Imperial College London

Charged Lepton Flavour Violation —COMET & PRISM—

25 September 2015 Yoshi Uchida PPAP Community Meeting at Imperial

The First PPAP Community Meeting

The experimental intense muon physics session at the 2009 meeting:

16:00	LHC Upgrades (LHCb + GPDs) 15' Speaker: Chris Parkes (Glasgow) Material: Slides
16:15	e+e- machines (SuperB, Belle upgrade, tau- Speaker: Adrian Bevan (QMUL) Material: Slides
16:30	 NA62 + kaons at JPARC/FNAL 10' Speaker: Cristina Lazzeroni (Birmingham) Material: Slides in more information in the second s
16:40	COMET + muons at FNAL 10' Speaker: Yoshi Uchida (Imperial College) Material: Slides 🔂
Discu	ssion <i>50'</i>
end	

nity Meeting on physics

-charm factories) 15'

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-		

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charm factories) 15'

Nuclear Capture of Mesons and the Meson Decay

B. PONTECORVO National Research Council, Chalk River Laboratory, Chalk River Ontario, Canada June 21, 1947

...Returning to the actual decay of the meson, an experiment suggests itself which might answer the following question: Is the electron emitted by the meson with a mean life of about 2.2 microseconds accompanied by a photon of about 50 Mev? This experiment is being attempted at the present time, since it is felt that the available analysis¹⁰ of the soft component in equilibrium with its primary meson component is probably insufficient to decide definitely whether the meson decays into either an electron plus neutral particle(s) or electron plus photon.

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1945



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1947



Charged Lepton Flavour Violation

90% C.L. upper limit on branching ratios Yoshi.Uchida@imperial.ac.uk



Muon CLFV: Global Activities

- MEG
 MEG Upgrade
 Mu2e
 COMET
 DeeMee
 Mu3e
- PRISM

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Lepton Flavour Conservation in the Standard Model

- In the Standard Model:
 - Lepton flavour is conserved absolutely not through a fundamental principle, but through the choice of fields • an accidental symmetry
- Deviations from the SM can introduce **Lepton Flavour Violation**

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Charged Lepton Flavour Violation Beyond-the-Standard Model Physics can cause CLFV

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• but this is **GIM-suppressed**:

 $B(\mu \to e + \gamma) = \frac{3\alpha}{32\pi} \left| \sum_{\ell} V^*_{\mu\ell} V_{e\ell} \frac{\Delta m^2_{\nu_{\ell}}}{m^2_W} \right|$

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• but this is **GIM-suppressed**: $B(\mu \to e + \gamma) \sim 10^{-54} \times \frac{\sin^2 2\theta_{13}}{2}$

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• but this is **GIM-suppressed**: $B(\mu \to e + \gamma) \sim 10^{-54}$

 if CLFV seen, unambiguous signal for new physics • without such cancellations, CLFV signal can be much larger

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 $\sim rac{m_{\mu}}{30m_{\oplus}}$

(beyond Dirac $m_{\nu} > 0$)

Charged Lepton Flavour Violation

90% C.L. upper limit on branching ratios Yoshi.Uchida@imperial.ac.uk



MEG HOME



Switzerland **PSI, ETH-Z**







MEG Collaboration

some 65 Physicists 5 Countries, 14 Institutes

Italy INFN + Univ. : Pisa, Genova, Pavia, Roma I & Lecce



USA

University of **California** Irvine UCI



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Russia **BINP**, Novosibirsk, JiNR, Dubna

Japan

Univ.Tokyo, KEK Waseda Univ., Kyushu Univ.



Muon-to-Electron Conversion Search for the process $\mu^- + N(A, Z) \rightarrow e^- + N(A, Z)$ muonic atom mono-energetic electron $(\boldsymbol{E_e} \leq 105 \text{ MeV})$

 Time available after formation of muonic atom: up to about 1 microsecond (Z-dependent)

• $E_e = m_\mu$ $-E_{\text{bind}} - E_{\text{recoil}}$

 observed signal is smeared because of detector effects

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Charged Lepton Flavour Violation

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"MELC" Search for Lepton-Flavor-Violating Rare Muon Processes proposal R. M. Djilkibaev* and V. M. Lobashev** from Institute for Nuclear Research, Russian Academy of Sciences, 1980s pr. Shestidesyatiletiya Oktyabrya 7a, Moscow, 117312 Russia

Received March 26, 2010; in final form, July 12, 2010

Abstract—A new approach to seeking three lepton-flavor-violating rare muon processes ($\mu \rightarrow e$ conversion, $\mu \to e + \gamma$, and $\mu \to 3e$) on the basis of a single experimental facility is proposed. This approach makes it possible to improve the sensitivity level of relevant experiments by factors of 10^5 , 600, and 300 for, respectively, the first, the second, and the third of the above processes in relation to the existing experimental level. The approach is based on employing a pulsed proton beam and on combining a muon source and the detector part of the facility into a unified magnetic system featuring a nonuniform field. A new detector design involving separate units and making it possible to study all three muonic processes at a single facility that admits a simple rearrangement of the detectors used is discussed.



Fig. 1. Central horizontal cut of the MELC facility: (1) proton target, (2) superconductor solenoid, (3) shield of the solenoid, (4) steel yoke, (5) transport solenoid and collimator, (6) detector target, (7) coordinate detector, (8) calorimeter, and (9) detector shield and beam trap.







COMET Phase-I

 COMET is the first experiment of its type to be built strong desire to study pion/muon production in detail Decision in 2012 to build and run the first 90 degrees of the beam line

 to study beam characteristics and perform CLFV physics 1/100th of the sensitivity of the full COMET configuration

> **Detector systems:** Cylindrical Drift Chamber, **Triggering Hodoscopes** Straw Tracker, ECAL

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MUSIC Facility at Osaka

Pion capture solenoid Max. Bsol: 3.5 T

-Muon transport solenoid (36deg Bsol: 2.0 T Bdipole: 0.04 T

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The COMET Collaboration

- Belarus
- Canada
- China
- Czech Republic
- France
- Georgia
- India
- Japan
- Malaysia
- Russia
- Saudi Arabia
- South Korea
- United Kingdom
- Vietnam

14 countries, 32 institutes, about 170 collaborators











Recent new participation from Germany (Kai Zuber)

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Markovská starovská starovska st

COMET Construction Status

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Phase-I Drift Chamber Status Four prototypes built, beam-tested All Front-end Boards built and undergoing testing (IHEP) Full detector ready Summer 2016 Wires about 60% strung ALCONTRA . to be completed later this year • Triggering algorithm being developed in UK



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Phase-I Straw Tracker Status

 Straw Tubes produced at JINR (based on NA62) • all 2,500 Phase-I straws delivered (20 µm walls) • Phase-II straws (thinner 12 µm design) being developed Testing and enclosure development underway



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Proton Beam Beam extinction (critical for μ -e conversion) reduced to O(10⁻¹²) at 8 GeV Graphite target for Phase-I, Tungsten for Phase-II (UK concept) Shielding designed for 56 kW (Phase-II specs)

Septum magnets delivered

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Magnets aligned

π/μ Transport Solenoid Experimental hall completed March 2015, and first curved solenoid installed

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π/μ Transport Solenoid

Undergoing extensive testing this year

- Close R&D partnership with manufacturer (Toshiba)
 - proprietary technology to superimpose a vertical dipole field on curved solenoid
 - ongoing iterative testing and feedback with Toshiba
- Gain experience towards rapid production of Phase-II solenoids

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year facturer (Toshiba) oose a vertical

Trigger/Fast Control/DAQ UK-designed Fast Control/Trigger/DAQ systems to be

- used across all subdetectors
- Prototype custom "FCT" boards and FC7 boards (developed here for CMS upgrade) distributed to/purchased by Japan, China, India, Russia, South Korea
 - interfaces with subdetectors being developed
- Performing radiation tests at Brunel and ISIS
 - Russian interface with FCT

Power for pre-trigger and mezzanin boards

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FC7 Testing in Korea FC7 under test

Software

- Decision made in 2013 to build dedicated software framework ICEDUST
 - based on T2K ND280 software, re-organised from the bottom up
 - truly UK-led effort, coordinating work across the collaboration
 - highly active participation in software effort from other nations
 - has become a very substantial and high-quality software suite

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Timeline

 J-PARC radiation accident (May 2013) COMET, being under construction, impacted directly

- J-PARC/KEK Stakeholders (all experiments) agreement to prioritise funding for COMET budget profile for "baseline" funding scenario shown below
 - more optimistic scenarios move data-taking forward

 \Rightarrow Data-taking in 2018 or 2019

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Agreed COMET Budget Profile (including B-Line construction)

COMET Budget Profile

Phase-I to Phase-II

- Rapid transition from Phase-I to Phase-II (about two years)
 - same pion capture solenoid
 - same detector solenoid, with extension
 - re-use of Straw Tracker components
 re-use of ECAL
- Main new component: electron spectrometer solenoid
 - build on Phase-I experience; superconducting cable can be produced in advance
- UK-led Phase-II physics studies ongoing
 - to be reflected in Phase-I running and analysis
 - sensitivity improvements

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COMET & Mu2E Comparison

- COMET has tunable, vertical dipole fields to select particle charge and momentum
 - "C-shaped" solenoids v "S-shaped" solenoids
 - muon beam momentum selection
 - electron spectrometer for background rejection
- Different beam power (pion production solenoid) design):
 - 56 kW for COMET, 8 kW for Mu2E
 - 3 years v 1 year data for full sensitivity
- "Phased approach" for COMET
 - to study partial beam line while preparing full experiment
- Also planning PRISM to allow precision measurements of CLFV

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Schedule of COMET Phase-I and Phase-II

	JFY	2014	2015	2016	2017	2018	2019	2020	2021	2022			
COMET Phase-I	construction												
	data taking												
COMET Phase-II	construction												
	data taking												

COMET Phase-I: 2018 ~ S.E.S. ~ 3x10⁻¹⁵ (for 110 days with 3.2 kW proton beam)

COMET Phase-II: 2021~ S.E.S. ~ 3x10⁻¹⁷ (for 2×10^7 sec with 56 kW proton beam)

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AlCap Support experiment at PSI to measure particle emissions when muons stop in Al Charged (Runs 1 & 3) Run-1: Setup and neutral emissions Vacuum Pump Joint COMET/Mu2E Lead Shielding effort



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Detector 11

Ben Krikler: bek07@imperial.ac.uk

UK involvement in all AlCap runs; publications

PRISM

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PRISM - Phase Rotated Intense Slow Muon beam



UK Timeline

- Participation since 2007; before the start of COMET Contributed to approval of Phase-I construction in 2012
- In 2013, had Subproject Manager roles in:
 - Trigger DAQ hardware/firmware
 - DAQ Readout
 - Software
 - Proton beam target
 - (Muon monitor detector)
- Inaugural Collaboration Board Chair
- Phase-I recommended by PPRP for funding in 2013 Science Board: "proposal fits well within the STFC Roadmap and that the Particle Physics Advisory Panel's 2012 report had concluded there was a strong science case"
- Science Board decision to not fund

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UK Status Reduced UK role significantly

- cannot reduce further
- Closely coordinating with other countries to cover **UK** responsibilities
 - UK-designed trigger hardware and readout, proton targetry etc, being built by other groups
- Current UK positions in collaboration:
 - (new) Collaboration Board Chair
 - Physics Analysis Coordinator
 - Software Coordinator
- UK GridPP expertise, working with CC-IN2P3

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UK Status



David Britton, University of Glasgow

PPAP

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LHC and More



UK Status

- Reduced UK role significantly as requested by STFC
 - cannot reduce further
- Closely coordinating with other countries to cover **UK** responsibilities
 - UK-designed trigger hardware and readout, proton targetry etc, being built by other groups
- Current UK positions in collaboration:
 - Collaboration Board Chair
 - Physics Analysis Coordinator
 - Software Coordinator
- UK GridPP expertise, working with CC-IN2P3
- Continued AlCap/MUSIC exploitation
 - Participation in AlCap runs in 2013 and 2015
- Leadership of PRISM Task Force

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The Current **State of** Muon Flavour Violation Searches

90% C.L. upper limit on branching ratios Yoshi.Uchida@imperial.ac.uk



The Mu3e Collaboration:

- beam and target (PSI)
- solenoidal magnet (PI-HD)
- pixel detector (PI-HD, KIT, Ma
- scintillating fibre detector (ETHZ, PSI, UniGe, UZH)
- scintillating tile detector (KIP-HD)
- detector readout and filter farm (Mainz)
- mechanics and cooling (PSI, PI-HD)
- experimental infrastructure (PSI)
- slow control (PSI)

		outer pixel layers
muon beam	inner pixe	Hayers
		17 cm
		¥



Future Muon Flavour Violation Searches

90% C.L. upper limit on branching ratios Yoshi.Uchida@imperial.ac.uk



Experimental **Limits on Lepton Flavour Violation**

90% C.L. upper limits on branching ratios

Reaction

$$\mu^{+} \rightarrow e^{+}\gamma$$

$$\mu^{+} \rightarrow e^{+}e^{+}e^{-}$$

$$\mu^{-}Ti \rightarrow e^{-}Ti$$

$$\mu^{-}Au \rightarrow e^{-}Au$$

$$\mu^{+}e^{-} \rightarrow \mu^{-}e^{+}$$

$$\tau \rightarrow e\gamma$$

$$\tau \rightarrow \mu\gamma$$

$$\tau \rightarrow \mu\mu\mu$$

$$\tau \rightarrow eee$$

$$\pi^{0} \rightarrow \mu e$$

$$K_{L}^{0} \rightarrow \mu e$$

$$K_{L}^{0} \rightarrow \mu e$$

$$K_{L}^{0} \rightarrow \pi^{0}\mu^{+}e^{-}$$

$$Z_{L}^{0} \rightarrow \tau e$$

$$Z_{L}^{0} \rightarrow \tau \mu$$

Present limit $< 5.7 \times 10^{-13}$ $< 1.0 \times 10^{-12}$ $< 6.1 \times 10^{-13}$ $< 7 \times 10^{-13}$ $e^{-}Au$ $< 8.3 \times 10^{-11}$ $< 3.3 \times 10^{-8}$ $< 4.4 \times 10^{-8}$ $< 2.1 \times 10^{-8}$ $< 2.7 \times 10^{-8}$ $< 3.8 \times 10^{-10}$ $< 4.7 \times 10^{-12}$ $< 1.3 \times 10^{-11}$ $< 7.6 \times 10^{-11}$ $< 1.7 \times 10^{-6}$ $< 9.8 \times 10^{-6}$ $< 1.2 \times 10^{-5}$

 $^{+}\mu^{+}e^{-}$ μ^+e^-

Complementarity of CLFV Channels

- The various CLFV channels complement each other
 - CLFV muons and tau decays
 - different muon-to-electron conversion nuclei
 - CLFV decays of heavier particles
 - mesons, vector bosons, Higgses....
- The relative rates can tell us the nature of the BSM physics

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Complementarity of CLFV Channels

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Relative dependences of the muonto-electron conversion branching ratio on the target nucleus

For different nuclei, different size of nucleus, radius of orbit, u- and d-quark composition



Conclusions

- An exciting time for muon CLFV experimentation
 - a thriving global effort to probe beyond the SM
- COMET on track to perform world's first new muonto-electron conversion experiment since SINDRUM-II (which last took data in 2000)
 - First example of novel, pulsed-beam measurement
 - Also considering other BSM physics channels
 - Phase-II to follow Phase-I within about two years
 - significant opportunities for UK leadership
 - based on Phase-I experience
 - if STFC agrees CLFV is worth it after all experimental design allows for further improvement
- PRISM presents unique future possibility within muon CLFV of a factor 1,000,000 improvement





PHYSICAL REVIEW D 74, 116002 (2006)

Lepton flavor violation from supersymmetric grand unified theories: Where do we stand for MEG, PRISM/PRIME, and a super flavor factory

L. Calibbi,¹ A. Faccia,¹ A. Masiero,¹ and S. K. Vempati^{2,3}

(Received 29 August 2006; published 13 December 2006)

We analyze the complementarity between lepton flavor violation (LFV) and LHC experiments in probing the supersymmetric (SUSY) grand unified theories (GUT) when neutrinos get a mass via the seesaw mechanism. Our analysis is performed in an SO(10) framework, where at least one neutrino Yukawa coupling is necessarily as large as the top Yukawa coupling. Our study thoroughly takes into account the whole renormalization group running, including the GUT and the right-handed neutrino mass scales, as well as the running of the observable neutrino spectrum. We find that the upcoming (MEG, SuperKEKB) and future (PRISM/PRIME, super flavor factory) LFV experiments will be able to test such SUSY framework for SUSY masses to be explored at the LHC and, in some cases, even beyond the LHC sensitivity reach.

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LEPTON FLAVOR VIOLATION FROM SUPERSYMMETRIC ...



 $\mu \rightarrow e$ in Ti as a probe of SUSY-GUT scenarios. The plots are obtained by scanning the LHC accessible parameter space. The horizontal lines are the present (SINDRUM II) bound and the planned (PRISM/PRIME) sensitivity to the process. We see that PRIME would be able to severely constrain the low tan β , low mixing angles case and to completely test the other scenarios.

PHYSICAL REVIEW D 74, 116002 (2006)

Profumo, Vempoti, and Yaguna H



Interplay of LFV and slepton mass splittings at the LHC as a probe of the SUSY seesaw

A. Abada^a, A. J. R. Figueiredo^b, J. C. Romão^b and A. M. Teixeira^c

We study the impact of a type-I SUSY seesaw concerning lepton flavour violation (LFV) both at low-energies and at the LHC. The study of the di-lepton invariant mass distribution at the LHC allows to reconstruct some of the masses of the different sparticles involved in a decay chain. In particular, the combination with other observables renders feasible the reconstruction of the masses of the intermediate sleptons involved in $\chi_2^0 \to \tilde{\ell} \,\ell \to \ell \,\ell \,\chi_1^0$ decays. Slepton mass splittings can be either interpreted as a signal of non-universality in the SUSY soft breakingterms (signalling a deviation from constrained scenarios as the cMSSM) or as being due to the violation of lepton flavour. In the latter case, in addition to these high-energy processes, one expects further low-energy manifestations of LFV such as radiative and three-body lepton decays. Under the assumption of a type-I seesaw as the source of neutrino masses and mixings, all these LFV observables are related. Working in the framework of the cMSSM extended by three right-handed neutrino superfields, we conduct a systematic analysis addressing the simultaneous implications of the SUSY seesaw for both high- and low-energy lepton flavour violation. We discuss how the confrontation of slepton mass splittings as observed at the LHC and low-energy LFV observables may provide important information about the underlying mechanism of LFV.

JHEP 1010:104,2010



One of the most interesting results region in parameter space associated with a $\tilde{e}_L - \tilde{\mu}_L$ mass splitting $\sim \mathcal{O}(1\%)$ is also within the future sensitivity of low-energy facilities, especially for $CR(\mu - e, Ti)$ (even without the expected) upgrade to $\mathcal{O}(10^{-18})$ for PRISM/PRIME)⁵. Also, any $\tilde{e}_L - \tilde{\mu}_L$ mass splitting above 4% would also be associated with a $\mu \to e\gamma$ signal within MEG sensitivity. A similar situation (albeit not so striking) is observed for $\tilde{\mu}_L - \tilde{\tau}_2$ mass differences: as an example, mass splittings above 3%, 4% and 6% would be associated to low-energy signals of LFV within PRISM/PRIME, SuperB, and MEG reach, respectively.

JHEP 1010:104,2010

consists in the fact that almost the entire

Probing the Randall-Sundrum geometric origin of flavor with lepton flavor violation

Kaustubh Agashe,¹ Andrew E. Blechman,² and Frank Petriello³

¹Department of Physics, Syracuse University, Syracuse, New York 13244, USA and School of Natural Sciences, Institute for Advanced Study, Princeton, New Jersey 08540, USA ²Department of Physics, The Johns Hopkins University, Baltimore, Maryland 21218, USA ³Department of Physics, University of Wisconsin, Madison, Wisconsin 53706, USA (Received 23 June 2006; published 28 September 2006)

The anarchic Randall-Sundrum model of flavor is a low energy solution to both the electroweak hierarchy and flavor problems. Such models have a warped, compact extra dimension with the standard model fermions and gauge bosons living in the bulk, and the Higgs living on or near the TeV brane. In this paper we consider bounds on these models set by lepton flavor-violation constraints. We find that loopinduced decays of the form $l \rightarrow l' \gamma$ are ultraviolet sensitive and incalculable when the Higgs field is localized on a four-dimensional brane; this drawback does not occur when the Higgs field propagates in the full five-dimensional space-time. We find constraints at the few TeV level throughout the natural range of parameters, arising from $\mu - e$ conversion in the presence of nuclei, rare μ decays, and rare τ decays. A tension exists between loop-induced dipole decays such as $\mu \rightarrow e\gamma$ and tree-level processes such as $\mu - e$ conversion; they have opposite dependences on the five-dimensional Yukawa couplings, making it difficult to decouple flavor-violating effects. We emphasize the importance of the future experiments MEG and PRIME. These experiments will definitively test the Randall-Sundrum geometric origin of hierarchies in the lepton sector at the TeV scale.

DOI: 10.1103/PhysRevD.74.053011

PACS numbers: 13.35.-r, 11.10.Kk

Probing the Randall-Sundrum geometric origin of flavor with lepton flavor violation

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PROBING THE RANDALL-SUNDRUM GEOMETRIC ...



Scan of the $\mu \rightarrow e\gamma$ and $\mu - e$ conversion predictions for $M_{\rm KK} = 3, 5, 10$ TeV and $\nu = 0$. The solid line FIG. 6 denotes the PDG bound on $BR(\mu \rightarrow e\gamma)$, while the dashed lines indicate the SINDRUM II limit on $\mu - e$ conversion and the projected MEG sensitivity to $BR(\mu \rightarrow e\gamma)$.

PHYSICAL REVIEW D 74, 053011 (2006)



Charged lepton flavour violation and $(g-2)_{\mu}$ in the Littlest Higgs model with T-Parity: a clear distinction from Supersymmetry

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ABSTRACT: We calculate the rates for the charged lepton flavour violating decays $\ell_i \to \ell_j \gamma$, $\tau \to \ell \pi, \tau \to \ell \eta, \tau \to \ell \eta', \mu^- \to e^- e^+ e^-$, the six three body leptonic decays $\tau^- \to \ell_i^- \ell_j^+ \ell_k^-$ and the rate for $\mu - e$ conversion in nuclei in the Littlest Higgs model with T-parity (LHT). We also calculate the rates for $K_{L,S} \to \mu e$, $K_{L,S} \to \pi^0 \mu e$ and $B_{d,s} \to \ell_i \ell_j$. We find that the relative effects of mirror leptons in these transitions are by many orders of magnitude larger than analogous mirror quark effects in rare K and B decays analyzed recently. In particular, in order to suppress the $\mu \to e\gamma$ and $\mu^- \to e^- e^+ e^-$ decay rates and the $\mu - e$ conversion rate below the experimental upper bounds, the relevant mixing matrix in the mirror lepton sector $V_{H\ell}$ must be rather hierarchical, unless the spectrum of mirror leptons is quasi-degenerate. We find that the pattern of the LFV branching ratios in the LHT model differs significantly from the one encountered in the MSSM, allowing in a transparent manner to distinguish these two models with the help of LFV processes. We also calculate $(g-2)_{\mu}$ and find the new contributions to a_{μ} below $1 \cdot 10^{-10}$ and consequently negligible. We compare our results with those present in the literature.

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Charged lepton flavour violation and $(g-2)_{\mu}$ in the Littlest Higgs model with T-Parity: a clear distinction from Supersymmetry

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ABSTRACT: We calculate the rates for the charged lepton flavour violating decays $\ell_i \to \ell_j \gamma$, $\tau \to \ell \pi, \tau \to \ell \eta, \tau \to \ell \eta', \mu^- \to e^- e^+ e^-$, the six three body leptonic decays $\tau^- \to \ell_i^- \ell_j^+ \ell_k^-$ and the rate for $\mu - e$ conversion in nuclei in the Littlest Higgs model with T-parity (LHT). We also calculate the rates for $K_{L,S} \to \mu e$, $K_{L,S} \to \pi^0 \mu e$ and $B_{d,s} \to \ell_i \ell_j$. We find that the relative effects of mirror leptons in these transitions are by many orders of magnitude larger than analogous mirror quark effects in rare K and B decays analyzed recently. In particular, in order to suppress the $\mu \to e\gamma$ and $\mu^- \to e^- e^+ e^-$ decay rates and the $\mu - e$ conversion rate below the experimental upper bounds, the relevant mixing matrix in the mirror lepton sector $V_{H\ell}$ must be rather hierarchical, unless the spectrum of mirror leptons is quasi-degenerate. We find that the pattern of the LFV branching ratios in the LHT model differs significantly from the one encountered in the MSSM, allowing in a transparent manner to distinguish these two models with the help of LFV processes. We also calculate $(g-2)_{\mu}$ and find the new contributions to a_{μ} below $1 \cdot 10^{-10}$ and consequently negligible. We compare our results with those present in the literature.

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LFV in tau and muon decays within SUSY seesaw S. Antusch et al. / Nuclear Physics B (Proc. Suppl.) 169 (2007) 155–165



Correlation between BR($\mu \to e \gamma$) and BR($\tau \to \mu \gamma$) as a function of m_{N_3} , for SPS 1a.

Flavor physics at large tan β with a binolike lightest supersymmetric particle

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The minimal supersymmetric extension of the standard model with large $\tan\beta$ and heavy squarks $(M_{\tilde{q}} \gtrsim 1 \text{ TeV})$ is a theoretically well-motivated and phenomenologically interesting extension of the standard model. This scenario naturally satisfies all the electroweak precision constraints and, in the case of not too heavy slepton sector ($M_{\tilde{\ell}} \leq 0.5$ TeV), can also easily accommodate the $(g-2)_{\mu}$ anomaly. Within this framework nonstandard effects could possibly be detected in the near future in a few lowenergy flavor violating observables, such as $\mathcal{B}(B \to \tau \nu)$, $\mathcal{B}(B_{s,d} \to \ell^+ \ell^-)$, $\mathcal{B}(B \to X_s \gamma)$, and $\mathcal{B}(\mu \to \tau \nu)$ $e\gamma$). Interpreting the $(g-2)_{\mu}$ anomaly as the first hint of this scenario, we analyze the correlations of these low-energy observables under the additional assumption that the relic density of a binolike lightest supersymmetric particle accommodates the observed dark-matter distribution.

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MSSM with Large $\tan\beta$ & Heavy Squarks

- Accommodates g-2 signal
- Red points satisfy **B-physics** constraints
- Predicted $\mu \rightarrow e +$ γ at 10-12
 - further factor O(100S) for $\mu + N$ $\rightarrow e + N$



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Isidori et al Phys Rev D75 (2007) 1150

"DNA" of flavour physics effects for the most interesting observables in a selection of SUSY and non-SUSY models $\star \star \star \star$ signals large effects, $\star \star \star$ visible but small effects and \star implies that the given model does not predict sizable effects in that observable.

	AC	RVV2	AKM	δLL	FBMSSM	LHT	RS
$D^0 - \overline{D}^0$	***	*	*	*	*	***	?
ϵ_K	\star	$\star\star\star$	$\star\star\star$	\star	\star	**	$\star\star\star$
$S_{\psi\phi}$	$\star\star\star$	$\star\star\star$	$\star\star\star$	\star	\star	$\star\star\star$	$\star\star\star$
$S_{\phi K_S}$	$\star \star \star$	**	\star	$\star\star\star$	$\star\star\star$	\star	?
$A_{\rm CP}(B \to X_s \gamma)$	\star	\star	\star	$\star\star\star$	$\star\star\star$	\star	?
$A_{7,8}(B\to K^*\mu^+\mu^-)$	\star	\star	\star	$\star\star\star$	$\star\star\star$	$\star\star$?
$A_9(B \to K^* \mu^+ \mu^-)$	\star	\star	\star	\star	\star	\star	?
$B \to K^{(*)} \nu \overline{\nu}$	\star	\star	\star	\star	\star	\star	\star
$B_s \to \mu^+ \mu^-$	$\star\star\star$	$\star\star\star$	$\star\star\star$	$\star \star \star$	$\star\star\star$	\star	\star
$K^+ \to \pi^+ \nu \bar{\nu}$	\star	\star	\star	\star	\star	$\star\star\star$	$\star\star\star$
$K_L \to \pi^0 \nu \bar{\nu}$	\star	\star	\star	\star	\star	$\star\star\star$	$\star\star\star$
$\mu \rightarrow e \gamma$	$\star \star \star$	$\star\star\star$	$\star\star\star$	$\star\star\star$	$\star\star\star$	$\star\star\star$	$\star\star\star$
$\tau \to \mu \gamma$	$\star \star \star$	$\star\star\star$	\star	$\star\star\star$	$\star\star\star$	$\star\star\star$	$\star\star\star$
$\mu + N \rightarrow e + N$	$\star\star\star$	$\star\star\star$	$\star\star\star$	$\star\star\star$	$\star\star\star$	$\star\star\star$	$\star\star\star$
d_n	***	***	***	**	***	\star	***
d_e	$\star\star\star$	$\star \star \star$	**	\star	$\star\star\star$	\star	$\star\star\star$
$(g - 2)_{\mu}$	$\star\star\star$	$\star\star\star$	$\star\star$	$\star\star\star$	$\star\star\star$	\star	?

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new physics →	AC	RVV2	AKM	δLL	FBMSSM	LHT	RS
$D^0 - \overline{D}^0$	***	*	*	*	\star	***	?
ϵ_K	\star	$\star\star\star$	$\star\star\star$	\star	\star	$\star\star$	$\star\star\star$
$S_{\psi\phi}$	$\star\star\star$	$\star\star\star$	$\star\star\star$	\star	\star	$\star\star\star$	$\star\star\star$
$S_{\phi K_S}$	$\star\star\star$	**	\star	$\star\star\star$	$\star\star\star$	\star	?
$A_{\rm CP}(B \to X_s \gamma)$	\star	\star	\star	$\star\star\star$	$\star\star\star$	\star	?
$A_{7,8}(B\to K^*\mu^+\mu^-)$	\star	\star	\star	$\star\star\star$	$\star\star\star$	$\star\star$?
$A_9(B\to K^*\mu^+\mu^-)$	\star	\star	\star	\star	\star	\star	?
$B \to K^{(*)} \nu \bar{\nu}$	\star	\star	\star	\star	\star	\star	\star
$B_s \to \mu^+ \mu^-$	$\star\star\star$	$\star\star\star$	$\star\star\star$	$\star\star\star$	$\star\star\star$	\star	\star
$K^+ \to \pi^+ \nu \bar{\nu}$	\star	\star	\star	\star	\star	$\star\star\star$	$\star\star\star$
$K_L \to \pi^0 \nu \bar{\nu}$	\star	\star	\star	\star	\star	$\star \star \star$	$\star\star\star$
$\mu \rightarrow e \gamma$	$\star\star\star$	$\star \star \star$	$\star \star \star$	$\star \star \star$	$\star\star\star$	$\star \star \star$	$\star\star\star$
$ au o \mu \gamma$	$\star\star\star$	$\star\star\star$	\star	$\star\star\star$	$\star\star\star$	$\star\star\star$	$\star\star\star$
$\mu + N \rightarrow e + N$	$\star\star\star$	$\star\star\star$	$\star\star\star$	$\star\star\star$	$\star\star\star$	$\star\star\star$	$\star\star\star$
d_n	***	***	***	**	***	\star	***
d_e	$\star\star\star$	$\star\star\star$	$\star\star$	\star	$\star\star\star$	\star	$\star\star\star$
$(g-2)\mu$	$\star\star\star$	$\star\star\star$	**	$\star\star\star$	$\star\star\star$	\star	?

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bservables	AC	RVV2	AKM	δLL	FBMSSM	LHT	RS
$D^0 - \overline{D}^0$	***	*	*	*	*	***	?
ϵ_K	\star	$\star\star\star$	$\star\star\star$	\star	\star	**	$\star\star\star$
$S_{\psi\phi}$	$\star\star\star$	$\star\star\star$	$\star\star\star$	\star	\star	$\star\star\star$	$\star\star\star$
$S_{\phi K_S}$	$\star \star \star$	**	\star	$\star\star\star$	$\star\star\star$	\star	?
$A_{\rm CP}(B \to X_s \gamma)$	\star	\star	\star	$\star\star\star$	$\star\star\star$	\star	?
$A_{7,8}(B\to K^*\mu^+\mu^-)$	\star	\star	\star	$\star\star\star$	$\star\star\star$	$\star\star$?
$A_9(B\to K^*\mu^+\mu^-)$	\star	\star	\star	\star	\star	\star	?
$B \to K^{(*)} \nu \bar{\nu}$	\star	\star	\star	\star	\star	\star	\star
$B_s \to \mu^+ \mu^-$	$\star\star\star$	$\star\star\star$	$\star\star\star$	$\star \star \star$	$\star \star \star$	\star	\star
$K^+ \to \pi^+ \nu \bar{\nu}$	\star	\star	\star	\star	\star	$\star\star\star$	$\star\star\star$
$K_L \to \pi^0 \nu \bar{\nu}$	\star	\star	\star	\star	\star	$\star\star\star$	$\star\star\star$
$\mu \rightarrow e \gamma$	$\star \star \star$	$\star\star\star$	$\star\star\star$	$\star \star \star$	$\star\star\star$	$\star\star\star$	$\star\star\star$
$ au o \mu \gamma$	$\star \star \star$	$\star\star\star$	\star	$\star \star \star$	$\star\star\star$	$\star\star\star$	$\star\star\star$
$\mu + N \rightarrow e + N$	$\star\star\star$	$\star\star\star$	$\star\star\star$	$\star\star\star$	$\star\star\star$	$\star\star\star$	$\star\star\star$
d_n	***	***	$\star \star \star$	**	***	\star	***
d_e	$\star\star\star$	$\star\star\star$	**	\star	$\star\star\star$	\star	$\star\star\star$
$(g - 2)_{\mu}$	$\star\star\star$	$\star\star\star$	**	$\star\star\star$	$\star\star\star$	\star	?

"DNA" of flavour physics effects for the most interesting observables in a selection of SUSY and non-SUSY models \star \star signals large effects, \star \star visible but small effects and \star implies that the given model does not predict sizable effects in that observable.

AC	RVV2	AKM	δLL	
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RS FBMSSM LHT • AC: Abelian model by Agashe and Carone based on a U(1) flavour symmetry • RVV2: the non-Abelian model by Ross, Velasco-Sevilla and • AKM: Antusch, King and Malinsky model based on the flavour symmetry SU(3) • δLL: flavour models predicting pure, CKM-like, left-handed currents • FBMSSM: flavour-blind MSSM • LHT: Littlest Higgs Model with T-Parity • RS: Randall–Sundrum model with custodial protection non-SUSY models SUSY models

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"DNA" of flavour physics effects for the most interesting observables in a selection of SUSY and non-SUSY models $\star \star \star \star$ signals large effects, $\star \star \star$ visible but small effects and \star implies that the given model does not predict sizable effects in that observable.

	AC	RVV2	AKM	δLL	FBMSSM	LHT	RS
$D^0 - \overline{D}^0$	***	*	\star	\star	\star	***	?
ϵ_K	\star	$\star\star\star$	$\star\star\star$	\star	\star	**	$\star\star\star$
$S_{\psi\phi}$	$\star\star\star$	$\star\star\star$	$\star\star\star$	\star	\star	$\star\star\star$	$\star\star\star$
$S_{\phi K_S}$	$\star \star \star$	**	\star	$\star\star\star$	$\star\star\star$	\star	?
$A_{\rm CP}(B\to X_s\gamma)$	\star	\star	\star	$\star\star\star$	$\star\star\star$	\star	?
$A_{7,8}(B\to K^*\mu^+\mu^-)$	\star	\star	\star	$\star\star\star$	$\star\star\star$	**	?
$A_9(B\to K^*\mu^+\mu^-)$	\star	\star	\star	\star	\star	\star	?
$B \to K^{(*)} \nu \overline{\nu}$	\star	\star	\star	\star	\star	\star	\star
$B_s \to \mu^+ \mu^-$	$\star \star \star$	$\star\star\star$	$\star\star\star$	$\star\star\star$	$\star\star\star$	\star	\star
$K^+ \to \pi^+ \nu \bar{\nu}$	\star	\star	\star	\star	\star	$\star\star\star$	$\star\star\star$
$K_L \to \pi^0 \nu \bar{\nu}$	\star	\star	\star	\star	\star	$\star\star\star$	$\star\star\star$
$\mu \rightarrow e \gamma$	$\star\star\star$	$\star\star\star$	$\star\star\star$	$\star\star\star$	$\star\star\star$	$\star\star\star$	$\star\star\star$
$\tau \to \mu \gamma$	$\star\star\star$	$\star\star\star$	\star	$\star\star\star$	$\star\star\star$	$\star\star\star$	$\star\star\star$
$\mu + N \rightarrow e + N$	$\star\star\star$	$\star\star\star$	$\star\star\star$	$\star\star\star$	$\star\star\star$	$\star\star\star$	$\star\star\star$
d_n	$\star \star \star$	$\star\star\star$	$\star\star\star$	**	$\star\star\star$	\star	***
d_e	$\star\star\star$	$\star\star\star$	**	\star	$\star\star\star$	\star	$\star\star\star$
$(g-2)_{\mu}$	$\star\star\star$	***	**	$\star\star\star$	$\star\star\star$	\star	?

W. Altmannshofer et al. / Nuclear Physics B 830 (2010) 17–94