

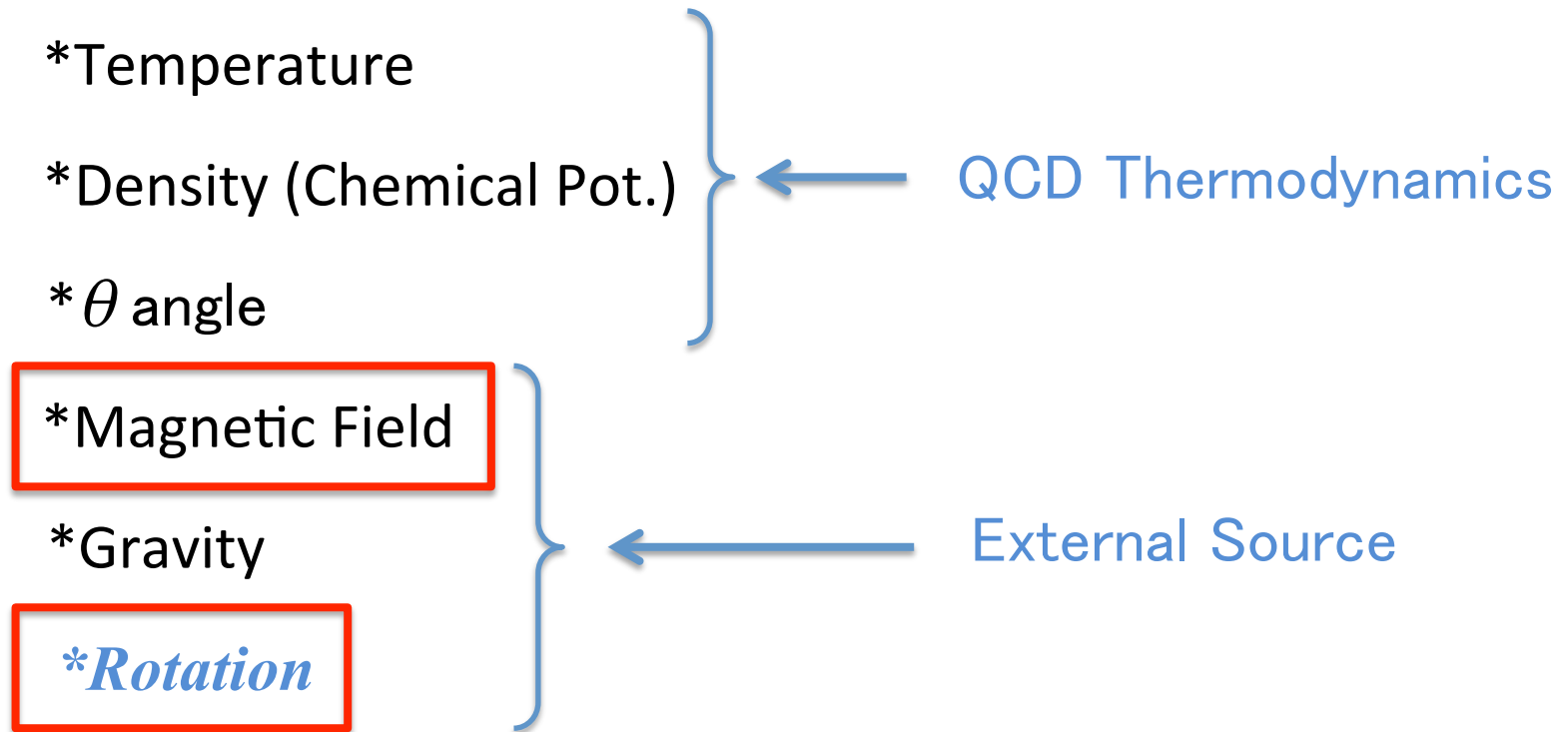
# *Novel Chiral Structure Realized by Rotation*

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Phys.Rev.D93(2016)104052, arXiv:1512.08974

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arXiv:1608.00336

# *QCD with Environmental Effects*



# Relativistic Systems with $eB$

Heavy-ion collisions  $eB \sim 10^{18}$  G

Skokov, Illarionov, Toneev (2009)

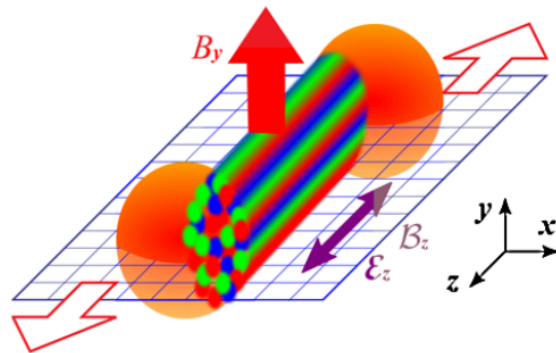
Magnetar(surface)  $eB \sim 10^{15}$  G

Duncan, Thompson (1992)

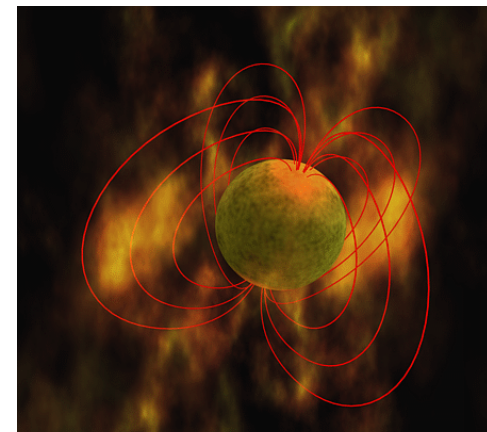
Magnetar(interior)  $eB \sim 10^{18-20}$  G

Lai, Shapiro (1991)

Ferrer et al. (2010)



10.1103/PhysRevLett.104.212001



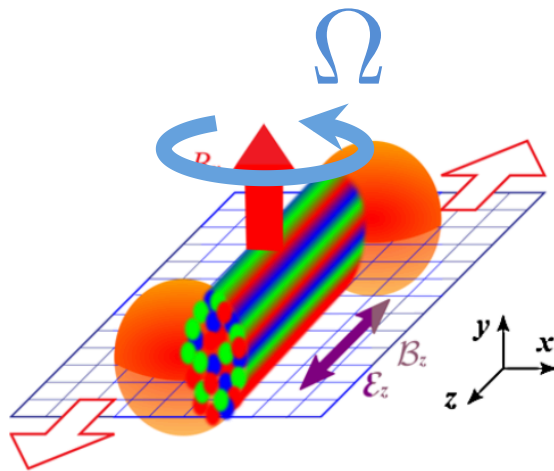
# Relativistic Systems with $\Omega$

Heavy-ion collisions  $\Omega \sim \Lambda_{\text{QCD}} ?$

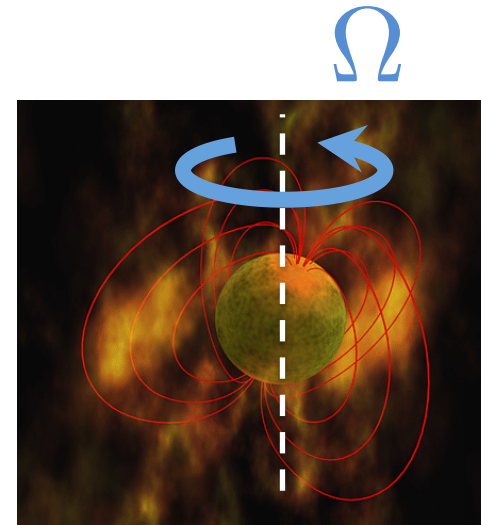
Csernai et al. (2011)

Magnetar  $R\Omega \sim 10^{-1}$

Marshall et al. (2004)



10.1103/PhysRevLett.104.212001

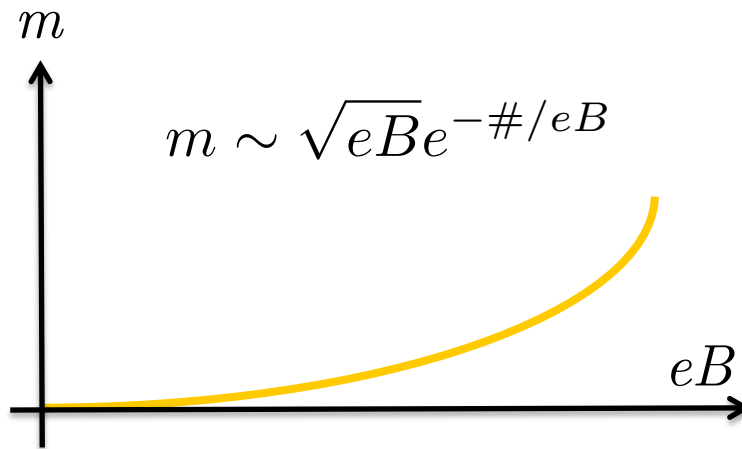


# *Magnetic Catalysis*

- \* Attractive interaction
- \* Large magnetic field

Klimenko (1992)

Gyusynin, Miransky, Shovkovy (1994)



Increasing  $eB$   $\longrightarrow$  Increasing  $m$

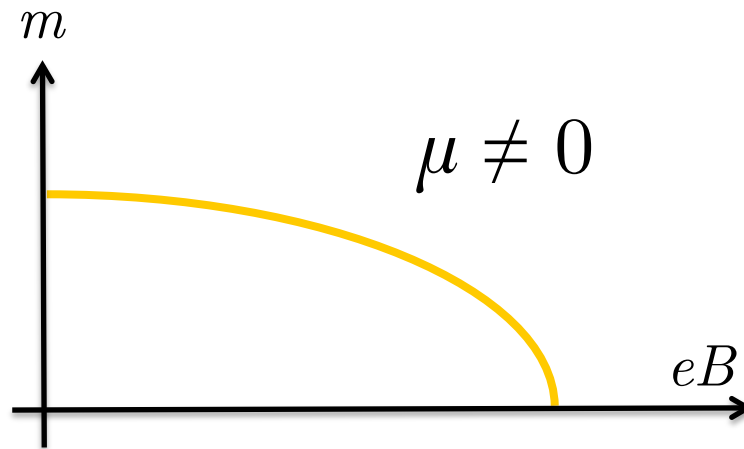
**Magnetic Catalysis**

# *Magnetic Inhibition*

- \* **Attractive interaction**
- \* **Large magnetic field**
- \* **Density**

Ebert, Klimenko (1999)

Preis, Rebhan, Schmitt (2012)



Increasing  $eB$   $\longrightarrow$  **Decreasing  $m$**

**Magnetic Inhibition / Inverse Magnetic Catalysis**

# “Rotational Magnetic Inhibition”

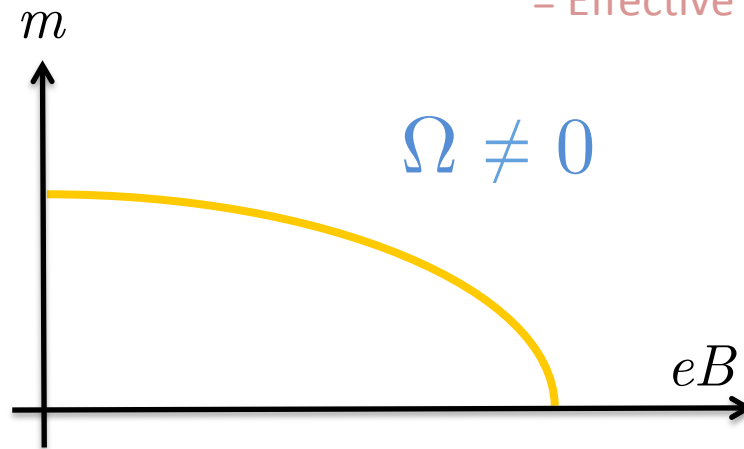
\* Attractive interaction

\* Large magnetic field

\* **Rotation**

$$E_{\text{Rest}} \rightarrow E_{\text{Rot.}} + \frac{(L_z + S_z)\Omega}{2}$$

Rotational Energy  
= Effective Chemical Potential



Increasing  $eB$   $\longrightarrow$  Decreasing  $m$

“Rotational Magnetic Inhibition”

# Fermions with $eB$ and $\Omega$

Dirac eq.  $\left[ i\gamma^\mu (D_\mu + \Gamma_\mu) - m \right] \psi = 0$

Rotation  $V^\mu = e_i^\mu V^i$

Magnetic field  $A_i = (0, By/2, -Bx/2, 0)$

Rotation  $g_{\mu\nu} = \begin{pmatrix} 1 - r^2\Omega^2 & y\Omega & -x\Omega & 0 \\ y\Omega & -1 & 0 & 0 \\ -x\Omega & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$

$$(E + \underbrace{\Omega\ell + \Omega s_z}_{\text{Rotational Energy}})^2 = eB(2n + 1 - 2s_z) + p_z^2 + m^2$$



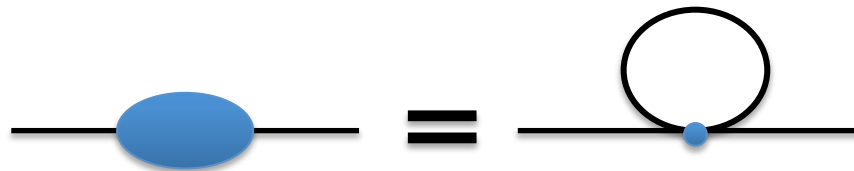
# *NJL Model with $eB$ and $\Omega$*

$$S = \int d^4x \sqrt{-\det(g_{\mu\nu})} \mathcal{L}(\bar{\psi}, \psi)$$

$$\mathcal{L} = \bar{\psi} [i\gamma^\mu (D_\mu + \Gamma_\mu)] \psi + \frac{G}{2} [(\bar{\psi}\psi)^2 + (\bar{\psi}\gamma_5\psi)^2]$$

Gap eq.  
(mean field approx.)

$$m = G \text{tr}[S(x, x)]$$

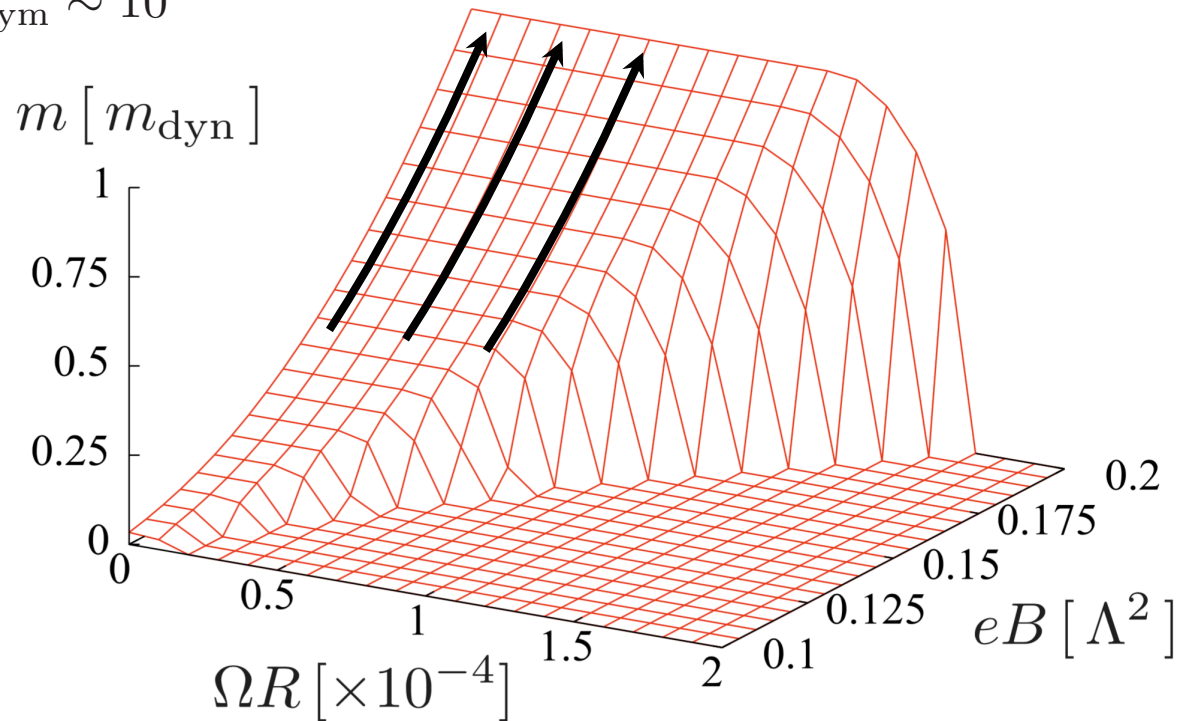


**NJL in finite density**  
Ebert, Klimenko (1999)

weak coupling : Catalysis  
strong coupling : Inhibition

# Weak Coupling ( $T = 0$ )

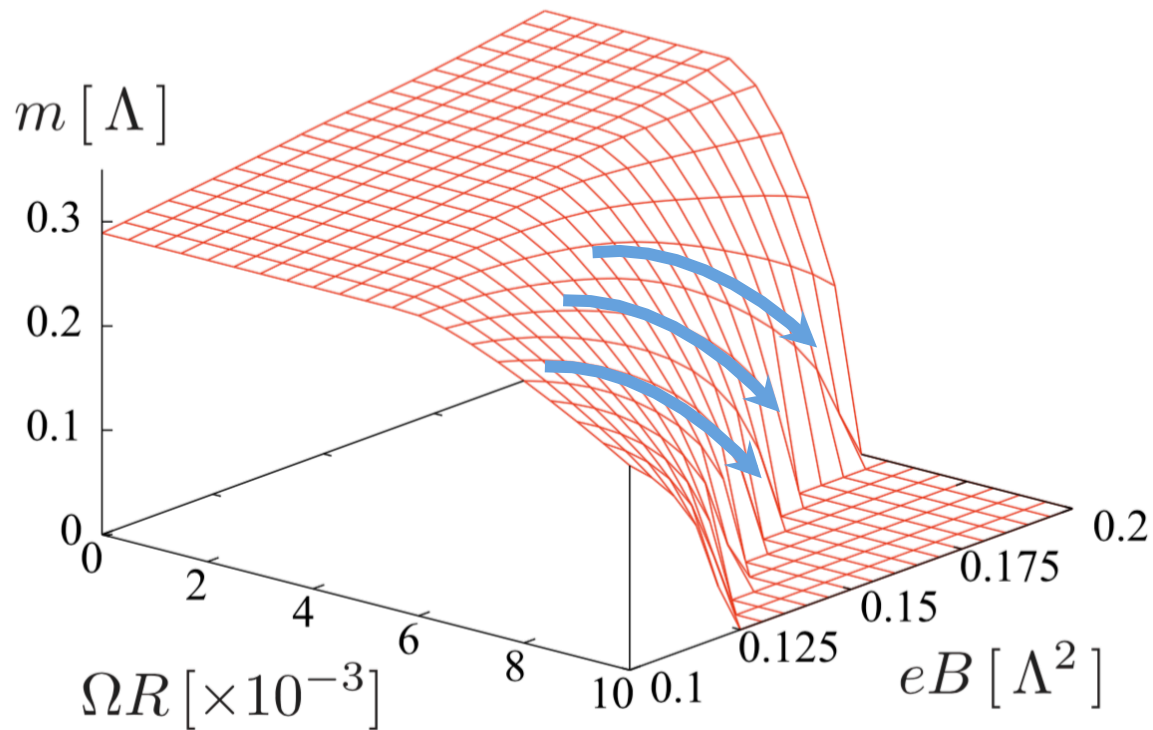
$$m_{\text{dym}} \sim 10^{-2}$$



Increasing  $eB$   $\longrightarrow$  Increasing  $m$

## Magnetic Catalysis

# Strong Coupling ( $T = 0$ )



Increasing  $eB$   $\longrightarrow$  **Decreasing  $m$**

**“Rotational Magnetic Inhibition”**

Chen, Huang, Fukushima, KM (2016)

# *How Rotation Affects Fermions*

Fermi dist.  $\frac{1}{1 - e^{-\beta\{\varepsilon - |\Omega(\ell + 1/2)|\}}}$   $\longrightarrow \theta\left(\varepsilon - |\Omega(\ell + 1/2)|\right)$

If there exists a mode that satisfies  $\varepsilon - \Omega|\ell + 1/2| < 0$   
rotation affects the fermion dynamics

$1/\sqrt{eB} \ll R$   $\longrightarrow$  The wave func. is localized  
 $\longrightarrow$   $R$  can be regarded as infinite

Eigen-mode  $\varepsilon = \sqrt{2neB + p_z^2} + m^2$

**Rotation with a strong magnetic field affects a system**

# Without Magnetic Field

The wave func. is *no longer* localized  $\longrightarrow$   $R$  should be finite

Eigen-mode  $\varepsilon = \sqrt{p_{\ell, k}^2 + p_z^2 + m^2} \quad \phi(r) \sim J_\ell(p_{\ell, k} r)$

IR mode  $J_\ell(p_{\ell, k} R) = 0 \longrightarrow p_{\ell, k} = \xi_{\ell, k} / R$

Causality  $\Omega R \leq 1$

$$\varepsilon - \Omega(\ell + 1/2) \geq \frac{\xi_{\ell, 1}}{R} - \Omega(\ell + 1/2)$$

$$= \frac{1}{R} \left[ \xi_{\ell, 1} - \Omega R(\ell + 1/2) \right]$$

$$\geq \frac{1}{R} \left[ \xi_{\ell, 1} - (\ell + 1/2) \right] > 0 \quad \text{e.g. } \xi_{0, 1} = 2.40483 > 1/2$$

**Rotation itself leads to NOTHING**

Ebihara, Fukushima, KM (2016)

# Summary

□ Rotation itself never changes dynamics (at zero temperature)

□  $eB$  with rotation  $\simeq$   $eB$  with density

□ A High  $eB$  or  $T$  allows us to ignore the boundary effect

$$j_A = \frac{T^2 \Omega}{12} + \frac{\Omega^2 \Omega}{48\pi^2} + O(\Omega R^{-2}) \quad \text{for } T \gg R^{-1} \geq \Omega$$

Vilenkin (1980)

□ For a small  $eB$  or  $T$ , the boundary effect should be taken into account