

dark matter from one-flavor SU(2) gauge theory

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outline

- motivation
- the basic theory
- preliminary lattice explorations
- coupling to the standard model

motivation

- a minimal non-Abelian dark sector
- dark matter stability exists naturally

SU(2) with one Dirac flavor

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}^a F^{a\mu\nu} + i\bar{Q}\gamma^\mu D_\mu Q + mQ^T C E Q$$

$$\text{where } Q = \begin{pmatrix} \chi_L \\ C\bar{\chi}_R^T \end{pmatrix}, \quad E = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$$

Q has no standard model quantum numbers.

\mathcal{L} has an unbroken global SU(2), $Q \rightarrow e^{i\sum_{i=1}^3 T_i \alpha_i} Q$, which is the generalization of baryon number.

For $m=0$, \mathcal{L} has an unbroken (but anomalous) global U(1), $Q \rightarrow e^{i\beta} Q$, like the axial U(1) in QCD. We expect the U(1) to be broken dynamically by a mass-like vev.

For $N_f > 1$, the global SU(2) would be SU($2N_f$).

the particle spectrum

The theory will have mesons, baryons and glueballs.

Simple operators for mesons: $\bar{Q}\Gamma Q$

Simple operators for baryons: $Q^T C \Gamma E Q$

$$\text{Recall } E = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$$

The spectrum forms multiplets of the global SU(2).

Examples:

$\bar{Q}\gamma_5 Q$ is a singlet. Let's name it η .

$\bar{Q}\gamma_\mu Q$ and baryon and antibaryon form a triplet $\rho^{\pm,0}$.

lattice ensembles

plaquette gauge action

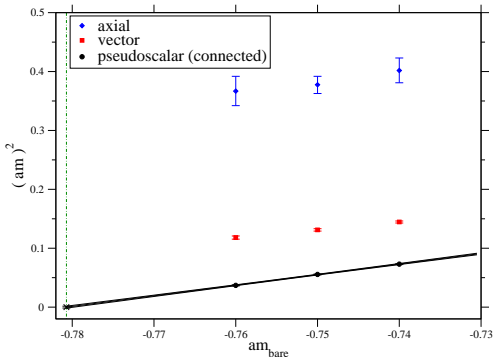
Wilson fermion action

HiRep code [Del Debbio, Patella, Pica, PRD81, 094503 (2010)]

β	2.2	2.309
lattice dimensions	$20^3 \times 56$	$28^3 \times 56$
number of configurations	2000	1540
acceptance	73%	74%
unitary m_{bare}	-0.865	-0.76
partially quenched m_{bare}	-0.845, -0.855	-0.74, -0.75
average plaquette	0.5989	0.6255
aw_0	1.430(5)	1.956(7)
$m_V L$	9.0(2)	8.8(2)

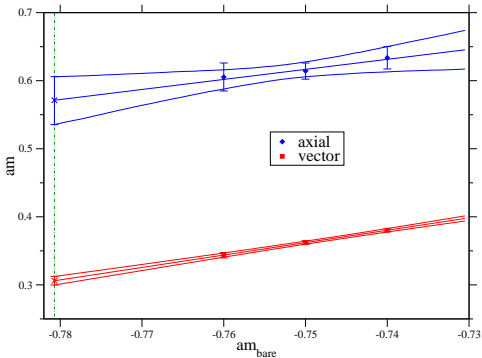
preliminary lattice spectrum

$\beta=2.309$



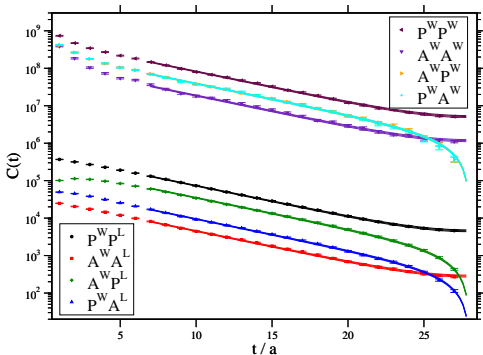
linear extrapolations

$$\beta=2.309$$



pseudoscalar correlators

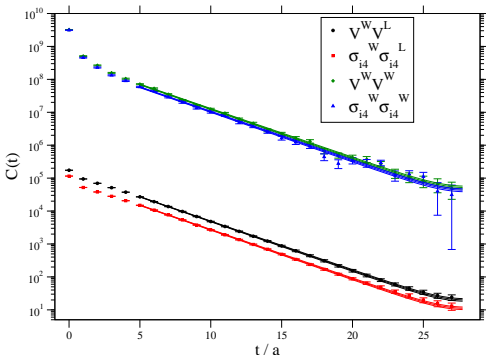
connected part only, $\beta=2.309$, $am_{\text{bare}} = -0.76$



W=wall (Coulomb gauge fixed), L=local

vector correlators

$$\beta=2.309, \quad am_{\text{bare}} = -0.76$$



W=wall (Coulomb gauge fixed), L=local

a large N_c limit

For $N_c > 2$,
one-flavor $SU(N_c)$ does *not* have the global $SU(2)$.

But recall $SU(2) = Sp(2)$.

The global $SU(2)$ is present for $Sp(N_c)$. (N_c is even.)

In $Sp(N_c)$,
the η becomes massless as $N_c \rightarrow \infty$ and $m_Q \rightarrow 0$.
In this double limit, the global $U(1)$ is only broken
dynamically and the η is its Goldstone boson.

particle spectrum for $N_c \rightarrow \infty$

How can the global SU(2) remain unbroken?

Mesons have this fermion content: $M = \sum_{i=1}^{N_c} \bar{Q}_i Q_i$

Baryons: ??? $X = \sum_{i,j,\dots,k=1}^{N_c} Q_i Q_j \dots Q_k$??? No.

Sp(N_c) baryons are $B = \sum_{i=1}^{N_c} \sum_{j=1}^{N_c} Q_i E_{ij} Q_j$

X operators reduce to a collection of B operators.

Thus the global SU(2) remains unbroken.

Baryons B and mesons M remain degenerate.

Higgs couplings to the dark sector

Can dim=5 BSM physics destabilize dark matter? No.

$$\text{dim 5 : } \delta\mathcal{L} \sim \bar{Q}\gamma_5 Q H^\dagger H \quad (\text{couples to } \eta)$$

$$\text{dim 6 : } \delta\mathcal{L} \sim \bar{Q}\gamma^\mu Q H^\dagger \nabla_\mu H \quad (\text{couples to } \rho)$$

Dark matter, ρ , is stable at dimension 5.

This feature is specific to the one-flavor theory.

Like SM proton decay, ρ decay is dimension 6.

Thus ρ can be stable for the life of the universe.

fermion mass with parity violation

The dark fermion gets mass from 2 sources:
BSM (dimension 4) and the SM Higgs (dimension 5):

$$\begin{aligned}\delta\mathcal{L} = & -m_4 \cos\theta_4 \bar{Q}Q - m_4 \sin\theta_4 \bar{Q}i\gamma_5 Q \\ & -\frac{v^2}{\Lambda} \cos\theta_5 \bar{Q}Q \left(1 + \frac{h}{v}\right)^2 - \frac{v^2}{\Lambda} \sin\theta_5 \bar{Q}i\gamma_5 Q \left(1 + \frac{h}{v}\right)^2\end{aligned}$$

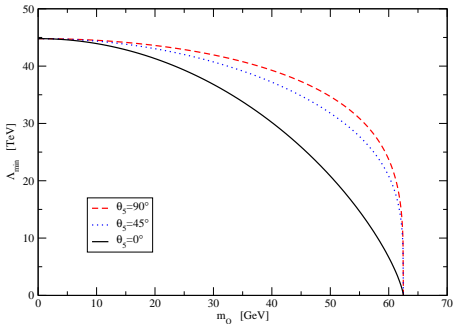
Λ =BSM scale. $v=246$ GeV. Parity violation allowed.

Mass terms: $\delta\mathcal{L} = m\bar{Q}_{\text{tw}}Q_{\text{tw}}$ where $Q_{\text{tw}} \equiv e^{i\gamma_5\alpha/2}Q$.

Vector and axial hadrons are invariant under $Q \rightarrow Q_{\text{tw}}$.

invisible Higgs decays for SU(2)

The experimental bound $\Gamma(h \rightarrow Q\bar{Q}) < 1.2 \text{ MeV}$ gives



decays from the dark sector

The only $\text{dark} \rightarrow \text{SM}$ decays are through Higgs bosons.
If $\theta_5 \neq 0$, then η can decay through a single Higgs.

Recall: ρ is essentially stable due to the global SU(2).

Example:

BBN \Rightarrow lifetime of $\eta < 1$ second.

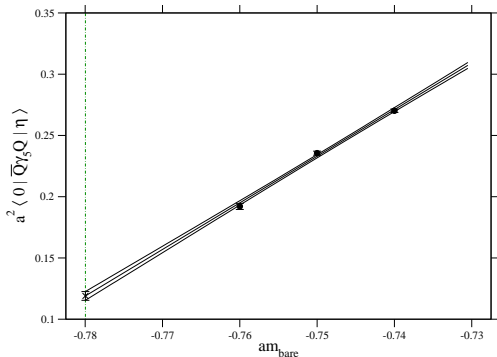
Use $\langle 0 | \bar{Q} \gamma_5 Q | \eta \rangle$ from lattice to bound $\frac{\sin \theta_5}{\Lambda}$.

For $m_\eta \ll m_H$,

$$\Gamma_\eta = |\langle 0 | \bar{Q} \gamma_5 Q | \eta \rangle|^2 \frac{m_\eta \sin^2 \theta_5}{2\pi \Lambda^2 m_H^4} \sum_{f \in \text{SM}} m_f^2 \left(1 - \frac{4m_f^2}{m_\eta^2} \right)^{3/2}$$

amplitude useful for η decay

connected part only, $\beta=2.309$



also in progress

1. direct detection: ρ scattering from a SM nucleon.

2. relic density:

Options for $m_\eta < m_\rho$ include

- $\rho\rho \rightarrow \eta\eta$
- asymmetric dark matter

Options for $m_\eta > m_\rho$ include

- $\rho\rho\rho \rightarrow \rho\rho$

Future lattice simulations will reveal $m_\eta - m_\rho$ ordering.
Disconnected diagrams are required.

For disconnected calculations in the 2-flavour theory, see [Arthur, Drach, Hietanen, Pica, Sannino, 1607.06654](#) and [Drach, Monday 14:15](#)

summary

1-flavor $SU(2)$ is a minimal non-Abelian dark sector.

It has a *global* $SU(2)$ to stabilize dark matter.

It has no dark matter decay at dimension 5.

1 Goldstone boson should emerge for $Sp(N_c \rightarrow \infty)$.

$SU(2)=Sp(2)$ lattice explorations are in progress.