

# Nucleon electromagnetic and axial form factors with $N_f=2$ twisted mass fermions at the physical point

<sup>a,b</sup>C. Alexandrou, <sup>b</sup>M. Constantinou, <sup>c</sup>K. Hadjiyiannakou, <sup>d</sup>K. Jansen,  
<sup>b</sup>C. Kallidonis, <sup>b</sup>G. Koutsou, <sup>e</sup>K. Ottnad, <sup>f</sup>A. Vaquero

<sup>a</sup>Department of Physics, University of Cyprus

<sup>b</sup>Computation-based Science and Technology Research Centre (CaSToRC),  
The Cyprus Institute

<sup>c</sup>Department of Physics, George Washington University, Washington DC

<sup>d</sup>DESY-Zeuthen

<sup>e</sup>Institut für Strahlen- und Kernphysik, Universität Bonn

<sup>f</sup>INFN Sezione Milano Bicocca



# Outline

## ★ Introduction and motivation

- ▶ Proton radius
- ▶ Axial matrix elements

## ★ Setup and method

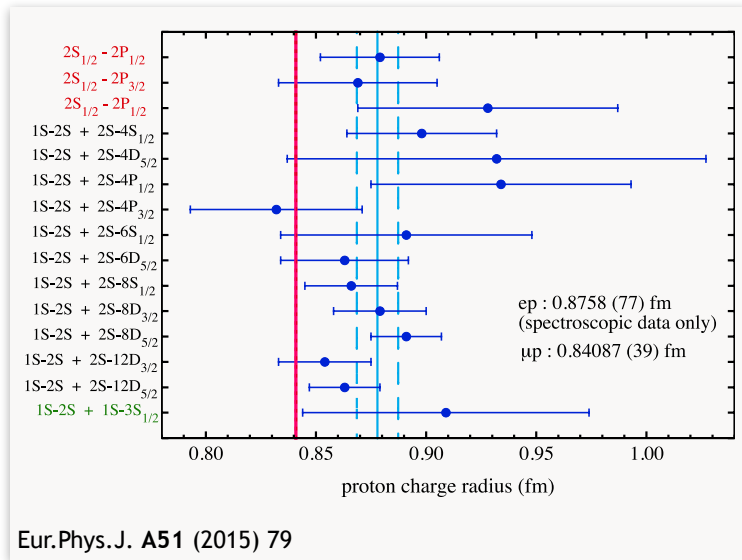
- ▶ Form factors from the lattice
- ▶ Lattice setup

## ★ Results

- ▶ Axial form factors
- ▶ Electromagnetic form factors

## ★ Summary and outlook

# Introduction and motivation



## ★ Axial and Electromagnetic Form Factors

### Insight on structure of nucleon

- ▶ Slope at  $Q^2 \rightarrow 0$ : Electric and magnetic radii (or Dirac and Pauli)
- ▶  $G_A(Q^2 = 0) = g_A$  Nucleon axial charge
- ▶  $G_M(Q^2 = 0) = \mu_N$  Nucleon magnetic moment
- ▶ Input to determination of  $M_A$
- ▶  $G_p(Q^2)$  test pion pole dominance expectation

## ★ Proton spin puzzle

- ▶ Discrepancy between electron scattering and muonic hydrogen Lamb shifts
- ▶ Need  $\sim 2\%$  accuracy on  $\langle r_p^2 \rangle$  to contact experiment
  - Large separations for suppressing excited state effects major challenge
  - Disconnected contributions to obtain up- and down-quark contributions – or equivalently, proton and neutron contributions

# Form factor decomposition

## ★ Axial form factors

$$\langle N(p', s') | A_\mu^3 | N(p, s) \rangle = i \sqrt{\frac{m_N^2}{E_N(\vec{p}') E_N(\vec{p})}} \bar{u}(p', s') [\gamma_\mu \gamma_5 G_A(q^2) + \frac{q_\mu \gamma_5}{2m_N} G_P(q^2)] \frac{1}{2} u_N(p, s)$$

with:

$$A_\mu^3(x) = \bar{\psi}(x) \gamma_\mu \gamma_5 \frac{\tau^3}{2} \psi(x), \text{ and } \psi = \begin{pmatrix} u \\ d \end{pmatrix}$$

## ★ Electromagnetic

$$\langle N(p', s') | j_\mu | N(p, s) \rangle = \sqrt{\frac{m_N^2}{E_N(\vec{p}') E_N(\vec{p})}} \bar{u}(p', s') [\gamma_\mu F_1(q^2) + \frac{i \sigma_{\mu\nu} q^\nu}{2m_N} F_2(q^2)] u_N(p, s)$$

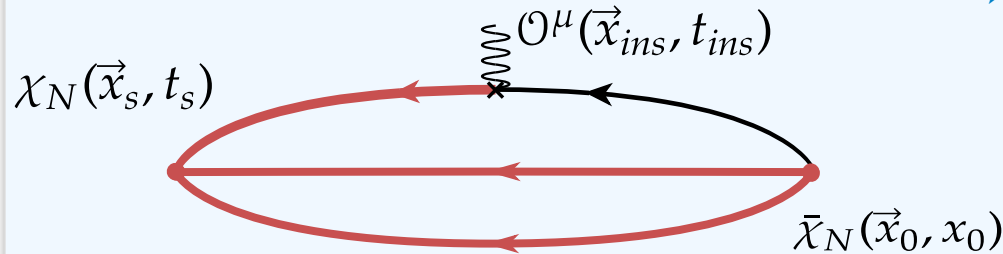
$j_\mu$ : lattice conserved current

$F_1$  and  $F_2$ , Dirac and Pauli form factors, alternatively, define Sachs electromagnetic form factors:

$$G_E(q^2) = F_1(q^2) + \frac{q^2}{(2m_N)^2} F_2(q^2)$$

$$G_M(q^2) = F_1(q^2) + F_2(q^2)$$

# Form factors from the lattice



## ★ Three-point correlation function

- ▶ Inversion through sink
- ▶ Fixed sink coordinates and momentum
- ▶ Free insertion indices (momentum, coordinate, operator)

$$G^\mu(\Gamma; \vec{q}; t_s, t_{ins}) = \sum_{\mathbf{x}_s \vec{x}_{ins}} e^{-i\vec{p}' \cdot \vec{x}_s} e^{-i(\vec{p}' - \vec{p}) \cdot \vec{x}_{ins}} \Gamma^{\alpha\beta} \langle \bar{\chi}_N^\beta(\vec{x}_s; t_s) | \mathcal{O}^\mu(\vec{x}_{ins}; t_{ins}) | \chi_N^\alpha(\vec{x}_0; t_0) \rangle$$

## ★ Ratio with two-point functions:

$$R^\mu(\Gamma; \vec{q}; t_s; t_{ins}) = \frac{G^\mu(\Gamma; \vec{q}; t_s; t_{ins})}{G(\vec{0}; t_s)} \sqrt{\frac{G(\vec{0}; t_s) G(\vec{q}; t_s - t_{ins}) G(\vec{0}; t_{ins})}{G(\vec{q}; t_s) G(\vec{0}; t_s - t_{ins}) G(\vec{q}; t_{ins})}}$$

## ★ Plateau method:

$$R^\mu \xrightarrow[t_s - t_{ins} \gg]{} \Pi^\mu(\Gamma; \vec{q})$$

Fit to constant when  $t_s - t_{ins} \gg$   
 $t_{ins} \gg$

## ★ Summation method:

$$\sum_{t_{ins}} R^\mu \xrightarrow[t_{ins} \gg]{} C + \Pi^\mu(\Gamma; \vec{q}) t_s$$

Slope from linear fit

# Form factor extraction

## ★ Four projectors:

- ▶ One unpolarised:  $\Gamma_0 = \frac{1 + \gamma_0}{4}$
- ▶ Three polarised:  $\Gamma_k = i\gamma_5\gamma_k\Gamma_0$

## ★ Final state at rest: $\vec{p}' = 0, \quad \vec{q} = -\vec{p}$

## ★ Electromagnetic form factors

$$\Pi^0(\Gamma_0; \vec{q}) = C \frac{E_N + m_N}{2m_N} G_E(Q^2)$$

$$\Pi^i(\Gamma_0; \vec{q}) = C \frac{q_i}{2m_N} G_E(Q^2)$$

$$\Pi^i(\Gamma_k; \vec{q}) = C \frac{\epsilon_{ijk} q_j}{2m_N} G_M(Q^2)$$

$$C = \sqrt{\frac{2m_N^2}{E_N(E_N + m_N)}}$$

## ★ Axial form factors

$$\Pi^i(\Gamma_k; \vec{q}) = \frac{iC}{4m_N} \left[ \frac{q_k q_i}{2m_N} G_p(Q^2) - (E_N + m_N) \delta_{ik} G_A(Q^2) \right]$$

$$\sum_n \frac{(\sum_{m=1}^2 D_{nm} G_m - \Pi_n)^2}{w_n^2}$$

- ▶ Minimise, via the singular value decomposition of  $D$ , where  $G$  is the vector of wanted form factors

# Lattice parameters

## ★ Lattice ensemble

- ▶  $N_f=2$  twisted mass with clover term
- ▶  $48^3 96$  lattice sites
- ▶  $a = 0.093(1)$  fm determined from nucleon mass
- ▶ Appx. 3000 independent configurations
- ▶  $m_\pi=0.132(1)$  GeV

## ★ Three-point functions

- ▶  $t_s = 10a, 12a, 14a \approx 0.9$  fm, 1.1 fm, 1.3 fm
  - 580 configs.  $\times$  16 randomised source locations
  - Four polarisations:  $\Gamma_0, \Gamma_1, \Gamma_2, \Gamma_3$

$G_E, G_M, G_A, G_p$

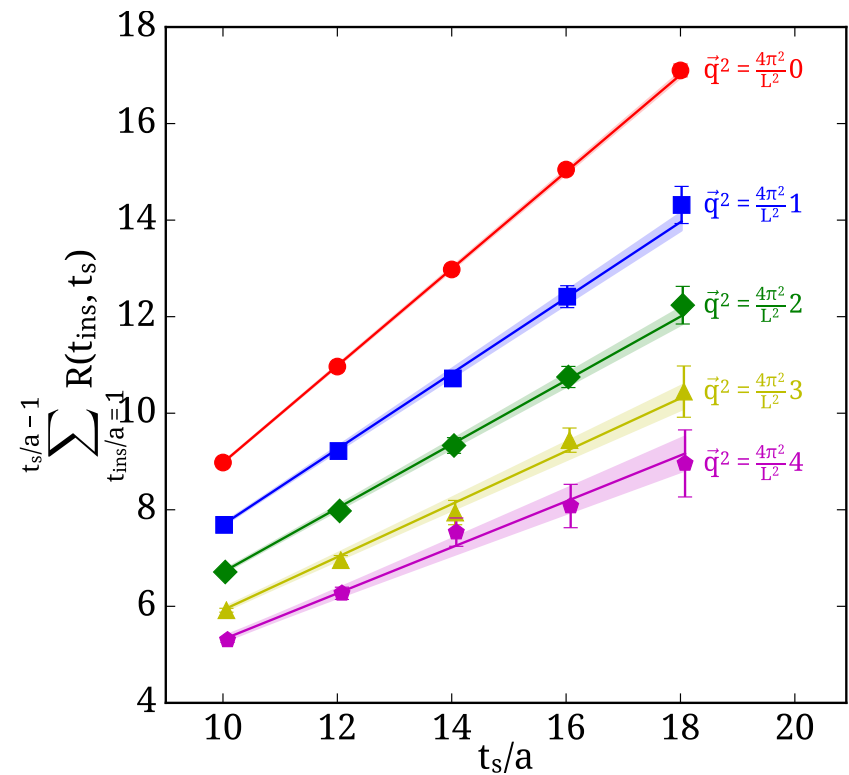
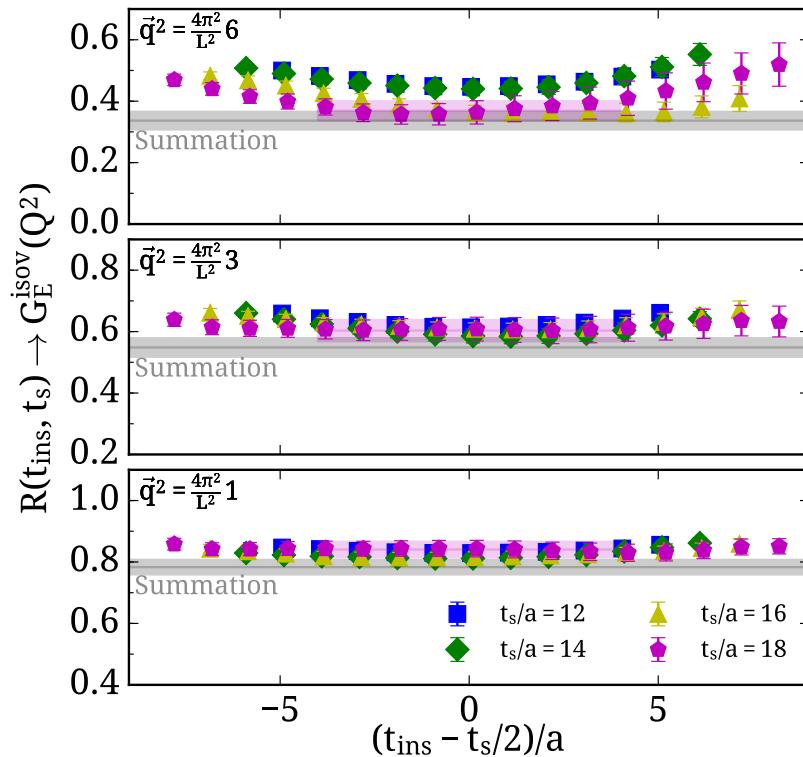
- ▶  $t_s = 16a, 18a \approx 1.5$  fm, 1.7 fm
  - 530 and 725 configs. resp.  $\times$  88 randomised source locations
  - One polarisation:  $\Gamma_0$

$G_E$

# Plateaus and summation fits

## ★ Plateaus and summation method fits example

- ▶ Contact points omitted in summation
- ▶ Different statistics for different sink – source separations → fits in bootstrap throughout

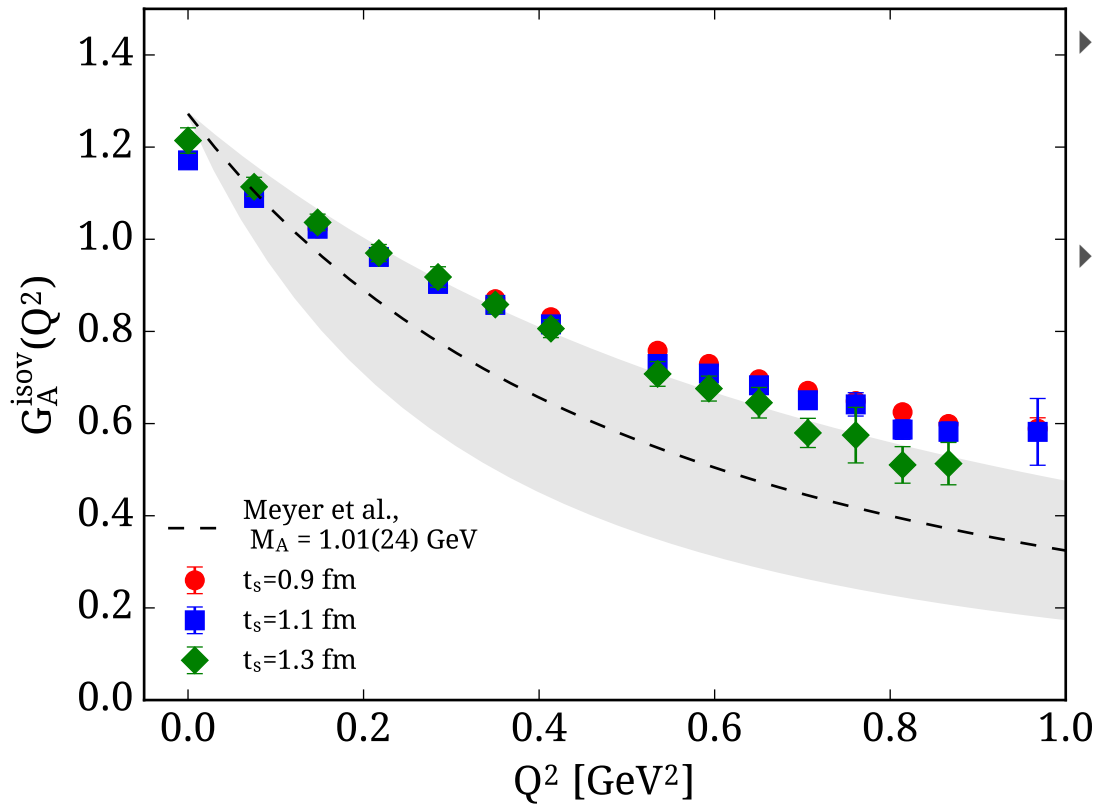




**Axial form factors:  
 $G_A(Q^2)$ ,  $G_p(Q^2)$**

# $G_A(Q^2)$

## ★ Axial form factor



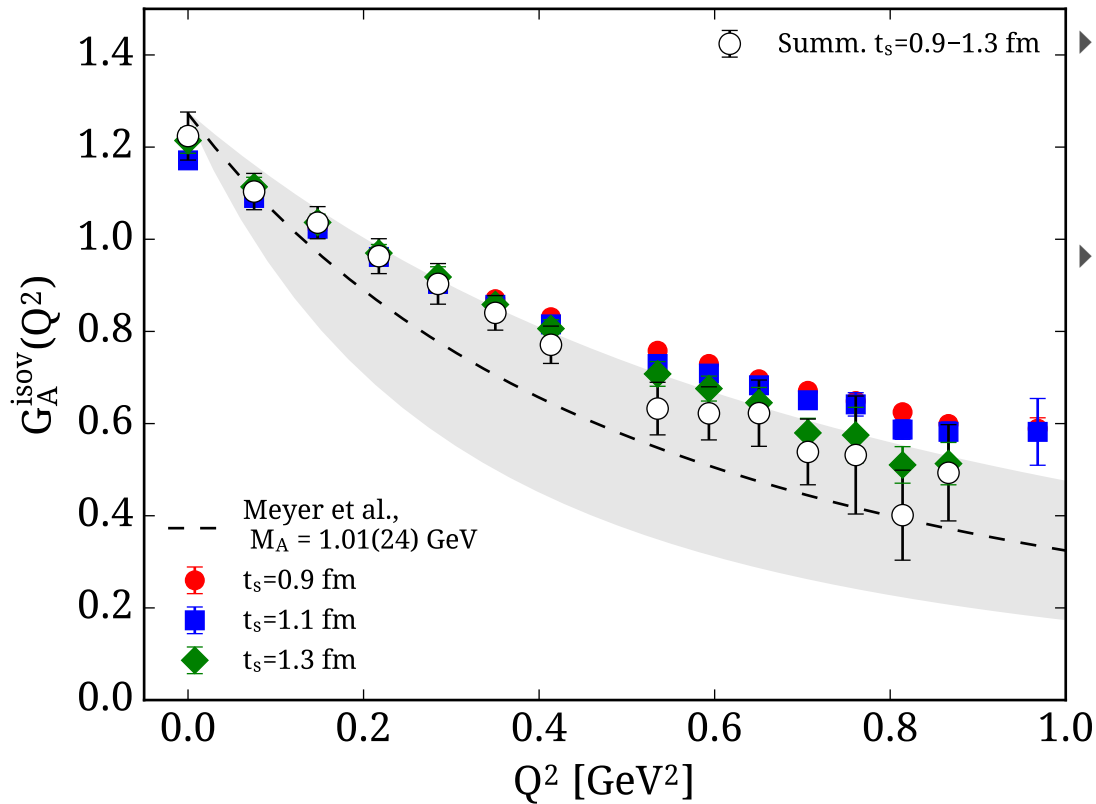
▶ Experimental parametrisation

$$G_A(Q^2) = \frac{g_A}{\left(1 + \frac{Q^2}{M_A^2}\right)^2}$$

▶  $M_A^{\text{exp}} = 1.01(24)$  GeV

# $G_A(Q^2)$

## ★ Axial form factor



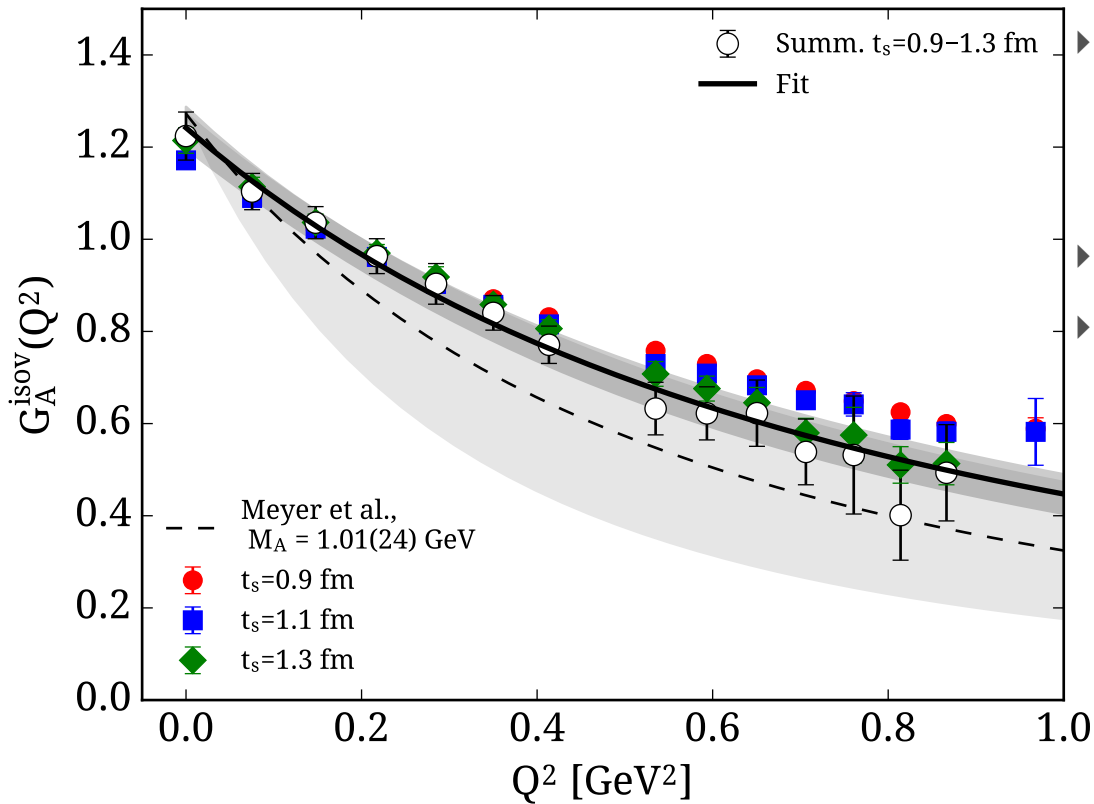
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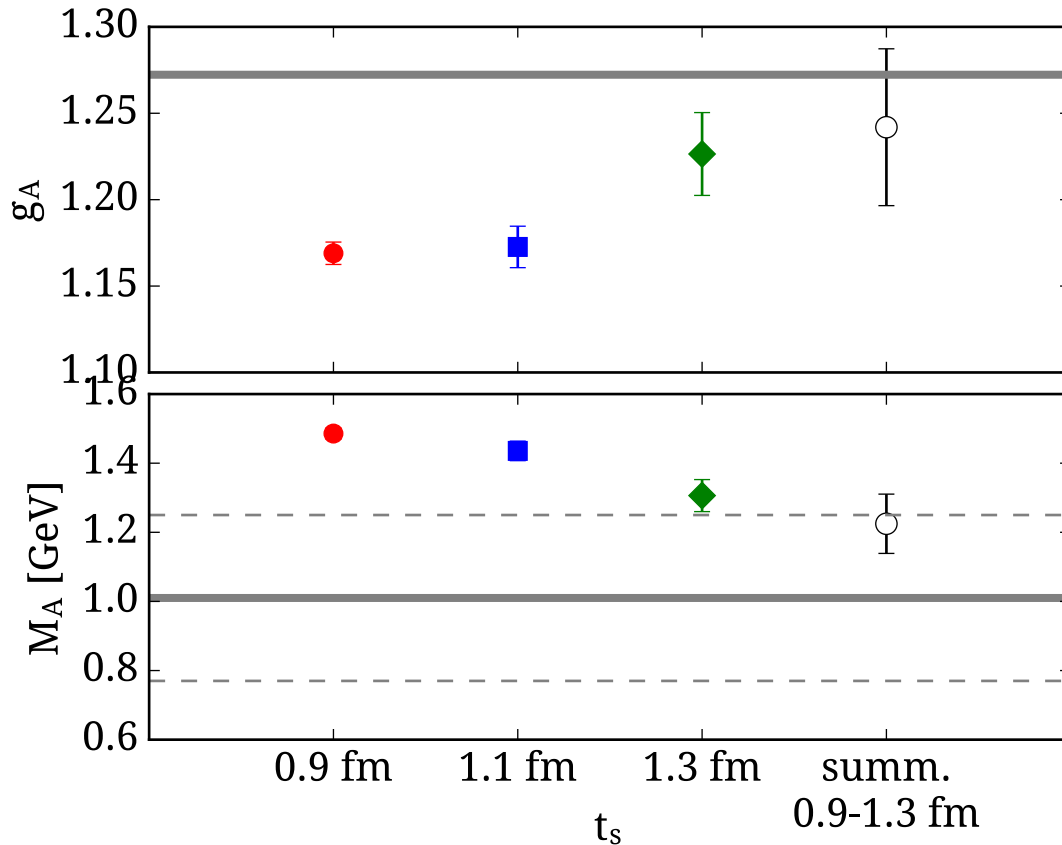
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▶  $M_A^{\text{summ.}} = 1.24(8)$  GeV

$g_A^{\text{summ.}} = 1.24(4)$

# $G_A(Q^2)$

## ★ Axial form factor



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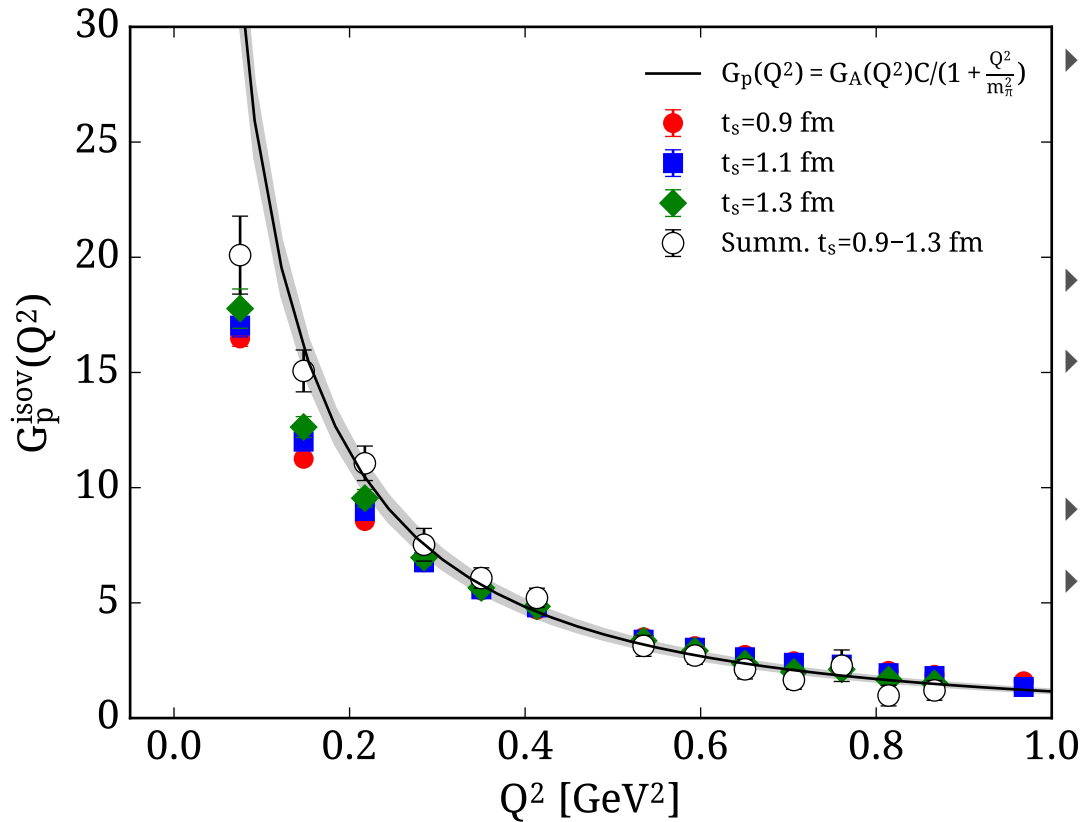
▶  $M_A^{\text{summ.}} = 1.24(8)$  GeV

$g_A^{\text{summ.}} = 1.24(4)$

▶ Suppression of excited states  
→ experimental values

# $G_p(Q^2)$

## ★ Induced pseudo-scalar form factor



▶ Expected pole dependence:

$$G_p(Q^2) = G_A(Q^2)C / (1 + \frac{Q^2}{m_\pi^2})$$

▶ Single parameter fit for  $C$

▶  $C = 138(9)$  (excl. lowest  $Q^2$ )

$$C = 4 \left[ \frac{m_N^0}{m_\pi} \right]^2 \Rightarrow$$

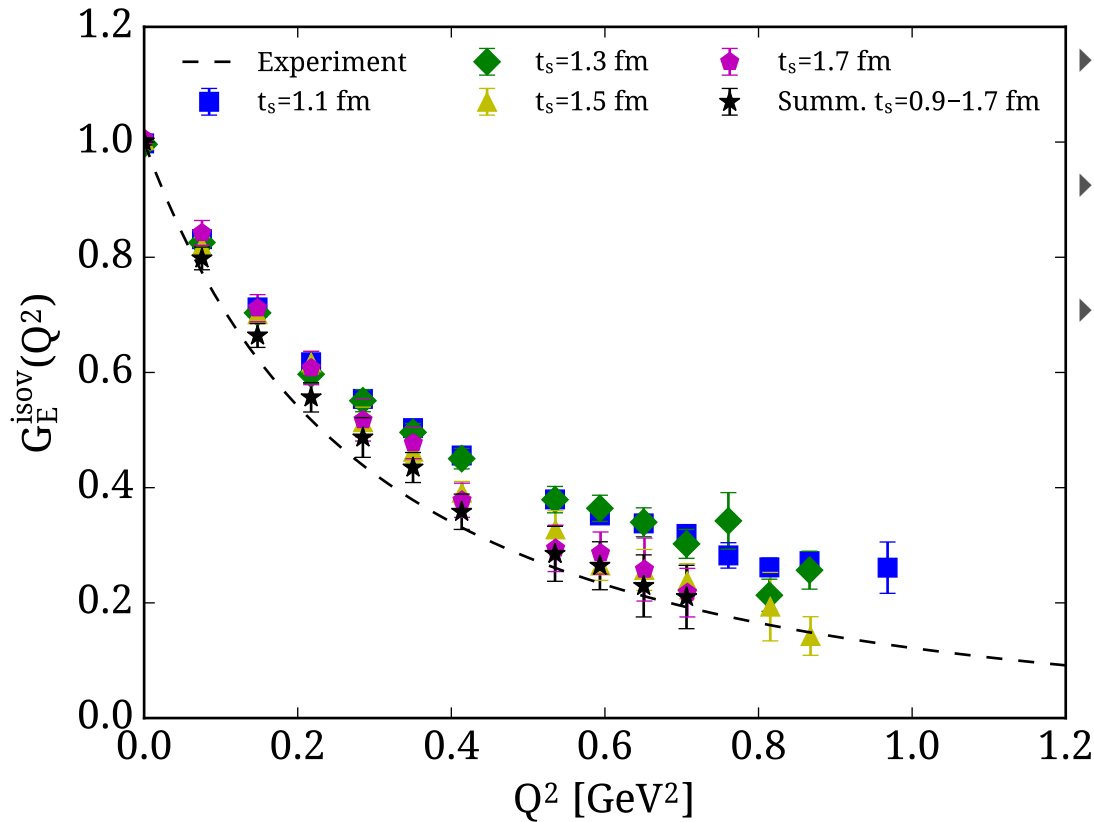
▶  $m_N^0 = 0.77(2)$  GeV

▶ Compare to  $\sim 0.88$  GeV from HB $\chi$ PT

# Electromagnetic form factors

# $G_E(Q^2)$

## ★ Isvector electric Sachs form factor

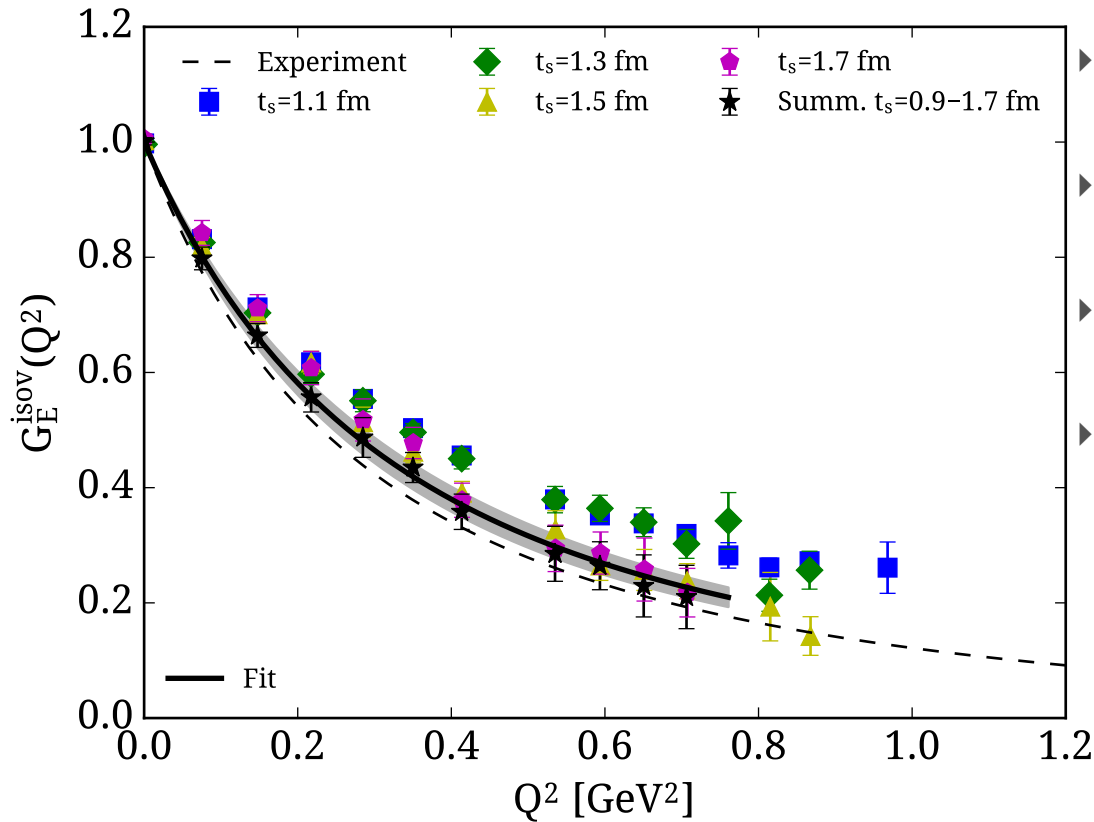


- ▶ Five sink-source separations, between 0.9 – 1.7 fm
- ▶ ~60000 statistics at largest separation
- ▶ Consistent approach to experiment



# $G_E(Q^2)$

## ★ Isovector electric Sachs form factor



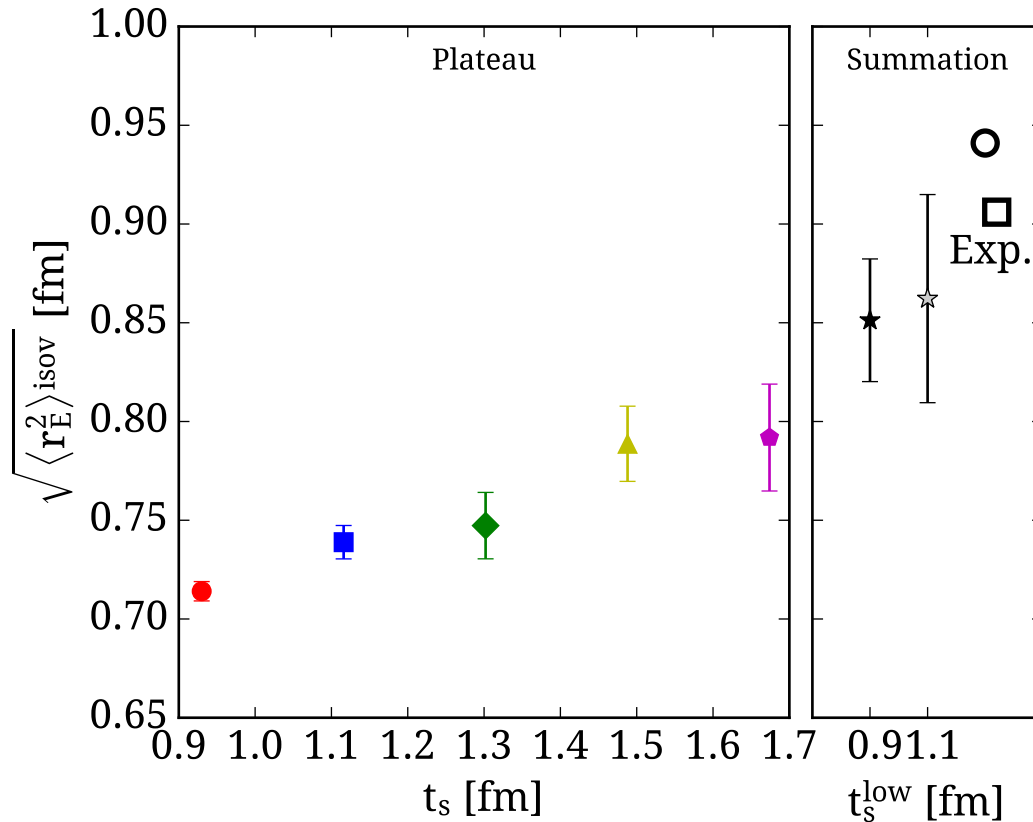
- ▶ Five sink-source separations, between 0.9 – 1.7 fm
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- ▶ Consistent approach to experiment
- ▶ Dipole fit:

$$G_E(Q^2) = \frac{1}{\left(1 + \frac{Q^2}{M_E^2}\right)}$$

$$\langle r_E^2 \rangle^{\text{isov}} = \frac{12}{M_E^2}$$

# $G_E(Q^2)$

## ★ Isovector electric Sachs form factor



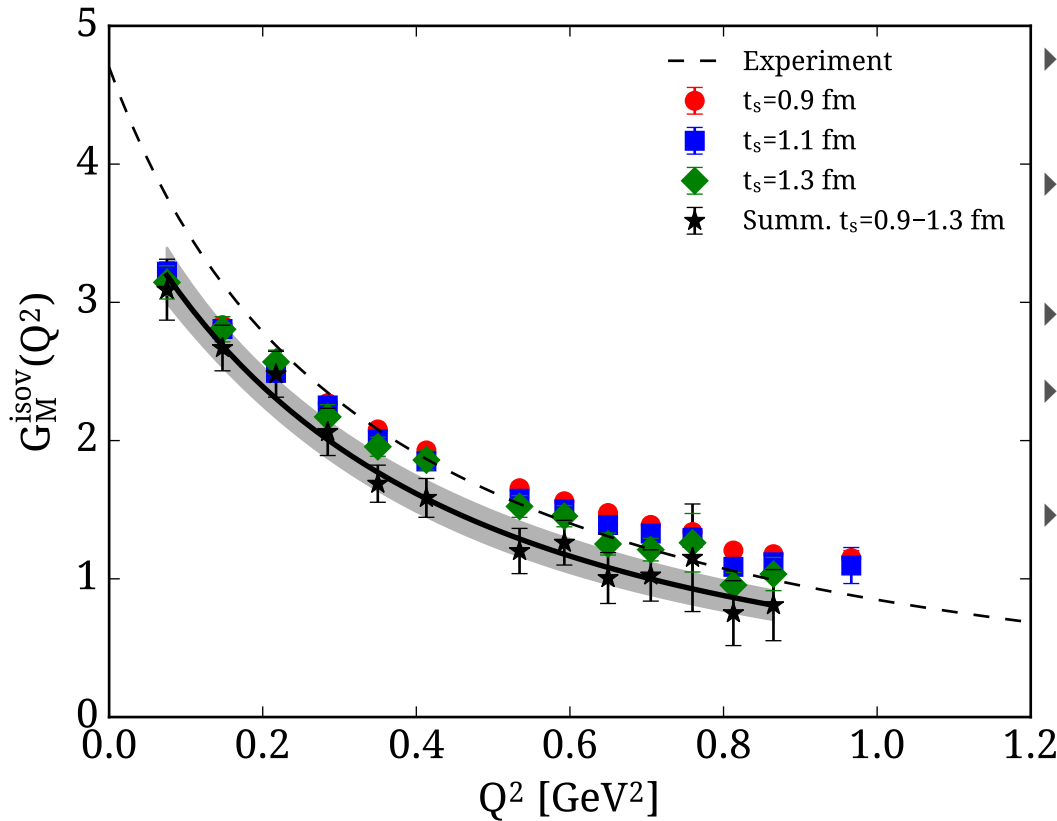
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# $G_M(Q^2)$

## ★ Isovector magnetic Sachs form factor



▶ Three sink-source separations, between 0.9 – 1.3 fm

▶ ~9300 statistics at largest separation

▶ Underestimated  $G_M(0)$

▶ Similar radius behaviour as electric case

▶ Two-parameter fit:

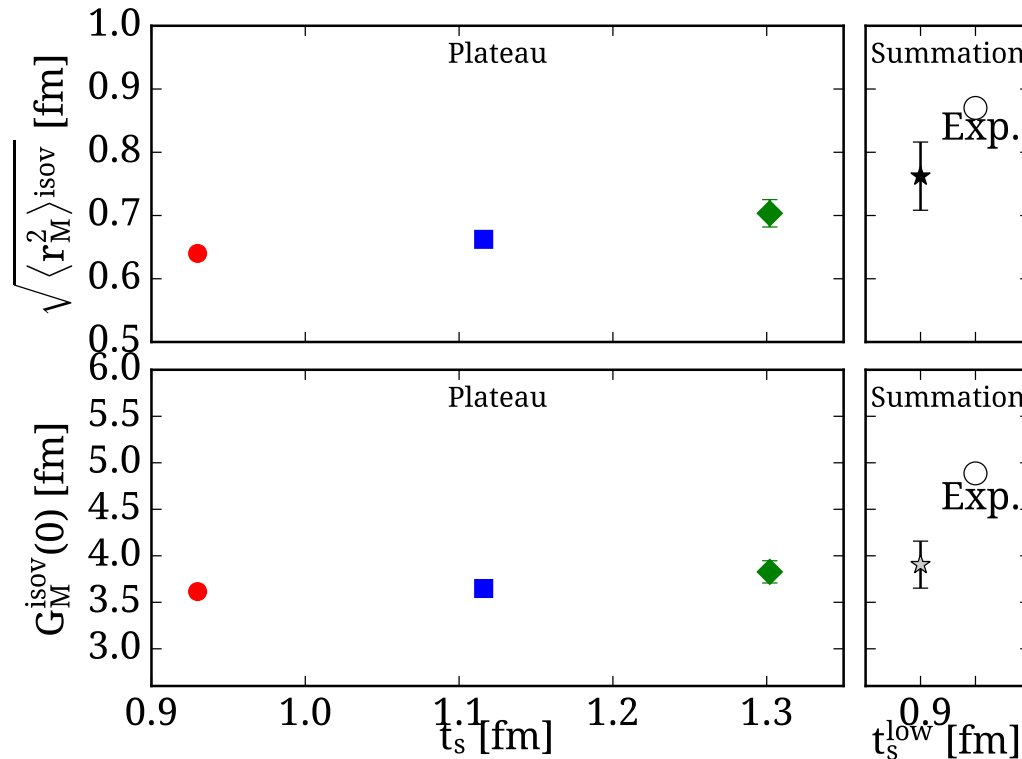
$$G_M(Q^2) = \frac{G_M(0)}{\left(1 + \frac{Q^2}{M_M^2}\right)^2}$$

$$\langle r_M^2 \rangle^{\text{isov}} = \frac{12}{M_M^2}$$

$$\mu_N^p - \mu_N^n = G_M(0)$$

# $G_M(Q^2)$

## ★ Isovector magnetic Sachs form factor



- ▶ Three sink-source separations, between 0.9 – 1.3 fm
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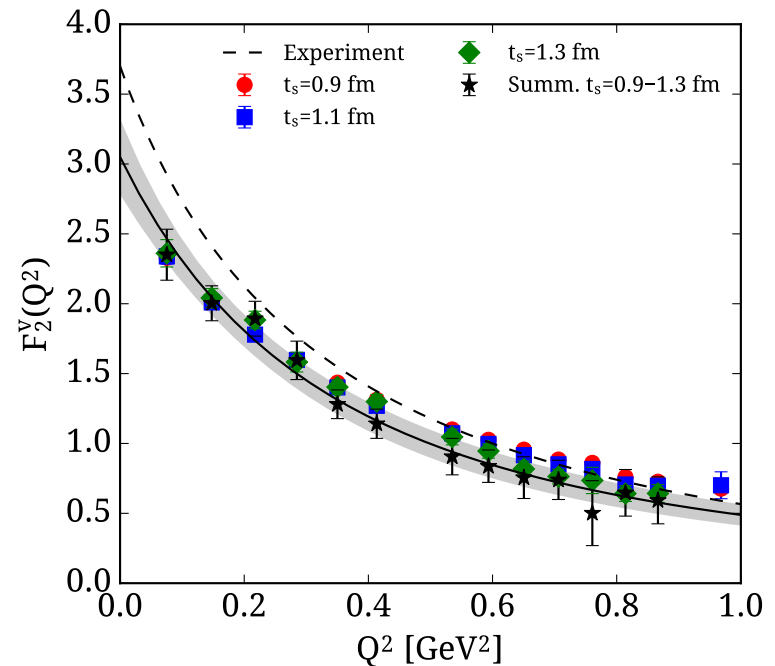
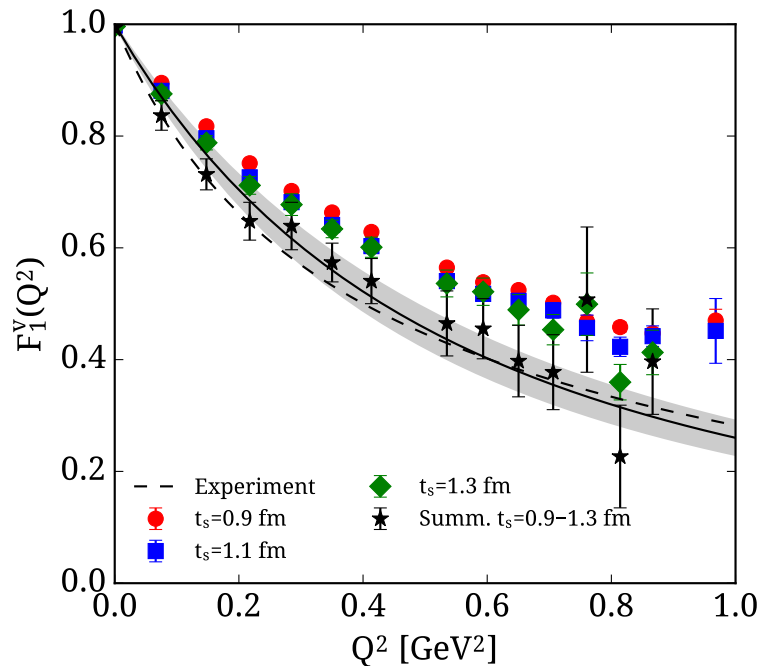
$$\langle r_M^2 \rangle^{\text{isov}} = \frac{12}{M_M^2}$$

$$\mu_N^p - \mu_N^n = G_M(0)$$

# Dirac and Pauli isovector form factors

Only for common sink – source separations between  $G_E$  and  $G_M$

$$F_1(Q^2) = \frac{G_E(Q^2) + \tau G_M(Q^2)}{1 + \tau} \quad F_2(Q^2) = \frac{G_M(Q^2) - G_E(Q^2)}{1 + \tau} \quad \tau = \frac{Q^2}{(2m_N)^2}$$

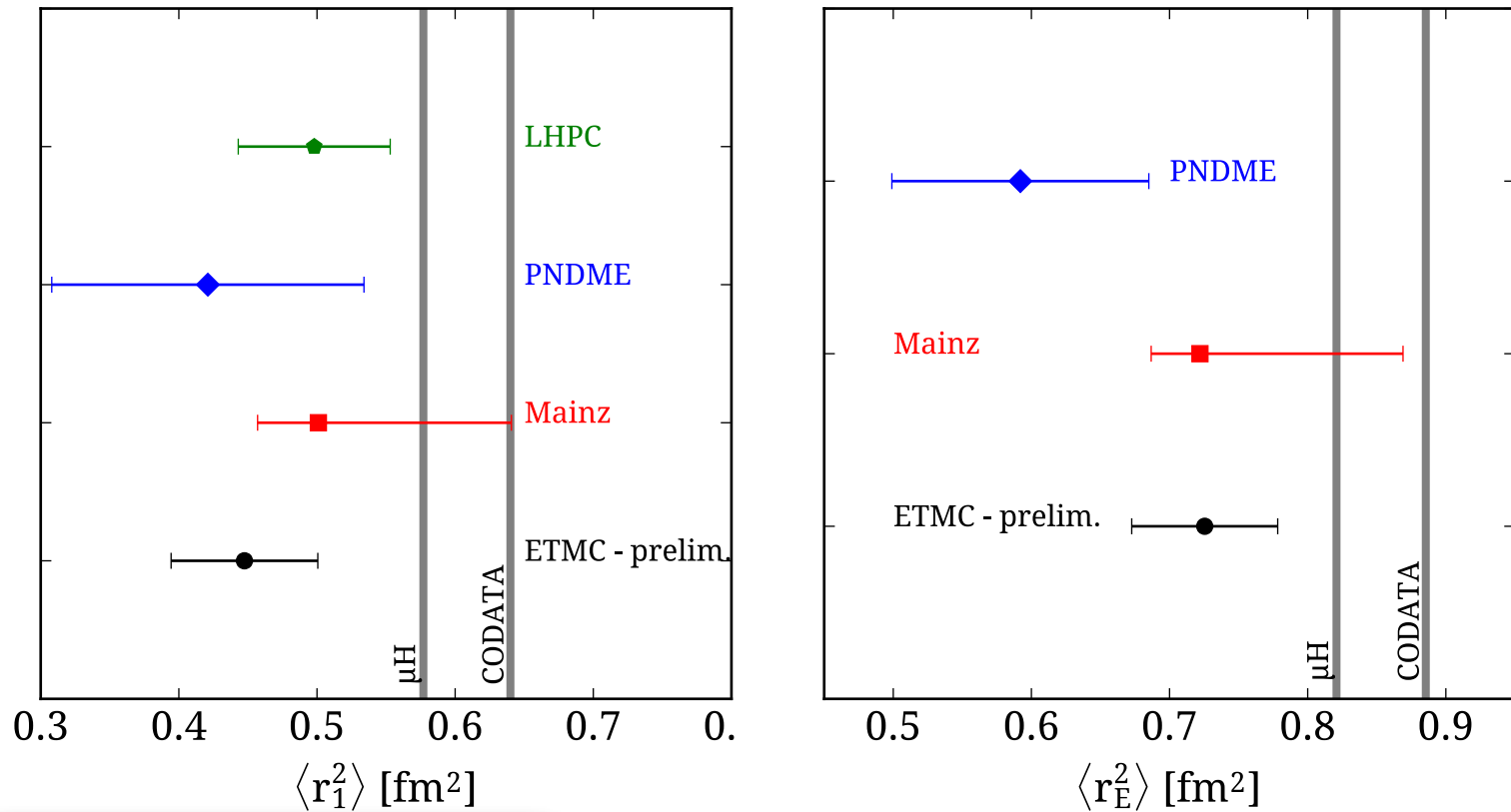


Fit to dipole form:

$$F_i(Q^2) = \frac{F_i(0)}{\left(1 + \frac{Q^2}{M_i^2}\right)^2} \quad F_1(0) = 1, F_2(0) \text{ allowed to vary}$$

# Electric and Dirac radii

Comparison of recent results near or at physical pion mass



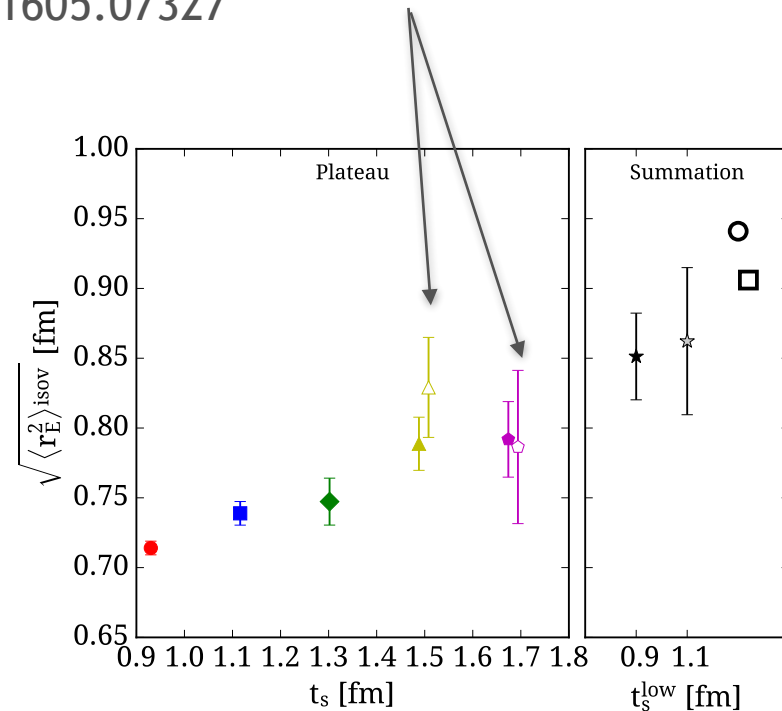
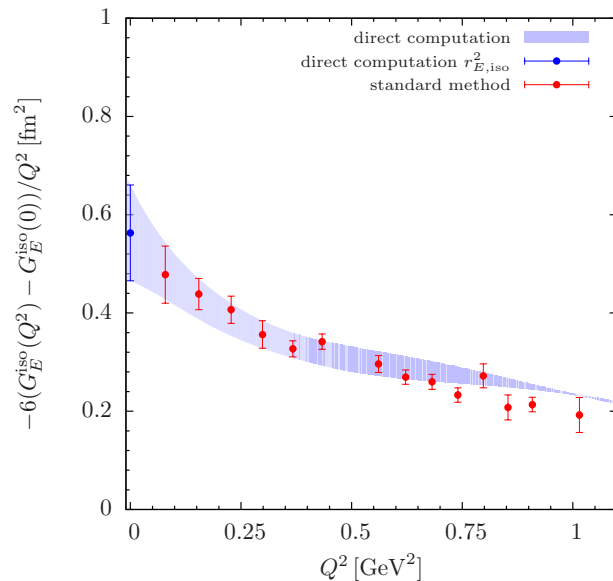
- ◆ LHPC, arXiv:1404.4029
- Mainz, arXiv:1504.04628
- ◆ PNDME, arXiv:1306.5435

Consistent trend towards larger radii with increasing sink-source separations

# Conclusions – outlook

## ★ Outlook: position space methodS for radius

- ▶ Preliminary results using few momenta
- ▶ Radius consistent within large error
- ▶ Applied to 370 MeV in arXiv:1605.07327



# Conclusions – outlook

## ★ Direct physical point calculation of axial and EM form factors

- ▶ Separations up to 1.3 or 1.7 fm
- ▶ Still observe excited state effects in radii

## ★ Axial form factors

- ▶ Axial pole mass within error of latest analysis of experimental data
- ▶ Excited state continuation → larger pole mass

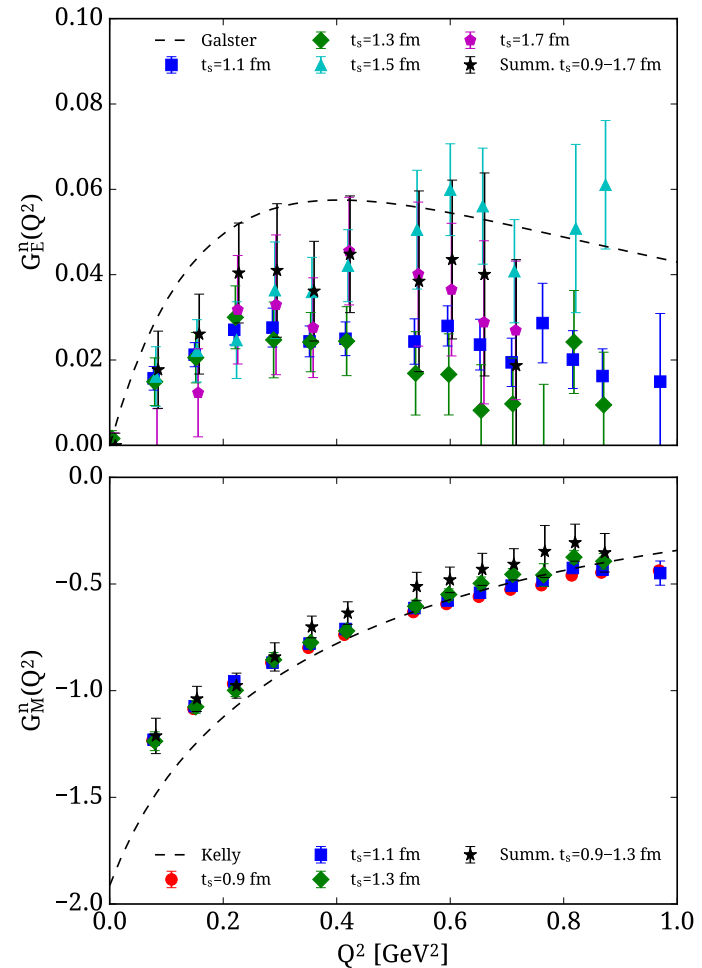
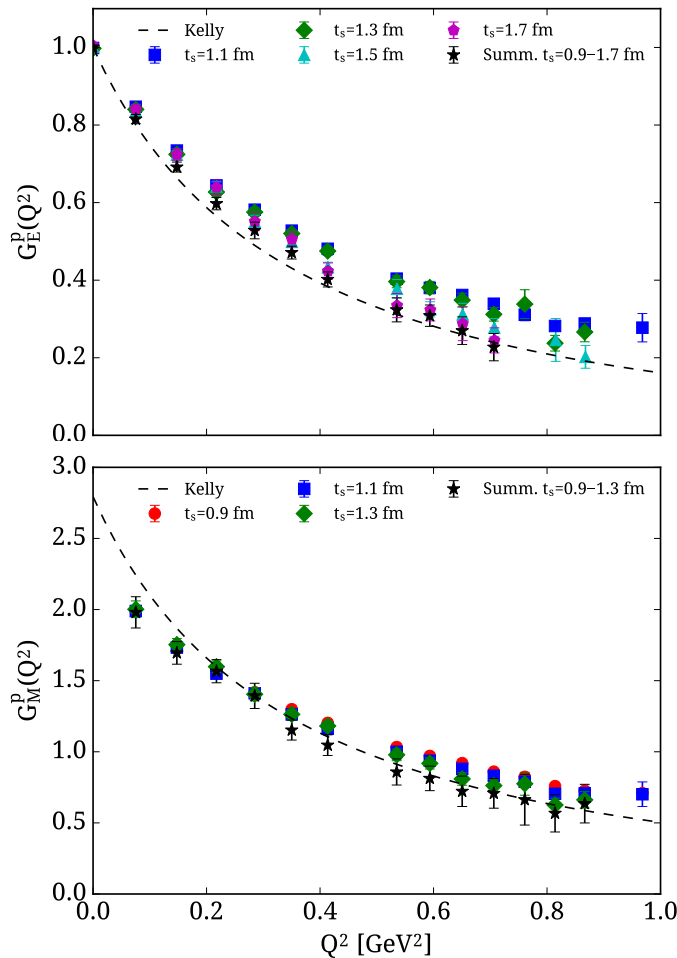
## ★ Electromagnetic form factors

- ▶ Radii approaching experimental values with increasing separations
- ▶ Still more statistics and larger separations needed though
- ▶ May need ~2 fm separation and  $O(10^5)$  statistics
- ▶  $G_M(Q^2 \rightarrow 0)$  still puzzling, excited state effects seem mild



# BONUS

# Proton and neutron [ignoring disconnected]



Disconnected diagram contributions being assessed

[see talk by A. Vaquero, Thu. 17:30]

# Error of form factors and radii

## ★ Variance estimate

- ▶ Note that different sink-source separations have different statistics
- ▶ Exponential suppression:  $(m_N - 3/2m_\pi)$  consistent with nucleon mass at all momentum transfers

