# New polynomially exact integration rules on $U(N)$ and $S U(N)$ 

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We consider the $0+1$ dimensional QCD with chemical potential.
Dirac operator for a quark of mass $m$ at chemical potential $\mu$ :

$$
\mathfrak{D}(U)=\left(\begin{array}{ccccc}
m & \frac{e^{\mu}}{2} U_{1} & & & \frac{e^{-\mu}}{2} U_{n}^{*} \\
-\frac{e^{-\mu}}{2} U_{1}^{*} & m & \frac{e^{\mu}}{2} U_{2} & & \\
& \ddots & \ddots & \ddots & \\
& & -\frac{e^{-\mu}}{2} U_{n-2}^{*} & m & \frac{e^{-\mu}}{2} U_{n-1} \\
-\frac{e^{\mu}}{2} U_{n} & & & -\frac{e^{-\mu}}{2} U_{n-1}^{*} & m
\end{array}\right)
$$

$N_{f}$ flavor partition function:

$$
Z(m, \mu, G, n)=\int_{G^{n}} \operatorname{det} \mathfrak{D}(U)^{N_{f}} d h_{G}^{n}(U)
$$

with $G \in\{S U(N), U(N)\}$ and $h_{G}$ the (normalized) Haar measure on $G$.

Choice of Gauge: $U_{j}=1$ except $U=U_{n}$. Then,

$$
\operatorname{det} \mathfrak{D}=\operatorname{det}\left(\prod_{j=1}^{n} \tilde{m}_{j}+2^{-n} e^{-n \mu} U^{*}+(-1)^{n} 2^{-n} e^{n \mu} U\right)
$$

with $\tilde{m}_{1}:=m$,

$$
\forall j \in[2, n-1] \cap \mathbb{N}: \quad \tilde{m}_{j}:=m+\frac{1}{4 \tilde{m}_{j-1}},
$$

and

$$
\tilde{m}_{n}:=m+\frac{1}{4 \tilde{m}_{n-1}}+\sum_{j=1}^{n-1} \frac{(-1)^{j+1} 2^{-2 j}}{\tilde{m}_{j} \prod_{k=1}^{j-1} \tilde{m}_{k}^{2}}
$$

$0+1$ dimensional QCD for a quark at non-zero chemical potential

$$
\int_{G} \operatorname{det}\left(\prod_{j=1}^{n} \tilde{m}_{j}+2^{-n} e^{-n \mu} U^{*}+(-1)^{n} 2^{-n} e^{n \mu} U\right)^{N_{f}} d h_{G}(U)
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is a highly oscillating integral.

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The sign problem is very problematic for $m$ small and $n \mu$ large.

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## Idea

Use polynomially exact quadrature rules over $U(N), S U(N)$.

We will consider ( $N_{f}=1$ )

$$
Z(m, \mu, G, n)=\int_{G} \operatorname{det} \mathfrak{D} d h_{G}
$$

and the chiral condensate

$$
\partial_{m} \ln Z(m, \mu, G, n)=\frac{\int_{G} \partial_{m} \operatorname{det} \mathfrak{D} d h_{G}}{\int_{G} \operatorname{det} \mathfrak{D} d h_{G}}
$$

with

$$
\operatorname{det} \mathfrak{D}=\operatorname{det}\left(\prod_{j=1}^{n} \tilde{m}_{j}+2^{-n} e^{-n \mu} U^{*}+(-1)^{n} 2^{-n} e^{n \mu} U\right)
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## Solution: Symmetrization?

## Bloch, Bruckmann, Wettig (2013)

Center Symmetry of $S U(3)$ : Choosing quadrature rules $Q$ satisfying

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\forall U \in Q: U^{*}, e^{\frac{2 \pi i}{3}} U, e^{\frac{4 \pi i}{3}} U \in Q
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Can we construct polynomially exact quadrature rules?

We want completely symmetric quadrature rules.

- $U(N) \cong S U(N) \rtimes U(1)$ and $h_{U(N)}=h_{S U(N)} \times h_{U(1)}$ where $\rtimes$ denotes the semidirect product and $h$ the normalized Haar measure.
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- $\left\{e^{\frac{2 \pi i k}{t+1}+\varphi_{0}} ; k \in\{0, \ldots, t\}\right\} \subseteq U(1)$ integrates all polynomials of degree $\leq t$ exactly ( $\varphi_{0}$ arbitrary).
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- polynomials in $S U(N)$ of degree $\leq k$ are mapped to "unitary polynomials" of degree $\leq k$ in $\times_{j=1}^{N-1} S^{2 j+1}$


## Genz (2003)

(Randomized) Completely symmetric (and thus polynomially exact) quadrature rules on spheres $S^{n}$ can be constructed explicitly.
computing the partition function $Z(m, \mu, G, n)=\int_{G} \operatorname{det} \mathfrak{D} d h_{G}$ $\operatorname{det} \mathfrak{D}(U)=\operatorname{det}\left(\prod_{j=1}^{n} \tilde{m}_{j}+2^{-n} e^{-n \mu} U^{*}+(-1)^{n} 2^{-n} e^{n \mu} U\right)$

computing the partition function $Z(m, \mu, G, n)=\int_{G} \operatorname{det} \mathfrak{D} d h_{G}$


Note that the symmetrized (polynomially exact) quadrature rule yields double precision results for all values of $m$.
computing the partition function $Z(m, \mu, G, n)=\int_{G} \operatorname{det} \mathfrak{D} d h_{G}$
$n=20, \mu=1.0, U(3)$, number sequences $=10,1024$ bit ext. floats

computing the chiral condensate $\partial_{m} \ln Z(m, \mu, G, n)=\partial_{m} Z / Z=\int_{G} \partial_{m} \operatorname{det} \mathfrak{D} d h_{G} / \int_{G} \operatorname{det} \mathfrak{D} d h_{G}$

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$n=8, \mu=1.0$, number of sequences $=50, U(2)$


## Conclusion

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arXiv:1607.05027 [hep-lat]
$0+1$ dimensional QCD and Polynomially Exact Rules
Numerical Results


