Leptonic decay-constant ratio f_K/f_π from clover-improved $N_f = 2 + 1$ QCD

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34th International Symposium on Lattice Field Theory July, 24-30, 2016 Southampton, UK • in collaboration with (Wuppertal U., JSC, CNRS Marseille, UR)

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• preprint:

arXiv:1601.05998[hep-lat]

Leptonic decay-constant ratio f_K/f_{π} from lattice QCD using 2+1 clover-improved fermion flavors with 2-HEX smearing





Cabibbo-Kobayashi-Maskawa (CKM)-matrix elements

- (leptonic) decay constants $f_{\pi^{\pm}}$, $f_{K^{\pm}} \rightarrow \ell \nu_{\ell}$ of (charged) π -, K-meson $\longrightarrow V_{us}$, V_{ud}
- leptonic decay ($\ell=e^{\pm}$, μ^{\pm}), experimentally measured: decay widhts

$$\frac{\Gamma(K^{\pm} \to \ell \nu_{\ell})}{\Gamma(\pi^{\pm} \to \ell \nu_{\ell})} = \frac{V_{\rm us}^2}{V_{\rm ud}^2} \frac{f_{K^{\pm}}^2}{f_{\pi^{\pm}}^2} \frac{M_{K^{\pm}}^2}{M_{\pi^{\pm}}^2} \frac{\left(1 - \frac{m_{\ell}^2}{M_{K^{\pm}}^2}\right)^2}{\left(1 - \frac{m_{\ell}^2}{M_{\pi^{\pm}}^2}\right)^2} \left(1 + \delta_{\rm em}\right)$$

MARCIANO, 2004

• measured on lattice (w/o isospin-splitting/EM)

$$rac{f_K}{f_\pi}$$

• $V_{\rm us}/V_{\rm ud}$, with further input: first-row unitarity, $V_{\rm us}$





ensembles

- generated by the Budapest-Marseille-Wuppertal (BMW)-Collaboration
- $N_f = 2 + 1$ tree-level clover-improved fermions, 2-HEX smearing
- tree-level Symanzik gluon action
- 47 ensembles in total
- five β -values: 3.31, 3.50, 3.61, 3.70, 3.80 ($1/a \simeq 1.7$ –3.7 GeV, $a \simeq 0.12$ –0.05 fm)
- M_{π} : 130–680 MeV, m_s physical (except two ensembles)



• have been used in other studies: quark masses, SU(2)-LECs and f_{π} , B_K , nucleon σ -term





extrapolation to the physical world

- \bullet interpolation to physical quark-masses $m_{\rm ud}~\to~m_{\rm ud}^{\rm phys}$, $m_{\rm s}~\to~m_{\rm s}^{\rm phys}$
- extrapolate to infinite volume $V~\rightarrow~\infty$
- extrapolate to the continuum limit $a \rightarrow 0$
- scale-setting 1/a

variation over fits, ranges, . . . ("Wuppertal-method") for systematic uncertainty

• at least two different functional forms for each extrapolation





extrapolation to $M_{\pi}^{\rm phys}$, $M_{K}^{\rm phys}$

• NLO-SU(3): 2 fit-parameters

$$\frac{f_K}{f_\pi} = 1 + \frac{c_0}{2} \left\{ \frac{5}{4} M_\pi^2 \log\left(\frac{M_\pi^2}{\mu^2}\right) - \frac{1}{2} M_K^2 \log\left(\frac{M_K^2}{\mu^2}\right) - \frac{3}{4} M_\eta^2 \log\left(\frac{M_\eta^2}{\mu^2}\right) + c_1 [M_K^2 - M_\pi^2] \right\}$$
$$M_\eta^2 = \frac{1}{3} (4M_K^2 - M_\pi^2), \quad c_0 = \frac{1}{(4\pi F_0)^2}, \quad c_1 = 128\pi^2 L_5(\mu)$$

• 3-,4-,6-parameter polynomial in M_{π}^2 (dominant) and $[M_K^2 - M_{\pi}^2]$ (flavor-symmetry constr.):

$$\begin{aligned} \frac{f_K}{f_\pi} &= 1 + [M_K^2 - M_\pi^2] \Big(c_0^{3\text{-par}} + c_1^{3\text{-par}} [M_K^2 - M_\pi^2] + c_2^{3\text{-par}} M_\pi^2 \Big) \\ \frac{f_K}{f_\pi} &= 1 + [M_K^2 - M_\pi^2] \Big(c_0^{4\text{-par}} + c_1^{4\text{-par}} [M_K^2 - M_\pi^2] + c_2^{4\text{-par}} M_\pi^2 + c_3^{4\text{-par}} M_\pi^4 \Big) \\ \frac{f_K}{f_\pi} &= 1 + [M_K^2 - M_\pi^2] \Big(c_0^{6\text{-par}} + c_1^{6\text{-par}} [M_K^2 - M_\pi^2] + c_2^{6\text{-par}} M_\pi^2 + c_3^{6\text{-par}} M_\pi^4 \\ &+ c_4^{6\text{-par}} M_\pi^2 [M_K^2 - M_\pi^2] + c_5^{6\text{-par}} [M_K^2 - M_\pi^2]^2 \Big) \end{aligned}$$

• all fit-functions satisfy flavor-symmetry constraint: $\frac{f_K}{f_\pi}\Big|_{m_{\rm ud}=m_{\rm s}}=1$





continuum limit extrapolation

• cut-off effects

(clover-improved fermions, tree-level Symanzik gluons)

$\propto \alpha a$

with strong coupling $\alpha = g^2/(4\pi)$ at scale 1/a

- α : logarithmic function of a
- two ansätze: a^2 and $a\alpha$ -scaling

β	$a^{-1}/{\sf GeV}$	$lpha_{N_f=3}$	$\alpha_{N_f=4}$
3.31	1.670(07)	0.327	0.333
3.50	2.134(15)	0.286	0.295
3.61	2.576(28)	0.262	0.271
3.70	3.031(32)	0.244	0.254
3.80	3.657(37)	0.227	0.237

$$\frac{f_{K}}{f_{\pi}} = 1 + \left(\frac{f_{K}}{f_{\pi}}(M_{\pi}, M_{K}) - 1\right) \left(1 + c^{\mathsf{disc}}a^{2}\right)$$
$$\frac{f_{K}}{f_{\pi}} = 1 + \left(\frac{f_{K}}{f_{\pi}}(M_{\pi}, M_{K}) - 1\right) \left(1 + c^{\mathsf{disc}}\alpha a\right)$$

1 fit-parameter, flavor-symmetry constraint

scale-setting: mass-independent via M_Ω for each β or per-ensemble







infinite volume extrapolation

• ChPT: finite volume $M_{\pi,K}(L)$, $f_{\pi,K}(L) \leftrightarrow \text{infinite volume } M_{\pi,K}$, $f_{\pi,K}$

$$\begin{aligned} \frac{f_K(L)}{f_\pi(L)} &= \frac{f_K}{f_\pi} \Big(1 + c^{\mathsf{FV}} \left[\frac{5}{8} \tilde{g}_1(M_\pi L) - \frac{1}{4} \tilde{g}_1(M_K L) - \frac{3}{8} \tilde{g}_1(M_\eta L) \right] \Big) \\ \tilde{g}_1(z) &= \frac{24}{z} K_1(z) + \frac{48}{\sqrt{2}z} K_1(\sqrt{2}z) + \frac{32}{\sqrt{3}z} K_1(\sqrt{3}z) + \frac{24}{2z} K_1(2z) + \dots \\ K_1(z) &= \sqrt{\frac{\pi}{2z}} e^{-z} \left\{ 1 + \frac{3}{8z} - \frac{3 \cdot 5}{2(8z)^2} + \frac{3 \cdot 5 \cdot 21}{6(8z)^3} - \frac{3 \cdot 5 \cdot 21 \cdot 45}{24(8z)^4} + \dots \right\} \end{aligned}$$

- either full $ilde{g}_1(z)$ or first term of expansion in z
- 1 fit-parameter
- flavor-symmetry constraint satisfied





functional forms for extrapolation — summary

- 4 functional forms for physical mass (NLO-SU(3) ChPT, 3-, 4-, 6-par. polynomial)
 2, 3, 4, or 6 fit-parameters
- 2 functional forms for continuum limit $(a^2, a\alpha)$ 1 fit-parameter
- 2 functional forms for infinite volume limit (full \tilde{g}_1 , leading terms) 1 fit-parameter
- 2 method for scale-setting (mass-independent, per-ensemble)

 $4 \times 2 \times 2 \times 2 = 32$ combinations

4, 5, 6 or 8 fit-parameters





fit ranges

- ensembles: $2^{47} \simeq \mathcal{O}(10^{14})$ combinations • 47 neither feasible nor sensible
- M_{π}^{max} : none, 350 MeV, 300 MeV, 250 MeV
- M_K^{max} : none, 600 MeV, 550 MeV, 500 MeV
- $(M_{\pi}L)^{\min}$: none, 3.85, 4.05
- β^{\min} : 3.31, 3.50, 3.61
- ≥ 5 data-points: 63 combinations
- only "true" fits

- theor. $32 \times 63 = 2016$ fits
- effective: **1368 fits** $(n_{d.o.f.} \geq 1)$















dependence on fit types, fit ranges





E.E. Scholz — Leptonic decay-constant ratio f_K/f_π



${f_{K^\pm}}/{f_{\pi^\pm}}$, $V_{ m us}$, $V_{ m ud}$

 $m_{\rm up} \neq m_{\rm down}$, EM-effects: correction from ChPT [Gasser, Leutwyler (1985); Cirigliano, Neufeld (2011)]

$$\delta_{
m SU(2)} = -0.0061(61)
ightarrow rac{f_K \pm}{f_{\pi \pm}} = rac{f_K}{f_{\pi}} \sqrt{1 + \delta_{
m SU(2)}} = 1.178(10)_{
m stat}(26)_{
m syst}$$

 $\delta_{SU(2)}$: -0.0043(12) [CIRIGLIANO, NEUFELD (2011)] and FLAG-II, -0.0078(7) [DIVITIIS ET AL. (2012)]

• $V_{\rm us}/V_{\rm ud}$ from exp. measured decay widths

$$\frac{V_{\rm us}}{V_{\rm ud}} \frac{f_{K^{\pm}}}{f_{\pi^{\pm}}} = 0.27599(38)_{\rm exp} \rightarrow \frac{V_{\rm us}}{V_{\rm ud}} = 0.2343(20)_{\rm stat}(52)_{\rm syst}(03)_{\rm exp}$$

• using super-allowed nuclear β -decay $V_{\rm ud} = 0.97417(21)_{\rm nuc}$

[HARDY, TOWNER (2015)]

[Moulson (2014), Rosner et al. (2015)]

 $V_{\rm us} = 0.2282(19)_{\rm stat}(51)_{\rm syst}(03)_{\rm exp\&nuc}$

• first-row unitarity with $|V_{\mathsf{ub}}| = 4.12(37)(06) \cdot 10^{-3}$ [Rosner et al. (2015)]

$$V_{\rm ud}^2 + V_{\rm us}^2 + |V_{\rm ub}|^2 - 1 = 0.0011(09)_{\rm stat}(23)_{\rm syst}(05)_{\rm exp\&nuc} = 0.0011(25)_{\rm comb}$$

• alternatively: w/o V_{ud} from nuclear β -decay using $|V_{ub}|$ and first-row unitarity

$$V_{\rm ud} = 0.9736(04)_{\rm stat}(11)_{\rm syst}(01)_{\rm exp}$$
 $V_{\rm us} = 0.2281(18)_{\rm stat}(48)_{\rm syst}(03)_{\rm exp}$







comparison with FLAG-2016



our result: $f_{K^\pm}/f_{\pi^\pm} = 1.178(28)$



FLAG-2016 $N_f = 2 + 1$ avr.: $V_{\rm us} = 0.2243(10)$, $V_{\rm ud} = 0.97451(23)$

our results: (first row unit.+ $V_{
m ub}$) $V_{
m us}=0.2281(51)$, $V_{
m ud}=0.9736(12)$





BACKUP





NLO SU(3)-ChPT fit



mass-independent scale, full FV, a^2 -discretization $M_\pi \leq 300 \, {
m MeV}$, $M_K \leq 600 \, {
m MeV}$, $(M_\pi L) \geq 3.85$, $eta \geq 3.50$

 $\chi^2=7.6$, $n_{
m d.o.f.}=6$, p=0.27





NLO SU(3)-ChPT fit



mass-independent scale, full FV, $a\alpha$ -discretization $M_{\pi} \leq 350 \text{ MeV}, M_K \leq 600 \text{ MeV}$

 $\chi^2=47$, $n_{
m d.o.f.}=23$, $p=2\cdot 10^{-3}pprox 0$







mass-independent scale, full FV, a^2 -discretization $M_\pi \leq 300 \ {
m MeV}$, $M_K \leq 600 \ {
m MeV}$, $(M_\pi L) \geq 3.85$, $eta \geq 3.50$

 $\chi^2=5.8$, $n_{
m d.o.f.}=5$, p=0.33







mass-independent scale, full FV, a^2 -discretization $M_{\pi} \leq 350$ MeV, $M_K \leq 600$ MeV

 $\chi^2=54$, $n_{
m d.o.f.}=22$, $p=1.7~\cdot~10^{-4}~pprox~0$







mass-independent scale, full FV, a^2 -discretization $M_{\pi} \leq 300$ MeV, $M_K \leq 600$ MeV, $(M_{\pi}L) \geq 3.85$, $\beta \geq 3.50$

 $\chi^2=5.8$, $n_{
m d.o.f.}=4$, p=0.21







mass-independent scale, full FV, $a^2\text{-discretization}$ $\beta \geq 3.61$

 $\chi^2=72$, $n_{
m d.o.f.}=18$, $p=2~\cdot~10^{-8}pprox 0$







mass-independent scale, full FV, a^2 -discretization $m_K \leq 600 \text{ MeV}, (M_{\pi}L) \geq 4.05$

 $\chi^2=11$, $n_{ ext{d.o.f.}}=10$, p=0.32











flavor-symmetry constraint

$$m_{
m ud}~=~m_{
m s}$$
 implies $M_{\pi}~=~M_{K}\,,~f_{\pi}~=~f_{K}$

$$\left.\frac{f_K}{f_\pi}\right|_{m_{\rm ud}=m_{\rm s}} = 1$$

also at (any) finite volume and lattice spacing

only functional forms satisfying this constraint





extrapolation to physical masses

• we use M_{π} , M_K instead of (bare, renormalized) quark masses

$$\begin{split} M_{\pi}^{2}|_{\text{LO-ChPT}} &= 2B_{0}m_{\text{ud}} \\ M_{K}^{2}|_{\text{LO-ChPT}} &= B_{0}(m_{\text{ud}} + m_{\text{s}}) \\ (2M_{K}^{2} - M_{\pi}^{2})|_{\text{LO-ChPT}} &= 2B_{0}m_{\text{s}} \end{split}$$

• correction in isospin-limit (FLAG-recommendation)

$$M_{\pi}^{\text{phys}} = 134.8(0.3) \text{ MeV} \quad M_{K}^{\text{phys}} = 494.2(0.4) \text{ MeV}$$





NLO-SU(3)-ChPT extrapolation

$$\frac{f_K}{f_\pi} = 1 + \frac{c_0}{2} \left\{ \frac{5}{4} M_\pi^2 \log\left(\frac{M_\pi^2}{\mu^2}\right) - \frac{1}{2} M_K^2 \log\left(\frac{M_K^2}{\mu^2}\right) - \frac{3}{4} M_\eta^2 \log\left(\frac{M_\eta^2}{\mu^2}\right) + c_1 [M_K^2 - M_\pi^2] \right\}$$

$$M_{\eta}^2 = \frac{1}{3} (4M_K^2 - M_{\pi}^2), \quad c_0 = \frac{1}{(4\pi F)^2}, \quad c_1 = 128\pi^2 L_5(\mu)$$

2 fit-parameters, flavor-symmetry constraint satisfied





polynomial forms

expansion in

- M_{π}^2 : dominant behaviour
- $[M_K^2 M_\pi^2]$: flavor-symmetry constraint

3-, 4-, or 6-parameter fit:

$$\begin{aligned} \frac{f_K}{f_\pi} &= 1 + [M_K^2 - M_\pi^2] \Big(c_0^{3\text{-par}} + c_1^{3\text{-par}} [M_K^2 - M_\pi^2] + c_2^{3\text{-par}} M_\pi^2 \Big) \\ \frac{f_K}{f_\pi} &= 1 + [M_K^2 - M_\pi^2] \Big(c_0^{4\text{-par}} + c_1^{4\text{-par}} [M_K^2 - M_\pi^2] + c_2^{4\text{-par}} M_\pi^2 + c_3^{4\text{-par}} M_\pi^4 \Big) \\ \frac{f_K}{f_\pi} &= 1 + [M_K^2 - M_\pi^2] \Big(c_0^{6\text{-par}} + c_1^{6\text{-par}} [M_K^2 - M_\pi^2] + c_2^{6\text{-par}} M_\pi^2 + c_3^{6\text{-par}} M_\pi^4 \\ &+ c_4^{6\text{-par}} M_\pi^2 [M_K^2 - M_\pi^2] + c_5^{6\text{-par}} [M_K^2 - M_\pi^2]^2 \Big) \end{aligned}$$





real world leptonic decay constant ratio

• measured in experiment: ratio of **charged** leptonic decay constants with 6 quark flavors (u,d,s,c,b,t)

$$rac{f_{K^{\pm}}}{f_{\pi^{\pm}}}$$

- suppressed in $N_f=2+1$: dynamical charm, bottom, top, but $(m_{
 m charm}/m_{
 m strange})^2~\simeq~140~\dots$
- (hard) isospin-breaking $m_{\rm up} \neq m_{\rm down}$, EM-effects (charges)

not taken into account in this lattice QCD simulation

• ChPT-analysis of these effects [GASSER, LEUTWYLER (1985); CIRIGLIANO, NEUFELD (2011)]

$$rac{f_{K^{\pm}}}{f_{\pi^{\pm}}} \;=\; rac{f_{K}}{f_{\pi}} \, \sqrt{1 \;+\; \delta_{\mathrm{SU}(2)}}$$

• $\delta_{SU(2)}$: -0.0043(12) [C.,N. (2011)] and FLAG-II, -0.0078(7) [DIVITIIS ET AL. (2012)]

$$\delta_{
m SU(2)} \;=\; -0.0061(61) \;
ightarrow \; rac{f_K \pm}{f_\pi \pm} \;=\; 1.178(10)_{
m stat}(26)_{
m syst}$$







CKM-matrix elements V_{ud} , V_{us}

$$\frac{V_{\rm us}}{V_{\rm ud}} \frac{f_{K^{\pm}}}{f_{\pi^{\pm}}} = 0.27599(38)_{\rm exp}$$

[Moulson (2014), Rosner et al. (2015)]

$$rac{V_{
m us}}{V_{
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• using super-allowed nuclear β -decay $V_{ud} = 0.97417(21)_{nuc}$ [HARDY, TOWNER (2015)]

 $V_{\rm us} = 0.2282(19)_{\rm stat}(51)_{\rm syst}(03)_{\rm exp\&nuc}$

• first-row unitarity ($|V_{ub}| = 4.12(37)(06) \cdot 10^{-3}$ [ROSNER ET AL. (2015)])

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