

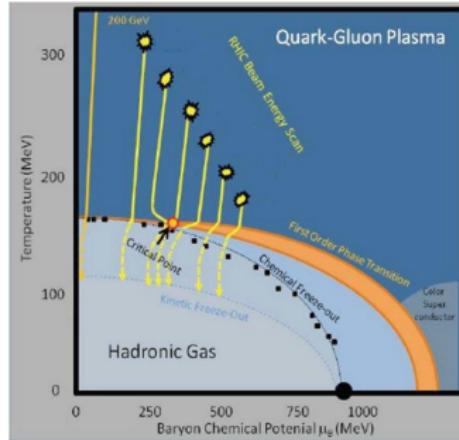
# The QCD equation of state at finite density from analytical continuation

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for the Wuppertal-Budapest-Collaboration

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# The $(T, \mu_B)$ -phase diagram of QCD



Our observables:

Last Year:  $T_c$

R. Bellwied et al., Phys. Lett. B751, 559 (2015), arXiv:1507.07510

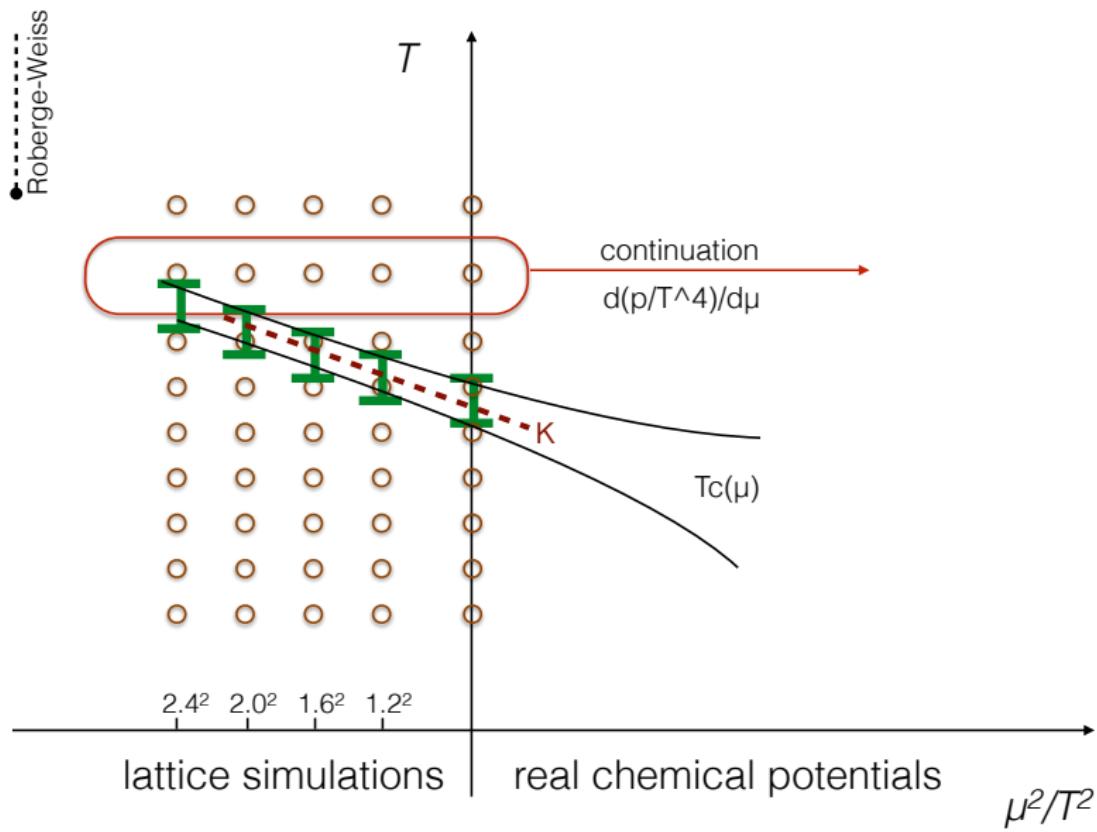


This year: The Equation of State along trajectories of constant  $\frac{S}{N_B}$  and its Taylor coefficients determined by the method of analytical continuation

J. Günther et al., arXiv:1607.02493

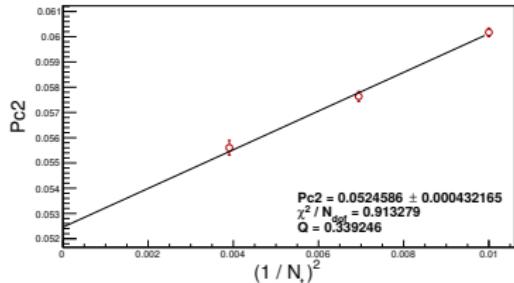
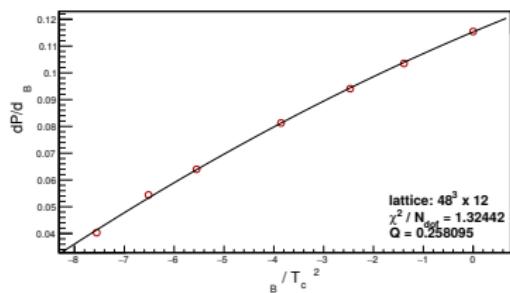
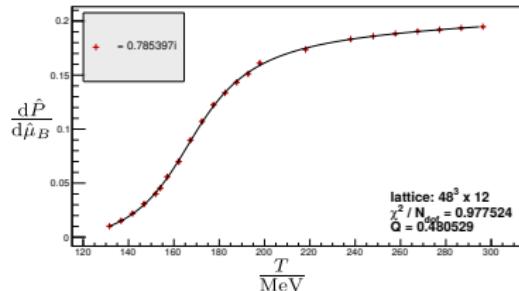


# Analytic continuation

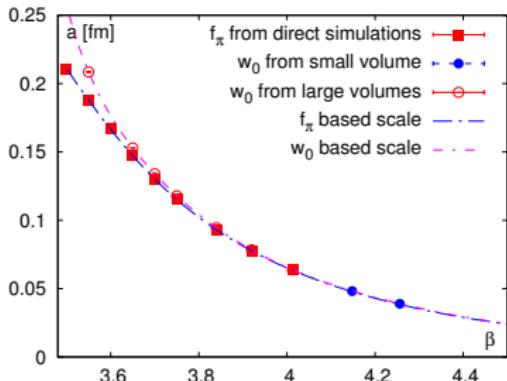


# Overview over the Analysis

1. Do the simulations at  $\langle n_s \rangle \approx 0$
2. Extrapolate to  $\langle n_s \rangle = 0$  and  $\langle n_Q \rangle = 0.4\langle n_B \rangle$
3. Make a fit in the  $T$  direction
  
4. Determine everything you need for the observables
5. Make a fit in the  $\mu_B$  direction
  
6. Make a fit in the  $\frac{1}{N_t^2}$  direction
7. Determine the observables



# Simulation details



- ▶ Action: tree-level Symanzik improved gauge action, with four times stout smeared staggered fermions
- ▶ 2+1+1 flavour, on LCP with pion and kaon mass
- ▶ Simulation at  $\langle n_S \rangle = 0$  (as for heavy ion collisions, in contrast to simulations with  $\mu_s = 0$  or  $\mu_S = 0$  where  $\mu_S = \frac{1}{3}\mu_B - \mu_s$ )
- ▶ Lattice sizes:  $40^3 \times 10$ ,  $48^3 \times 12$  and  $64^3 \times 16$
- ▶  $\frac{\mu_B}{T} = i \frac{j\pi}{8}$  with  $j = 0, 3, 4, 5, 6, 6.5$  and  $7$
- ▶ Two methods of scale setting:  $f_\pi$  and  $w_0$ ,  $Lm_\pi > 4$

## Tuning to $\langle n_S \rangle = 0$

Aim: For a given  $\mu_B$  determine  $\mu_S$  so that  $\langle n_S \rangle = 0$ . This means solving the differential equation

$$\langle n_S \rangle = 0 \Leftrightarrow \frac{\partial \log Z}{\partial \mu_S} = 0$$

Notation:

$$\chi_{udsc} = -\frac{1}{T^4} \frac{\partial^4}{\partial(\mu_u/T) \partial(\mu_d/T) \partial(\mu_s/T) \partial(\mu_c/T)} \frac{T}{V} \log Z$$

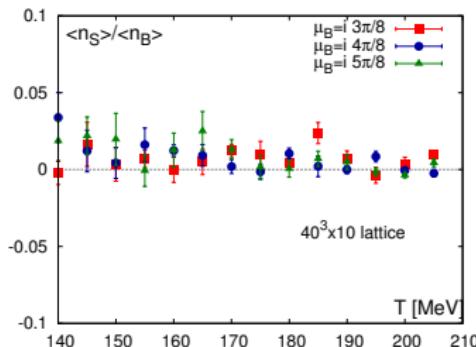
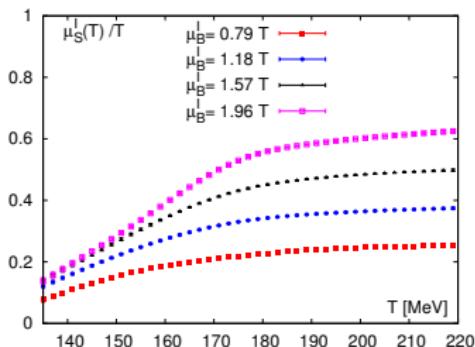
Assuming we know the value for  $\mu_S(\mu_B)$  so that  $\langle n_s \rangle = 0$  for  $\mu_S(\mu_B^0)$  and  $\mu_S(\mu_B^0 - \Delta\mu_B)$  with all the derivatives. Then (Runge-Kutta):

$$\mu_S(\mu_B^0 + \Delta\mu_B) = \mu_S(\mu_B^0 - \Delta\mu_B) + 2\Delta\mu_B \frac{d\mu_S}{d\mu_B}(\mu_B^0).$$

In the simulations with  $\mu_B^0$  and  $\mu_B^0 - \Delta\mu_B$ ,  $\mu_S$  might not precisely tuned. There we want to extrapolate to a better value. We assume that correct value of  $\mu_S$  is  $\tilde{\mu}_S = \mu'_S + \Delta\mu'_S$ . Then:

$$\langle n_S \rangle = \frac{\partial \log Z}{\partial \tilde{\mu}_S} = \frac{\partial \log Z}{\partial \mu'_S} + \frac{\partial^2 \log Z}{\partial \mu'^2_S} \Delta\mu'_S = 0$$

# Tuning to $\langle n_S \rangle = 0$



This yields

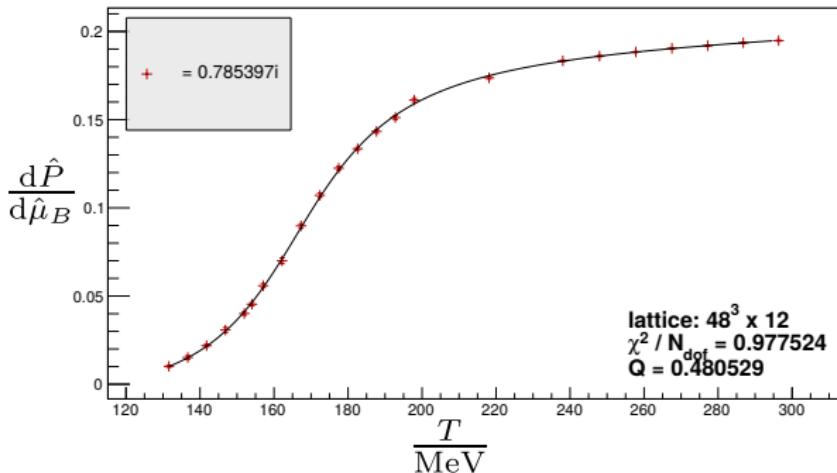
$$\Delta\mu'_S = -\frac{\chi_S}{\chi_{SS}}$$

Similar for the derivative we get:

$$\frac{d\tilde{\mu}_S}{d\mu_B} = -\frac{\tilde{\chi}_{SB}}{\tilde{Z}_{SS}}|_{\langle n_S \rangle=0} = -\frac{\chi_{SB}}{\chi_{SS}} - \frac{\chi_{SSB}\chi_{SS} - \chi_{SSS}\chi_{SB}}{(\chi_{SS})^2} \Delta\mu'_S + \mathcal{O}(\Delta\mu'^2_S)$$

With similar arguments we get to extrapolate to  $\langle n_Q \rangle = 0.4\langle n_B \rangle$

## Fit in the $T$ direction



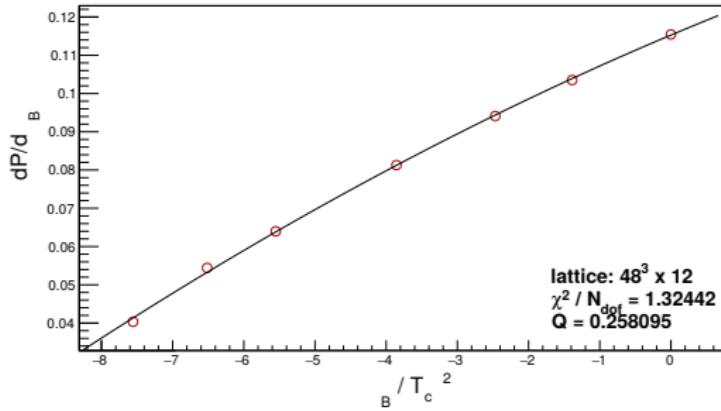
$$A_1(T) = a + bT + c/T + d \arctan(e(T - f))$$

$$A_2(T) = a + bT + c/T + d/(1 + e(T - f)^g)^{1/g},$$

$$A_3(T) = a + bT + cT^2 + d \arctan(e(T - f))$$

$$A_4(T) = a + bT + cT^2 + d/(1 + e(T - f)^g)^{1/g}.$$

# Fit in the $\mu_B$ direction

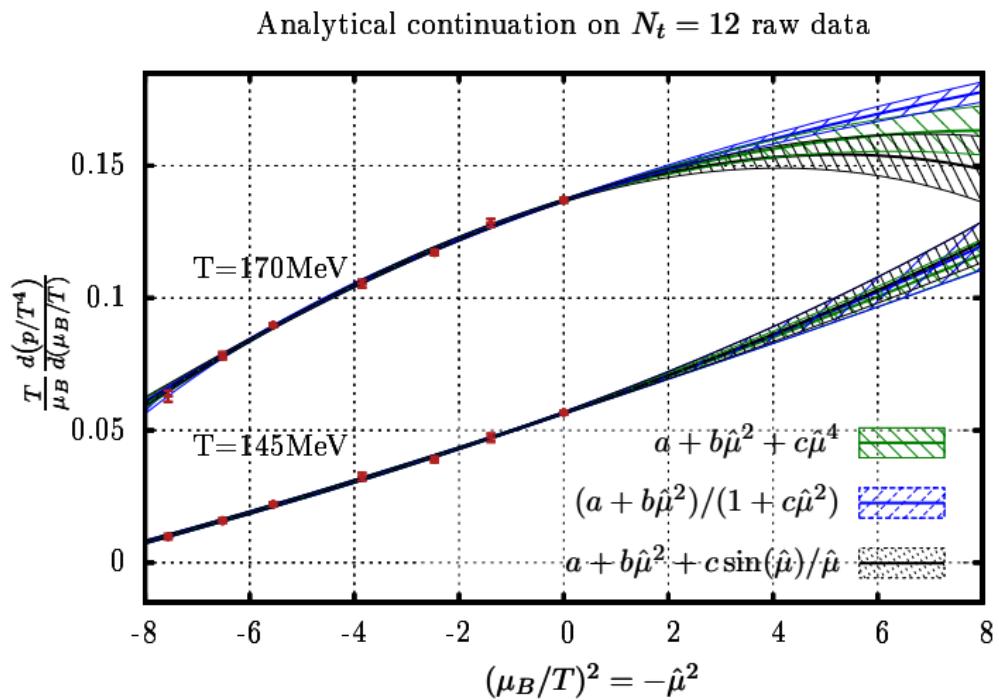


$$B_1(\hat{\mu}) = a + b\hat{\mu}^2 + c\hat{\mu}^4$$

$$B_2(\hat{\mu}) = (a + b\hat{\mu}^2)/(1 + c\hat{\mu}^2)$$

$$B_3(\hat{\mu}) = a + b\hat{\mu}^2 + c \sin(\hat{\mu})/\hat{\mu}$$

# Extrapolation from different fit functions



# Entropy

- ▶  $S = \frac{\partial P}{\partial T} \Big|_{\mu_i}$
- ▶ We have  $\frac{d\hat{P}}{d\hat{\mu}_B}$
- ▶ And we can only do a total derivative in  $T$

$$\frac{d\hat{P}}{dT} = \frac{\partial \hat{P}}{\partial T} \Bigg|_{\mu_i} + \frac{d\mu_B}{dT} \frac{\hat{n}_B}{T} + \frac{d\mu_Q}{dT} \frac{\hat{n}_Q}{T}$$

How we correct this:

$$\frac{d\mu_B}{dT} = \frac{d(\hat{\mu}_B T)}{dT} = \hat{\mu}_B$$

$$\frac{d\mu_Q}{dT} = \frac{d(\hat{\mu}_Q T)}{dT} = \hat{\mu}_Q + T \frac{d\hat{\mu}_Q}{dT}$$

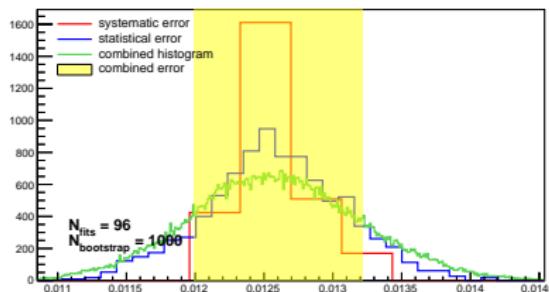
$$\begin{aligned}\hat{S} &= 4\hat{P} + T \frac{d\hat{P}}{dT} - \hat{\mu}_B \hat{n}_B - \hat{\mu}_Q \hat{n}_Q - T \frac{d\hat{\mu}_Q}{dT} \\ &= 4\hat{P} + T \frac{d\hat{P}}{dT} - \hat{n}_B (\hat{\mu}_B + 0.4\hat{\mu}_Q) - T \frac{d\hat{\mu}_Q}{dT}\end{aligned}$$

New terms we have to determine

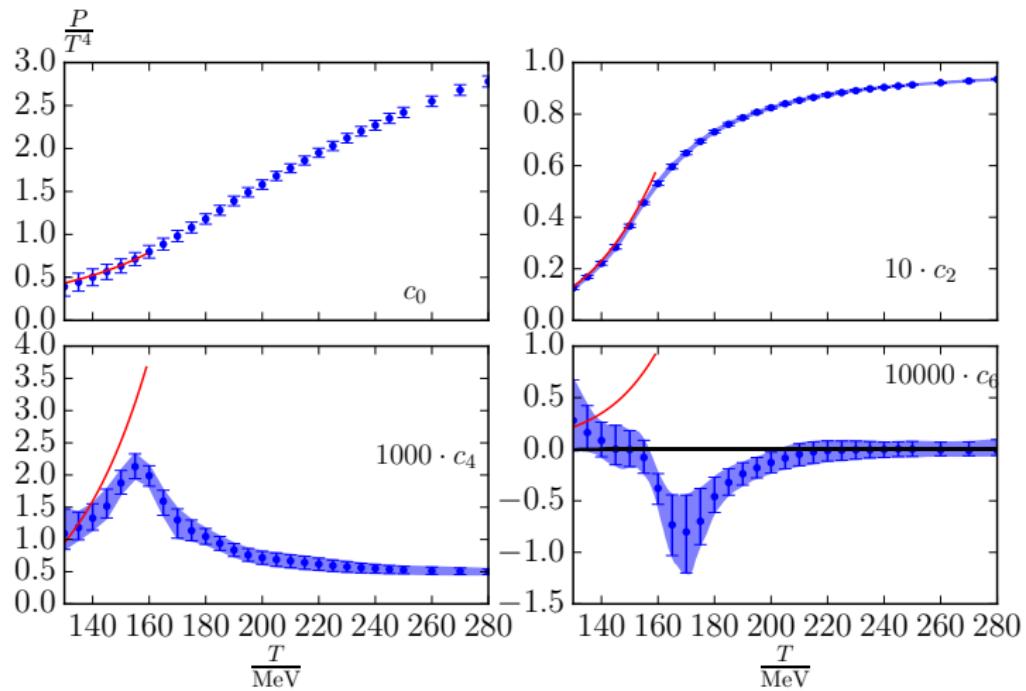
# Error estimation

- ▶ Statistical error:  
Bootstrap method
- ▶ Systematic error:  
Using different way of analysis, combining them in a histogram:
  - ▶ 4 fit functions for the  $T$  direction
  - ▶ 3 fit functions in the  $\mu_B$  direction
  - ▶ Doing continuum extrapolation and  $\mu_B$ -fit in one or two steps
  - ▶ 2 methods of scale setting:  $f_\pi$  and  $w_0$
  - ▶ 2 temperatures from where we use the extrapolated data

This adds up to 96 ways of analysis



# Taylor coefficients



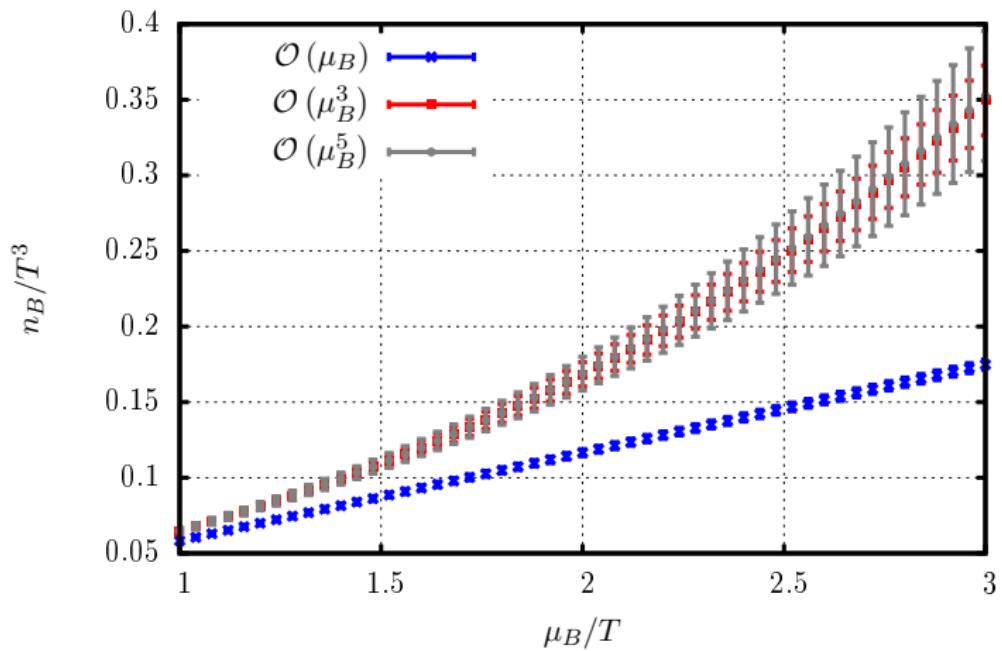
The Taylor coefficients of  $\frac{P}{T^4} = c_0 + c_2 \left(\frac{\mu_B}{T}\right)^2 + c_4 \left(\frac{\mu_B}{T}\right)^4 + c_6 \left(\frac{\mu_B}{T}\right)^6$

$c_0$  is from S. Borsanyi et al, Phys. Lett. B730, 99 (2014), arXiv:1309.5258



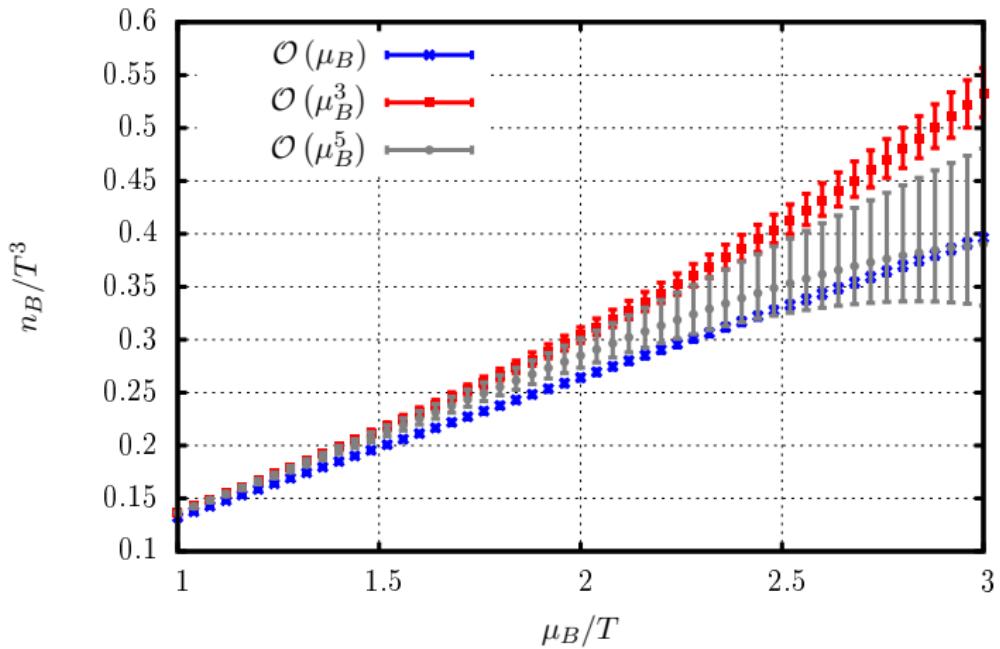
# Influence of different orders

$T = 145\text{MeV}$

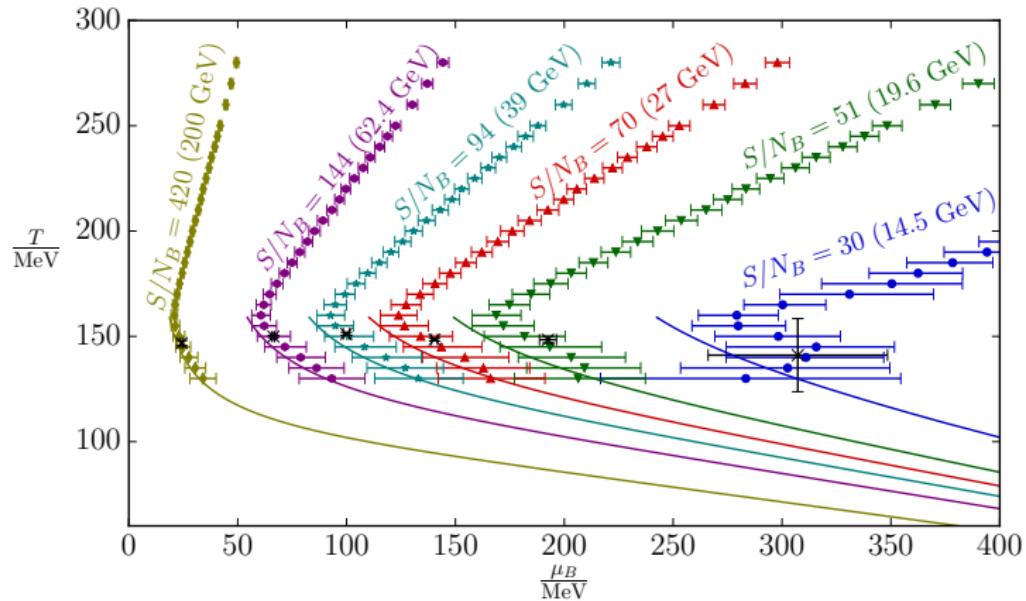


# Influence of different orders

$T = 170\text{MeV}$



# Trajectories



# Equation of state

