A first look at Staggered Domain Wall Fermions

Christian Hoelbling, Christian Zielinski

Bergische Universität Wuppertal NTU Singapore

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Christian Hoelbling (Wuppertal)

Introduction

Lattice fermions



Introduction

Lattice fermions



Introduction

Staggered Wilson term construction

(Golterman, Smit, 1984; Adams, 2010; C.H. 2010; deForcrand et.al. 2010,2012; Durr 2012)

$$A^{\mu_1...\mu_{2n}} = i^n \eta_{\mu_1} \dots \eta_{\mu_{2n}} \Gamma_{\mu_1...\mu_{2n}} (C_{\mu_1} \dots C_{\mu_{2n}})_{sym}$$

•
$$\{D_{st}, \epsilon\} = 0$$
 and $D_{st}^{\dagger} = -D_{st}$

- Mass term: $[\mathbf{A}, \epsilon] = \mathbf{0}, \, \mathbf{A}^{\dagger} = \mathbf{A}$
- $\rightarrow D_{\mathsf{A}}\epsilon = \epsilon D_{\mathsf{A}}^{\dagger}$
- → $\lambda_i = \lambda_{i^*}^*$, real determinant



$$egin{aligned} \mathcal{C}_{\mu} &:= rac{1}{2} \left(V_{\mu} + V_{\mu}^{\dagger}
ight) \ (V_{\mu})_{xy} &:= U_{\mu}(x) \delta_{x+\hat{\mu},y} \end{aligned}$$

•
$$\Gamma_{\mu_{1}...\mu_{2n}} = \epsilon_{\mu_{1}...\mu_{2n}} (-1)^{\sum_{i} x_{\mu_{i}}}$$

 $\sim (\gamma_{\mu_{1}} \dots \gamma_{\mu_{2n}} \otimes \xi_{\mu_{1}} \dots \xi_{\mu_{2n}})$
• $\eta_{\mu} = (-1)^{\sum_{\nu < \mu} x_{\nu}} \sim (\gamma_{\mu} \otimes 1)$
• $\epsilon = (-1)^{\sum_{\mu} x_{\mu}} \sim (\gamma_{5} \otimes \xi_{5})$
• $\{C_{\mu}, \epsilon\} = 0$
• $A^{\mu_{1}...\mu_{2n}} \sim (1 \otimes \xi_{\mu_{1}} \dots \xi_{\mu_{2n}}) + O(a)$

Domain wall operator (Kaplan 1992; Shamir 1993; Furman, Shamir 1994)

$$\bar{\Psi} D_{\mathsf{DW}} \Psi = \sum_{s=1}^{N_s} \bar{\Psi}_s \left(D_W^+ \Psi_s - P_- \Psi_{s+1} - P_+ \Psi_{s+1} \right)$$

with $P_{\pm} = \frac{1}{2}(1 \pm \gamma_5), D_W^{\pm} = D_W(-M_0) \pm 1$

Boundary conditions with mass term:

$$P_{+} (\Psi_{0} + m\Psi_{N_{s}}) = 0 \qquad P_{-} (\Psi_{N_{s}+1} + m\Psi_{1}) = 0$$

modification (Boriçi, 1999) Optimal DWF (Chiu, 2002)

$$\mathcal{P}_{\pm}\Psi_{m{s}\mp1}
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- $D_W^{\pm}
 ightarrow D_W^{\pm}(s) = \omega_s D_W(-M_0) \pm 1$
 - Staggered versions (Adams, 2011) $D_W \rightarrow D_A$ $\gamma_5 \rightarrow \epsilon \sim (\gamma_5 \otimes \xi_5)$

Borici's

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$$P_{\pm}\Psi_{s\mp1} \rightarrow -D_W^-P_{\pm}\Psi_{s\mp1}$$





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Boric

Numerical setup

- Schwinger model (QED₂)
- Wilson gauge action
- Unsmeared and 3-APE ($\alpha = 0.5$)

| β | $N_x \times N_t$ | #conf |
|---------|--------------------------|-------|
| 0.8 | 8 × 8 | 1000 |
| 1.8 | 12×12 | 1000 |
| 3.2 | 16 	imes 16 | 1000 |
| 5 | 20 	imes 20 | 1000 |
| 7.2 | 24 imes 24 | 1000 |
| 9.8 | 28 imes 28 | 1000 |
| 12.8 | 32 	imes 32 | 1000 |

Observables:

- Normality violation $\Delta_N = ||[D, D^{\dagger}]||$
- GW violation $\Delta_{GW} = ||\gamma_5 D D\hat{\gamma}_5||$
- Residual mass $m_{\rm eff} = |\lambda_{\rm min}|$











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Staggered Domain Wall Fermions

Wilson kernel

Staggered Wilson kernel



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Bulk operators

Wilson kernel

Staggered Wilson kernel



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Wilson kernel

Staggered Wilson kernel



Wilson kernel

Staggered Wilson kernel















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Continuum behaviour



Smeared continuum behaviour



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Why the good chiral properties?



QCD outlook

QCD, $6^4 \times 8$, $\beta = 6$, APE smeared



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QCD outlook

QCD, $6^4 \times 8$, $\beta = 6$, APE smeared



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Summary

Could staggered DWF be useful?

- ✓ Construction is straightforward
- ✓ Good chiral symmetry properties
- Easily parallelizable

- Counterterms for N_f = 1 (Sharpe 2012)
- X Staggered spin structure





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