

### Measurement and interpretation of the Fermi-LAT angular power spectrum of gamma-ray anisotropies Mattia Fornasa

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### The Diffuse Gamma-Ray Background

http://svs.gsfc.nasa.gov/vis/ a010000/a011300/a011342/

## Anisotropies in the Diffuse Gamma-Ray Background



- cumulative emission of unresolved sources
- guaranteed components from unresolved astrophysical sources

 $|a_{\ell m}|^2$ 

 constraints on additional contributors (Dark Matter)

$$I(\psi) = \sum_{\ell m} a_{\ell,m} Y_{\ell,m}(\psi) \qquad C_{\ell} = \frac{1}{2\ell + 1} \sum_{m=1}^{N} \frac{1}{2\ell + 1} \sum_{m=1}^{N}$$

- measure  $C_{\ell}$  (update the 2012 detection by Fermi-LAT)
- develop a model of  $C_{\ell}$  in terms of astrophysical sources to fit the data

### New APS measurement



Fornasa et al. (2016)	Ackermann et al. (2012)
81 months	22 months
Pass 7 reprocessed (ULTRACLEAN_v15) front	Pass 6 (DIFFUSE_v3) front and back
13 energy bins between 0.5-500 GeV	4 energy bins between 1-50 GeV
masking sources in 3FGL	masking sources in 1FGL

### **APS** estimator



### Binned APS measurement



- contamination of Galactic foreground at low  $\ell$  and effect of the beam window function at large  $\ell$
- fitting the data with a Poissonian APS:  $\chi^2$ /dof = 0.91, *p*-value=0.96
- fits with A(ℓ/ℓ<sub>0</sub>)<sup>α</sup> leads to a best-fit α of -0.06±0.08, with the same p-value of the Poissonian best fit

#### Anisotropy energy spectrum



$$C_{\ell} = \frac{1}{2\ell + 1} \sum_{m = -\ell}^{\ell} |a_{\ell m}|^2$$

$$I(\psi) = \sum_{\ell m} a_{\ell,m} Y_{\ell,m}(\psi)$$

• 1 population of source: factor out the energy dependence

$$C_{\mathbf{P}}^{i,j} = I(E_i) \, I(E_j) \, \tilde{C}_{\mathbf{P}}$$

• 2 populations of sources (contributions sum up linearly):

$$C_{\rm P}^{i,j} = I_{\rm A}(E_i) I_{\rm A}(E_j) \tilde{C}_{{\rm P},{\rm A}} + I_{\rm B}(E_i) I_{\rm B}(E_j) \tilde{C}_{{\rm P},{\rm B}}$$

 features in the anisotropy energy spectrum indicates multiple populations of sources

## Cross-correlation APS <sup>10<sup>2</sup></sup> Multipole

$$C_{\ell}^{ij} = \frac{1}{2\ell + 1} \sum_{m=-\ell}^{\ell} a_{\ell m}^{i} a_{\ell m}^{j\star}$$

- 91 independent combination of en. bins: 91 Poissonian  $C_{P^{i,j}}$
- cross correction coefficients

$$C_{\mathrm{P}}^{i,j}/\sqrt{C_{\mathrm{P}}^{i,i}C_{\mathrm{P}}^{j,j}}$$

• one source class:

$$C_{\mathbf{P}}^{i,j} = I(E_i)I(E_j)\tilde{C}_{\mathbf{P}}$$

• multiple source classes:

$$C_{\mathbf{P}}^{i,j} = \sum_{\alpha} C_{\mathbf{P},\alpha}^{i,j} = \sum_{\alpha} I(E_i) I(E_j) \tilde{C}_{\mathbf{P},\alpha}^{i,j}$$

 cross-correlation coefficients different than 1.0 hint at multiple components



10<sup>3</sup>

### Cross-correlation APS



### Interpretation in terms of multiple populations

Fitting the data with one or more populations, assuming specific energy spectra:

$$I(E) \propto E^{-\alpha} \qquad I(E) \propto \begin{cases} (E/E_0)^{-\alpha} & \text{if } E \leq E_b \\ (E_0/E_b)^{-\alpha+\beta} (E/E_0)^{-\beta} & \text{otherwise} \end{cases}$$



### Best-fit interpretation

Two populations of sources with broken-power-law spectra has the lowest  $\chi^2$  ( $\chi^2$ /dof=1.10, *p*-value=0.24)



### Parametric model for blazars

Luminosity-dependent density evolution:

• double power-law gamma-ray luminosity function

$$\Phi(L_{\gamma}, z=0, \Gamma) = \frac{dN}{dL_{\gamma}dVd\Gamma} = \frac{A}{\ln(10)L_{\gamma}} \left[ \left(\frac{L_{\gamma}}{L_{0}}\right)^{\gamma_{1}} + \left(\frac{L_{\gamma}}{L_{0}}\right)^{\gamma_{2}} \right]^{-1} \times \exp\left[ -\frac{(\Gamma - \mu(L_{\gamma}))^{2}}{2\sigma^{2}} \right]$$

 spectral index Γ depends on luminosity Lγ (modelling FSRQs and BL Lacs at the same time)

$$\mu(L_{\gamma}) = \mu^* + \beta \left( \log(L_{\gamma}) - 46 \right)$$

• evolutionary factor:  $\Phi(L_{\gamma}, z, \Gamma) = \Phi(L_{\gamma}, 0, \Gamma) \times e(z, L_{\gamma})$ 

$$e(z, L_{\gamma}) = \left[ \left( \frac{1+z}{1+z_c(L_{\gamma})} \right)^{-p_1(L_{\gamma})} + \left( \frac{1+z}{1+z_c(L_{\gamma})} \right)^{-p_2(L_{\gamma})} \right]^{-1}$$

$$z_c(L_{\gamma}) = z_c^* \left(\frac{L_{\gamma}}{10^{48} \text{erg s}^{-1}}\right)^{\alpha} \qquad \qquad p_1(L_{\gamma}) = p_1^0 + \tau \left(\log L_{\gamma} - 44\right) \\ p_2(L_{\gamma}) = p_2^0 + \delta \left(\log L_{\gamma} - 44\right)$$

# $Data (3F_1 GL \frac{dN}{dF} and APS) 3$

1

49r



• best fit from Ajello et al. (2015):  $\gamma_2 = 1.83$ ,  $\mu^{*0} = 2.22$ ,  $\beta = 0.1$ ,  $\sigma = 0.28$ ,  $z_c^* = 1.25$ ,  $\alpha = 7.23$ ,  $p_1^{0} = 3.39$ 



#### Evidence for new source class



- blazars alone cannot explain the measured APS below 1 GeV: χ<sup>2</sup>=1.59 (*p*-value=0.005) and χ<sup>2</sup>=1.70 (*p*-value=0.0003)
- include an additional class of sources, assuming  $C_{\rm P}(E) \propto E^{-2\Gamma \rm new}$
- new fits including new source class (Γ<sub>new</sub> and C<sub>P<sup>0,0</sup></sub>): χ<sup>2</sup>=1.11 (p-value=0.28) and χ<sup>2</sup>= (p-value=). New class is preferred at more than 99%CL.
- best-fit  $\Gamma_{\text{new}}$  is  $\Gamma_{\text{new}}=3.00^{0.27}_{0.22}$  and  $\Gamma_{\text{new}}=xxxx$



### Residuals

#### modelgll\_iem\_v05\_rev1.fit



### Validation of foreground cleaning



- significant less power at low multipoles (large angular scales)
- region with *l*>143 (bin #5 in the signal region) is not affected by Galactic foreground, even without cleaning
- foreground cleaning also reduces correlation between multipole bins

### Correlation between multipole bins





### Validation of mask (3FGL)



• 1°-mask: disc of 1 deg around all sources

- 2°-mask: disc of 2 deg around 500 brightest sources and disc of 1 deg around all the others
- 3°-mask: disc of 3.5 deg around 500 brightest sources, disc of 2.5 deg around following 500 brightest sources, disc of 2.0 deg around 1000 following brightest sources and disc of 1.5 deg around all the others

### Front only and front+back



- two data sets are not independent
- possibility of leakage outside of point sources in front+back data set because of the larger PSF

default default mask

#### Pass 7 vs. Pass 8



### How to bin APS

- produce 200 Monte Carlo realisations of the gamma-ray sky with a fixed nominal C<sub>P</sub>
- analytical expression for the error is

$$\sigma_{\ell} = \sqrt{\frac{2}{\left(2\ell+1\right)f_{\rm sky}}} \left(C_{\ell} + \frac{C_{\rm N}}{W_{\ell}^2}\right)$$

• to bin  $C_{\ell}$  in one multipole bin, you can compute:

#### A. unweighted average

- B. weighted average with weight =  $1/\sigma_{\ell}$
- C. weighted average with weight =  $1/\sigma_{\ell}$  and only photon noise
- Monte Carlo simulations prove that method B underestimates the APS



### How to estimate the error on the binned APS

- method A: unweighted average of the analytical expression for the error  $\sigma_\ell$ 

$$\sigma_{\ell} = \sqrt{\frac{2}{\left(2\ell+1\right)f_{\rm sky}}} \left(C_{\ell} + \frac{C_{\rm N}}{W_{\ell}^2}\right)$$

- method B: weighted average of  $\sigma_\ell$
- method C: average of the variances and covariances computed by PolSpice
- 40 Unweighted average Weighted average with weights  $w_1 = \sigma_1^{-2}$ they agree on the MC data 35 Nominal C<sub>P</sub> Number of MC realizations Gaussian distribution with 30  $\sigma$  estimated by PolSpice 25 20 15 10 <u></u>∃≍10<sup>-18</sup> 0 5 3 6 7 2  $\overline{C}$  [cm<sup>-4</sup> s<sup>-2</sup> sr<sup>-2</sup> sr]

How to fit the binned APS

• method A: minimise the  $\chi^2$ 

$$\chi^2(C_{\rm P}) = \sum_{\ell} \frac{(\overline{C_{\ell}} - C_{\rm P})^2}{\overline{\sigma_{\ell}}^2}$$

method B: maximise the log(L)

$$\log \mathcal{L}(C_{\rm P}) = -\sum_{\ell} \log(\overline{\sigma_{\ell}}) - \frac{1}{2} \sum_{\ell} \frac{(\overline{C_{\ell}} - C_{\rm P})^2}{\overline{\sigma_{\ell}}^2}$$



#### Comparison with old measurement



#### **Cross-correlation**



### Gamma-ray anisotropies from Dark Matter



 $E=4 \text{ GeV}, M_{\text{min}}=10^{-6} \text{ M}_{\odot}, b \text{ quarks}$ 

 $m_{\chi}$ =200 GeV,  $\sigma v$ =3×10<sup>-26</sup>cm<sup>3</sup>s<sup>-1</sup> (annihilation),  $m_{\chi}$ =2 TeV,  $\tau$ =2×10<sup>27</sup> s (decay)

## DM-induced emission

- repetition of the Millennium-II simulation box to cover a large portion of the Universe
- extrapolation below the mass resolution of the Millennium-II (assuming low-mass halos trace the smallest halos in Millennium-II)
- unresolved subhalos accounted for through an analytic fit to  $P(\rho, r)$
- Milky Way smooth halo and Galactic subhalos from Aquarius (carved in the centre)



10

26/13

# Effect of an uncertain MW mass on GAL-AQ

- uncertainty of a factor 4 on the mass of the Milky Way (MW)
- 16 bins in M<sub>MW</sub> accounting for a correspondent depletion in the amount of Galactic subhalos
- including uncertainty on the position of the observer



### Effect of an too-bright subhalos on GAL-AQ

- for certain combination of ( $m_{\chi}$ ,  $\sigma_{ann}v$ ) and ( $m_{\chi}$ ,  $\tau$ ), some subhalos are brighter than the 3FGL sensitivity
- those structures should be masked



### Effect of an too-bright subhalos on GAL-AQ



### **DM-induced APS**



### Conservative exclusion limits

 $\langle C_{\ell,\rm DM}^{i,j} \rangle < C_{\rm P}^{i,j} + 1.64 \, \sigma_{C_{\rm P}^{i,j}}$ 



 REF, MIN and MAX encompass the uncertainties considered in the distribution of the DM (for a fixed M<sub>min</sub>)

31/13

• shaded bands describe the uncertainty on the value of  $M_{\min}$ 



New measurement of the angular power spectrum of anisotropies using the data of the Fermi-LAT

32/13

### Conservative limits

10<sup>-25</sup>



10<sup>-20</sup> 10<sup>-20</sup> Measurement and interpretentiation of Spirit fairly 2017 Fermi LAT Oward Spheroidals (Ackermann et al., 2015) Fermi LAT IGRB intensity (Ackermann et al., 2015) Fermi LAT IGRB intensity (Ackermann et al., 2015) Fermi LAT IGRB intensity (Ackermann et al., 2015)

33/13

### 2-component fit to the binned APS



- TS = -2 ln[ $\chi^2(no \mathcal{D}_{g})$ ] + 2 ln[ $\chi^2(m_\chi, \sigma v)$ ]
- annihilation best  $f_{35\%}^{S} T_{35\%} T_{5\%} T_{75\%} T_{75$
- decay best fit: TS=xx,  $m_{\chi}$ =1743.2 GeV,  $\tau$ =1.15×10<sup>26</sup> s <sup>10<sup>27</sup></sup> 10 <sup>10<sup>2</sup></sup> 10<sup>3</sup>



# <sup>10<sup>-24</sup></sup> Exclusion limits for the 2-component fit

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35/13

### TS for 2-component fit

