Where do we go from here?

Ben Gripaios

Cambridge

February 2016

A cautionary tale
cf. Observational Geography, c. 1953

Newcastle Journal
North Mail
No. 33,336 TUESDAY JUNE 2 1953 A KEMSLEY NEWSPAPER 2d
6 AM

Britain on top of the world on this Elizabeth II day

EVEREST IS CONQUERED

Queen awakened to hear of climbers' triumph

4-PAGE TV AND RADIO GUIDE INSIDE

COMMENTARY

Queen of all hearts

THIS is a great day. Once more we British make plain to the world that we hold firmly to our way of life.

And, most appropriately, last night brought great news to back the claim that our way usually takes us, in the end, to the top.

At the third attempt, the sixth British expedition, under Col. John Hunt, has conquered Everest. It is the happiest of omens.

Skill, tenacity, courage, and great endurance have overcome the terribly powerful defences of the mightiest mountain on earth.

Right instinct

THE significance of Everest may be clearer to the mass of our people than is the symbolism of the Coronation ceremony.

But they know instinctively that the underlying concern is always for them; for their welfare, their protection, their rights and liberties.

They know, too, something which baffles other peoples; that the Crown has gained in influence far more than has been surrendered in direct power.

Yet, ordinarily, they make no parade whatever of their

2 REACH SUMMIT ON 11th EXPEDITION

THE QUEEN WAS WAKENED AT BUCKINGHAM PALACE LATE LAST NIGHT TO BE TOLD THAT THE BRITISH EXPEDITION HAS CONQUERED MOUNT EVEREST.

This great news, on the eve of the Coronation, reached London last night in a message to "The Times" from Col. John Hunt in charge of the expedition.

MR. E. P. HILLARY
He reached the summit.

The climb was made on Friday and Col. Hunt has reported that "all is well."

The successful assault was made by Mr. E. P. Hillary, a New Zealander, and the Sherpa porter named Tensing Bhutia.

This great feat of the new Elizabethans flashed round the world adding still further joy to the heightening Coronation fever.

Mr. Hillary, aged 34, is a bee-keeper in New Zealand. His climbing experience was gained in the Southern Alps in the South Island, a range that has attracted mountaineers from all over the world because of the difficulty of the climbs.

New route

He was an originator of winter ski mountaineering in New Zealand. During the war he served in the Royal New Zealand Air Force.

He had experience in the Himalayas two years ago, when he was a member of the expedition which, led by Eric Shipton, found a way into the Western Cwm.

This discovery opened up a new route for attacks on Everest, and it has been by following up the pioneering work done by Shipton's and other expeditions of the last two years that success has been achieved this time.

Tensing, the "King" of the Sherpa porters of Mount

1 a.m. AND ROYAL ROUTE BARRIERS ARE CLOSED

CROWD barriers were closed hours before schedule early this morning as tens of thousands of sightseers camped along the Coronation route.

Fifty thousand people jammed in Trafalgar-

CORONATION GOWN

GIRL, 16, STABBED, FRIEND MISSING

AN attractive girl of 16, found dead in the Thames at Richmond yesterday, was murdered and her body flung in the river near Teddington Lock on Sunday night, police decided after a post-mortem examination last night.

There was a gash on her forehead, but death was due to three stab wounds in the back.

Blood-stains were found on the grassy and wooded verge of the river towpath near the lock a mile and a half away.

WENT CYCLING

Police later feared that they had a double murder to contend with.

The dead girl was identified by her mother as Barbara Songhurst, of Princess Road, Teddington. She was an assistant in a chemist's shop.

On Sunday, she went out with her friend, Christine Reed (18), of Roy Crescent, Hampton Hill, for a cycle ride. Christine Reed has now been reported as missing.

Earlier, detectives had been puzzled by two pairs of women's shoes beside the towpath.

Neither cycle has been found.

Duke will watch

1953 - The high energy frontier (of gravitational potential).

1953 - Where do we go from here?

Nowhere.

2016 - Where do we go from here?

2016 - The high energy frontier (of particle physics).

Where is ‘here’?

UNIVERSITY OF CAMBRIDGE

NATURAL SCIENCES TRIPOS **Part IB & II (General)**

Tuesday, 31 May, 2011 9:00 am to 12:00 pm

Exercise 1.1.1.1.1a: Given locality, causality, Lorentz invariance, and known physical data since 1860, show that the Lagrangian describing all observed physical processes (sans gravity) can be written:

$$
\begin{aligned}
&-\tfrac{1}{2}\partial_\nu g^a_\mu \partial_\nu g^a_\mu - g_s f^{abc}\partial_\mu g^a_\nu g^b_\mu g^c_\nu - \tfrac{1}{4}g_s^2 f^{abc}f^{ade}g^b_\mu g^c_\nu g^d_\mu g^e_\nu + \\
&\tfrac{1}{2}ig_s^2(\bar{q}^\sigma_i\gamma^\mu q^\sigma_j)g^a_\mu + \bar{G}^a\partial^2 G^a + g_s f^{abc}\partial_\mu \bar{G}^a G^b g^c_\mu - \partial_\nu W^+_\mu \partial_\nu W^-_\mu - \\
&M^2 W^+_\mu W^-_\mu - \tfrac{1}{2}\partial_\nu Z^0_\mu \partial_\nu Z^0_\mu - \tfrac{1}{2c_w^2}M^2 Z^0_\mu Z^0_\mu - \tfrac{1}{2}\partial_\mu A_\nu \partial_\mu A_\nu - \tfrac{1}{2}\partial_\mu H \partial_\mu H - \\
&\tfrac{1}{2}m_h^2 H^2 - \partial_\mu\phi^+\partial_\mu\phi^- - M^2\phi^+\phi^- - \tfrac{1}{2}\partial_\mu\phi^0\partial_\mu\phi^0 - \tfrac{1}{2c_w^2}M\phi^0\phi^0 - \beta_h[\tfrac{2M^2}{g^2} + \\
&\tfrac{2M}{g}H + \tfrac{1}{2}(H^2 + \phi^0\phi^0 + 2\phi^+\phi^-)] + \tfrac{2M^4}{g^2}\alpha_h - igc_w[\partial_\nu Z^0_\mu(W^+_\mu W^-_\nu - \\
&W^+_\nu W^-_\mu) - Z^0_\nu(W^+_\mu\partial_\nu W^-_\mu - W^-_\mu\partial_\nu W^+_\mu) + Z^0_\mu(W^+_\nu\partial_\nu W^-_\mu - \\
&W^-_\nu\partial_\nu W^+_\mu)] - igs_w[\partial_\nu A_\mu(W^+_\mu W^-_\nu - W^+_\nu W^-_\mu) - A_\nu(W^+_\mu\partial_\nu W^-_\mu - \\
&W^-_\mu\partial_\nu W^+_\mu) + A_\mu(W^+_\nu\partial_\nu W^-_\mu - W^-_\nu\partial_\nu W^+_\mu)] - \tfrac{1}{2}g^2 W^+_\mu W^-_\mu W^+_\nu W^-_\nu + \\
&\tfrac{1}{2}g^2 W^+_\mu W^-_\nu W^+_\mu W^-_\nu + g^2c_w^2(Z^0_\mu W^+_\mu Z^0_\nu W^-_\nu - Z^0_\mu Z^0_\mu W^+_\nu W^-_\nu) + \\
&g^2 s_w^2(A_\mu W^+_\mu A_\nu W^-_\nu - A_\mu A_\mu W^+_\nu W^-_\nu) + g^2 s_w c_w[A_\mu Z^0_\nu(W^+_\mu W^-_\nu - \\
&W^+_\nu W^-_\mu) - 2A_\mu Z^0_\mu W^+_\nu W^-_\nu] - g\alpha[H^3 + H\phi^0\phi^0 + 2H\phi^+\phi^-] - \\
&\tfrac{1}{8}g^2\alpha_h[H^4 + (\phi^0)^4 + 4(\phi^+\phi^-)^2 + 4(\phi^0)^2\phi^+\phi^- + 4H^2\phi^+\phi^- + 2(\phi^0)^2H^2] - \\
&gMW^+_\mu W^-_\mu H - \tfrac{1}{2}g\tfrac{M}{c_w^2}Z^0_\mu Z^0_\mu H - \tfrac{1}{2}ig[W^+_\mu(\phi^0\partial_\mu\phi^- - \phi^-\partial_\mu\phi^0) - \\
&W^-_\mu(\phi^0\partial_\mu\phi^+ - \phi^+\partial_\mu\phi^0)] + \tfrac{1}{2}g[W^+_\mu(H\partial_\mu\phi^- - \phi^-\partial_\mu H) - W^-_\mu(H\partial_\mu\phi^+ - \\
&\phi^+\partial_\mu H)] + \tfrac{1}{2}g\tfrac{1}{c_w}(Z^0_\mu(H\partial_\mu\phi^0 - \phi^0\partial_\mu H) - ig\tfrac{s_w^2}{c_w}MZ^0_\mu(W^+_\mu\phi^- - W^-_\mu\phi^+) +
\end{aligned}
$$

$$
\begin{gathered}
igs_wMA_\mu(W_\mu^+\phi^- - W_\mu^-\phi^+) - ig\frac{1-2c_w^2}{2c_w}Z_\mu^0(\phi^+\partial_\mu\phi^- - \phi^-\partial_\mu\phi^+) + \\
igs_wA_\mu(\phi^+\partial_\mu\phi^- - \phi^-\partial_\mu\phi^+) - \frac{1}{4}g^2W_\mu^+W_\mu^-[H^2 + (\phi^0)^2 + 2\phi^+\phi^-] - \\
\frac{1}{4}g^2\frac{1}{c_w^2}Z_\mu^0Z_\mu^0[H^2 + (\phi^0)^2 + 2(2s_w^2-1)^2\phi^+\phi^-] - \frac{1}{2}g^2\frac{s_w^2}{c_w}Z_\mu^0\phi^0(W_\mu^+\phi^- + \\
W_\mu^-\phi^+) - \frac{1}{2}ig^2\frac{s_w^2}{c_w}Z_\mu^0H(W_\mu^+\phi^- - W_\mu^-\phi^+) + \frac{1}{2}g^2s_wA_\mu\phi^0(W_\mu^+\phi^- + \\
W_\mu^-\phi^+) + \frac{1}{2}ig^2s_wA_\mu H(W_\mu^+\phi^- - W_\mu^-\phi^+) - g^2\frac{s_w}{c_w}(2c_w^2-1)Z_\mu^0A_\mu\phi^+\phi^- - \\
g^1s_w^2A_\mu A_\mu\phi^+\phi^- - \bar{e}^\lambda(\gamma\partial + m_e^\lambda)e^\lambda - \bar{\nu}^\lambda\gamma\partial\nu^\lambda - \bar{u}_j^\lambda(\gamma\partial + m_u^\lambda)u_j^\lambda - \\
\bar{d}_j^\lambda(\gamma\partial + m_d^\lambda)d_j^\lambda + igs_wA_\mu[-(\bar{e}^\lambda\gamma^\mu e^\lambda) + \frac{2}{3}(\bar{u}_j^\lambda\gamma^\mu u_j^\lambda) - \frac{1}{3}(\bar{d}_j^\lambda\gamma^\mu d_j^\lambda)] + \\
\frac{ig}{4c_w}Z_\mu^0[(\bar{\nu}^\lambda\gamma^\mu(1+\gamma^5)\nu^\lambda) + (\bar{e}^\lambda\gamma^\mu(4s_w^2 - 1 - \gamma^5)e^\lambda) + (\bar{u}_j^\lambda\gamma^\mu(\frac{4}{3}s_w^2 - \\
1 - \gamma^5)u_j^\lambda) + (\bar{d}_j^\lambda\gamma^\mu(1 - \frac{8}{3}s_w^2 - \gamma^5)d_j^\lambda)] + \frac{ig}{2\sqrt{2}}W_\mu^+[(\bar{\nu}^\lambda\gamma^\mu(1+\gamma^5)e^\lambda) + \\
(\bar{u}_j^\lambda\gamma^\mu(1+\gamma^5)C_{\lambda\kappa}d_j^\kappa)] + \frac{ig}{2\sqrt{2}}W_\mu^-[(\bar{e}^\lambda\gamma^\mu(1+\gamma^5)\nu^\lambda) + (\bar{d}_j^\kappa C_{\lambda\kappa}^\dagger\gamma^\mu(1 + \\
\gamma^5)u_j^\lambda)] + \frac{ig}{2\sqrt{2}}\frac{m_e^\lambda}{M}[-\phi^+(\bar{\nu}^\lambda(1-\gamma^5)e^\lambda) + \phi^-(\bar{e}^\lambda(1+\gamma^5)\nu^\lambda)] - \\
\frac{g}{2}\frac{m_e^\lambda}{M}[H(\bar{e}^\lambda e^\lambda) + i\phi^0(\bar{e}^\lambda\gamma^5e^\lambda)] + \frac{ig}{2M\sqrt{2}}\phi^+[-m_d^\kappa(\bar{u}_j^\lambda C_{\lambda\kappa}(1-\gamma^5)d_j^\kappa) + \\
m_u^\lambda(\bar{u}_j^\lambda C_{\lambda\kappa}(1+\gamma^5)d_j^\kappa] + \frac{ig}{2M\sqrt{2}}\phi^-[m_d^\lambda(\bar{d}_j^\lambda C_{\lambda\kappa}^\dagger(1+\gamma^5)u_j^\kappa) - m_u^\kappa(\bar{d}_j^\lambda C_{\lambda\kappa}^\dagger(1 - \\
\gamma^5)u_j^\kappa] - \frac{g}{2}\frac{m_u^\lambda}{M}H(\bar{u}_j^\lambda u_j^\lambda) - \frac{g}{2}\frac{m_d^\lambda}{M}H(\bar{d}_j^\lambda d_j^\lambda) + \frac{ig}{2}\frac{m_u^\lambda}{M}\phi^0(\bar{u}_j^\lambda\gamma^5u_j^\lambda) - \\
\frac{ig}{2}\frac{m_d^\lambda}{M}\phi^0(\bar{d}_j^\lambda\gamma^5d_j^\lambda) + \bar{X}^+(\partial^2 - M^2)X^+ + \bar{X}^-(\partial^2 - M^2)X^- + \bar{X}^0(\partial^2 - \\
\frac{M^2}{c_w^2})X^0 + \bar{Y}\partial^2Y + igc_wW_\mu^+(\partial_\mu\bar{X}^0X^- - \partial_\mu\bar{X}^+X^0) + igs_wW_\mu^+(\partial_\mu\bar{Y}X^- - \\
\partial_\mu\bar{X}^+Y) + igc_wW_\mu^-(\partial_\mu\bar{X}^-X^0 - \partial_\mu\bar{X}^0X^+) + igs_wW_\mu^-(\partial_\mu\bar{X}^-Y - \\
\partial_\mu\bar{Y}X^+) + igc_wZ_\mu^0(\partial_\mu\bar{X}^+X^+ - \partial_\mu\bar{X}^-X^-) + igs_wA_\mu(\partial_\mu\bar{X}^+X^+ - \\
\partial_\mu\bar{X}^-X^-) - \frac{1}{2}gM[\bar{X}^+X^+H + \bar{X}^-X^-H + \frac{1}{c_w^2}\bar{X}^0X^0H] + \\
\frac{1-2c_w^2}{2c_w}igM[\bar{X}^+X^0\phi^+ - \bar{X}^-X^0\phi^-] + \frac{1}{2c_w}igM[\bar{X}^0X^-\phi^+ - \bar{X}^0X^+\phi^-] + \\
igMs_w[\bar{X}^0X^-\phi^+ - \bar{X}^0X^+\phi^-] + \frac{1}{2}igM[\bar{X}^+X^+\phi^0 - \bar{X}^-X^-\phi^0]
\end{gathered}
$$

$$\mathscr{L}_{SM}$$

$$\mathscr{L}_{?} = \mathscr{L}_{SM} + \sum \frac{\mathscr{O}_n}{\Lambda^n}$$

- Effects of $\mathscr{O}_n$, $\sim (\frac{E}{\Lambda})^n$.
- What is Λ?

- LHC, all $\mathscr{O}_n$
- LEP & al., $\mathscr{O}_6 = (H^\dagger D_\mu H)^2, \ldots$
- flavour mixing, $\mathscr{O}_6 = (\bar{s}\gamma_\mu d)^2, \ldots$
- proton decay, $\mathscr{O}_6 = qqql, u^c u^c d^c e^c, \ldots$

Probes of generic new physics:

- LHC, $\Lambda \gtrsim$ TeV
- LEP & al., $\Lambda \gtrsim 1 - 10$ TeV
- flavour mixing, $\Lambda \gtrsim 10^{3-5}$ TeV
- proton decay, $\Lambda \gtrsim 10^{13}$ TeV

$\exists$ 1 measurement of Λ:

- ν masses, $\mathscr{O}_5 = (LH)^2$
- $\implies \Lambda \sim 10^{10}$ TeV

This is evidence for, not against, the SM!

Other ‘evidence’ for Λ:

- Dark Energy $\implies \Lambda \sim 10^{-3}$ eV!
- Dark Matter: $\frac{\Delta\Lambda}{\Lambda} \sim 10^{80}$!
- Baryogenesis $\implies \Lambda \lesssim M_P$!

So why are we bothering to look at all?!

$\exists$ 1 troublesome operator

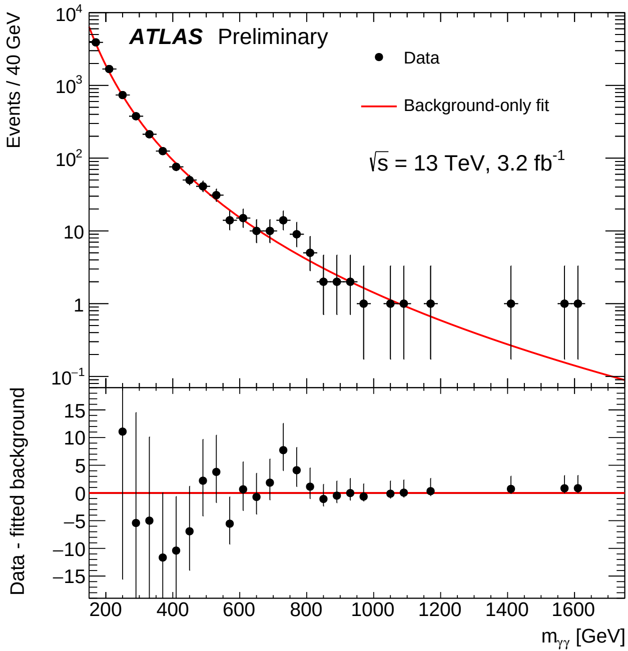
- $\mathscr{O}_2 = H^\dagger H$
- $\mathscr{L} \supset \Lambda^2 H^\dagger H \implies \Lambda \sim 100$ GeV
- naturalness vs. fine-tuning/anthropics ...

Running out of good solutions to this problem!

Anomalies . . .

The thing that went bump in the night . . .

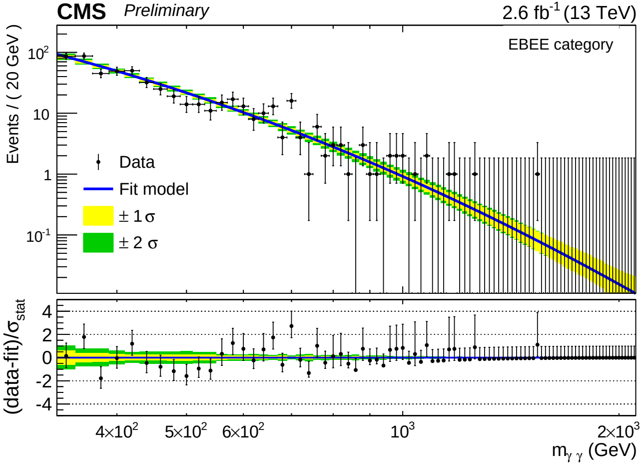
ATLAS 13 TeV 3.2 /fb: 14 events at 750 GeV



ATLAS 3.2/fb: 3.9σ local, 2.3σ global

ATLAS-CONF-2015-081

CMS 2.6/fb: 10 events at 760 GeV



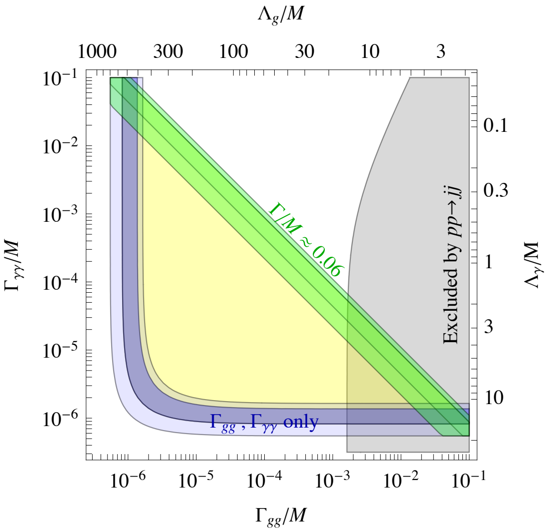
CMS-EXO-15-004

Qualitatively

- Big $\sigma \times BR$
- Excess in 2 bins $\implies$ wide
- $\implies$ strong interactions?
- $\implies$ inconsistent with 8 TeV?
- $\times 5$ pdf gain for 2σ compatibility $\implies$ gg or QQ production modes

Franceschini et al. et al. et al., 1512.04933

$$\frac{\Gamma_{\gamma\gamma}}{M}\frac{\Gamma_{gg}}{M} \approx 1.1 \times 10^{-6}\ \frac{\Gamma}{M} \approx 6 \times 10^{-8}$$



Franceschini et al. et al. et al., 1512.04933

Couplings from EW invariants:

$$\frac{g_3^2}{\Lambda_3}\eta G^{\mu\nu}\tilde{G}_{\mu\nu} + \frac{g_2^2}{\Lambda_2}\eta W^{\mu\nu}\tilde{W}_{\mu\nu} + \frac{g_1^2}{\Lambda_1}\eta B^{\mu\nu}\tilde{B}_{\mu\nu}$$

BB: $\frac{\Gamma(S\to Z\gamma)}{\Gamma(S\to\gamma\gamma)} = 2\tan^2\theta_W \approx 0.6, \qquad \frac{\Gamma(S\to ZZ)}{\Gamma(S\to\gamma\gamma)} = \tan^4\theta_W \approx 0.08.$

$$\frac{\Gamma(S\to WW)}{\Gamma(S\to\gamma\gamma)} = \frac{2}{\sin^4\theta_W} \approx 40,$$

WW: $\frac{\Gamma(S\to ZZ)}{\Gamma(S\to\gamma\gamma)} = \frac{1}{\tan^4\theta_W} \approx 12, \qquad \frac{\Gamma(S\to Z\gamma)}{\Gamma(S\to\gamma\gamma)} = \frac{2}{\tan^2\theta_W} \approx 7.$

Franceschini et al. et al. et al., 1512.04933

A bit premature, but let's imagine it's true . . .

1. Naturalness is back!

‘To do one tuning may be regarded as a misfortune. To do two looks a little careless.’

2. Strong coupling: SUSY unlikely.

3. Composite Higgs.

Composite Higgs overview

Composite Higgs $\equiv$ modern incarnation of natural EWSB via strong dynamics.

A solution to the hierarchy problem that is (almost) literally natural.

A rhetorical question: What if $\nexists$ Higgs?

What if $\nexists$ Higgs?

- An ‘almost perfect’ rendition of EWSB!
- QCD has a natural scale $\sim$ GeV
- Global χSB: $SU(2)_L \times SU(2)_R \to SU(2)_V$
- Gauge $\supset SU(2)_L \times U(1)_Y \to U(1)_{em}$
- But $m_{W,Z} \sim$ GeV :-(

QCD Colour $\rightarrow$ Technicolour

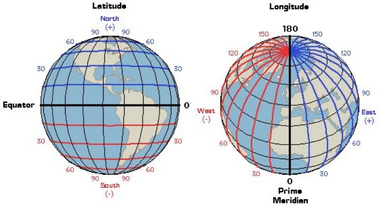
- natural scale $\sim$ 100 GeV
- Global $SU(2)_L \times SU(2)_R \rightarrow SU(2)_V$
- Gauge $\supset SU(2)_L \times U(1)_Y \rightarrow U(1)_{em}$
- A perfect, natural rendition of EWSB
- But no Higgs, flavour, EWPT, ...

Technicolour $\rightarrow$ Composite Higgs

Kaplan & Georgi, 84 ...

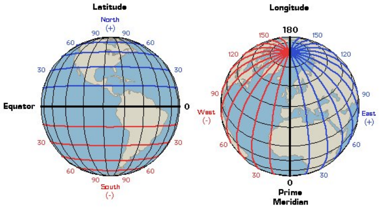
- $SU(2)_L \times SU(2)_R \rightarrow SU(2)_V$ is equivalent to $SO(4) \rightarrow SO(3)$
- Generalize to $SO(n+1) \rightarrow SO(n)\ldots$

$SO(n+1) \to SO(n)$ is in fact rather mundane.



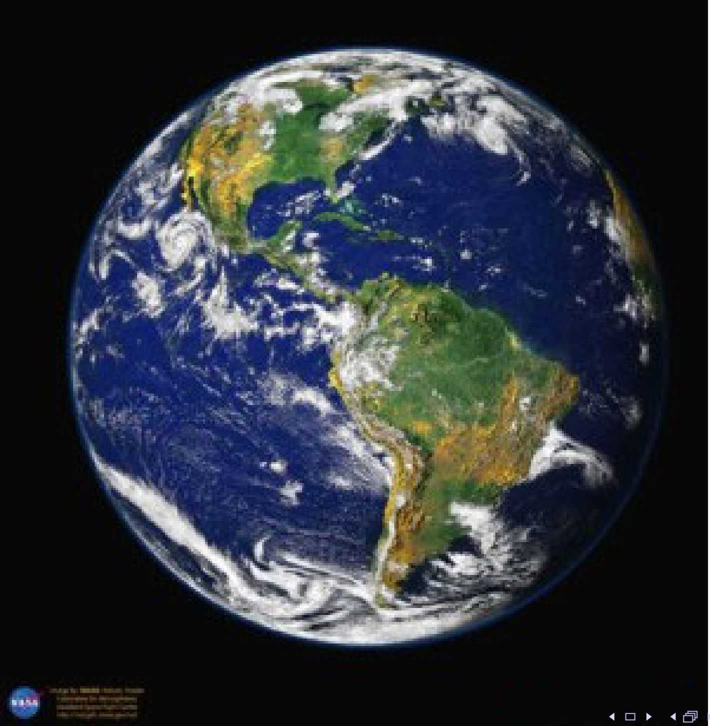
Consider $SO(3) \rightarrow SO(2)$:

- There are 2 Goldstone bosons: latitude and longitude.



Now gauge $SO(2) \subset SO(3) \ldots$

- Goldstone boson $\rightarrow$ pseudo-GB
- Gets potential and coupling to gauge fields
- cf. potential on Earth



Consider $SO(5) \to SO(4)$:

- There are 4 Goldstone bosons: angles of S^4
- they are a $2_{\frac{1}{2}}$ of $SU(2) \times U(1)_Y \subset SO(4)$, viz. the Higgs field, H
- Gauging $SU(2) \times U(1)_Y$ plus coupling to t generates $V(H)$ and $HWW, H\gamma\gamma$ etc
- a.k.a. the Minimal Composite Higgs model

Agashe, Contino, & Pomarol, 0412089

The minimal composite Higgs model

- $\Delta S \propto \theta^2 \implies \gtrsim 10 - 20$ % tuning

Phenomenology of composite Higgs models: bad news

- Departures from SM in e.g. H couplings $\propto \theta^2 \lesssim 10 - 20$ %

  Giudice, Grojean, Pomarol & Rattazzi,0703164

  Falkowski, 0711.0828

  Low, Rattazzi & Vichi, 0907.5413

- Generic resonance masses $\sim (4\pi)v/\theta \gtrsim$ few TeV

Phenomenology of composite Higgs models: good news

- $m_h = 125 GeV \implies$ light, fermionic top partner

  Contino, da Rold & Pomarol, 0612048

- $\exists$ search strategies for these now

  Contino & Servant, 0801.1679

  de Simone & al., 1211.5663

  BMG, Muller, Parker & Sutherland, 1406.5957

- Flat $\mathbb{R}^4$ vs. curved S^4: h self-couplings at high E
- Effects in flavour physics. Tough to call.

But this is true for the minimal model with just a ‘curved Higgs’.
What about the 750 GeV anomaly?

Are extra scalars plausible?
Of course: just change G/H!

e.g. the ‘next-to-minimal’ model based on $SO(6)/SO(5)$ has the Higgs plus 1 scalar.

BMG, A. Pomarol, F. Riva, J. Serra, 0902.1483

$SO(6)/SO(5)$ is unique: $SO(6) \simeq SU(4)$.
Can get this from a chiral (ergo natural) gauge theory.

Another desideratum: extra scalars in composite models have mass $\gtrsim \theta \times m_h$.

$\implies$ Singlet mass is expected to be $\gtrsim$ few 100 GeV.

Another

desiderratum: composite models feature anomalies
cf. $\pi^0 \to \gamma\gamma$ in QCD.

BMG, 0803.0497

Physics agenda

- With $g_{SM} = 0$, the structure is fixed by group theory.
- hh, $h\eta$, $\eta\eta$ couplings fixed.
- With $g_{SM} \neq 0, \eta$ couples to everything (like h)
- couplings to fermions scale like Higgs Yukawas

Physics agenda II

LHC run II

- Confirm excess
- Look for couplings to WW, $Z\gamma$, ZZ; check $SU(2)_L \times U(1)_Y$
- (In general, get ηBB and ηWW.)
- Measure another channel $gg \to gg$ and pin down couplings

Physics agenda III

Future Collider(s)

- Decide on future facility(ies): η-strahlung from e^+e^- at 850 GeV? QQ or gg via pp?
- What $\eta - SM$ couplings can be probed?
- Can we measure hh, $h\eta$, or $\eta\eta$?
- Are top partners/other resonances within reach?
- What about flavour?

Suffice to say plenty to do . . .

But what if it goes away again?

“So many centuries after the Creation, it is unlikely that anyone could find hitherto unknown lands of any value.” Spanish Royal Commision, 1490

“The more important fundamental laws and facts of physical science have all been discovered, and these are now so firmly established that the possibility of their ever being supplanted in consequence of new discoveries is exceedingly remote.... Our future discoveries must be looked for in the sixth place of decimals.” Michelson, 1894