

B_c decays from highly improved staggered quarks and NRQCD

Andrew Lytle (HPQCD Collaboration)
University of Glasgow

Lattice 2016
Southampton, UK
27.07.16

Intro & Motivation

- Obtain $|V_{cb}|$ from $b \rightarrow c$ transitions in semileptonic decays.
- Treatment of c and especially b quarks challenging in lattice simulations due to lattice artifacts which grow as $(am_q)^n$.
- We use two complementary approaches:
 - ▶ Highly improved relativistic action at small a , extrapolate $m_h \rightarrow m_b$.
 - ▶ Improved non-relativistic formalism (NRQCD) at m_b .
- First study:
 - ▶ $B_c \rightarrow \eta_c$
 - ▶ $B_c \rightarrow J/\psi$ [accessible at LHCb]
- More precise $b \rightarrow c$ currents used in $B \rightarrow D, B \rightarrow D^*$.

Outline

1. Intro & Motivation.
2. Calculation Framework.
 - ▶ HISQ action.
 - ▶ Improved NRQCD.
3. Semileptonic Decays.
 - ▶ Correlation functions.
 - ▶ $B_c \rightarrow \eta_c$ and results.
 - ▶ $B_c \rightarrow J/\psi$ and results.
4. Discussion & Future Work.

DiRAC II computing

Computations carried out on the Darwin cluster at Cambridge.

Includes:

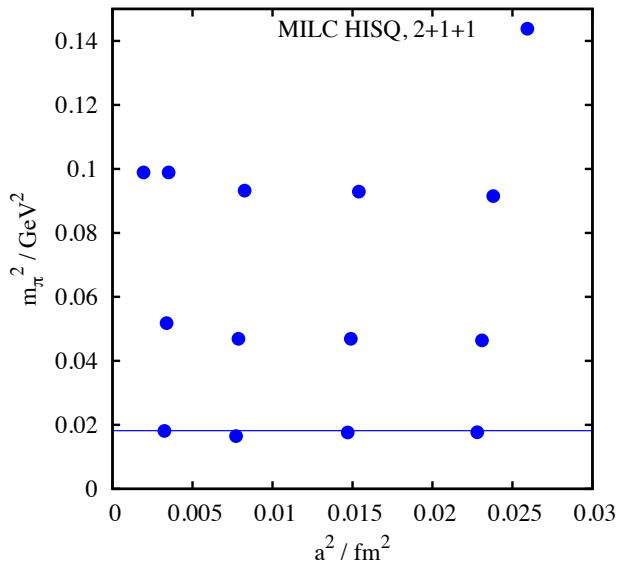
- 9600 Intel Sandy Bridge cores
- 2.6 GHz, 4 GB RAM/core
- 2 PB storage



Part of STFC's HPC facility for theoretical particle physics and astronomy.

- HISQ fermion action.
- Symanzik-improved gauge action, takes into account $\mathcal{O}(N_f \alpha_s a^2)$ effects of HISQ quarks in sea. [0812.0503]
- Multiple lattice spacings down to ~ 0.045 fm.
- Effects of u/d , s , and c quarks in the sea.
- Multiple light-quark input parameters down to physical pion mass.
 - ▶ Chiral fits.
 - ▶ Reduce statistical errors.

MILC ensemble parameters



Heavy quark propagators are calculated using a non-relativistic formalism.

Improved Non-relativistic QCD action

- Accurate through $\mathcal{O}(\alpha_s v^4)$.
- Discretisation corrections through $\mathcal{O}(\alpha_s v^2 a^2 p^2)$.
- $v^2 \sim 0.1$ bottomonium, ~ 0.3 charmonium.
- $am > 1 \rightarrow b$ quarks on $a = 0.15 - 0.06$ fm (down to $m_b/2$ on $a = 0.15$ fm).

Propagators constructed via an evolution equation,

$$G(\mathbf{x}, t + a) = e^{-aH_{\text{eff}}} G(\mathbf{x}, t) .$$

$$aH_{\text{NRQCD}} = aH_0 + a\delta H$$

$$aH_0 = -\frac{\Delta^{(2)}}{2am_b}$$

$$\begin{aligned} a\delta H = & -c_1 \frac{(\Delta^{(2)})^2}{8(am_b)^3} + c_2 \frac{i}{8(am_b)^2} (\nabla \cdot \mathbf{E} - \mathbf{E} \cdot \nabla) \\ & - c_3 \frac{1}{8(am_b)^2} \sigma \cdot (\nabla \times \mathbf{E} - \mathbf{E} \times \nabla) \\ & - c_4 \frac{1}{2am_b} \sigma \cdot \mathbf{B} + c_5 \frac{\Delta^{(4)}}{24am_b} \\ & - c_6 \frac{(\Delta^{(2)})^2}{16n(am_b)^2} \end{aligned}$$

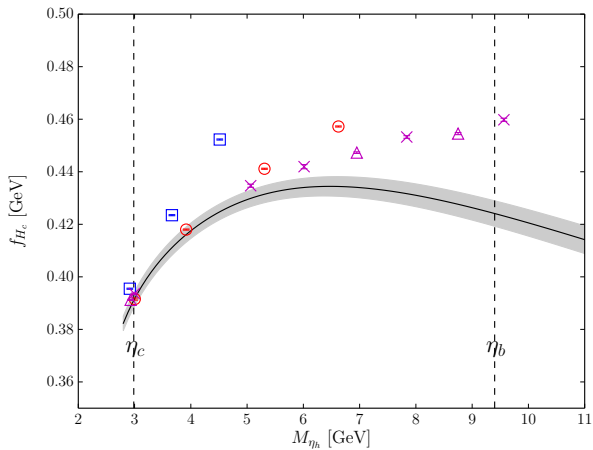
General strategy

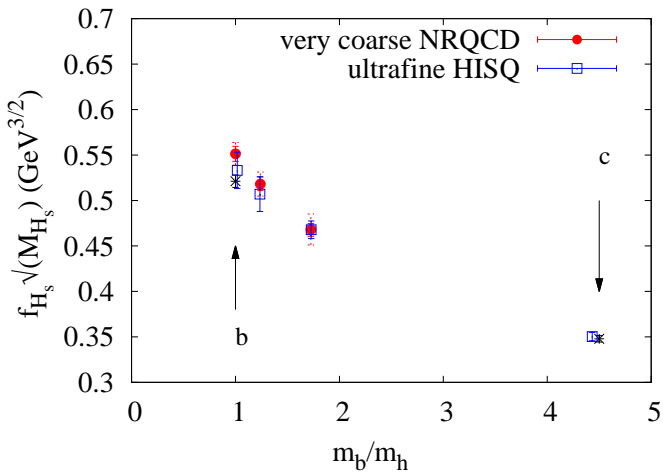
Strategy pursued by HPQCD collaboration:

- Staggered quarks \rightarrow small a , physical pions, multiple lattice spacings..
- Highly improved action
 - \rightarrow discretisation effects under control at m_c
 - \rightarrow reduced taste-splittings
 - \rightarrow physical point ensembles with dynamical u/d , s , and c quarks.
- Compute heavy quarks using (improved) NRQCD.

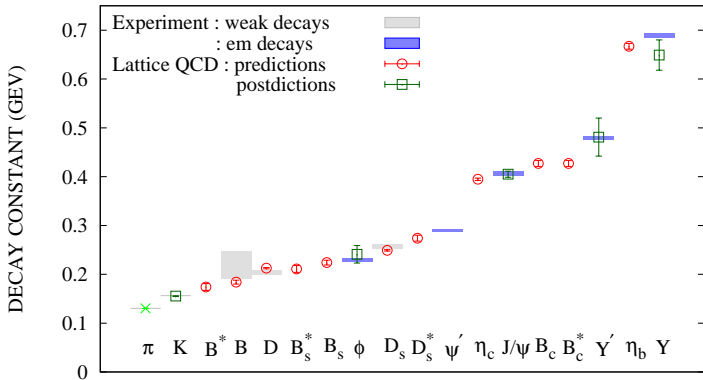
These two approaches are complementary. Ideally there is a range of overlap in applicability to check the approaches are mutually consistent.

f_{H_c} from HISQ.





Decay constants – summary plot.



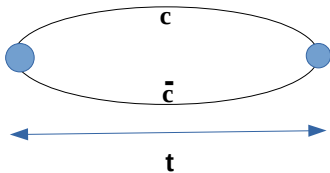
B_c semileptonic decays

Semileptonic decays

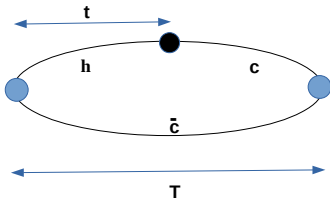
- Study of $B_c \rightarrow \eta_c$, $B_c \rightarrow J/\psi$ decay matrix elements.
- We work in the frame where the B_c is at rest.
- The form factors which parametrise the matrix elements are functions of q^2 , where q is the four-momentum transferred to the leptons.
 - ▶ $q_{\max}^2 = (M - m)^2$, zero recoil of decay hadron.
 - ▶ $q^2 = 0$, maximum recoil of decay hadron.
- Matrix elements are determined by simultaneous fitting of three-point and two-point functions.

Semileptonic decays – meson correlators

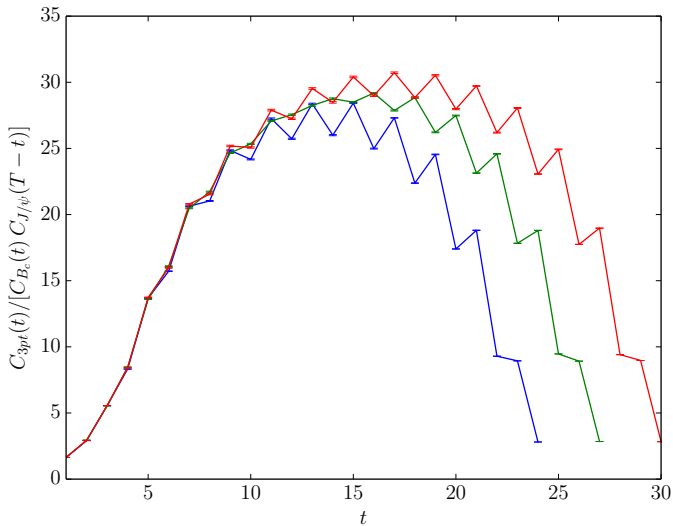
Two-point functions:



Three-point functions:



Semileptonic decays



$$Z \langle \eta_c(p) | V^\mu | B_c(P) \rangle = f_+(q^2) \left[P^\mu + p^\mu - \frac{M^2 - m^2}{q^2} q^\mu \right] + f_0(q^2) \frac{M^2 - m^2}{q^2} q^\mu,$$

From PCVC,

$$\langle \eta_c(p) | S | B_c(P) \rangle = \frac{M^2 - m^2}{m_{b0} - m_{c0}} f_0(q^2)$$

Find Z by calculating both matrix elements at q_{\max}^2 .

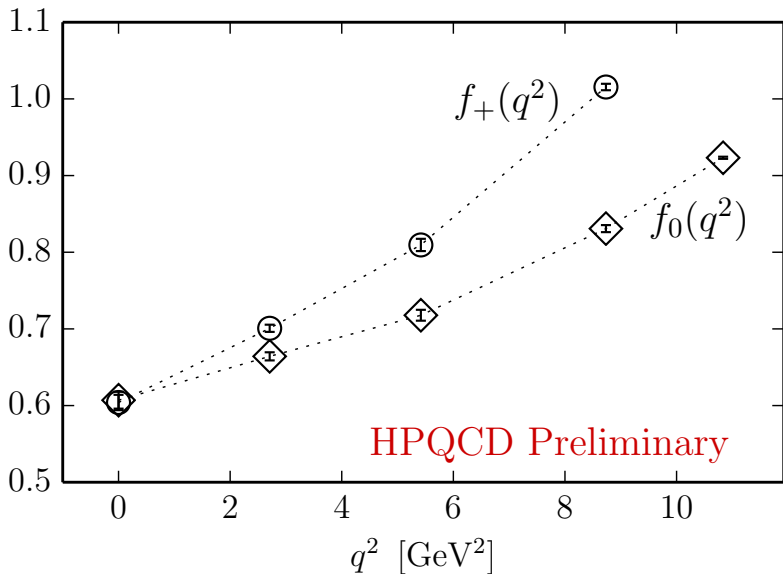
f_0 and f_+ are determined in the NRQCD formalism from matrix elements of the vector current $\langle V_\mu^{\text{nrqcd}} \rangle$, where

$$V_0^{\text{nrqcd}} = (1 + \alpha_s z_0^{(0)}) \left[V_0^{(0)} + (1 + \alpha_s z_0^{(1)}) V_0^{(1)} + \alpha_s z_0^{(2)} V_0^{(2)} \right]$$

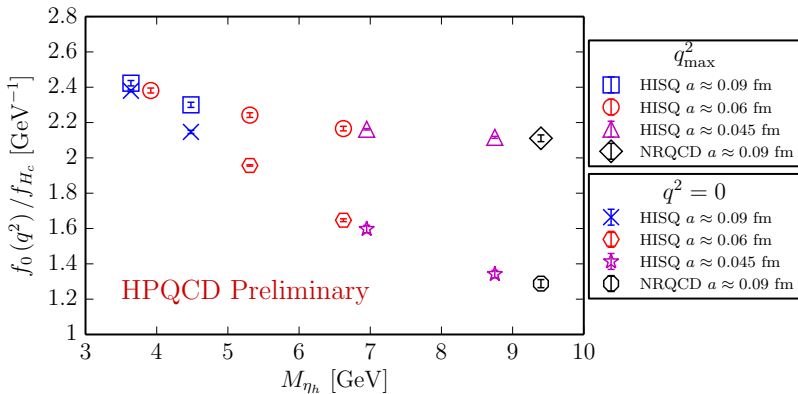
$$V_k^{\text{nrqcd}} = (1 + \alpha_s z_k^{(0)}) \left[V_k^{(0)} + (1 + \alpha_s z_k^{(1)}) V_k^{(1)} + \alpha_s z_k^{(2)} V_k^{(2)} + \alpha_s z_k^{(3)} V_k^{(3)} + \alpha_s z_k^{(4)} V_k^{(4)} \right].$$

One goal of the present work is to constrain the coefficients entering V_μ^{nrqcd} using fully relativistic HISQ data.

NRQCD form factors.



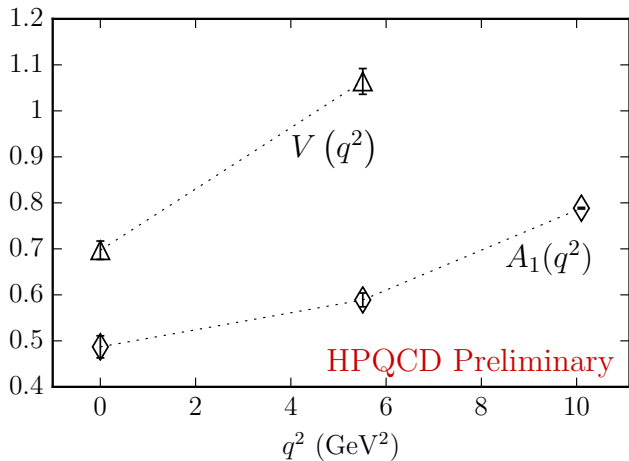
f_0 from HISQ.



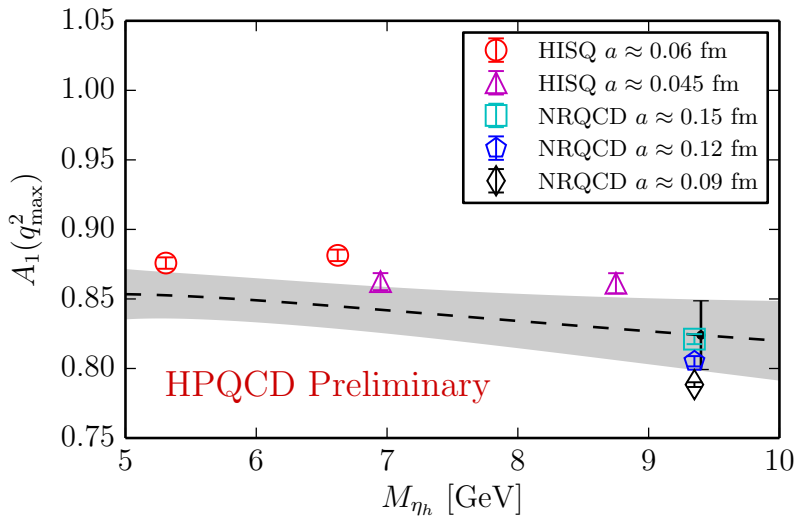
$B_c \rightarrow J/\psi$

$$\begin{aligned} \langle J/\psi(p, \varepsilon) | V^\mu - A^\mu | B_c(P) \rangle = & \\ & \frac{2i\epsilon^{\mu\nu\rho\sigma}}{M+m} \varepsilon_\nu^* p_\rho P_\sigma V(q^2) - (M+m) \varepsilon^{*\mu} A_1(q^2) + \\ & \frac{\varepsilon^* \cdot q}{M+m} (p+P)^\mu A_2(q^2) + 2m \frac{\varepsilon^* \cdot q}{q^2} q^\mu A_3(q^2) - 2m \frac{\varepsilon^* \cdot q}{q^2} q^\mu A_0(q^2) \end{aligned}$$

$B_c \rightarrow J/\psi$.



A_1 from HISQ.



Summary

- A promising approach to study of $b \rightarrow c$ transitions:
 - ▶ Lattice NRQCD with HISQ quarks, plus
 - ▶ Fully relativistic formulation, extrapolate m_h to m_b .
- Proof-of-principle demonstrated for f_0 .
 - ▶ Controlled calculation over full q^2 range.
 - ▶ Good agreement seen with NRQCD results.
- Outputs:
 - ▶ B_c to $J/\Psi \rightarrow$ new possible determination of $|V_{cb}|$.
 - ▶ Improved understanding of NRQCD currents feeds into additional calculations (B to D , B to D^* , ...).

Thank you!