

# A high-statistics lattice QCD study of nucleon sigma terms

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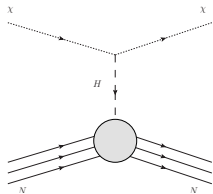
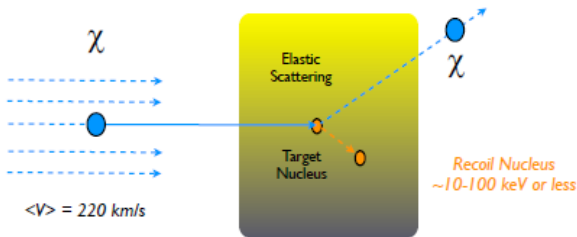
Budapest-Marseille-Wuppertal collaboration (BMWc)

Special thanks to Christians Hoelbling and Torrero and Lukas Varnhorst

(Preliminary results)



# Direct WIMP dark matter detection



$$\mathcal{L}_{q\chi} = \sum_q \lambda_q^{\Gamma} [\bar{q}\Gamma q][\bar{\chi}\Gamma\chi] \rightarrow \mathcal{L}_{N\chi} = \lambda_N^{\Gamma} [\bar{N}\Gamma N][\bar{\chi}\Gamma\chi]$$

Quarks are confined within nucleons  
→ nonperturbative QCD tool

# WIMP-nucleus spin-independent cross section

In low- $E$  limit

$$\frac{d\sigma_{\chi Z^A X}^{\text{SI}}}{dq^2} = \frac{1}{\pi v^2} [Z f_p + (A - Z) f_n]^2 |F_X(q^2)|^2$$

w/  $F_X(\vec{q} = 0) = 1$  nuclear FF and  $\chi N$  couplings ( $N = p, n$ )

$$\frac{f_N}{M_N} = \sum_{q=u,d,s} f_{qN} \frac{\lambda_q}{m_q} + \sum_{Q=c,b,t} f_{QN} \frac{\lambda_Q}{m_Q}$$

such that ( $f = u, \dots, t$  and  $\langle N(\vec{p}') | N(\vec{p}) \rangle = (2\pi)^3 \delta^{(3)}(\vec{p}' - \vec{p})$ )

$$f_{udN} M_N = \sigma_{\pi N} = m_{ud} \langle N | \bar{u}u + \bar{d}d | N \rangle, \quad f_{fN} M_N = \sigma_{fN} = m_f \langle N | \bar{f}f | N \rangle$$

For heavy  $Q = c, b, t$ , can use (Shifman et al '78)

$$f_{QN} \approx \frac{2}{27} \left( 1 - \sum_{q=u,d,s} f_{qN} \right)$$

# Feynman-Hellmann method

- Use Feynman-Hellmann (FH) theorem

$$\langle N | m_q \bar{q} q | N \rangle = m_q \left. \frac{\partial M_N}{\partial m_q} \right|_{m_q^\Phi}$$

- FH vs direct matrix element (ME) method

- ✓ Only simpler and less noisy 2pt-fn is needed
- ✓ No difficult quark-disconnected contributions
- ✓ No difficult singlet renormalization w/ Wilson fermions

✗  $m_{ud}$ -dependence not very large

$$M_N \simeq 939 \text{ MeV} \simeq 850 \text{ MeV} + O(m_{ud})$$

→ must extract smallish correction

✗  $m_s$  dependence even smaller

# Lattice details

- $N_f = 1 + 1 + 1 + 1$
- 3HEX clover-improved Wilson fermions on tree-level improved Symanzik gluons
- 29 ensembles w/ total  $\sim 155000$  trajectories
- $\sim 500$  measurements per configuration
- 4  $a \in [0.064, 0.102]$  fm;
- $M_\pi \in [195, 450]$  MeV w/  $LM_\pi > 4$

## Improvements over BMWc, PRL '16

- ✓ Charm in sea
- ✓  $\gtrsim \times 100$  in statistics (reduced to  $\gtrsim 35$  by later plateaux)
- ✓  $\gtrsim \times 2$  lever arm in  $m_s$
- ✓ Like PRL '16 FH in terms of quark and not meson masses
- ✗ No physical  $m_{ud}$ , but small enough and know  $M_N$

# Analysis strategy

- Dependence of 4 hadron masses analyzed simultaneously, vs  $m_{ud}$ ,  $m_s$  and  $m_{d-u} = m_d - m_u$ :
  - $M_\Omega$ ,  $M_{\pi^+}^2$  and  $M_{K^X}^2 = (M_{K^+}^2 + M_{K^0}^2 - M_{\pi^+}^2)/2$  to fix  $a$ ,  $m_{ud}^\phi$  and  $m_s^\phi$
  - $M_N$  to determine sigma terms (and sometimes scale)
- Quark masses obtained w/ *ratio-difference method* (BMWc, JHEP 1108)
- Since  $Z_S(\beta)$  is not computed, write functional dependence in terms of ( $q = ud, s$ , or  $d - u$ )

$$c_q \left[ \frac{(am_q)}{aZ_s(\beta)} - m_q^\phi \right] \longrightarrow \tilde{c}_q \left[ \frac{(am_q)}{aZ_s(\beta)m_q^\phi} - 1 \right] \longrightarrow \tilde{c}_q \left[ \frac{(am_q)}{a \cdot \tilde{m}_q^\phi(\beta)} - 1 \right]$$

- Cross checks:
  - $\tilde{m}_s^\phi / \tilde{m}_{ud}^\phi = m_s^\phi / m_{ud}^\phi$ ?
  - Values of  $Z_S(\beta)$ ?

# Analysis details (cont'd)

- Minimize  $\chi^2 = V^T C^{-1} V$ , with  $C$  full covariance matrix,

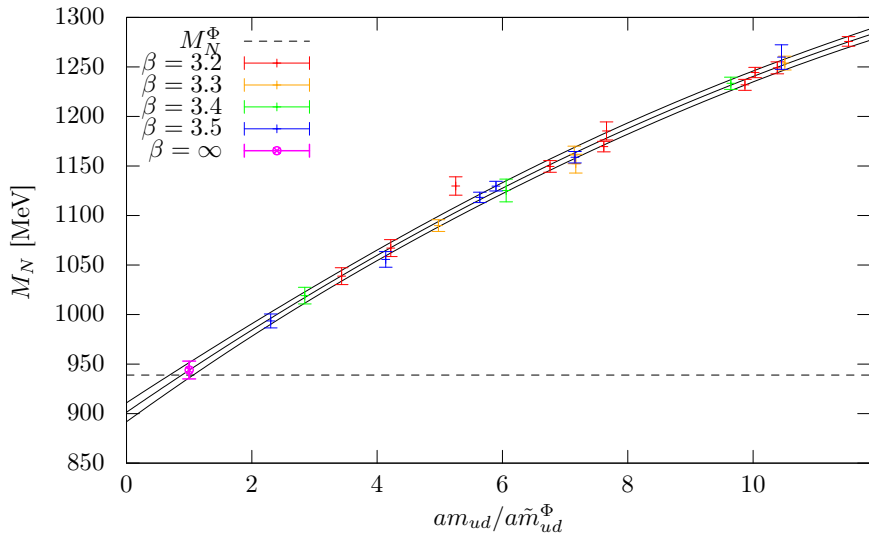
$$V^T = \left[ aM_{[\Omega,1]} - f_{\Omega}(\mathbf{c}, X_{[q,1]}), \dots, aM_{[N,N_{\text{ens}}]} - f_N(\mathbf{c}, X_{[q,N_{\text{ens}}]}), \right. \\ \left. m_{[q,1]} - X_{[q,1]}, \dots, m_{[q,N_{\text{ens}}]} - X_{[q,N_{\text{ens}}]} \right]$$

and e.g.

$$f_N(\mathbf{c}, X) = a(\beta) M_N^{\Phi} [1 + \text{FV}(aM_{\pi}, L/a)] \left\{ 1 + \sum_i^2 c_{ud,i} \left[ \frac{(am_{ud})}{a \cdot \tilde{m}_{ud}^{\Phi}(\beta)} - 1 \right]^i + c_{s,1} \left[ \frac{(am_s)}{a \cdot \tilde{m}_s^{\Phi}(\beta)} - 1 \right] \right\}$$

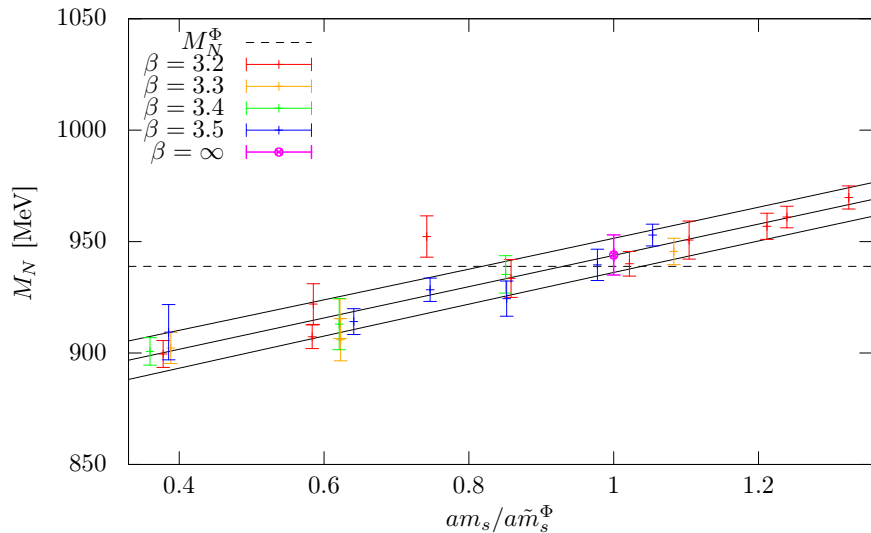
- Statistical error with  $N_{\text{boot}} = 1000$
- To determine fit ranges, use Kolmogorov-Smirnov test with different  $t_{\text{min}}$  (BMWc, Science '15)

# Example fits (preliminary)

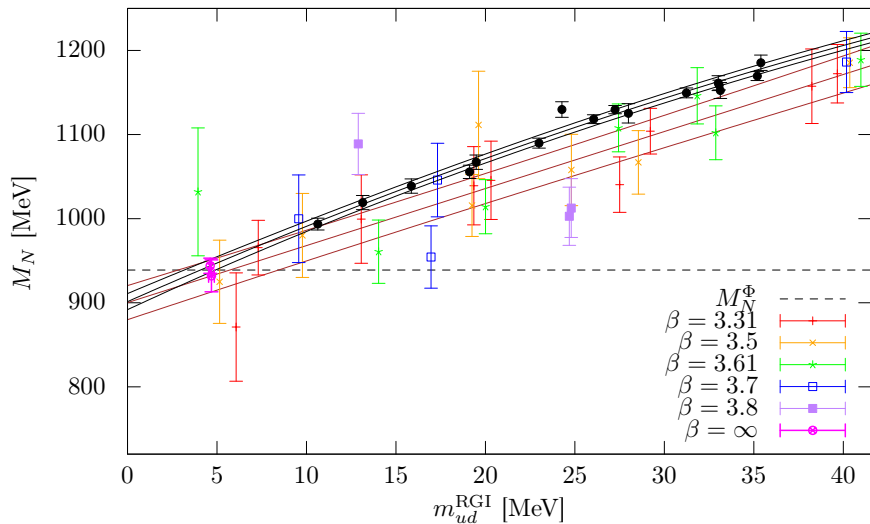




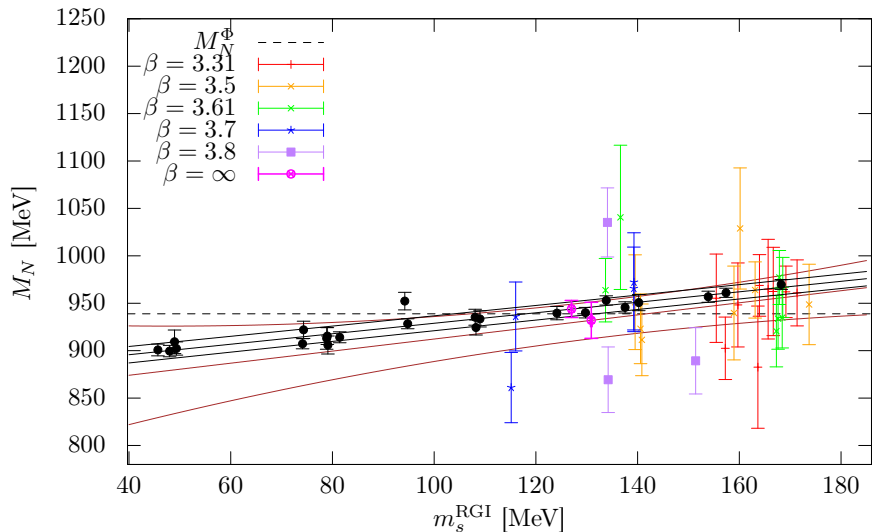
# Example fits (preliminary)



# Comparison with BMWc, PRL '16 (preliminary)



# Comparison with BMWc, PRL '16 (preliminary)



# New method for obtaining $f_{u/d}^{p/n}$

[BMWc, PRL 116 (2016)]

- Input:  $f_{ud}^N$  and  $\Delta_{\text{QCD}} M_N = M_n - M_p$  (from BMWc, Science '15)
- SU(2) relations w/  $\delta m = m_d - m_u$

$$H = H_{\text{ISO}} + H_{\delta m}, \quad H_{\delta m} = \frac{\delta m}{2} \int d^3x (\bar{d}d - \bar{u}u)$$

$$\Delta_{\text{QCD}} M_N = \delta m \langle p | \bar{u}u - \bar{d}d | p \rangle$$

lead to, w/  $r = m_u/m_d$ ,

$$f_u^{p/n} = \left( \frac{r}{1+r} \right) f_{ud}^N \pm \frac{1}{2} \left( \frac{r}{1-r} \right) \frac{\Delta_{\text{QCD}} M_N}{M_N}$$

$$f_d^{p/n} = \left( \frac{1}{1+r} \right) f_{ud}^N \mp \frac{1}{2} \left( \frac{1}{1-r} \right) \frac{\Delta_{\text{QCD}} M_N}{M_N}$$

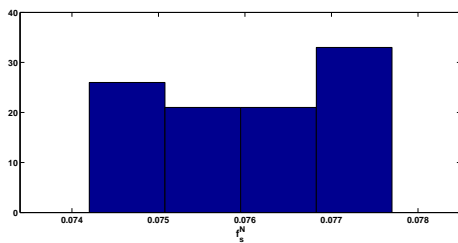
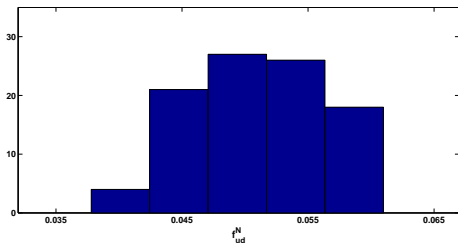
- Huge improvement on usual SU(3)-flavour approach

$$\text{systematic: } \left( \frac{m_s - m_{ud}}{\Lambda_{\text{QCD}}} \right)^2 \approx 10\% \quad \longrightarrow \quad \left( \frac{m_d - m_u}{\Lambda_{\text{QCD}}} \right)^2 \approx 0.01\% .$$

# Systematic error assessment (preliminary)

Estimated using extended frequentist approach (BMWc, Science '08)

- Excited state contamination: 2 time intervals ( $t_{\min} = 1.3$  or  $1.4$  fm)
- Mass interpolation/extrapolation errors
  - $M_{\pi} \leq 400$  MeV &  $450$  MeV
  - different  $m_q$  dependences (polynomials & Padés)
- lattice spacing uncertainty:  $M_{\Omega}$  vs  $M_N$
- continuum extrapolation:  $O(a^2)$  vs  $O(\alpha_S a)$
- AIC weight (BMWc, Science '15)



# Preliminary results

- Checks

- $Z_S(\beta) \sim 0.7$  and expected  $\beta$ -dependence
- Compare

$$M_N = 940(9)(6) \text{ MeV} \qquad m_s/m_{ud} = 26.33(37)(2)$$

w/  $M_N = 939 \text{ MeV}$  (PDG '15) and  $m_s/m_{ud} = 27.53(20)(08)$  (BMWc, JHEP 08 (2011))

- Predictions

$$f_{ud}^N = 0.0517(49)(70) \qquad f_s^N = 0.0760(43)(13)$$

$$f_u^p = 0.0174(16)(23) \qquad f_d^p = 0.0329(34)(48)$$

$$f_u^n = 0.0151(16)(22) \qquad f_d^n = 0.0379(34)(48)$$

# Comparison

