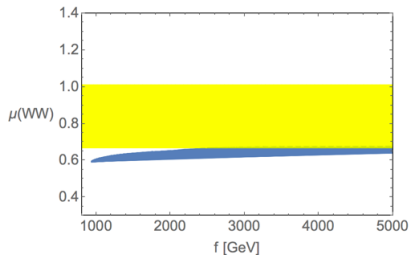
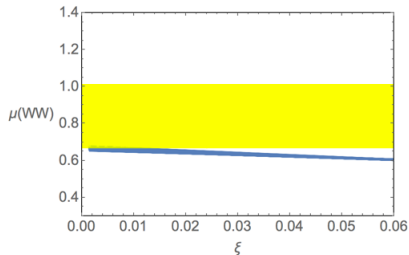
[Panico & Wulzer 15]

$$\xi < 0.15, \quad \xi \rightarrow 0 \iff \text{SM}$$

Phenomenology II - exotica

$M/\text{TeV} \in [1.5, 3.5]$, $\lambda_t \in [0, 4\pi]$, using Feynrules, MadEvent

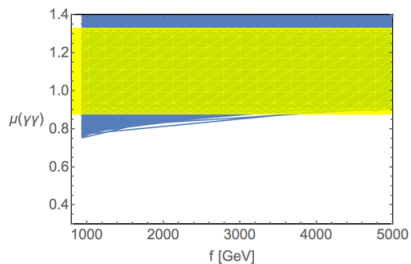
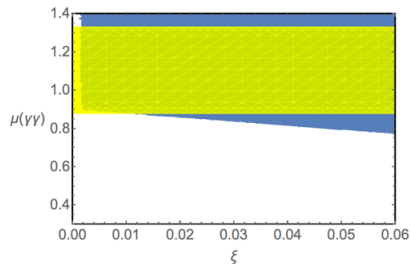
Higgs decay to WW



data from [Atlas & CMS 16]

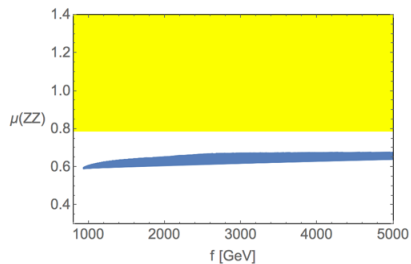
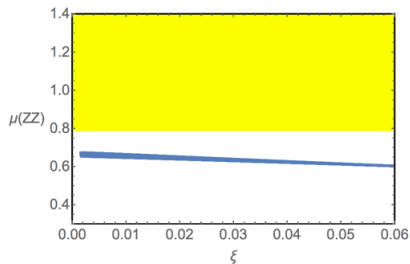
Phenomenology III - exotica

loop-induced Higgs decay to $\gamma\gamma$



Phenomenology IV - exotica

Higgs decay to ZZ



some tension in this channel

Scope for lattice simulations

Ferretti's model is not ruled out by LHC Run1 data

Strong constraints from a lattice computation of α and β

Estimate the precision required in order to have an impact on pheno

The lattice simulation is a difficult computation...

[Svetitsky, Jay later today]

... and needs to be extrapolated to the chiral limit

Ferretti? or someone else? ... theoretical tools, new observables



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