

Critical endline of the finite temperature phase transition for 2+1 flavor QCD around the SU(3)-flavor symmetric point

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in collaboration with

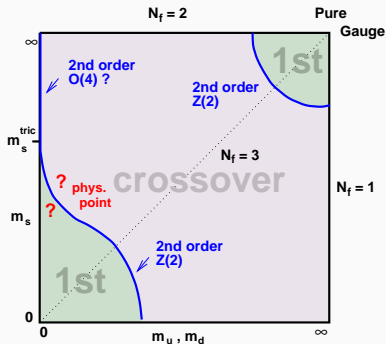
Y. Kuramashi, S. Takeda & A. Ukawa
(arXiv:1605.04659)

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XQCD 2016 in Plymouth

Motivation

- the critical endline at zero chemical potential for small quark mass region has not determined



[Laermann, Philipsen, '03]

Taylor expansion of kurtosis at critical endpoint around $m^{\text{sym}} (=m_l=m_s)$

$$K_E + c(\delta m_u + \delta m_d + \delta m_s) + \mathcal{O}(\delta m^2)$$

when changing quark mass as $\delta m_s = -2\delta m_l$, K_E remains unchanged up to $\mathcal{O}(\delta m^2)$

So, slope for the critical endline at m^{sym} should be -2

- we determine the critical endline around m^{sym} with clover fermions

Slope -2 for the critical endline at sym. point

- Taylor expansion of flavor symmetric operator at symmetric point
 $m_u = m_d = m_s = \bar{m}$

$$O(\delta m_u, \delta m_d, \delta m_s) = O(\bar{m}, \bar{m}, \bar{m}) + \mathbf{c}(\delta m_u + \delta m_d + \delta m_s) + \mathcal{O}(\delta m_q^2)$$

where $\delta m_i = m_i - \bar{m}$ and

$$\mathbf{c} = \frac{\partial O(\bar{m}, \bar{m}, \bar{m})}{\partial m_u} = \frac{\partial O(\bar{m}, \bar{m}, \bar{m})}{\partial m_d} = \frac{\partial O(\bar{m}, \bar{m}, \bar{m})}{\partial m_s}$$

- changing $\delta m_u + \delta m_d + \delta m_s = 0$ ($\delta m_s = -2\delta m_l$)

$$O(\delta m_u, \delta m_d, \delta m_s) = O(\bar{m}, \bar{m}, \bar{m}) + \mathcal{O}(\delta m_q^2)$$

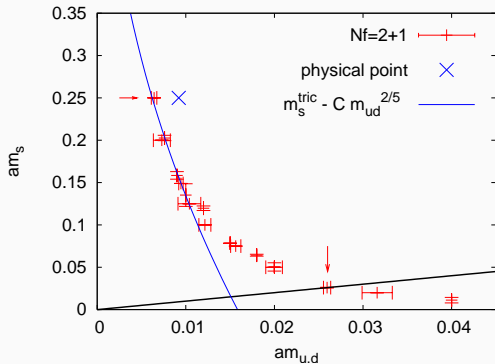
- $\chi, \mathbf{S}, \mathbf{K}$ with $\delta m_s = -2\delta m_l$

$$\chi(\bar{m}, \bar{m}, \bar{m}) = \chi(\delta m_u, \delta m_d, \delta m_s) + \mathcal{O}(\delta m_q^2)$$

$$\mathbf{S}(\bar{m}, \bar{m}, \bar{m}) = \mathbf{S}(\delta m_u, \delta m_d, \delta m_s) + \mathcal{O}(\delta m_q^2)$$

$$\mathbf{K}(\bar{m}, \bar{m}, \bar{m}) = \mathbf{K}(\delta m_u, \delta m_d, \delta m_s) + \mathcal{O}(\delta m_q^2)$$

Previous study with staggered fermions



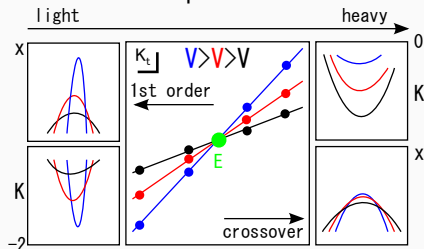
[de Forcrand, Philipsen, '06]

- $N_T = 4$, $a \approx 0.3$ fm
- data exhibits that slope at m^{sym} is not - 2
- $am_s^{\text{crit}} \approx 0.7$ (roughly 5 times larger than m_s^{phy})

Simulations

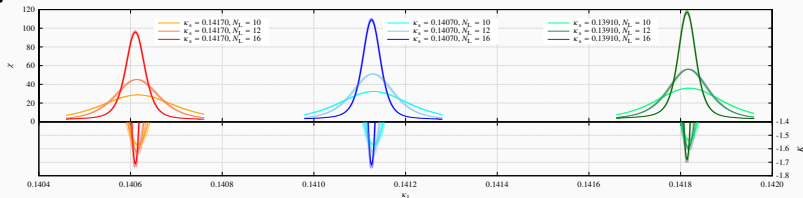
- Iwasaki gauge + NP O(a) improved Wilson fermions
- chiral condensate (10 - 20 noises for $\text{Tr}D^{-1,-2,-3,-4}$)
- kurtosis intersection method to determine the critical endpoint
- reweighting method to obtain more critical endpoints
- $N_t = 6$ ($a \approx 0.19\text{fm}$)
- $N_l = 10, 16, 20, 24$

β	κ
1.715	0.140900 – 0.141100
1.73	0.140420 – 0.140450
1.75	0.139620 – 0.139700

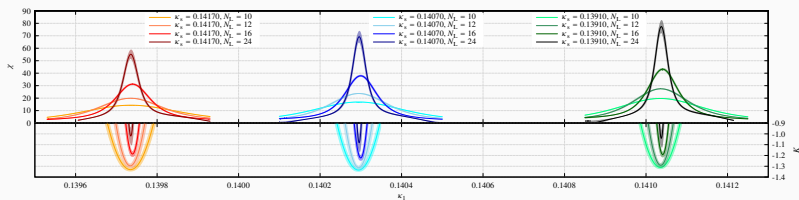


Susceptibility and kurtosis

$\beta = 1.715$



$\beta = 1.73$



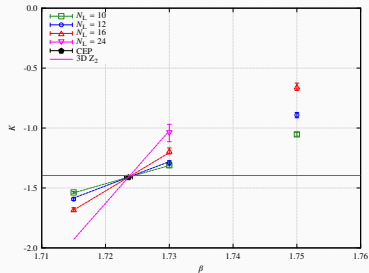
↑
 $\kappa_S = 0.1417$

↑
 $\kappa_S = 0.1407$

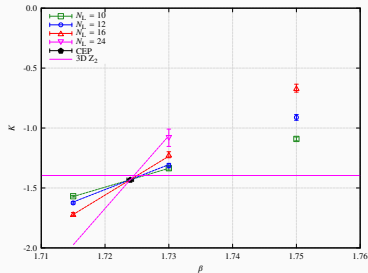
↑
 $\kappa_S = 0.1391$

Kurtosis intersection plot

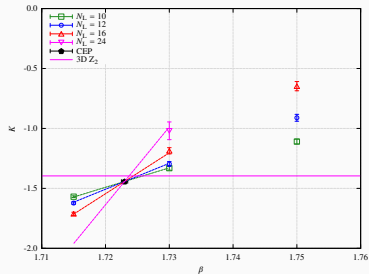
$\kappa_S = 0.1391$



$\kappa_S = 0.1407$

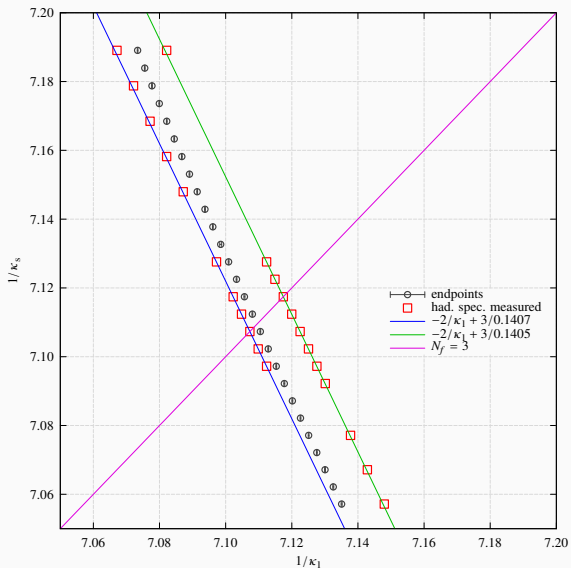


$\kappa_S = 0.1417$

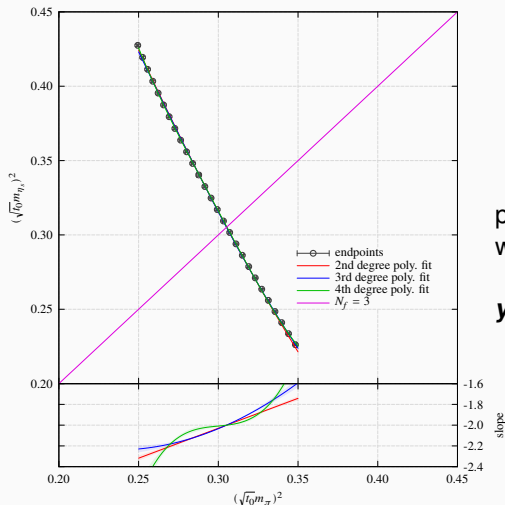


$$\text{fit} : \kappa_E + aN_L^{1/\nu}(\beta - \beta_E)$$

Critical endpoints in bare parameter plane



Critical endline



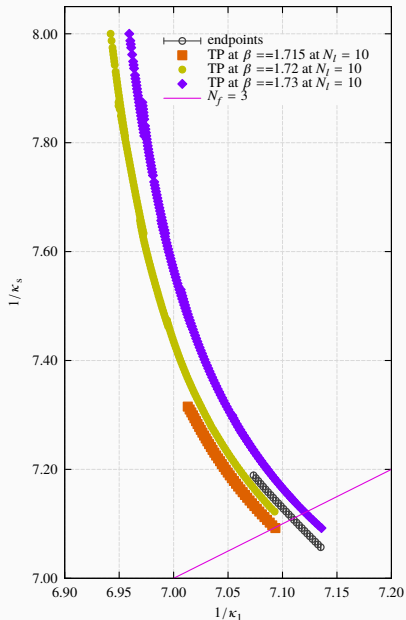
polynomial fits
with slope -2 at m^{sym}

$$y = c - 2(x - c) + a_1(x - c)^2$$

- $\chi^2/\text{d.o.f.} \sim 1$
- positive 2nd derivative

(cf. $1/\sqrt{t_0} = 1.347(30)$ GeV [Borsanyi et al. '12])

Critical endpoints away from m^{sym} (preliminary)



We have determined the critical endpoints (black points) by using $N_f = 3$ configurations

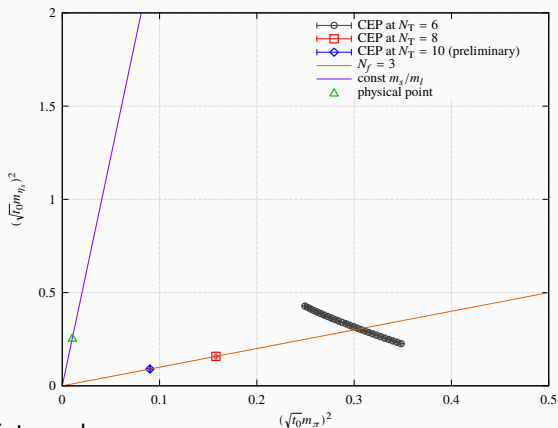
We are doing same analysis by using $N_f = 2 + 1$ configurations, so far only transition points at $N_f = 10$ are determined

Question : $m_s - m_s^{\text{tric}} \sim m_1^{2/5}$?

[Rajagopal '95]

Summary

We have determined the critical endline around the SU(3)-flavor symmetric point at $N_f = 6$ with NP O(a) improved Wilson fermions



- confirmed slope -2
- found positive 2nd derivative

Future plans

- critical endpoints away from symmetric point
- larger N_f for the continuum limit