

# Angular and chiral content of the $\rho$ and $\rho'$ mesons

Christian Rohrhofer, Markus Pak, Leonid Y. Glozman

based on 1603.04665 [hep-lat]

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*Physical mesons*



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*Chiral representations*

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$$^3S_1$$



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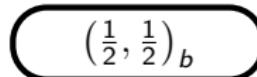
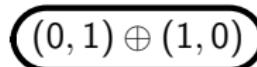
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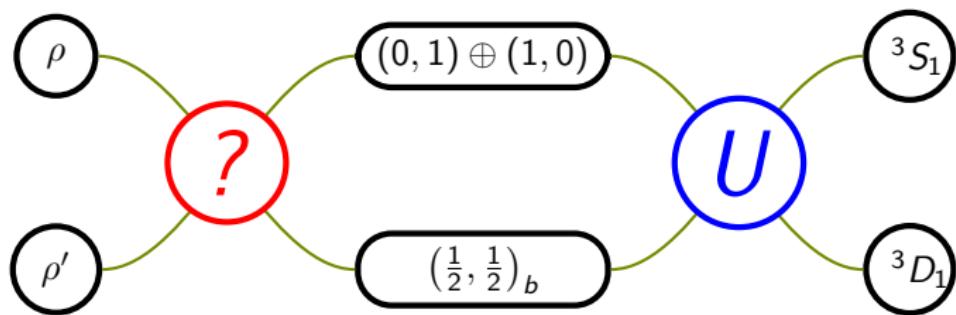


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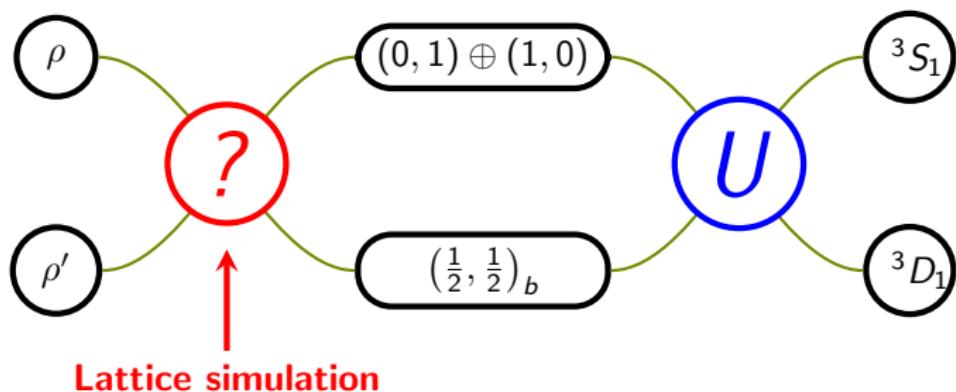


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*Angular momentum*



# The quark model

K.A. Olive et al. (PDG), Chin. Phys. C, 38, 090001 (2014)

$n^{2s+1}\ell_J$	$J^{PC}$	type	$ l=1$	$ l=\frac{1}{2} $	$ l=0 $
$1^1 S_0$	$0^{--}$	pseudoscalar	$\pi$	$K$	$\eta, \eta'$
$1^3 S_1$	$1^{--}$	vector	$\rho(770)$	$K^*$	$\omega, \phi$
$1^1 P_1$	$1^{+-}$	pseudovector	$b_1(1235)$	$K_{1B}$	$h_1$
$1^3 P_0$	$0^{++}$	scalar	$a_0(1450)$	$K_0^*$	$f_0$
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$1^3 P_2$	$2^{++}$	tensor	$a_2(1320)$	$K_2^*$	$f_2$
⋮					
$2^3 S_1$	$1^{--}$	vector	$\rho(1450)$		
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$1^3 D_1$	$1^{--}$	vector	$\rho(1700)$		

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"The physical vector mesons listed under  $1^3D_1$  and  $2^3S_1$  may be mixtures of  $1^3D_1$  and  $2^3S_1$ , or even have hybrid components."

# Angular momentum content of the $\rho$ meson



$$|11^{--}\rangle = \textcolor{blue}{x} |(0,1) + (1,0)\rangle + \textcolor{red}{y} |(\frac{1}{2}, \frac{1}{2})_b\rangle$$

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$$|11^{--}\rangle = \left( \sqrt{\frac{2}{3}}\textcolor{blue}{x} + \sqrt{\frac{1}{3}}\textcolor{cyan}{y} \right) |1, {}^3S_1\rangle + \left( \sqrt{\frac{1}{3}}\textcolor{blue}{x} - \sqrt{\frac{2}{3}}\textcolor{cyan}{y} \right) |1, {}^3D_1\rangle$$

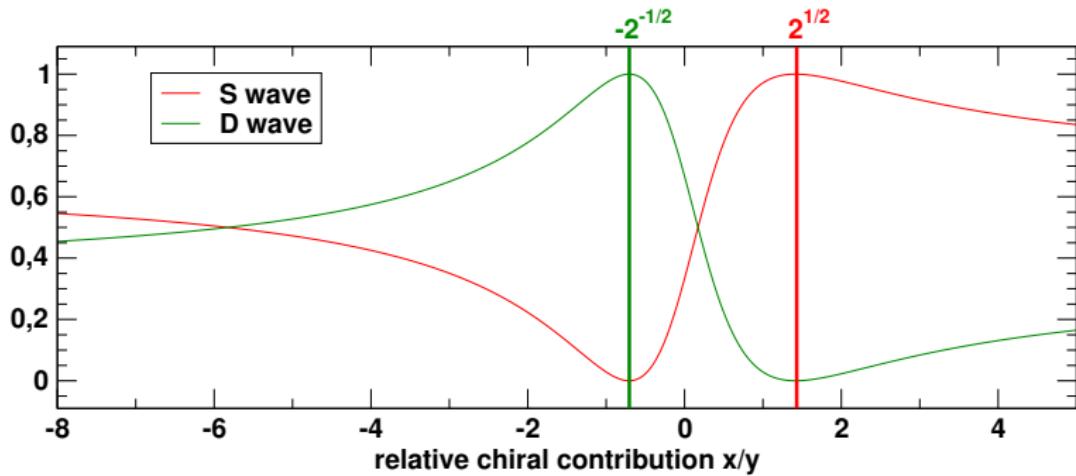
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# Gauge configurations and lattice details

## JLQCD $n_f = 2$ configurations

- Ensemble: 100 configurations
- Lattice size  $16^3 \times 32$
- Spacing  $a = 0.1184(3) fm$
- Topological charge  $Q = 0$
- Dynamical overlap fermions
- Pion mass  $m_\pi = (289 \pm 1.8) MeV$
- Lowest Energy for zero-momentum  $\pi\pi$ -state:  
 $\approx 1400 MeV \longrightarrow \rho$  is stable particle

## Interpolators for a 8x8 correlation matrix

	chiral representation
	$(0, 1) \oplus (1, 0)$ "vector"
	$(\frac{1}{2}, \frac{1}{2})_b$ "tensor"
	$\bar{q}(\vec{\tau} \otimes \gamma^k)q$ $\bar{q}(\vec{\tau} \otimes \gamma^0 \gamma^k)q$

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super narrow	$O_{sn}^V$	$O_{sn}^T$
narrow	$O_n^V$	$O_n^T$
wide	$O_w^V$	$O_w^T$
ultra wide	$O_{uw}^V$	$O_{uw}^T$

$\longleftrightarrow$   
probes chiral contribution

$\longleftrightarrow$   
probes different resolutions

## Variational analysis & interpolator contribution



Construct a cross correlation matrix

$$C_{ij}(t) = \langle O_i(t) O_j^\dagger(0) \rangle = \sum_n \langle 0 | O_i | n \rangle \langle n | O_j^\dagger | 0 \rangle e^{-E_n t}$$

where the overlap of interpolator  $O_i$  with physical state  $|n\rangle$  is given by

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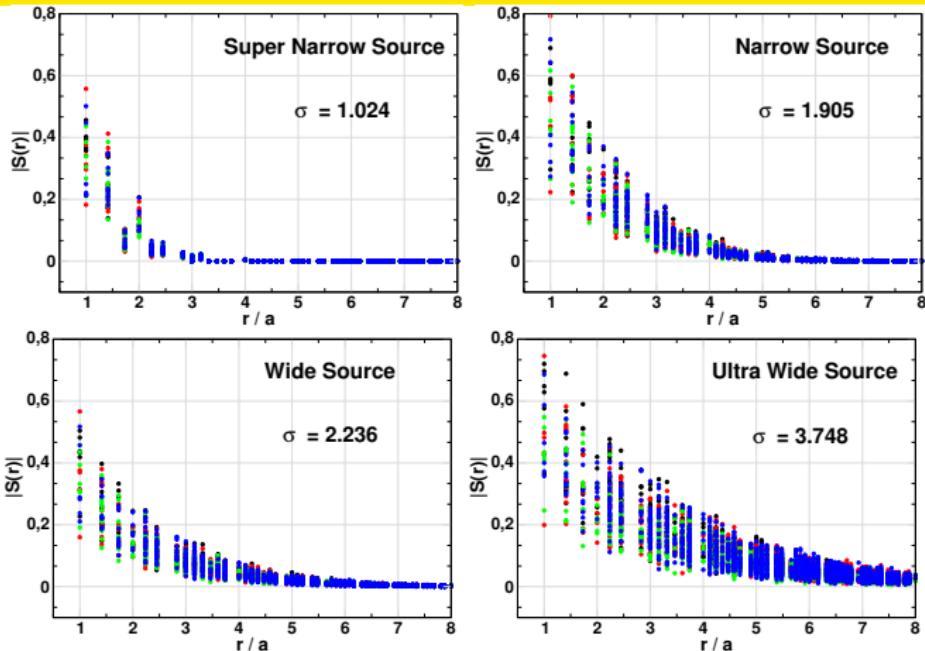
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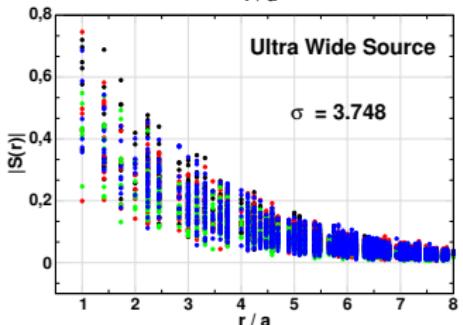
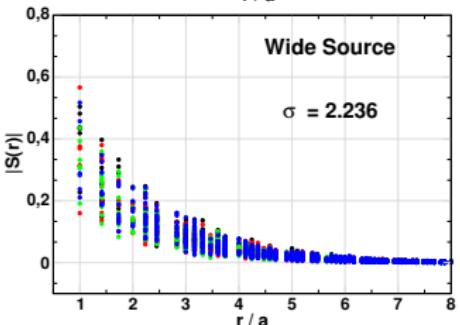
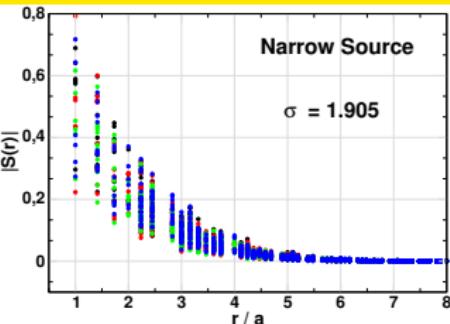
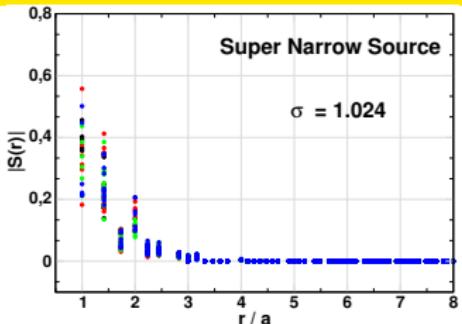
→ use ratio of overlaps as relative chiral contribution!

## Source smearing and resolution scale

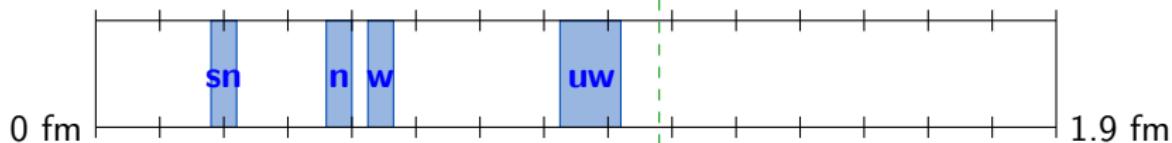
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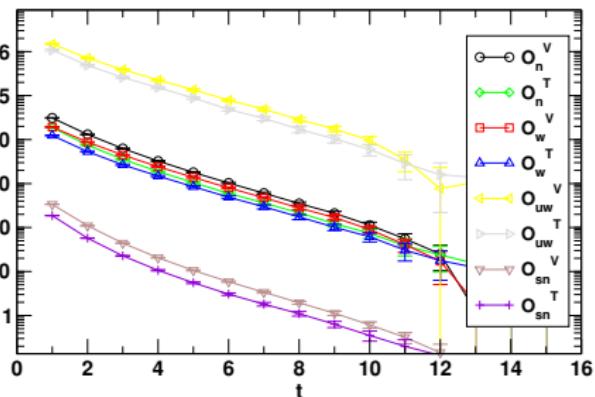
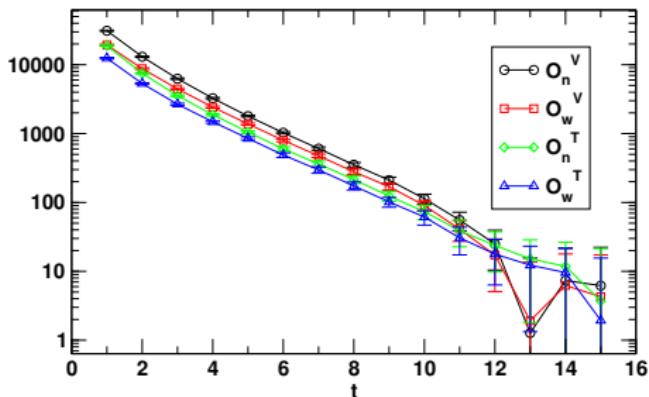
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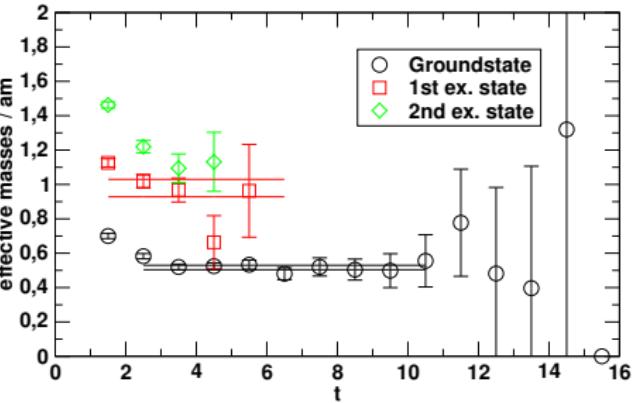
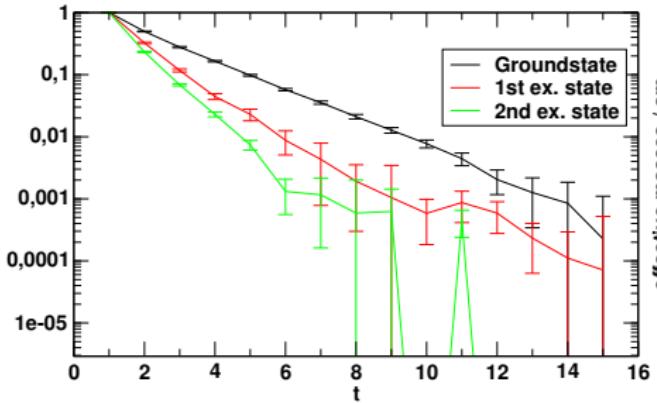
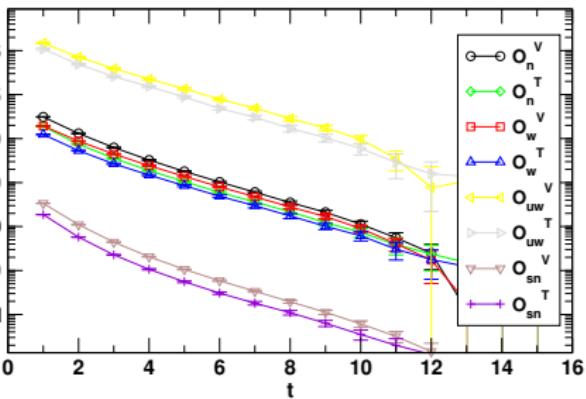
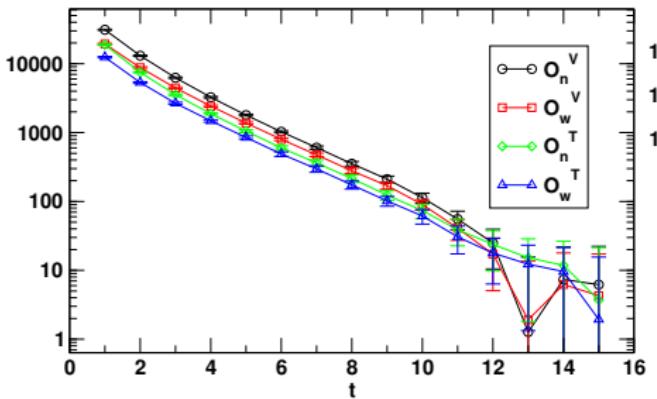
1fm



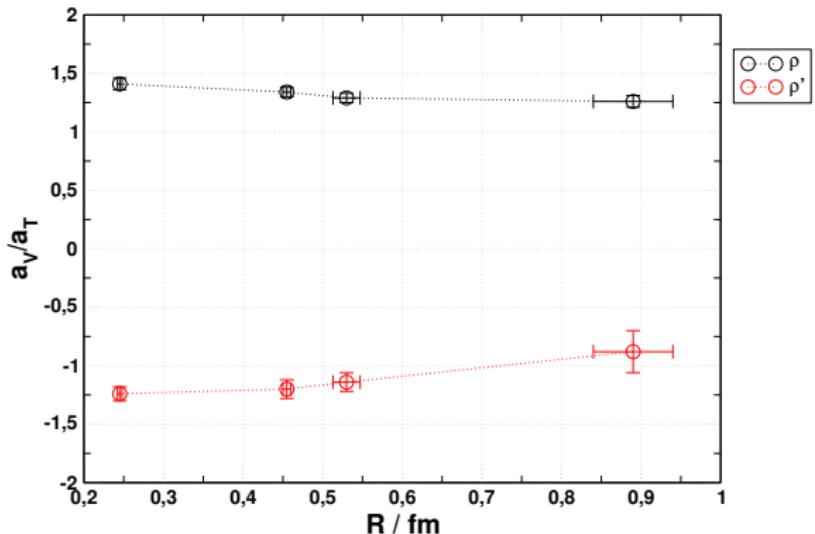
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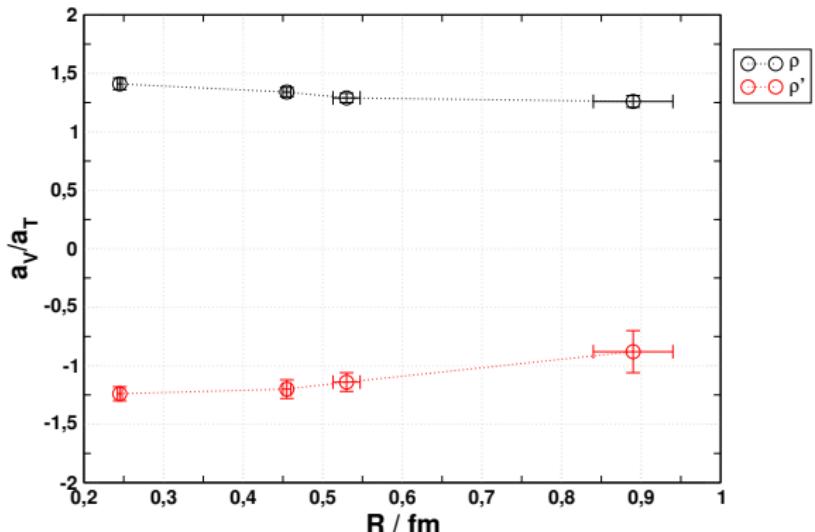
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# Angular and chiral content



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$$|\rho\rangle = (0.998 \pm 0.002) |^3S_1\rangle - (0.05 \pm 0.025) |^3D_1\rangle$$

$$|\rho'\rangle = -(0.106 \pm 0.09) |^3S_1\rangle - (0.994 \pm 0.005) |^3D_1\rangle$$

$$|\rho''\rangle = (0.99 \pm 0.1) |^3S_1\rangle - (0.01 \pm 0.12) |^3D_1\rangle$$

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- In nature chiral symmetry is *spontaneously* broken
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- Remove lowest modes to *unbreak* chiral symmetry:

$$D_{restored}^{-1} = D_{full}^{-1} - \sum_{i=1}^k \frac{1}{\lambda_i} |v_i\rangle \langle v_i|$$

- $k$  denotes the number of removed eigenmodes
- previous work e.g. (*PhysRevD*)
  - arXiv:1107.5195
  - arXiv:1205.4887
  - arXiv:1410.8751 (isoscalar mesons)
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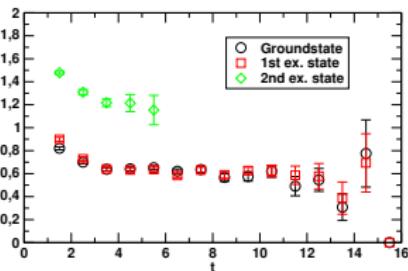
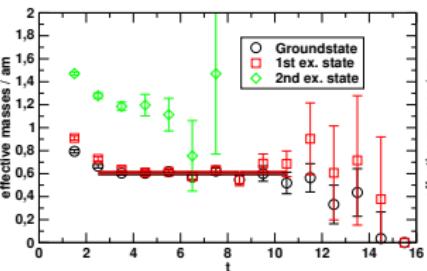
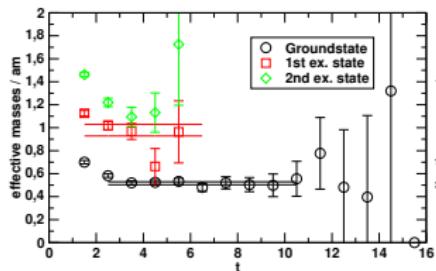
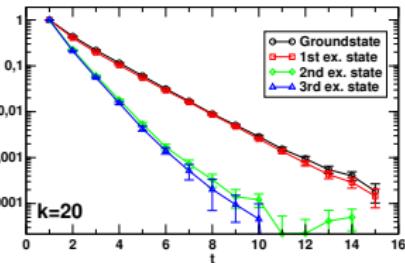
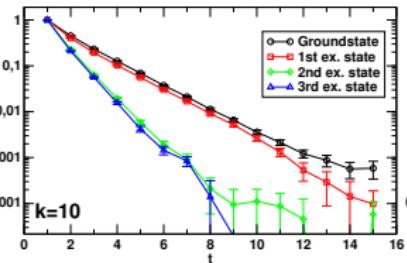
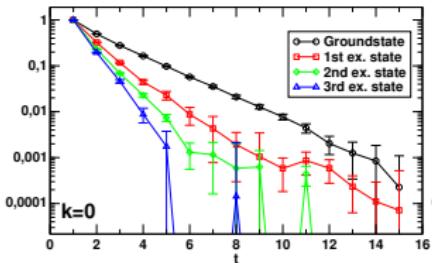
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→ *chiral symmetry indeed gets restored in the hadron spectrum*

# Effective masses after unbreaking $\chi$ symmetry



# Chiral structure



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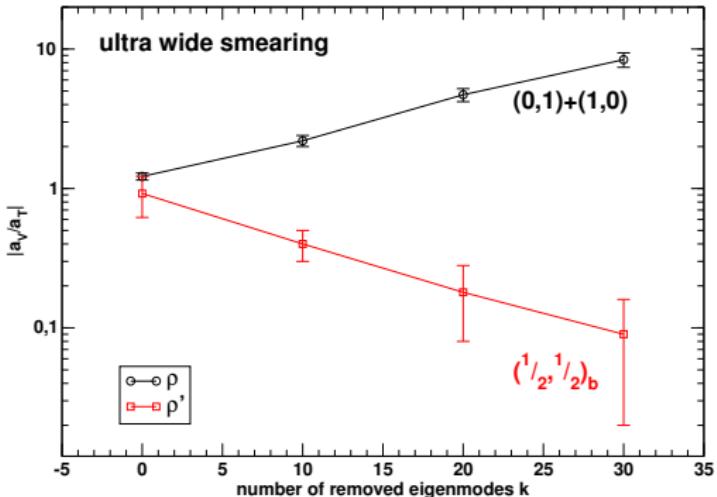
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In a scenario, where chiral symmetry is restored in the hadron spectrum:

- $\rho, \rho''$  live in  $(0, 1) \oplus (1, 0)$  representation
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