

# BSM ON THE LATTICE

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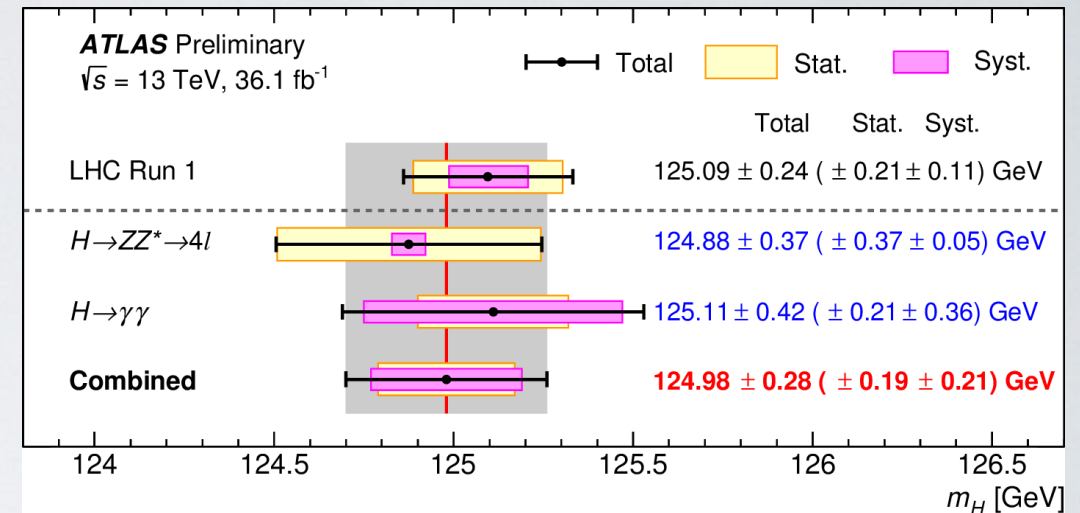
# Higgs: the story so far

- Spin 0, scalar...
- Coupling to the other SM particles proportional to their masses?
- Quantum effect consistent with SM Higgs?
- Is the Higgs Elementary or Composite?

ATLAS and CMS are providing precise measurements of the properties of the Higgs and strong constraints on Beyond SM physics

# THE SM works pretty well: Why BSM?

- Effective description / not UV complete
- EW scale stability / naturalness
- EW vacuum meta-stability
- Flavour problem / undetermined fermion masses & mixing angle / Flavour physics anomalies?
- Matter–antimatter asymmetry
- Dark matter / Dark energy
- New particles?
- The Higgs itself!



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Don't bite off too much:

- Look at a specific phenomenon
- ... in a class of model
- ... that explains something but not everything.



# Why lattice BSM?

- Most models for BSM traditionally based on weakly-coupled/calculable extensions. Experimental bounds are now constraining many models to tight corners of parameter space
- Strongly coupled model requiring non-perturbative dynamics are becoming more popular. In particular for the Lattice: Walking Technicolor and pNGB Composite Higgs
- Phenomenology needs non-perturbative input for strongly coupled models: quantum symmetries, spectrum, low energy constants, ...

# Why lattice BSM?

This talk follows the recent reviews by:

- B. Svetitsky at Lattice 2017 1708.04840
- C. Pica at Lattice 2016 1701.07782
- D. Negradi and A. Patella IJMPA 1607.07638
- T. DeGrand in Rev. Mod. Phys. 1510.05018

I will mostly talk of:

Technicolor / “Dilatonic Higgs”

Higgs as a pseudo-Goldstone boson / “partially Composite top”

Apologies, although lattice investigations exist for all the topics below I will not be able to address (time)

- SUSY
- Extradimensions
- Gauge/gravity duality
- Asymptotic Safety
- Dark Matter
- Lattice constraints to BSM physics

# Composite Higgs: how does it look like?

$\mathcal{L}_{SM-Higgs}$  Just the SM with no Higgs  $SU(3) \times SU(2) \times U(1)$

gauge group  $SU(N)/SP(N)/SO(N)$

$N_f$  fermions in some representation

$$\mathcal{L}_{SD} = -\frac{1}{4}F_{\mu\nu}^2 + i\bar{\psi}\not{D}\psi$$

Higgs impostor: scalar composite

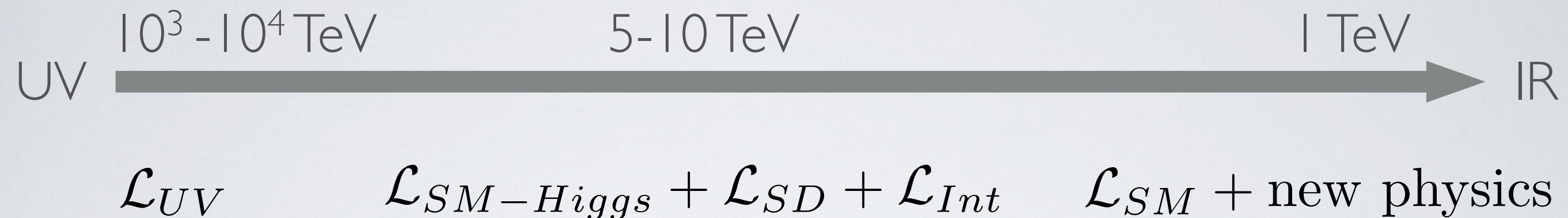
W/Z mass generation, new resonances

$\mathcal{L}_{Int}$  Effective interactions with the SM.  
Generates masses for the fermions

$$\frac{1}{\Lambda_{UV}^2}\bar{q}qO_B, \frac{1}{\Lambda_{UV}^2}qO_F$$



# Composite Higgs: how does it look like?



On the Lattice we only study the SD in isolation

Realistic phenomenology requires taking into account back-reactions from SM and other interactions.

## Our plan:

- Identify quantities which only depend on the new SD
- Compute (small?) corrections from other sectors
- Anyway ruling in or out a realistic model always requires to consider the full setting!

# Composite Higgs: how does it look like?

- If EWSB is due to a new of strongly-interacting sector with fermions, one would expect, in general, composite scalar particles
- To not be excluded by experiments, this scalar states should mimic a SM-like Higgs boson: correct mass and couplings
- This could happen if the composite scalar is a light pseudo-Goldstone boson of some broken symmetry:

Scale invariance symmetry  
(dilaton)  
Walking Technicolor

Flavour symmetry  
(pNGB)  
pNGB Higgs



# Walking technicolor

Break the Standard Model's  $SU(2)_L \times U(1)$  dynamically, without a scalar field

Original prototype: a copy of QCD with  $\Lambda \sim f_\pi \sim v \approx 245 \text{ GeV}$

- Chiral symmetry  $SU(2)_L \times SU(2)_R \rightarrow SU(2)_V$ 
  - ... leaving 3 Nambu–Goldstone bosons  $\pi^\pm, \pi^0$
  - ... which get eaten to give mass to  $W^\pm, Z^0$ .
  - ... The Higgs is the lightest scalar of the sector

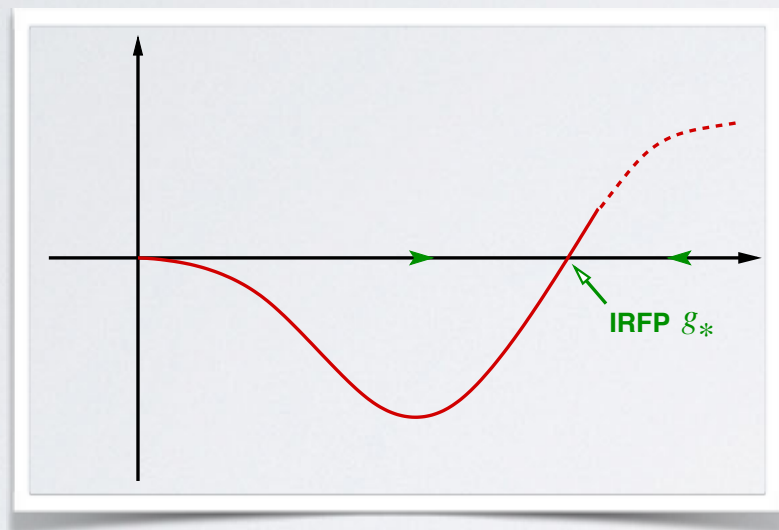
It simply doesn't work:  
You can generate FCNC

$$\mathcal{L}_{Int} = \dots + \frac{\gamma_{ab}}{\Lambda_{UV}^2} \bar{q} T^a q \bar{q} T^b q$$

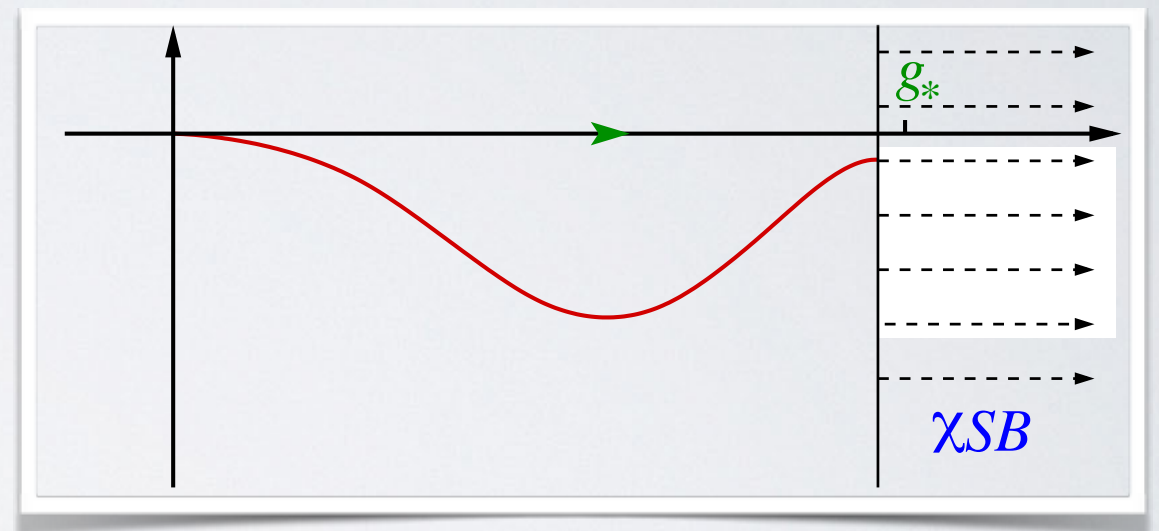
See [1005.5727], [1104.1255]

# Walking technicolor

For a light scalar suppose **APPROXIMATE SCALE INVARIANCE**  
For example by increasing  $N_f$



IR fixed point  $\Rightarrow$  scale invariance



**Approximate** scale invariance

**Light scalar** emerges as pseudo-Goldstone boson of approximate dilatation symmetry.

$m_H$  is protected from UV, like any PGB (and Yukawa couplings  $\propto m_q$ )

Scale separation driven by the anomalous dimension of the FP ( $\gamma$ )

# pNGB Higgs and partial Composite

Higgs is a **pseudo Goldstone boson**, hence naturally light!

Easily recover the **correct SM-like coupling** between the composite Higgs and the electroweak gauge bosons.

Pattern of symmetry breaking  $G_F / H_F$  should be such that the custodial symmetry of the SM is preserved

$$G_F \rightarrow H_F \supseteq SU(2)_L \times SU(2)_R \times U(1)_X$$

To give the correct hyper-charge to all SM fields we need a NGB with

$$\text{Higgs} = (2, 2)_0 \in G_F / H_F$$

4 $(\psi_\alpha, \tilde{\psi}_\alpha)$ <b>Complex</b>	$SU(4) \times SU(4)' / SU(4)_D$
4 $\psi_\alpha$ <b>Pseudoreal</b>	$SU(4) / Sp(4)$
5 $\psi_\alpha$ <b>Real</b>	$SU(5) / SO(5)$

**SU(3) Nf=4 Fund Dirac**

**SU(2) Nf=2 Fund Dirac**

**SU(4) Nf=2.5 2-A Dirac**



# pNGB Higgs and partial Composite

EW interactions break the global symmetry of SD and generate a mass for the composite Higgs, as for charged pions in massless QCD.

EW symmetry **must be broken** via radiative corrections (top)

Finally the top could be partially composite

$G_{HC}$	$\psi$	$\chi$	$G/H$
$SO(7, 9)$	$5 \times \mathbf{F}$	$6 \times \mathbf{Spin}$	$\frac{SU(5)}{SO(5)} \frac{SU(6)}{SO(6)} U(1)$
$SO(7, 9)$	$5 \times \mathbf{Spin}$	$6 \times \mathbf{F}$	
$Sp(4)$	$5 \times \mathbf{A}_2$	$6 \times \mathbf{F}$	$\frac{SU(5)}{SO(5)} \frac{SU(6)}{Sp(6)} U(1)$
$SU(4)$	$5 \times \mathbf{A}_2$	$3 \times (\mathbf{F}, \bar{\mathbf{F}})$	$\frac{SU(5)}{SO(5)} \frac{SU(3) \times SU(3)'}{SU(3)_D} U(1)$
$SO(10)$	$5 \times \mathbf{F}$	$3 \times (\mathbf{Spin}, \bar{\mathbf{Spin}})$	
$Sp(4)$	$4 \times \mathbf{F}$	$6 \times \mathbf{A}_2$	$\frac{SU(4)}{Sp(4)} \frac{SU(6)}{SO(6)} U(1)$
$SO(11)$	$4 \times \mathbf{Spin}$	$6 \times \mathbf{F}$	
$SO(10)$	$4 \times (\mathbf{Spin}, \bar{\mathbf{Spin}})$	$6 \times \mathbf{F}$	$\frac{SU(4) \times SU(4)'}{SU(4)_D} \frac{SU(6)}{SO(6)} U(1)$
$SU(4)$	$4 \times (\mathbf{F}, \bar{\mathbf{F}})$	$6 \times \mathbf{A}_2$	
$SU(5, 6)$	$4 \times (\mathbf{F}, \bar{\mathbf{F}})$	$3 \times (\mathbf{A}_2, \bar{\mathbf{A}}_2)$	$\frac{SU(4) \times SU(4)'}{SU(4)_D} \frac{SU(3) \times SU(3)'}{SU(3)_D} U(1)$

As in TC, **large anomalous dimensions** are advocated to suppress FCNC

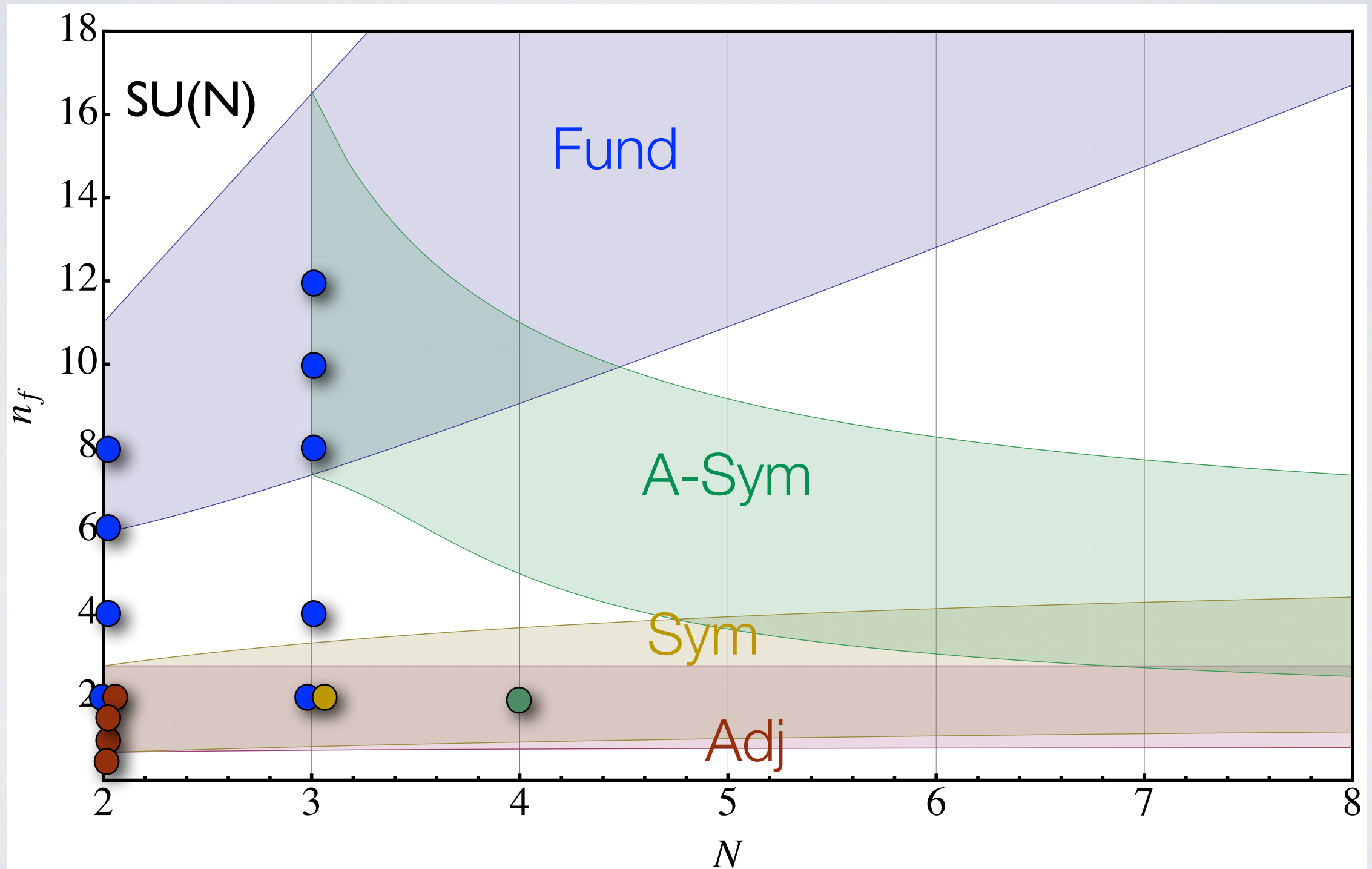
## QUESTIONS FOR DIRECT CALCULATIONS:

- Where is the exact location of the conformal window?
- Inside the conformal window:  
What is the anomalous dimension of the mass/baryonic operators?
- Outside the conformal window:  
 $\sigma$  resonance is light and narrow? Is it a dilaton? What are the couplings to NGBs?
- Generally, how does the spectrum change with the number of fermion fields or when changing the number of colors or fermion representation?
- How big is the S-parameter for models just below the conformal window?

How can I investigate: Many approaches and as many caveats:

- Mass spectrum (finite volume/lattice artefacts)
- Finite-T transition (lattice artefacts)
- Dirac eigenvalues (fit dependant)
- Beta function (finite volume/definition of the coupling)

# SU(N) phase diagram on the lattice





# SU(3) phase diagram on the lattice

**FINDING  $N_f^*$**  — SU(3)/fund (2 loops:  $N_f^* = 8.05$ )

- $N_f = 4n$  popular because of staggered fermions
- $N_f = 12$  — a long-running controversy
- many approaches: scaling of spectrum, finite- $T$  transition, Dirac eigenvalues —
- **Danger:** slow running (spontaneous  $\chi$ SB) is very similar to *no* running (fixed point)
  - possible fixed point at IR both scale  $L$  and  $a$  could be at strong coupling far from the continuum limit in UV
- **RG** was born for this purpose — compare  $L_1$  to  $L_2$  and obtain the beta function
- no scale  $\Lambda$  ( $1/L < \Lambda < 1/a$ )  $\Rightarrow$  everything is a function of  $a/L$ 
  - $\Rightarrow$  continuum extrapolation equivalent to  $L \rightarrow \infty$

# SU(3) phase diagram on the lattice

**Latest on**  $N_f = 12$  —  $\beta$  function from gradient flow

[Hasenfratz & Schaich 1610.10004, LatHiggs Fodor et al. 1607.06121, Lin, Ogawa, Ramos 1510.05755]

- **Existence** (or not) of fixed point  $g_*$  is universal together with its critical exponents
- **Location** is not. Neither is the  $\beta$  function.

– Varies with **scheme** for defining  $g$ :

e.g.: Schrodinger functional vs. gradient flow; parameter  $c$  in gradient flow

$$g_{GF}^2 = \frac{128\pi^2}{3(N^2 - 1)} \langle t^2 E(t) \rangle \quad \text{at} \quad \sqrt{8t} = cL$$

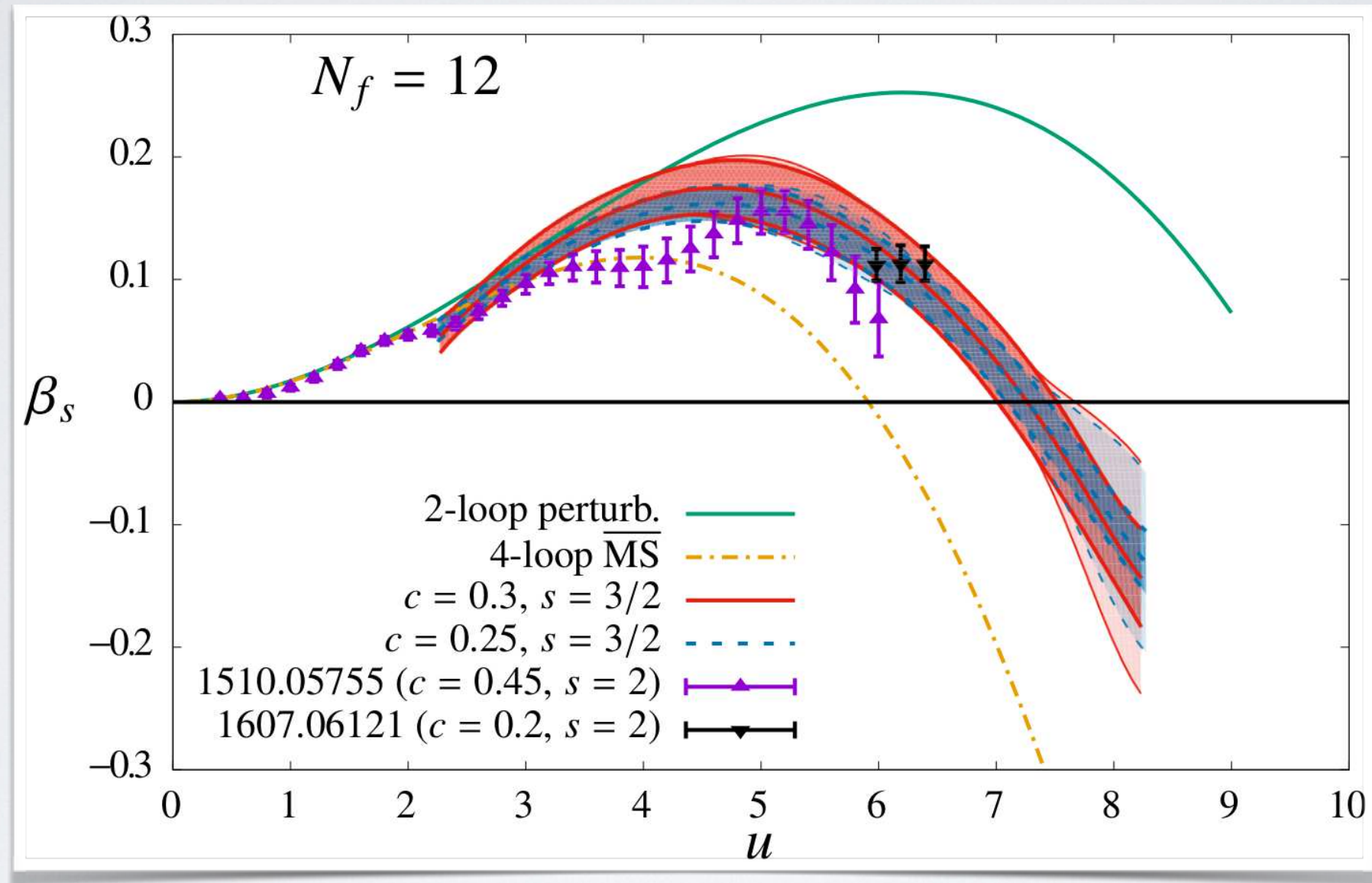
– **Shouldn't** depend on:

discretization/improvement of action (extrap.  $a \rightarrow 0$ );

disc./imp. of  $E$ ;

scale factor  $s \equiv L_2/L_1$

# SU(3) phase diagram on the lattice



The controversy has a direct consequences also on smaller  $N_F$



# SU(3) phase diagram on the lattice

Going down to  $N_f = 8$

Recall two-loops:  $N_f^* = 8.05$ .

[LSD Appelquist et al. 1601.04027 + recent][LatKMIY.Aoki et al. 1610.07011]

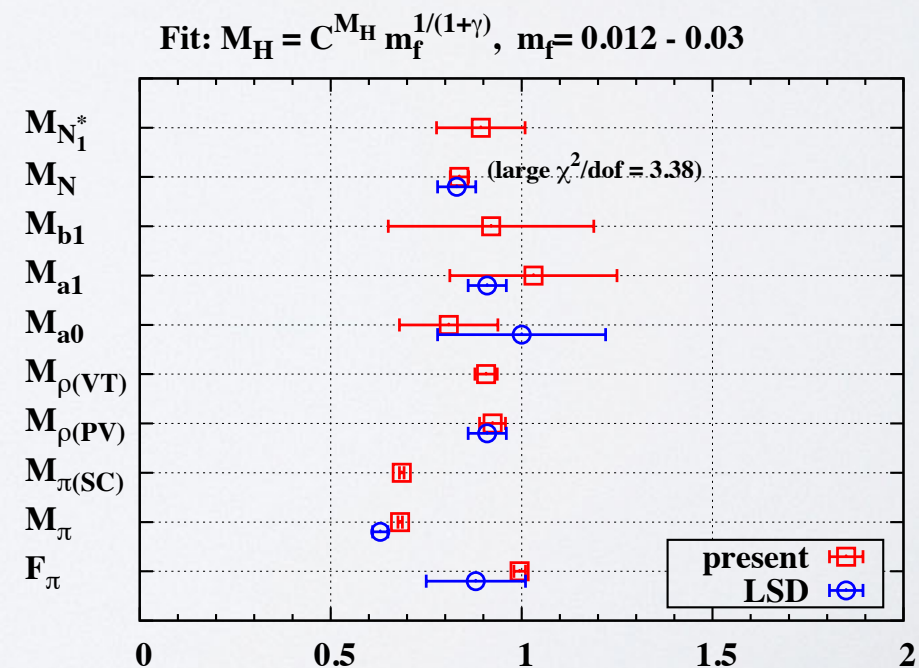
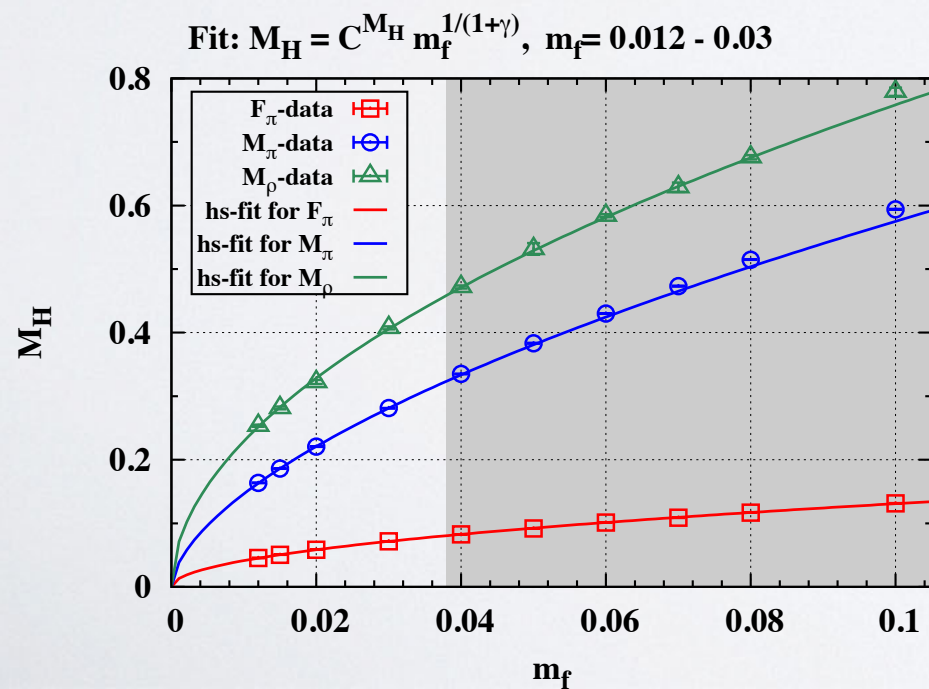
The claim is that  $N_f = 8$  *walks*.

- Study  $m_f$  dependence of various quantities:
- As  $m_f \rightarrow 0$ , find  $\chi_{SB}$ :  $f_{\pi\pi} \rightarrow const$  and massless  $\pi$
- Hyperscaling of all masses (except the  $\pi$ ) for a large range of  $m_f$ :

Near a fixed point  $M_H \sim m^{1/(1+\gamma)}$ ,  $\gamma \approx 1$

- Light scalar — degenerate with  $\pi$  over large range of  $m_f$

As  $m_{\pi\pi} \rightarrow 0$ , where does the scalar end up?



# Forcing a theory to walk

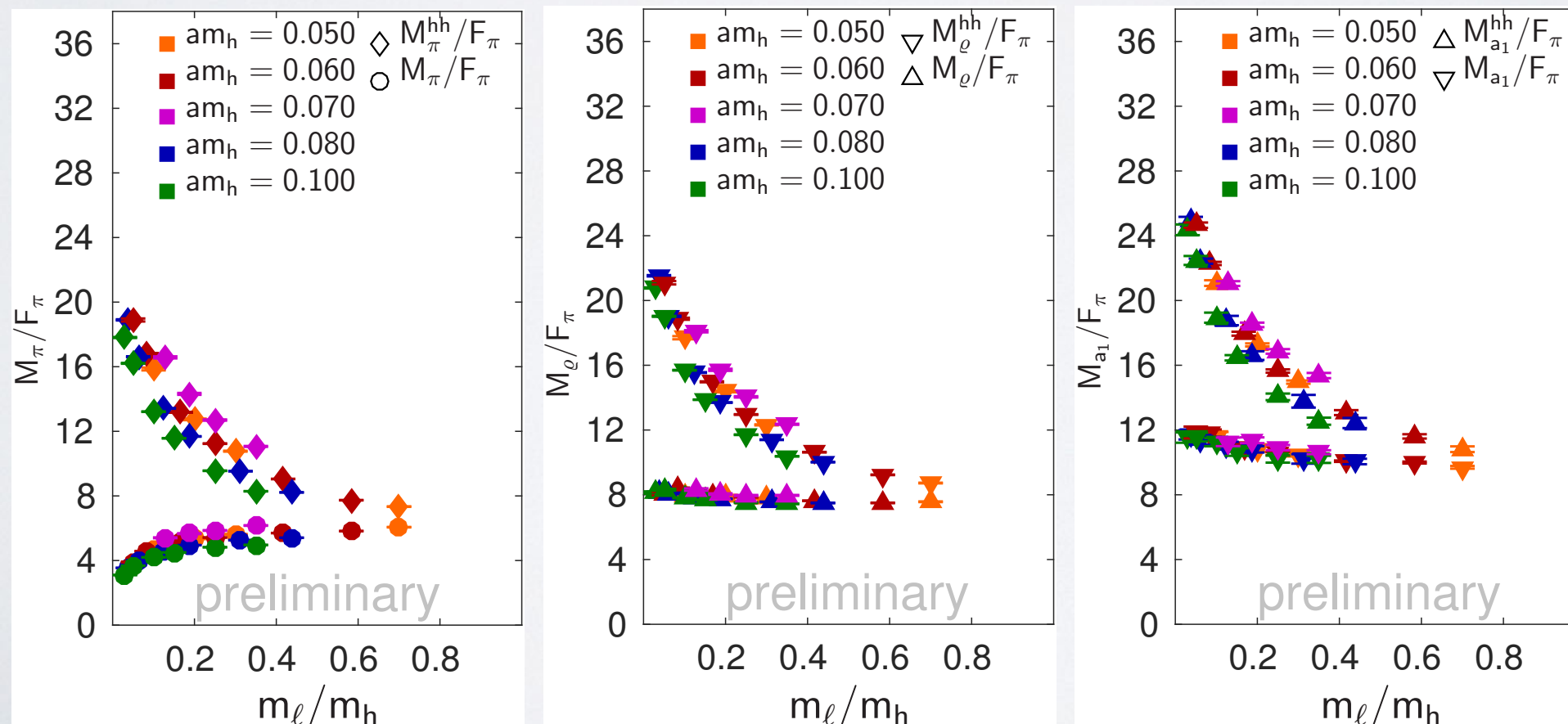
**And up again**  $N_f = 8$  heavy + 4 light

[Hasenfratz, Rebbi, Witzel 1609.01401 + recent,]

- SU(3) with  $N_f = 12$ , nearly conformal at scale  $\Lambda$ .
- Lift 8 flavors with  $m_h < \Lambda$ .
- Low-energy theory sees 4 flavors.
- Appeal to near-conformality to give light Higgs as above.

Note  $m_h$  is artificial.

Extra Goldstones ( $N_f = 4$ ) still to be addressed, maybe *composite Higgs* [Ma & Cacciapaglia].



A summary of the results of last lattice conference.

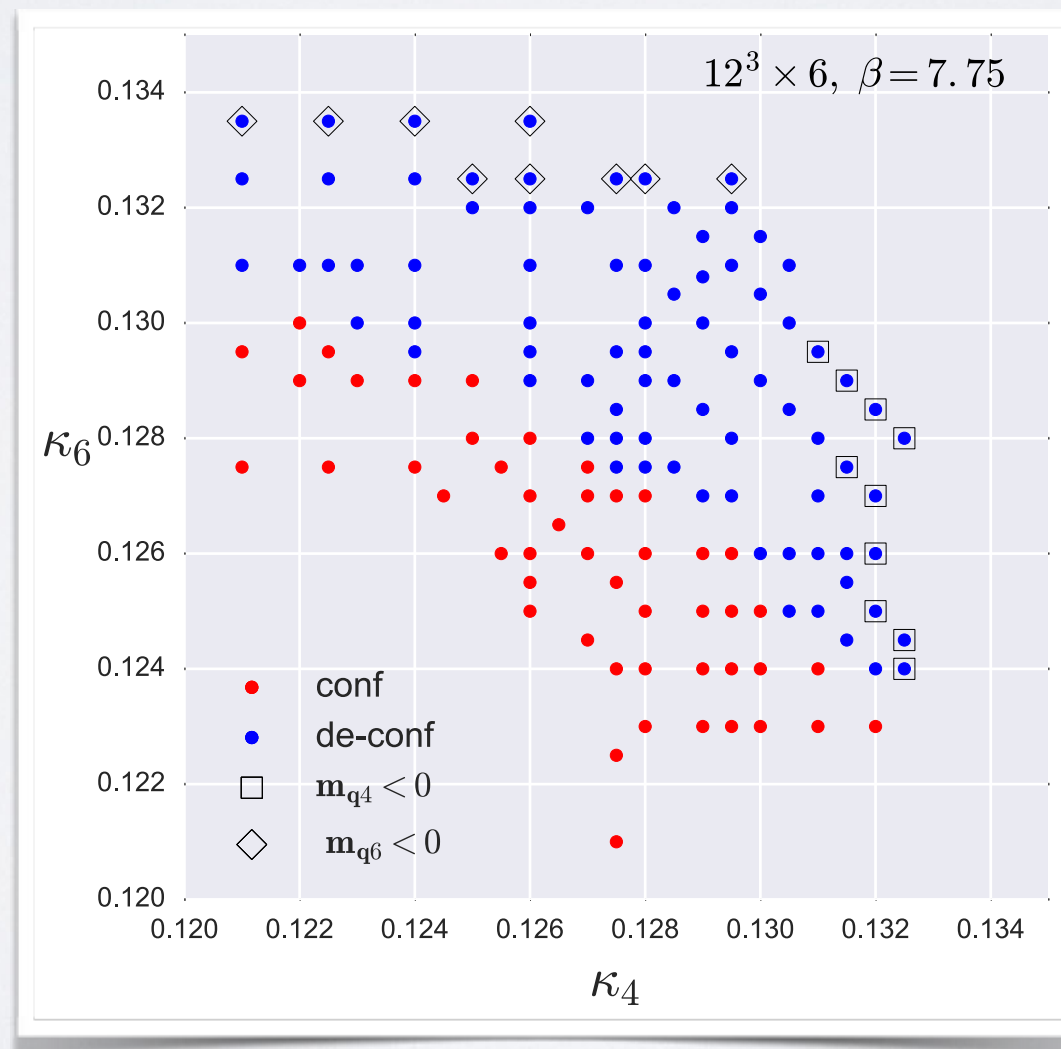
- $SU(2)/\text{fund}, N_f = 8$  — in the conformal window [Leino et al. 1701.04666]
- $SU(2)/\text{adj}, N_f = 1, 2$  — straddling the sill [Athenodorou et al. 1605.04298][Bergner et al. 1610.01576]
- $SU(2)/\text{adj}, N_f = 2 + NJL$  — induce walking with NJL term [Rantaharju, Pica, Sannino 1704.0397]
- $SU(3)/\text{sextet}, N_f = 2$  — may walk, with light Higgs  
[Fodor et al. 1506.06599, 1601.03302 (staggered)][Hansen, Drach, Pica 1705.11010 (unimp. Wilson)]



# Starting on the partial composite top

Current Work on Composite Higgs / Partially Composite Top  
TACO [DeGrand et al.]

$SU(4)$  gauge with  $\{N_f = 2 \times 6 \text{ and } 2 \times 4\}$  on the way to The Real Thing:  $\{5 \times 6 \text{ (Majorana) and } 3 \times 4\}$



## 1. Technicolor

1. The effort to nail down the sill of the conformal window continues,
2. If you know you're below the sill, it makes sense to look for walking as a mechanism for a light scalar that might be the Higgs boson.
3. Has your walking theory really produced a light, dilatonic Higgs and a low scale for the Higgs vev?

## 2. Composite Higgs and partially composite top quark

1. Multi-representation models are a whole new area for lattice simulations.
2. There are many opportunities for lattice calculations of low-energy constants. Unfortunately, they will always depend on unknown mixing parameters that come from yet higher energies.

## 3. Phenomenological constraints from these theories might be premature, as the models aren't perfect.