Numerical simulation of Dirac semimetals

V. V. Braguta, M. I. Katsnelson, <u>A. Yu. Kotov</u>, A. A. Nikolaev, M. A. Zubkov

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• Condensed matter analog of QGP

- strongly coupled
- massless fermions
- 3D analog of graphene (2D)

Dirac semimetals. Elementary structure





Na₃Bi arXiv:1310.0391, Z. K. Liu et al.

 $Cd_3As_2 \\ arXiv:1309.7892, \ M. \ Neupane \ et \ al. \\ arXiv:1309.7978, \ S. \ Borisenko \ al. \ al.$

Band structure



Take into account all energy levels?

Effective theory approach

- Relevant energy scales
- Relevant degrees of freedom



Dirac points



arXiv:1310.0391, Z. K. Liu et al. N = 2 Dirac points in electronic dispersion

$$E\sim \sqrt{v_\perp^2(k_x^2+k_y^2)+v_\parallel^2k_z^2}$$

- N = 2 Massless Dirac fermions
- Fermi velocity is $v_f \ll c$
- ${\scriptstyle \bullet}$ Anisotropy in Fermi velocity $\textit{\textbf{v}}_{\parallel} \neq \textit{\textbf{v}}_{\perp}$
- Electromagnetic interaction
- Effective charge is $\alpha_F = rac{lpha}{\epsilon_{V_F}} \sim 1$, $lpha = rac{1}{137}$

$$Z = \int D\bar{\psi}D\psi DA \exp\left(-\frac{1}{8\pi\alpha_F}\int d^4x [\partial_j A_4]^2 + i\int d^4x \bar{\psi}\left(\gamma^4 [\partial_4 + iA^4] + \xi_k \gamma^k \partial_k\right)\psi\right)$$

Semimetal - Insulator transition

At large enough α gapped phase?



Is there a gap? What is the value of α_c ?

Lattice discretization

Staggered fermions ($N = 4 \rightarrow N = 2$ via rooting):

$$S_{\Psi int.} = \sum_{x} m \bar{\Psi}_{x} \Psi_{x} + \frac{1}{2} \sum_{i=1,2,3} \xi_{i} (\bar{\Psi}_{x} \alpha_{x,i} \Psi_{x+\hat{1}} - \bar{\Psi}_{x} \alpha_{x,i} \Psi_{x-\hat{1}}) + \frac{1}{2} \bar{\Psi}_{x} \alpha_{x,4} \exp(i\theta_{I,4}(x)) \Psi_{x+\hat{4}} - \frac{1}{2} \bar{\Psi}_{x} \alpha_{x,4} \exp(-i\theta_{I,4}(x-\hat{4})) \Psi_{x-\hat{4}}$$

Noncompact gauge action:

$$S_g = \frac{\beta}{2} \sum_{i=1,2,3} \theta_{p,\hat{1}4}^2(x), \beta = \frac{1}{4\pi\alpha_F}$$
$$\theta_{p,\hat{1}4} = \theta_{l,4}(x+\hat{1}) - \theta_{l,4}(x)$$

- Hybrid Monte-Carlo
- Lattice size 20⁴
- N = 2 dynamical staggered fermions
- Chiral limit $m \to 0$

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$$\xi_1 = \xi_2 = 1$$
, different ξ_3

$$S_{\Psi int.} = \sum_{x} m \bar{\Psi}_{x} \Psi_{x} + \frac{1}{2} \sum_{i=1,2,3} \xi_{i} (\bar{\Psi}_{x} \alpha_{x,i} \Psi_{x+\hat{1}} - \bar{\Psi}_{x} \alpha_{x,i} \Psi_{x-\hat{1}}) + \frac{1}{2} \bar{\Psi}_{x} \alpha_{x,4} \exp(i\theta_{l,4}(x)) \Psi_{x+\hat{4}} - \frac{1}{2} \bar{\Psi}_{x} \alpha_{x,4} \exp(-i\theta_{l,4}(x-\hat{4})) \Psi_{x-\hat{4}}$$

• Order parameter for semimetal-insulator phase transition (chiral condensate):

$$\sigma = \langle \bar{\psi}\psi\rangle = \frac{1}{V}\langle \operatorname{tr} D^{-1}\rangle$$

• Susceptibility:

$$\chi = \frac{\partial \sigma}{\partial m} = \frac{1}{V} (\langle \operatorname{tr}^2 D^{-1} \rangle - \langle \operatorname{tr} D^{-2} \rangle - \langle \operatorname{tr} D^{-1} \rangle^2)$$

• Logarithmic derivative:

$$R = \frac{\partial \log \sigma}{\partial \log m} = \frac{m}{\sigma} \chi$$

χS	Broken	Restored
σ	\sim const	$\sim m ightarrow 0$
χ	Peak at the transition	
R	0	1



$\xi = 1$. Fisher plot

0.002 β=0.04 $\beta = 0.04$ $\beta = 0$ β=0 0.0015 $\beta = 0.0$ 92 0.001 0.0005 0 0.05 0.1 0.15 0.25 0.3 0.35 0.2 0.4 0 m/σ

Equation of State (motivated by QED): $m(X_0 + X_1(1 - \beta/\beta_c)) = \sigma^3 + Y_1(1 - \beta/\beta_c)\sigma.$

$\xi = 1$. Fisher plot

0.002 β=0.04 $\beta = 0.04$ $\beta =$ $\beta = 0$ β=0 0.0015 $\beta =$ 92 0.001 0.0005 0 0.05 0.1 0.15 0.2 0.25 0.3 0.35 0.4 0 m/σ

 $\begin{array}{l} \beta_c\sim0.045\rightarrow\alpha_c\sim1.76\\ \alpha_c\approx1.8660, \mbox{ J.Gonzalez, arXiv:1509.00210, Ladder approximation} \end{array}$

$\xi = 1$. Susceptibilities



$\xi = 1$. Logarithmic derivative



$\xi = 1$



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$\xi = 0.5$



$\xi = 0.2$



$\xi = 0.1$





Conclusions

- First results of study of Dirac semimetals by means of LFT
- Semimetal-insulator transition in Dirac semimetals
- α_c decreases if we increase asymmetry
- Dimensional reduction
- Systematical uncertanties ?



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Within EFT one can study:

- Phase diagram
- Magnetic field
- Temperature
- Impurities, deformations
- (Anomalous) transport phenomena and properties
- Chiral effects



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