### Hunting dark Z' and dark Higgs bosons.

#### Fantastic WIMPs and Where to Find Them

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#### Where are the WIMPs?



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## In the Dark about Dark Matter

Recent disappointments have physicists looking beyond WIMPs for dark matter particles



#### **New mediators**

- Experiments are closing in on the simplest WIMP models, in which the DM particle interacts via Standard Model bosons.
- A possible conclusion could be to abandon the WIMP idea and focus on different DM candidates, such as axions or sterile neutrinos.
  - Many great proposals discussed at this workshop!
- Nevertheless, the paradigm of thermal freeze-out is very general and does not require the DM particle to directly couple to any of the SM bosons.
- Instead, there may be a new mediator responsible for communicating the interactions of DM.



This opens up new parameter space – but also leads to new constraints.



#### **Searches for new resonances**



- Crucially, the mediator of the DM interactions will also lead to new interactions between Standard Model states,
- There may be observable signals from processes involving no DM particles at all.
- For example, if the mediator can be produced at the LHC, it can also decay back into quarks.
- One should therefore consider dedicated searches for the mediator particles themselves, such as searches for dijet resonances.





#### **Searches for dijet resonances**

- Results from dijet resonance searches at the LHC are presented in a way that they can be applied to a wide range of different models.
  - In particular, it is possible to directly compare simulated distributions of the dijet invariant mass to the observed spectra.
  - The SM expectation for these spectra are obtained by fitting a smooth function to the observed distribution:  $f(z) = p_1(1-z)^{p_2}z^{p_3}$
  - It is possible with this approach to reproduce experimental bounds (e.g. for the case of an RS graviton) to very good accuracy.





#### **Combination of dijet resonance searches**

- The great advantage of using actual dijet invariant mass spectra rather than published bounds is that one can not only reproduce the 95% CL bound, but in fact reconstruct the full likelihood (using e.g. a χ<sup>2</sup> test statistic).
- This makes it possible to combine dijet resonance searches from both ATLAS and CMS, as well as searches at both 8 TeV and 13 TeV.

	$m_{jj}$	$ \Delta \eta_{jj} $	additional
ATLAS 13 TeV	> 1.1  TeV	< 1.2	$p_{T,j_1} > 440 \text{ GeV} \text{ and } p_{T,j_2} > 50 \text{ GeV}$
$\mathrm{CMS}~13~\mathrm{TeV}$	> 1.2  TeV	< 1.3	$p_{T,j_1} > 500 \text{ GeV or } H_T > 800 \text{ GeV}$
ATLAS 8 TeV	> 250  GeV	< 1.2	-
$\mathrm{CMS}~8~\mathrm{TeV}$	$> 890 { m GeV}$	< 1.3	-
CMS 8 TeV $(low)$	$> 390 { m ~GeV}$	< 1.3	-



#### **Deriving model-independent bounds**

Previous analyses of LHC dijet searches: Focus either on resonances that decay exclusively into quarks or on certain DM models with specific choices of couplings.

An, Ji & Wang, arXiv:1202.2894 Dobrescu & Yu, arXiv:1306.2629 Chala, FK et al., arXiv:1503.05916

- Possible to analyse LHC dijet searches in a largely model-independent way
- Results can be applied to many different Z' models, including models with invisible or unobserved decay modes.
- Key idea: take the width of the Z' as a free parameter (rather than to calculate it from the assumed couplings).
- Since the shape of the dijet invariant mass distribution depends only on m<sub>z'</sub> and Γ<sub>z</sub>, changing g<sub>a</sub> then only changes the normalisation of the signal.
- For given m<sub>z</sub> and Γ<sub>z</sub>, we can thus easily construct an upper bound on the magnitude of the signal, leading to a bound on g<sub>a</sub>.



#### **Results**



- Narrow width:  $g_q \sim 0.1$
- Broad width:  $g_a \sim 0.3$

Sensitive to the typical weak interaction strength relevant for thermal freeze-out!



#### **Models of new mediators**

- Aim: combine all these LHC results with the information from other DM searches in order to understand whether thermal freeze-out can still be viable in this set-up.
- > This will typically only be possible within the context of a specific DM model.
- > At first sight, there is a huge number of DM models interacting via a new mediator.
- However, experimental constraints and theoretical considerations eliminate a large number of possibilities:
  - Direct detection experiments strongly disfavour a spin-1 mediator coupling to the DM vector current. Focusing on Majorana DM evades this constraint.
  - LHC dilepton resonance searches strongly constrain the coupling of the mediator to leptons. Consider models where the mediator couples dominantly to quarks.
  - The couplings of the mediator to SM fermions should respect the full symmetry of the SM gauge group before electroweak symmetry breaking.



### Mediators from a U(1)'

We consider a new U(1)' gauge group under which only quarks and the Majorana DM particle are charged. This gauge group is spontaneously broken by the vev of a SM singlet S.

$$\mathcal{L}_{\rm DM} = \frac{i}{2} \bar{\psi} \partial \!\!\!/ \psi - \frac{1}{2} g_{\rm DM}^A Z'^\mu \bar{\psi} \gamma^5 \gamma_\mu \psi - \frac{1}{2} y_{\rm DM} \bar{\psi} (P_L S + P_R S^*) \psi ,$$
  
$$\mathcal{L}_S = \left[ (\partial^\mu + i \, g_S \, Z'^\mu) S \right]^\dagger \left[ (\partial_\mu + i \, g_S \, Z'_\mu) S \right] + \mu_s^2 \, S^\dagger S - \lambda_s \left( S^\dagger S \right)^2$$

The Higgs singlet vev therefore generates both the Z' mass and the DM mass.

$$m_{\rm DM} = \frac{1}{\sqrt{2}} y_{\rm DM} w , \quad m_{Z'} \approx 2g_{\rm DM}^A w$$

This set-up is similar to DM models with gauged baryon number. For the purpose of this talk, I will not discuss the issue of anomaly cancellation.

Fileviez & Wise, arXiv:1002.1754, Duerr et al., arXiv:1304.0576



#### The two-mediator dark matter (2MDM) model

- The Z' is in fact not the only particle in this model that can mediate the interactions of DM.
- The dark Higgs will in general mix with the SM Higgs and thereby obtain couplings also to SM fermions.
- > We then obtain two dark mediators with the following interactions:

$$\mathcal{L}_{\chi} \supset -\frac{g_{\chi}}{2} \bar{\chi} \gamma^{\mu} \gamma^{5} \chi Z'_{\mu} - \frac{y_{\chi}}{2\sqrt{2}} \bar{\chi} \chi s , \qquad \qquad \frac{y_{\chi}}{m_{\chi}} = 2\sqrt{2} \frac{g_{\chi}}{m_{Z'}}$$
$$\mathcal{L}_{q} \supset -\sum_{q} \left( g_{q} \bar{q} \gamma^{\mu} q Z'_{\mu} + \sin \theta \frac{m_{q}}{v} \bar{q} q s \right)$$

The model is described by 6 independent parameters (three masses and three couplings).

Duerr, FK et al., arXiv:1606.07609



#### **Dark mediators and dark terminators**

- If one of the mediators is heavy and weakly-coupled, it will play no role in the phenomenology and one recovers a simpler model with only one mediator.
- If one of the two mediators is *light* and weakly coupled, it can provide a new final state for DM annihilation.





- We call such a light state a dark terminator, because it terminates rather than mediates the interactions of DM.
- Such a dark terminator can relax the relic density constraints on the interactions of the second mediator and thus extend the viable parameter space.



#### Saving WIMPs with dark terminators?



- Introducing a new weakly-coupled particle in the final state of DM annihilations is a drastic change of the traditional WIMP idea.
- It is certainly controversial whether such a model can still be called WIMP DM or if we need a different name (e.g. secluded DM, hidden sector freeze-out, ...).
- Nevertheless, in the context of the 2MDM model, we can continuously move from the traditional WIMP idea (Higgs exchange) to WIMP models with a new mediator (Z' exchange, dark Higgs exchange) to models with a dark terminator.
- This makes it possible to understand in general whether the simple WIMP idea is still viable or whether experiments favour more complex dark sectors.



#### **Global scans**

We can now perform a global analysis: For fixed masses we scan over all coupling combinations consistent with the relic density requirement.



# Red: All coupling combinations are excluded by at least one constraint.

White: At least one coupling combination is compatible with all constraints.

Orange: Large values of  $g_q$  cannot reliably be excluded due to the mediator width becoming large ( $\Gamma/m_{z'} > 0.3$ ).



#### A closer look



> Allowed parameter space for  $g_q \sim 0.04$  and  $\theta < 0.06$ .



#### **Global scans: Final results**



- Small DM masses are tightly constrained by the data.
  - The presence of a dark terminator is essential!
- Models with two dark terminators are in fact excluded by indirect detection constraints for m<sub>x</sub> < 100 GeV.</p>
  - For larger DM masses, there is larger viable parameter space (we find allowed solutions up to *m<sub>x</sub>* < 50 TeV).



### Hunting the dark Higgs

- At first sight, the case of a light dark Higgs terminator with tiny mixing with the SM Higgs seems impossible to probe experimentally.
- Nevertheless, the dark Higgs typically couples quite strongly to the DM particle and the Z'.
- If it is possible to produce these states at the LHC, they may emit Dark-Higgs-Strahlung.







#### Hunting the dark Higgs - predicted event rates

For large Z' masses, the rate of Dark-Higgs-Strahlung may actually exceed the rate of QCD radiation from the initial state.





Moreover, the dark Higgs is typically highly boosted and its decay therefore leads to a very characteristic signal: a fat jet containing two b-jets with an invariant mass close to the dark Higgs mass.



### Hunting the dark Higgs - projected sensitivity

- Looking specifically for this signature can suppress experimental backgrounds with high efficiency.
- > This makes it possible to probe the 2MDM model up to very large Z' masses.



Duerr, FK, et al., in preparation



#### Conclusions

- The non-observation of new physics at the LHC and null results in direct DM searches put significant pressure on the WIMP idea.
- A possible way forward: New mediators open up large new parameter space but also predict many new signatures (e.g. dijet resonances).
- The two-mediator DM model offers an attractive framework to study these ideas by allowing for a continuous transition from more conventional WIMPs (e.g. Higgs portal) to more exotic realizations of thermal freeze-out (e.g. dark terminators).
- We find significant experimental constraints on heavy mediators, unless DM is also rather heavy.
- Nevertheless, thermal freeze-out is certainly still viable if the dark sector has additional structure, for example if there is a new light terminator.
- Such models with light terminators can in fact be probed experimentally, for example with dedicated LHC searches for dark Higgs jets.

