

# Open charm correlators and spectral functions at high temperature

Jon-Ivar Skullerud, Aoife Kelly

National University of Ireland Maynooth  
FASTSUM collaboration

Lattice 2016, Southampton, 28 Jul 2016

# Outline

Background

Simulation and analysis

D meson results

Charmonium

Summary and outlook

## Why open charm?



- ▶ Heavy quarks are important probes of medium
- ▶ Long history of  $c\bar{c}$  studies: experiment, pheno, lattice
- ▶ Open charm still in its infancy

## Why D mesons?

### Open and hidden charm

Cannot study  $c\bar{c}$  in isolation from open charm

- ▶ Recombination at freeze-out
- ▶ Increased yield of D mesons relative to  $J/\psi$ ?
- ▶ Double ratio better measure than  $R_{AA}$ ?
- ▶ Thermal modifications of D mesons may be important
- ▶ Charm quark diffusion  $\leftrightarrow$  D meson flow

## Open charm — issues

### Open charm

- ▶ Increased experimental interest in open charm
- ▶ Suggestions of D meson survival in QGP?
- ▶ Modifications of yields of open charm states?
- ▶ Increased  $D_s/D$  ratio (strangeness enhancement)?

### Open charm from the lattice

Very few studies so far:

- ▶ Cumulants [Bazavov et al, Mukherjee et al (2015)]
- ▶ Screening correlators [Bazavov et al (2014)]

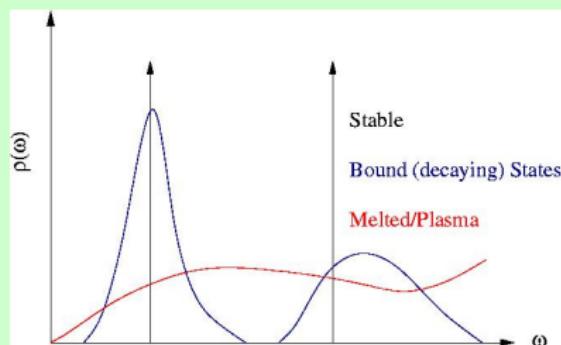
May contribute up to  $1.2 T_c$ ?

# Spectral functions

- ▶ contain information about the fate of hadrons in the medium
  - ▶ **stable states**  $\rho(\omega) \sim \delta(\omega - m)$
  - ▶ **resonances or thermal width**  $\rho(\omega) \sim$  lorentzian
  - ▶ **continuum** above threshold

## Spectral functions

- ▶ contain information about the fate of hadrons in the medium
  - ▶ **stable states**  $\rho(\omega) \sim \delta(\omega - m)$
  - ▶ **resonances or thermal width**  $\rho(\omega) \sim$  lorentzian
  - ▶ **continuum above threshold**



## Spectral functions

- ▶ contain information about the fate of hadrons in the medium
  - ▶ **stable states**  $\rho(\omega) \sim \delta(\omega - m)$
  - ▶ **resonances or thermal width**  $\rho(\omega) \sim$  lorentzian
  - ▶ **continuum** above threshold
- ▶  $\rho_\Gamma(\omega, \vec{p})$  related to **euclidean correlator**  $G_\Gamma(\tau, \vec{p})$  according to

$$G_\Gamma(\tau, \vec{p}) = \int \rho_\Gamma(\omega, \vec{p}) K(\tau, \omega) d\omega, \quad K(\tau, \omega) = \frac{\cosh[\omega(\tau - 1/2T)]}{\sinh(\omega/2T)}$$

- ▶ an **ill-posed problem** — requires a large number of time slices
  - ▶ Fit to physically motivated Ansatz
  - ▶ Use **Maximum Entropy Method** or other Bayesian methods
  - ▶ Other inversion methods, eg Cuniberti, Tikhonov–Morozov

## Reconstructed correlators

The systematic uncertainty of the MEM can be avoided by studying the **reconstructed correlator**, defined as

$$G_r(\tau; T, T_r) = \int_0^\infty \rho(\omega; T_r) K(\tau, \omega, T) d\omega$$

where  $K$  is the kernel

$$K(\tau, \omega, T) = \frac{\cosh[\omega(\tau - 1/2T)]}{\sinh(\omega/2T)}$$

If  $\rho(\omega; T) = \rho(\omega; T_r)$  then  $G_r(\tau; T, T_r) = G(\tau; T)$

Small changes in correlators is compatible with large changes in spectral function [Mocsy&Petreczky (2007)]

## Direct correlator reconstruction

[Ding et al (2012)]

With

$$T = \frac{1}{a_\tau N}, \quad T_r = \frac{1}{a_\tau N_r}, \quad \frac{N_r}{N} = m \in \mathbb{N}$$

and using

$$\frac{\cosh [\omega(\tau - N/2)]}{\sinh(\omega N/2)} = \sum_{n=0}^{m-1} \frac{\cosh [\omega(\tau + nN + mN/2)]}{\sinh(\omega mN/2)}$$

we have

$$G_r(\tau; T, T_r) = \sum_{n=0}^{m-1} G(\tau + nN, T_r)$$

## Simulation parameters

**FASTSUM Gen2 ensemble:**  $N_f = 2 + 1$  anisotropic clover  
 [HadSpec, PRD **79** 034502 (2009); FASTSUM, JHEP **1502** 186 (2015)]

$\xi$	3.5
$a_s$ (fm)	0.123
$a_\tau^{-1}$ (GeV)	5.63
$m_\pi$ (MeV)	380
$m_\pi/m_\rho$	0.45
$N_s$	24
$L_s$ (fm)	2.94

$N_\tau$	$T$ (MeV)	$T/T_c$	$N_{\text{cfg}}$
128	44	0.24	500
40	141	0.76	500
36	156	0.84	500
32	176	0.95	1000
28	201	1.09	1000
24	235	1.27	1000
20	281	1.52	576
16	352	1.90	1000

Charm action params from HadSpec: JHEP **1207** 126 (2012)

## Spectral function reconstruction

Spectral function  $\rho(\omega)$  is expressed in terms of default model  $m(\omega)$

$$\rho(\omega) = m(\omega) \exp\left[\sum_{k=1}^{N_b} b_k u_k(\omega)\right]$$

Singular value decomposition:

$$K(\omega, \tau) \rightarrow K(\omega_i, \tau_j) = K_{ij} = U \Xi V^T$$

Standard MEM (SVD basis):  $u_k$  are column vectors of  $U$ :

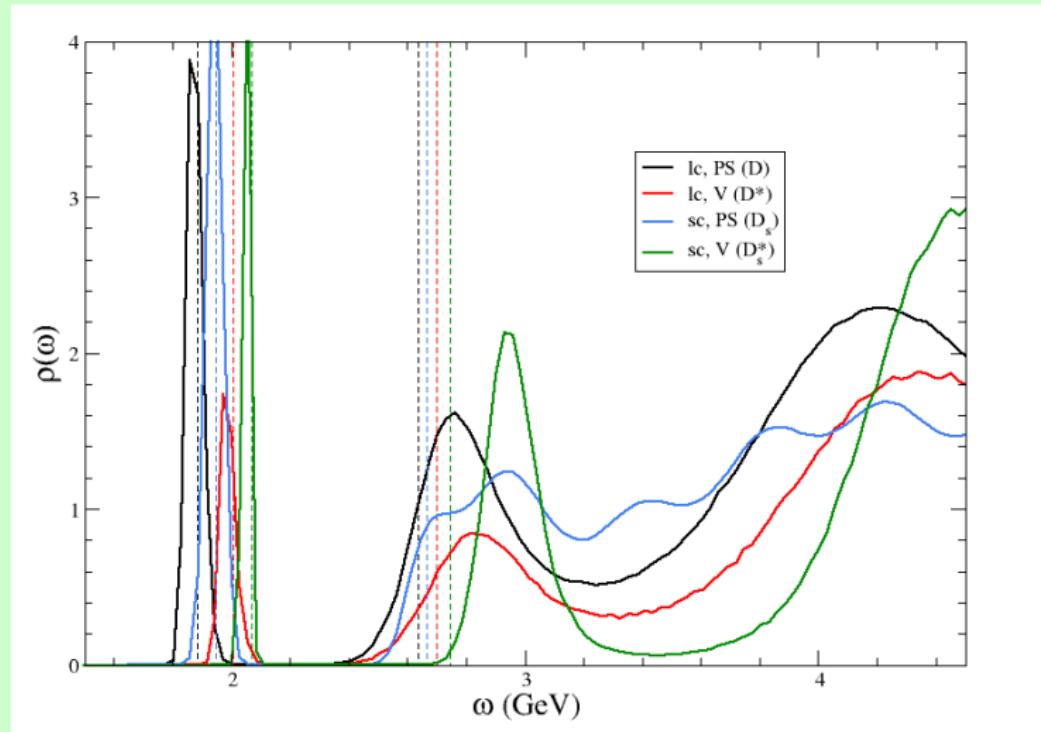
$$N_b = N_s \leq N_{\text{data}}$$

Extended basis: use  $N_{\text{ext}}$  additional column vectors of  $U$

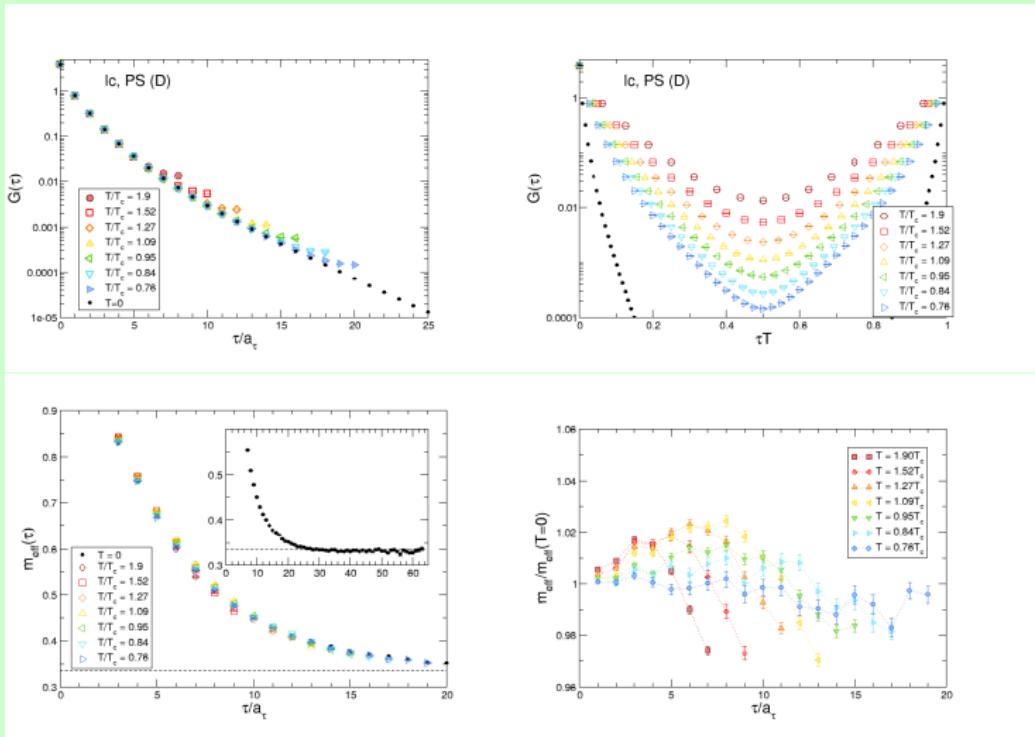
Fourier basis: use  $N_b$  Fourier modes as  $u_k$

Using MEM analysis code from Alexander Rothkopf

## Zero temperature spectral functions

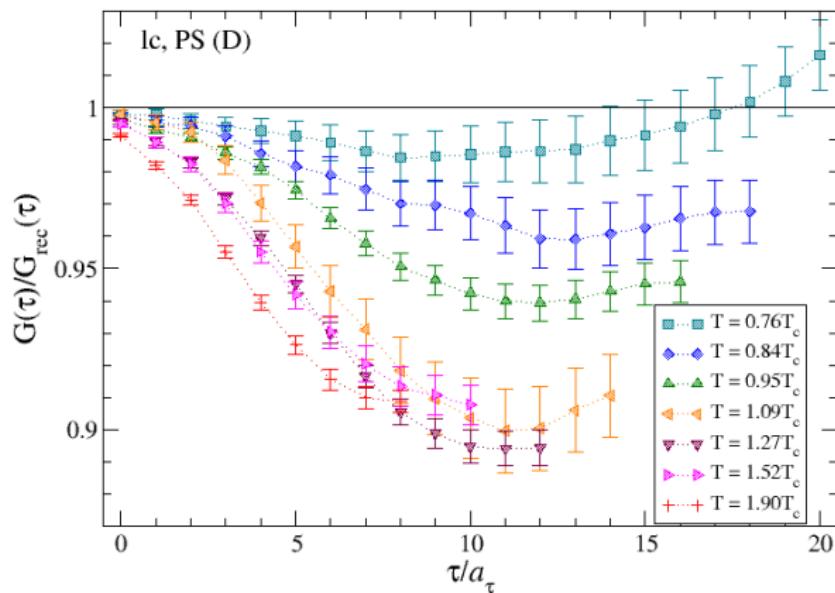


# D meson correlators



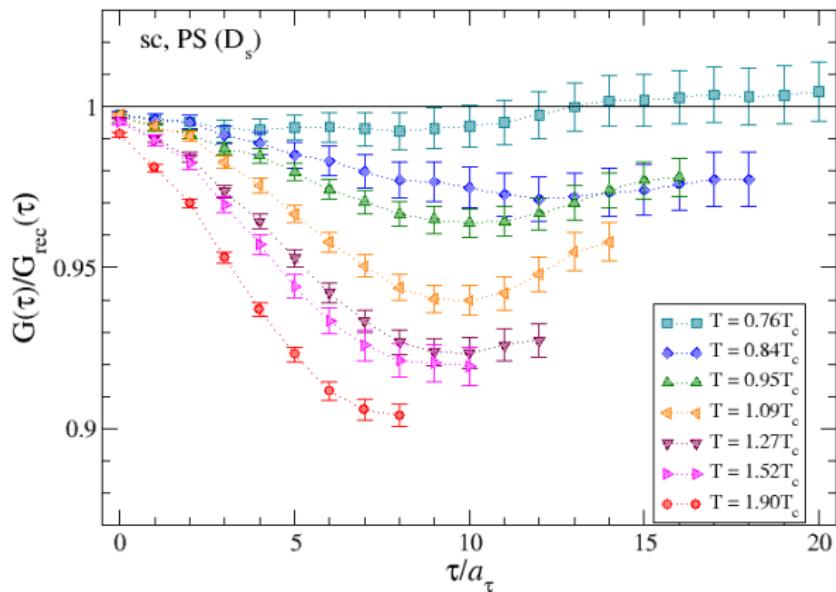
# Reconstructed correlators

*D* channel (lc, PS)



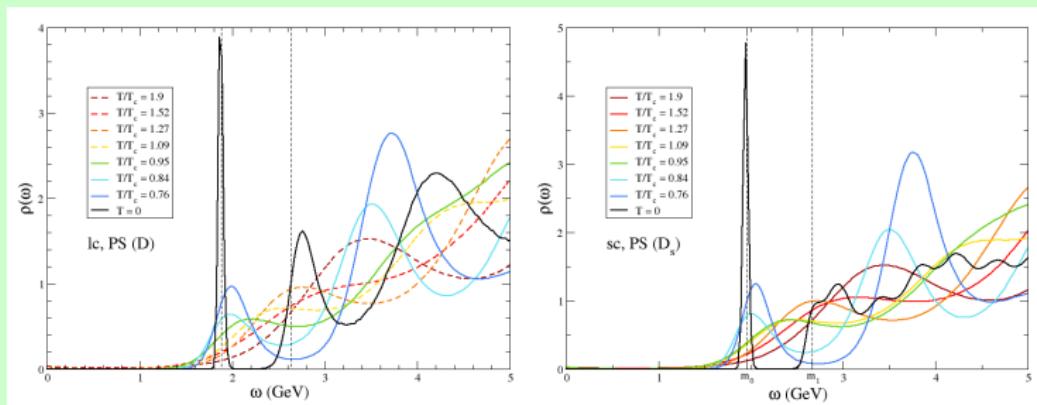
# Reconstructed correlators

$D_s$  channel (sc. PS)



# Open charm: spectral functions

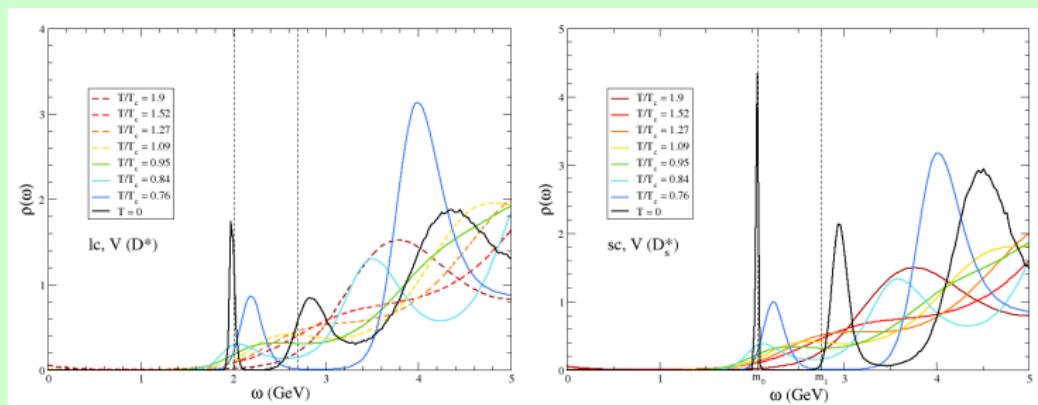
## Pseudoscalar channel



- ▶ Both  $D$  and  $D_s$  mesons dissociate close to  $T_c$
- ▶ Thermal mass shift below  $T_c$ ?

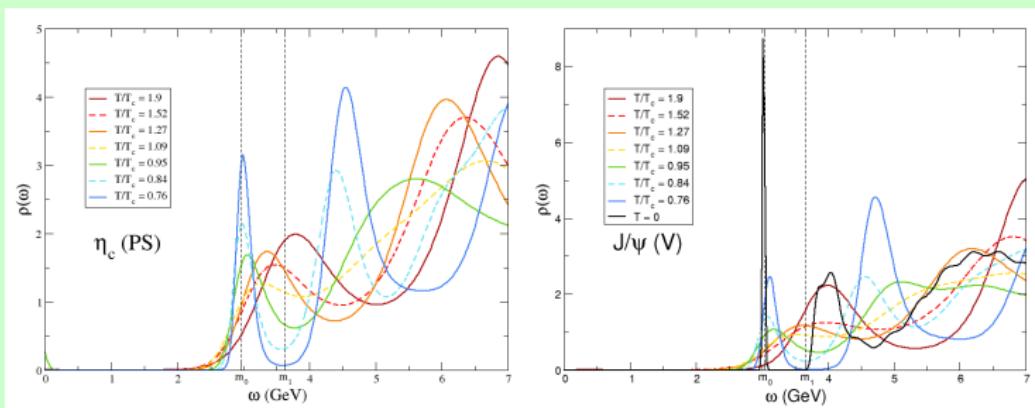
# Open charm: spectral functions

## Vector channel



- ▶ Thermal mass shift stronger in vector channel?

# Charmonium



# Summary and outlook

## Summary

- ▶ First lattice study of open charm temporal correlators and spectral functions
- ▶ Clear thermal modifications already below  $T_c$
- ▶ Possible thermal mass shift observed?
- ▶ No sign of surviving bound states above  $T_c$

# Summary and outlook

## Summary

- ▶ First lattice study of open charm temporal correlators and spectral functions
- ▶ Clear thermal modifications already below  $T_c$
- ▶ Possible thermal mass shift observed?
- ▶ No sign of surviving bound states above  $T_c$

## Outlook

- ▶ Complete study of MEM systematics
- ▶ Improved statistics (using multiple sources)
- ▶ Open beauty?