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# Importance of closed quark loops for lattice QCD studies of tetraquarks

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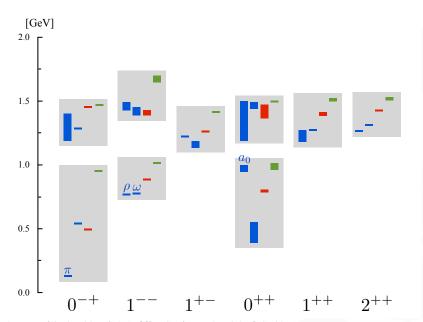




#### Outline

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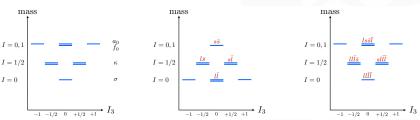
# Light mesons



#### Tetraquark interpretation

Nonet of light scalar mesons  $(J^P = 0^+)$  still poorly understood

- I=1 (two  $u/\bar{d}$  quarks) states  $(a_0,f_0)$  are heavier than the I=1/2  $(u/\bar{d}+s$  quark) states  $(\kappa)$
- $\,\blacksquare\,$  Tetraquark interpretation resolves the mass ordering of the  $0^+$  sector naturally
  - $a_0 \equiv \underline{u} s \bar{s} \bar{d}$  and  $\kappa \equiv \underline{u} (\underline{u} \bar{u} + d \bar{d}) \bar{s}$
  - $a_0(980) \longrightarrow K\bar{K}[\bar{s}u][\bar{d}s] \& a_0(980) \longrightarrow \eta_s\pi[\bar{s}s][\bar{d}u]$



experimental results

conventional  $qar{q}$  pairs

tetraquark interpretation

## Approach

Study of effective masses from mesonic two-quark and four-quark operators.

- Information about possible stable states around threshold
- Composition of states from the solution of the generalized eigenvalue problem
- Relies on large operator basis, in particular 2 meson states

#### Gauge configurations:

- 2+1 dyamical clover fermions and Iwasaki gauge action
- generated by the PACS-CS Collaboration

S. Aoki et al. [PACS-CS Collaboration], Phys. Rev. D 79 (2009) 034503 [arXiv:0807.1661 [hep-lat]].

- Lattice:  $32^3 \times 64$ ,  $a \approx 0.09$ fm
- ho pprox 500 configurations at  $M_\pi pprox 300 {
  m MeV}$

# Operator basis

In our study: 6 operators with the quantum numbers of  $a_0(980)$ .

$$\mathcal{O}^{qar{q}} = \sum_{\mathbf{x}} \left( ar{d}_{\mathbf{x}} \mathbf{u}_{\mathbf{x}} 
ight)$$

$$\mathcal{O}^{Kar{K},\; \mathsf{point}} = \sum_{\mathbf{x}} \left( ar{s}_{\mathbf{x}} \gamma_5 \mathbf{u}_{\mathbf{x}} 
ight) \left( ar{d}_{\mathbf{x}} \gamma_5 s_{\mathbf{x}} 
ight)$$

$$\mathcal{O}^{\eta_s\pi,\; \mathsf{point}} \;\; = \sum_{\mathbf{x}} \left( ar{s}_{\mathbf{x}} \gamma_5 s_{\mathbf{x}} 
ight) \left( ar{d}_{\mathbf{x}} \gamma_5 rac{\mathbf{u}_{\mathbf{x}}}{\mathbf{u}_{\mathbf{x}}} 
ight)$$

$$\mathcal{O}^{Q\bar{Q}} = \sum_{\mathbf{x}} \epsilon_{abc} \left( \bar{s}_{\mathbf{x},b} (C\gamma_5) \bar{d}_{\mathbf{x},c}^T \right) \epsilon_{ade} \left( \mathbf{u}_{\mathbf{x},d}^T (C\gamma_5) \mathbf{s}_{\mathbf{x},e} \right)$$

$$\mathcal{O}^{Kar{K},\; ext{2-part}} = \sum_{\mathbf{x},\mathbf{y}} \Big(ar{s}_{\mathbf{x}} \gamma_5 \mathbf{u}_{\mathbf{x}} \Big) \Big(ar{d}_{\mathbf{y}} \gamma_5 s_{\mathbf{y}} \Big)$$

$$\mathcal{O}^{\eta_s\pi,\; ext{2-part}} = \sum_{\mathbf{x},\mathbf{y}} \Big( ar{s}_{\mathbf{x}} \gamma_5 s_{\mathbf{x}} \Big) \Big( ar{d}_{\mathbf{y}} \gamma_5 rac{\mathbf{u}_{\mathbf{y}}}{\mathbf{v}} \Big)$$









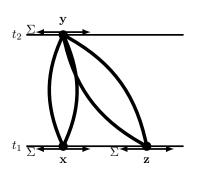


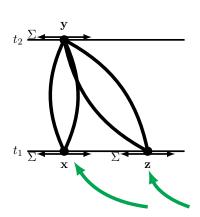


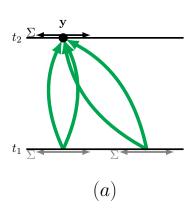
# $C_{jk}(t)$

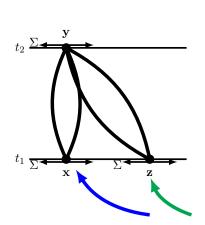
$$C_{jk} = \langle \mathcal{O}_j \mathcal{O}_k^{\dagger} \rangle$$

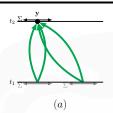
	$\mathcal{O}^{qar{q}^{\dagger}}$	$\mathcal{O}_{point}^{Kar{K}^{\dagger}}$	$\mathcal{O}_{point}^{\eta_s\pi}$ †	$\mathcal{O}^{Qar{Q}^\dagger}$	$\mathcal{O}_{2part}^{Kar{K}^{\dagger}}$	$\mathcal{O}_{2part}^{\eta_s\pi}$ †
$\mathcal{O}^{qar{q}}$	- 2 1	- L O	* D	- 1	-	+ 💆
$\mathcal{O}_{point}^{Kar{K}}$	- 2 <u>C</u>	+ <u>                                    </u>	- <u>O</u> + O	- 1 0 0	· D - S	- 1
$\mathcal{O}_{point}^{\eta_s\pi}$	+ (1)	- 🔯 + 💆	+ Q - Q	- 🔯 + 💆	- 1 - 2	+ 1 - 1
$\mathcal{O}^{Qar{Q}}$	- ( <u>)</u>	- 🔯 - 💆	- <u>O</u> + O	+ <u> </u>	- 1 - 2	- 1 - 1
$\mathcal{O}_{2part}^{Kar{K}}$	- 2	- 7	+ 1	- 7 + 7	-	- 🗘 +
$\mathcal{O}_{2part}^{\eta_s\pi}$	+ 2 1	- +	+ 00		- 📉 + 💆	- 0

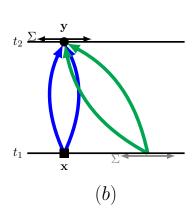


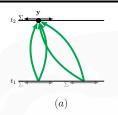


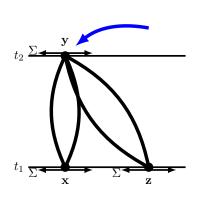


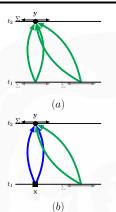


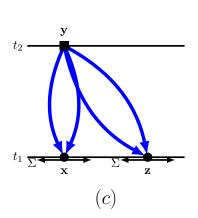


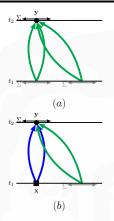


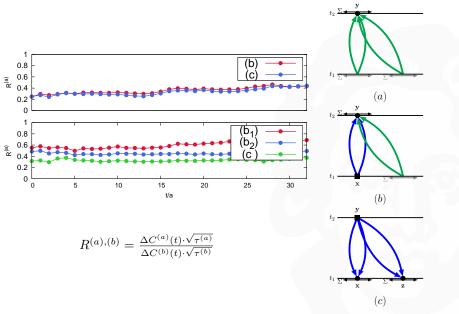






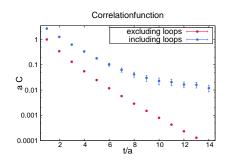


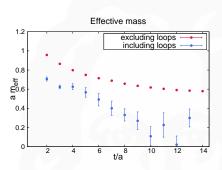




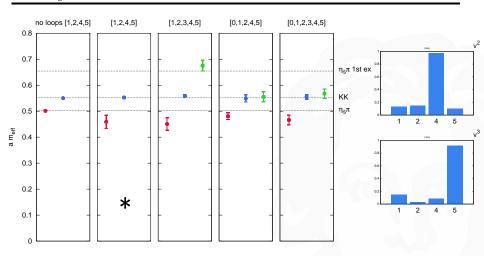
# Evidence for the relevance of closed quark loops

closed fermion loops not only required to include  $\mathcal{O}^{q\bar{q}}$  to the operator basis

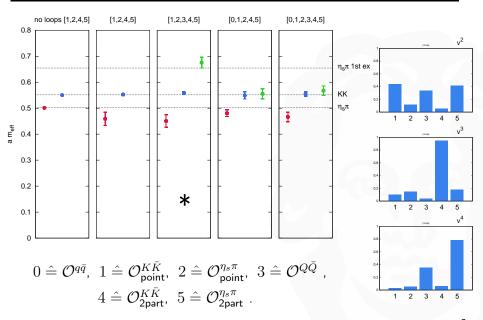


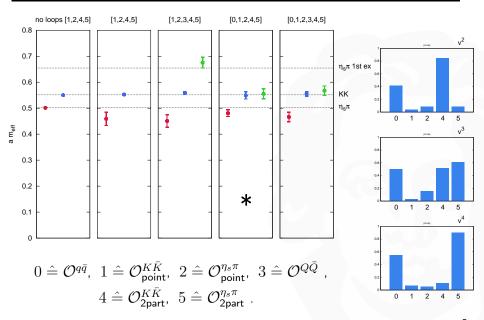


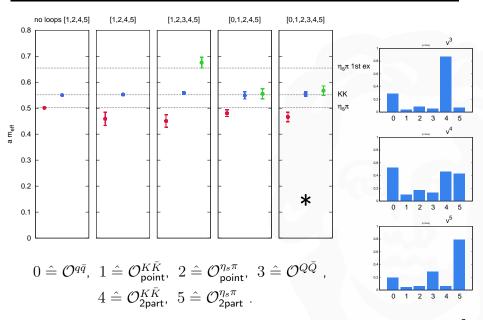
- several matrix elements experience distinct changes in their characteristics ( not only an addition of stochastic noise )
  - ullet e.g. the correlationfunction for  $\mathcal{O}^{Qar{Q}}$



$$\begin{split} 0 & \stackrel{\frown}{=} \mathcal{O}^{q\bar{q}}, \ 1 \stackrel{\frown}{=} \mathcal{O}^{K\bar{K}}_{\mathsf{point}}, \ 2 \stackrel{\frown}{=} \mathcal{O}^{\eta_s\pi}_{\mathsf{point}}, \ 3 \stackrel{\frown}{=} \mathcal{O}^{Q\bar{Q}}_{} \ , \\ 4 \stackrel{\frown}{=} \mathcal{O}^{K\bar{K}}_{\mathsf{2part}}, \ 5 \stackrel{\frown}{=} \mathcal{O}^{\eta_s\pi}_{\mathsf{2part}} \ . \end{split}$$







#### Summary

- Study of effective masses from mesonic two-quark and four-quark operators.
- Investigation of methods and combination of methods to find the optimal strategy to compute each diagram of the correlation matrix
- Computation of closed fermion loops expensive, but essential
- Analysis of states around the two particle threshold reveals **evidence for** an additional state  $(a_0$ , likely of  $q\bar{q}$  structure)