Dark Matter: Theory Overview from 10⁻²² eV to 10⁶⁷ eV

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work w/ Asimina Arvanitaki, Savas Dimopoulos, Sergei Dubovsky, Isabel Garcia Garcia, Lawrence Hall, Ed Hardy, Karsten Jedamzik, Nemanja Kaloper, Robert Lasenby, Chris McCabe Matthew McCullough, James Unwin, and Stephen West (and unpublished with Peter Graham, Jeremy Mardon, Doddy Marsh, Surjeet Rajendran)

The Dark Matter Landscape



~10⁶⁷eV (non-baryonic solar mass)

The Dark Matter Landscape



The Dark Matter Landscape



The Dark Matter Landscape (fermionic)









$$\rho_{\rm DM} \approx 0.3 \, \frac{\mathrm{crev}}{\mathrm{cm}^3} \approx (0.04 \, \mathrm{eV})^4$$

Generic Candidates: Light Pseudo-Nambu-Goldstones (Axions and Axion Like Particles — ALPs); Massive Hidden Vector Bosons



search for coherent effects of the entire field, not single hard particle scatterings

Early universe production? eg misalignment mechanism (other mechanisms possible)



initial value in inflationary patch displaced from minimum (during inflation spatial gradients become small)

$$\ddot{a} + 3H(t)\dot{a} + V'(a) = 0$$

Early universe production? eg misalignment mechanism (other mechanisms possible)



$$\ddot{a} + 3H(t)\dot{a} + V'(a) = 0$$

value of field frozen until $H(t) \leq m_a$ when it starts oscillating at frequency set by mass $a(t) \simeq a_0 \sin(m_a t)$

Early universe production? eg misalignment mechanism (other mechanisms possible)



Oscillating field redshifts as non-relativistic matter (due to slow evolution of a_0)

$$m_a^2 a_0^2 \simeq \rho_{DM}$$

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Oscillating field redshifts as non-relativistic matter (due to slow evolution of a_0)

$$m_a^2 a_0^2 \simeq \rho_{DM}$$

For scales longer than deBroglie wavelength acts just like cold DM

Today: a random field



 $\ell_c \sim 1/(m_a v)$

 $v={\rm Galactic}$ virial velocity $\sim 10^{-3}$

Today: a random field Coherence time for Earth experiments?



 $v = {\rm Galactic} \ {\rm virial} \ {\rm velocity} \sim 10^{-3}$

Today: a random field



Most important: the frequency spread of oscillation is *small*



 \rightarrow

resonant enhancement in detection is possible with $Q\sim 10^6$

Today: a random field



Most important: the frequency spread of oscillation is *small*



exactly how one detects depends on type of bosonic DM

Dark Matter



A New Particle (or Sector): Non gravitational interactions with SM?

What kind of bosons?

What kind of bosons?

Naturalness. Structure set by symmetries.

Spin 0 Axions and other goldstone bosons Easy to get in many UV theories





Strong CP Problem

$$S_{\theta} = \frac{\theta}{32\pi^2} \int d^4 x \epsilon^{\mu\nu\lambda\rho} \mathrm{Tr} \, G_{\mu\nu} G_{\lambda\rho}$$

Neutron edm bound implies

$$\bar{\theta} = \theta + \arg \det m_q \lesssim 10^{-10}$$

Like Cosmo Const and Electroweak hierarchy problems requires precise cancelation of apparently unrelated quantities

Unlike CC and EW problems NO anthropic reason

A clear call for new dynamics

QCD (Invisible) Axion

Peccei-Quinn-Weinberg-Wilczek Kim-Shifman-Vainshtein-Zakharov Dine-Fischler-Srednicki-Zhitnitsky

$$S_a = \int d^4x \left(\frac{1}{2} (\partial_\mu a)^2 + \frac{a}{32\pi^2 f_a} \epsilon^{\mu\nu\lambda\rho} \operatorname{Tr} G_{\mu\nu} G_{\lambda\rho} \right)$$

Non-pert: QCD gives potential V(a)

$$\sim m_a \sim \frac{\Lambda_{QCD}^2}{f_a} \sim 6 \times 10^{-10} \text{eV} \left(\frac{10^{16} \text{GeV}}{f_a}\right)$$

Minimum of potential leads to axion vev such that

$$\theta_{eff} \equiv \frac{\langle a(x) \rangle}{f_a} + \bar{\theta} = 0$$
 solves strong CP!

Axion vev only cancels $\bar{\theta}$ to required 10^{-10} accuracy if all other non-pert sources of axion mass smaller than QCD by 10^{-10}

$f_a \lesssim 10^9 { m GeV}$ is excluded by stellar and supernova 1987A physics



 $f_a \gg 10^{12} {
m GeV}$ is an especially interesting region: would be the evidence that Ω_{DM} is fixed anthropically

The Axiverse Axion Landscape

Arvanitaki, Dimopoulos, Dubovsky, Kaloper, JMR; arXiv:0905.4720

In the next decade cosmo and astro observations will be exploring 23 orders of magnitude in energy



Taking properties of axions in string theory seriously, there can exist a plenitude of axions with log-flat distribution of masses



String Axiverse

antisymmetric forms



compactification



many (100-10000) massless axions from topology, eg: $a_i = \int_{\Sigma_3^i} C_3$

axions can be removed from the spectrum by fluxes, branes, orientifold planes, but many survive....

Chern-Simons coupling _____ (Green-Schwarz anomaly cancelation)

axionic couplings

ALP masses in String Thy

Must suppress all possible non-pert string effects that contribute to the QCD axion mass $> 10^{-10} \times \text{QCD} \implies$ action (eg, of Euclidian wrapped D-brane) $S \gtrsim 200$

This QCD axion constraint on string model building impacts phys of all string axions

$$m_a \sim \frac{\mu_{UV}^2}{f_a} e^{-S/2}$$

Axion masses exponentially sensitive to precise S

$$f_a \sim \frac{M_{pl}}{S} \sim 10^{16} \mathrm{GeV}$$

Axion couplings only linearly sensitive to S

ALP properties in String Thy







Axion/ALP DM behaves just like ColdDM (despite being a BEC) except at "small" scales

Uncertainty Principle prevents density perturbation growth at

$$\frac{k_J}{a} > \sqrt{Hm}$$





What kind of bosons?

Naturalness. Structure set by symmetries.



Spin 0

Axions and other goldstone bosons

Easy to get in many UV theories



Naturalness. Structure set by symmetries.



Electromagnetism Nuclear Force Nuclear Spin

$$\left(\frac{a}{f_a}F\tilde{F}\right)$$
$$\left(\sim\frac{a}{f_a}\vec{E}.\vec{B}\right)$$

$$\left(\frac{a}{f_a}G\tilde{G}\right)$$

$$\left(rac{\partial_{\mu}a}{f_a}ar{N}\gamma^{\mu}\gamma_5 N
ight.$$

QCD Axion

General Axions



all interactions suppressed by (large) scale f_a

What kind of bosons?

Naturalness. Structure set by symmetries.






 $Hz \lesssim \omega \lesssim GHz$



Dark Matter



A New Particle (or Sector): Non gravitational interactions with SM? Need exquisite sensitivity to detect such tiny interactions



Massive Hidden Photons as DM

First consider the relatively unexplored possibility of massive Hidden Photon (ie, Z') DM Nelson, Scholtz arXiv:1105.2812

Arias et al arXiv:1201.5902

DM is classical A'_{μ} field oscillating at $\omega = m'_A$ (in random dir'n)

Acts as an electric field that is not shielded

$$E' \sim \sqrt{\rho_{DM}} \simeq 2000 \mathrm{V/m}$$

can excite an EM resonator that is shielded from normal fields

(The axion search experiment ADMX is sensitive in range 10^{-6} eV $\lesssim m_{A'} \lesssim 10^{-4}$ eV due to cavity size)

Hidden Photon Landscape



Arias et al arXiv:1201.5902

Hidden Photon DM



A powerful and flexible way is with a high-Q radio

S. Chaudhuri etal arXiv:1411.7382

a resonant tunable LC circuit inside a Faraday cage with SQUID readout

Iower and wider frequency range possible than with cavities
Metal box to shield backgrounds





Powerful ways to confirm signal is DM

- signal is narrow band at constant frequency
- $\frac{1}{2}$ directional & phase coherence over ~10³ wavelengths
- dependence on orientation characteristic of vector

another experiment in early stages...





New topic...

see, eg, Baer, Choi, Kim, Roszkowski arXiv:1407.0017

Search Non-Asymmetric Dark Matter

Axions & Other Light Bosons: Mis-alignment or thermal or non-thermal production

Non-Asymmetric Dark Matter

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Axions & Other Light Bosons: Mis-alignment or thermal or non-thermal production

WIMPs: Calculable thermal freeze-out at EW scale v

FIMPs: Calculable thermal freeze-in (possibly at EW scale v)

Asymmetric DM & Baryons

Sharing Co-genesis In both cases *final* DM state can often be composite ("atomic", "nuclear"...)

Search Non-Asymmetric Dark Matter

Axions & Other Light Bosons: Mis-alignment or thermal or non-thermal production

WIMPs: Calculable thermal freeze-out at EW scale v

FIMPs: Calculable thermal freeze-in (possibly at EW scale v)

Service Asymmetric DM & Baryons

Sharing

Co-genesis

Very Heavy Extended Objects (PBHs, Q-balls,...)

Many possibilities...

WIMP Freeze-out: DM starts with full T^3 density but interactions cannot track falling equilibrium density



But Baryon density set by particle-antiparticle asymmetry...

$$\eta_B = Y_B - Y_{\bar{B}}$$

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Unrelated origin, involving very different physics, of baryons & DM, makes it hard to understand $\Omega_{DM}/\Omega_B \simeq 4.86$

But Baryon density set by particle-antiparticle asymmetry...

$$\eta_B = Y_B - Y_{\bar{B}}$$



Unrelated origin, involving very different physics, of baryons & DM, makes it hard to understand $\Omega_{DM}/\Omega_B \simeq 4.86$

Freeze-out has dominated our thinking about DM candidates, detection, and LHC phenomenology



A Comment wrt SUSY LSP

Arvanitaki, Craig, Dimopoulos, Dubovsky, JMR: arXiv:0909.5440

In the context of String Theory highly unusual for SM to contain the LSP

- A plenitude of hidden sectors, each with their gauginos, sfermions, goldstini, moduli,...
- For Standard Model sector to have the LSP the SM must be most weakly coupled of all sectors to SUSY-breaking



A Comment wrt SUSY LSP

In the context of String Theory highly unusual for SM to contain the LSP

Can get long decay chains of the *many* neutralinos to true LSP via emission of visible or hidden states



eg, for the kinetically mixed "photini" we considered have basic vertices







Hall, Jedamzik, JMR, West, arXiv:0911.1120

Suppose

- $\overset{\circ}{\mathscr{S}}$ X only feebly coupled to visible-sector thermal bath particles V_i
- X never in thermal equilibrium with SM

as universe evolves, a tiny X abundance is produced

 $Y_X(t) \sim \Gamma_V t$





two calculable thermal mechanisms



Comments

FI yield is IR-dominated for renormalizable interactions



(heaviest particle in vertex)

The lightest ordinary-sector particle (LOSP) transforming under the X-stabilising symmetry is automatically long-lived

Even if LOSP is neutral so leading decay to X invisible, sub-dominant 3- or 4-body decays can involve charged SM states and allow measurement of lifetime and X mass

Solutions of Small Coupling?

The 'WIMP miracle' is that for $m' \sim v$ and $\lambda' \sim 1$

$$Y_{FO} \sim \frac{1}{\lambda'^2} \left(\frac{m'}{M_{Pl}}\right) \sim \frac{v}{M_{Pl}}$$

gives the observed value of $\Omega_{DM}h^2$

The 'FIMP miracle' is that for $m \sim v$ and $\lambda \sim v/M_{Pl}$

$$Y_{FI} \sim \lambda^2 \left(\frac{M_{Pl}}{m}\right) \sim \frac{v}{M_{Pl}}$$

Suggests that FIMPs occur where small couplings arise at linear order in the weak scale

For example...

moduli of the SUSY-breaking sector giving MSSM soft terms

$$m^{2}\left(1+\frac{T}{M}\right)\left(\phi^{\dagger}\phi+h^{\dagger}h\right) \quad \mu B\left(1+\frac{T}{M}\right)h^{2} \qquad Ay\left(1+\frac{T}{M}\right)\phi^{2}h$$
$$m_{\tilde{g}}\left(1+\frac{T}{M}\right)\tilde{g}\tilde{g} \qquad \mu y\left(1+\frac{T}{M}\right)\phi^{2}h^{*} \qquad \mu\left(1+\frac{T}{M}\right)\tilde{h}\tilde{h},$$

similarly for the modulini

$$\mu \frac{T}{M} \tilde{h}h \qquad \qquad \frac{m_{susy}}{M} \tilde{T}(q\tilde{q}^{\dagger}, l\tilde{l}^{\dagger}, \tilde{h}h^{\dagger})$$

For $M \sim M_{GUT}$ (natural value of compactification scale in realistic string theories) give renormalizable couplings $\lambda \sim 10^{-13}$

So far assumed FIMP mass close to weak-scale. For WIMPs this must be so as unitarity limits size of annihilation cross-section

FIMPs completely different:

DM with relic abundance Y and mass m leads to temperature for matter-rad'n equality of parametric form $T_e \sim Ym$

Remarkably for FI this is independent of mass $T_{e,FI} \sim \lambda^2 M_{Pl}$



Calculable thermal production of superheavy FIMP DM possible

w/ apologies to Rocky: FIMPzilla's!
Another possibility...

Asymmetric DM

Alternative: similar physics underlies both Ω_B and Ω_{DM}

(Nussinov '85; Gelmini, Hall, Lin '87; Barr '91; Kaplan '92; Thomas '95; Hooper, JMR, West '04; explosion in last few yrs esp work of Zurek etal; Sarkar etal; Sannino etal; now many others...)

Baryons: $U(1)_B$

u, d, s... p stable $\Omega_B \propto m_B \eta_B$

DM: $U(1)_X$ $X_0, X_1, X_2... X_0$ stable $\Omega_X \propto m_X \eta_X$

At some era

Interactions violate B and X to yield related values for η_B and η_X

$$\frac{\Omega_X}{\Omega_B} = \frac{\eta_X}{\eta_B} \frac{m_X}{m_B}$$

Ω_X =	$\underline{\eta_X} \underline{m_X}$
Ω_B	$\eta_B m_B$

only true if X density is determined by the asymmetric part otherwise

$$\frac{\Omega_X}{\Omega_B} = \frac{Y_X + Y_{\bar{X}}}{Y_B + Y_{\bar{B}}} \frac{m_X}{m_B}$$

need $Y_X + Y_{\bar{X}} = Y_X - Y_{\bar{X}} + \text{small corrections}$

non-trivial constraint as initially

$$Y_X + Y_{\bar{X}} = \frac{Y_X - Y_{\bar{X}}}{\epsilon}$$

where $\epsilon \ll 1$ measures CP-violation

\implies Must efficiently annihilate away symmetric part to light states

there has to be an efficient X-preserving freeze-out process

Options:

June of the second seco

 \Longrightarrow operators connecting X & SM sectors with strength bounded below

direct FO to light dark sector dof

 \implies (potentially) new "long-range" DM interactions

FO to dark sector dof which then late decay to SM

 \implies late-time energy injection in early universe

direct FO to light SM dof

limits from direct detection experiments and monojet etc searches at Tevatron and LHC are very constraining

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with slight exceptions if we want asymmetric DM in natural region $m_X < 10 \text{ GeV}$ then direct FO to SM is disfavoured

eliminating symm component likely implies new dark-sector dynamics



 $\eta_X \sim \eta_B$ by sharing

 $\eta_X \sim \eta_B$ by co-generation

Co-generation is more ambitious: attempts to explain simultaneous origin of B & X asymmetries (if at scale ~ TeV allows test at LHC...)

Sharing:

Assumes presence of some initial asymmetry in (at least) one of B, L & X

A "connector interaction" breaks a combination of B/L & X, such that



there is an era when only conserved U(I) is

$$B - L + X \implies \eta_B : \eta_L : \eta_X = N_1 : N_2 : N_3$$



Alternative view (either sharing or co-generation):

- incompatible with standard SUSY neutralino DM
- alters expected LHC signals of new physics
- Second change one or both direct/indirect DM detection

major issue is why is DM mass near that of baryon?

(but see, eg, Garcia Garcia, Lasenby, JMR; arXiv:1505.07410 for automatic explanation directly connected with naturalnessthe "Twin Higgs" mechanism: also see work on "mirror world" models, by Foot, Volkas, etal)

Consequences of ADM

For non-collider DM searches:

Light DM ~ few GeV is favoured

(but see, eg, JMR+McCullough, arXiv:1106.4319, and Sarkar etal for other possibilities)

indirect detection strongly modified - DM can't annihilate to only photons but can give rise to anti-B/L final states (Hall, JMR, Unwin, West, unpublished)

sharing allows normal direct detection, but co-generation can sometimes kill direct detection (but again see, eg, JMR+McCullough, arXiv:1106.4319 for co-generation theory with direct detection signals)

exotic possibilites, eg, DM-stimulated nucleon decay

(Hall, JMR, Unwin, West, unpublished; Huang & Zhao)

Consequences of ADM

Other generic astro signals

both sharing and co-generation generate an initially dominant symmetric $(X+\bar{X})$ component

$$\label{eq:eg} \mathbf{eg} \quad Y_X + Y_{\bar{X}} \sim \frac{Y_X - Y_{\bar{X}}}{\epsilon}$$

must be efficiently removed

$$\Rightarrow$$
 typically involves new light Dark-Sector states

$$\Longrightarrow$$

(work in progress w/West...)

Finally...

Macroscopic DM

There are rich possibilities for getting "macroscopic" DM

eg,

Scalar solitons like Q-balls...

Kusenko, Shaposhnikov, etal

"Nuclear" DM made in process of big-bang-DM-Hardy, Lasenby, JMR, & West: arXiv:1411.3739 & arXiv:1504.05419

🏺 Primordial BHs...

Carr; Bird etal; Garcia Bellido, etal

PBHs from T=0 (quantum) vacuum decay?

Garcia Garcia, Kripendorf, JMR; arXiv:1607.06813



Nobody has reliably computed resulting PBH mass spectrum...!

Conclusions Huge amount of exciting thy/expt awaits!



COMC ON.