Flavour Anomalies on the Eve of the Run-2 Verdict

Diego Guadagnoli LAPTh Annecy (France)



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0 A first qualitative observation

A whole range of $b \rightarrow s$ measurements involving a $\mu\mu$ pair display a consistent pattern: Exp < SM



Flavor anomalies ····· 0 A first qualitative observation A whole range of $b \rightarrow s$ measurements involving a $\mu\mu$ pair display a consistent pattern: Exp < SMLCSR Lattice - Data LCSR Lattice - Data Lattice - Data LCSR c^4/GeV^2 $dB/dq^2 [10^{-8} \times c^4/GeV^2]$ $B^+ \rightarrow K^+ \mu^+ \mu^-$ 5E $B^0 \rightarrow K^0 \mu^+ \mu^-$ LHCb $B^+ \rightarrow K^{*+} \mu^+ \mu^-$ Gev LHCb LHCb LHCb 2014 × dB/dq^{2} [10⁻ 10 dB/dq^{2} [10⁻⁸ 0 00 $q^{2} [GeV^{2/}c^{4}]$ $q^{20} q^{20} [\text{GeV}^2/c^4]$ 5 10 15 5 10 15 0 $q^{2} [GeV^{2}/c^{4}]$ 10 15 5 9E $\mathrm{dB}(B_s^0 \to \phi \mu \mu)/\mathrm{d}q^2 \ [10^{-8} \mathrm{GeV}^{-2} c^4]$ 1.6 LHCb $\frac{{\rm d} {\cal B}}{{\rm d} q^2} ~~[10^{-7}~{\rm GeV}^{-2}]$ 1.4 SM pred. ဖ 5 1.2 LHCb 2015 -Data 6 **Detmold+Meinel** 5 1.0 LHCb 2015 $\Lambda_{\rm b} \rightarrow \Lambda \ \mu^+ \mu^-$ 0.8 0.6 0.4 0.2 10 15 5

0.0

0

 $q^2 \,[{
m GeV}^2/c^4]$

20

15

10

 $q^2 \, \, [{
m GeV}^2]$

õ

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$$R_{K} = \frac{BR(B^{+} \rightarrow K^{+} \mu \mu)_{[1,6]}}{BR(B^{+} \rightarrow K^{+} e e)_{[1,6]}} = 0.745 \cdot (1 \pm 13\%)$$
(2.6 σ effect)

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2 $BR(B_s \rightarrow \varphi \mu\mu): >3\sigma$ below SM prediction. Same kinematical region $m_{\mu\mu}^2 \in [1, 6]$ GeV² Initially found in 1/fb of LHCb data, then confirmed by a full Run-I analysis (3/fb)



B

We know that BR measurements suffer from large f.f. uncertainties. However, here's a clean quantity:

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 $B \rightarrow K^* \mu \mu$ angular analysis: discrepancy in one combination of the angular expansion coefficients, known as P'_{5}

 $m{B}
ightarrow m{K}^* \, \mu\mu$ angular analysis: The P'₅ anomaly

B

- From LHCb's full angular analysis of the decay products in $B \rightarrow K^* \mu\mu$, one can construct observables with limited sensitivity to form factors.









b → c data

There are long-standing discrepancies in $b \rightarrow c$ transitions as well.

$$R(D^{(*)}) = \frac{BR(B \rightarrow D^{(*)}\tau\nu)}{BR(B \rightarrow D^{(*)}\ell\nu)} (\text{with } \ell = e,\mu)$$

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- R_{κ} significance fairly low. Yet interesting that all $b \rightarrow s \mu \mu$ modes go in a consistent direction
- Focusing for the moment on the $b \rightarrow s$ discrepancies
 - **Q1:** Can we (easily) make theoretical sense of data?
 - **Q2:** What are the most immediate signatures to expect ?

$\mathsf{B} ightarrow \mathsf{K}(*) \,\ell\ell$ decays: basic theory considerations

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$B \to K(\mbox{*}) \ \ell \ell$ decays: basic theory considerations



Caveat

• In practice, the short-distance part is dominated by the top loop, because of the large top mass:

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$$\frac{m_t^2}{m_W^2} = O(1)$$

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 \Longrightarrow "Hard" GIM breaking







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• Yes we can. Consider the following Hamiltonian

$$H_{\rm SM+NP}(\bar{b} \rightarrow \bar{s}\mu\mu) = -\frac{4G_F}{\sqrt{2}} V_{tb}^* V_{ts} \frac{\alpha_{\rm em}}{4\pi} \left[\bar{b}_L \gamma^\lambda s_L \cdot \left(C_9^{(\mu)} \bar{\mu} \gamma_\lambda \mu + C_{10}^{(\mu)} \bar{\mu} \gamma_\lambda \gamma_5 \mu \right) \right]$$

Concerning Q1: can we easily make theoretical sense of these data?
• Yes we can. Consider the following Hamiltonian

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- Advocating the same $(V A) \times (V A)$ structure also for the corrections to $C_{9,10}^{SM}$ (in the $\mu\mu$ -channel only!) would account for:
 - R_{κ} lower than 1
 - $b \rightarrow s \mu \mu$ BR data below predictions
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A fully quantitative test requires a global fit.

new physics contributions to the Wilson coefficients. We find that the by far largest decrease in the χ^2 can be obtained either by a negative new physics contribution to C_9 (with $C_9^{\text{NP}} \sim -30\% \times C_9^{\text{SM}}$), or by new physics in the $SU(2)_L$ invariant direction $C_9^{\text{NP}} = -C_{10}^{\text{NP}}$, (with $C_9^{\text{NP}} \sim -12\% \times C_9^{\text{SM}}$). A positive NP contribution to C_{10} alone would also improve the fit, although to a lesser extent. [Altmannshofer, Straub, EPJC '15]

For analogous conclusions, see also [Ghosh, Nardecchia, Renner, JHEP '14]

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 - $H_{\rm NP} = G \bar{b}'_{L} \gamma^{\lambda} b'_{L} \bar{\tau}'_{L} \gamma_{\lambda} \tau'_{L}$ with $G = 1/\Lambda_{\rm NP}^{2} \ll G_{F}$ expected e.g. in partial-compositeness frameworks

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 - They need to be rotated to the mass eigenbasis

$$mass \\ b'_{L} \equiv (d'_{L})_{3} = (U_{L}^{d})_{3i} (d_{L})_{i} \\ \tau'_{L} \equiv (\ell'_{L})_{3} = (U_{L}^{\ell})_{3i} (\ell_{L})_{i}$$

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 This rotation induces <u>LUV and LFV</u> effects *b* '_L = (*d* '_L)₃ = (*U*^{*d*}_L)_{3i} (*d*_L)_i *t* '_L = (*t* '_L)₃ = (*U*^{*d*}_L)_{3i} (*t*_L)_i

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$$\square \qquad \frac{BR(B^+ \to K^+ \mu e)}{BR(B^+ \to K^+ \mu \mu)} = \frac{|\delta C_{10}|^2}{|C_{10}^{SM} + \delta C_{10}|^2} \cdot \frac{|(U_L^{\ell})_{31}|^2}{|(U_L^{\ell})_{32}|^2} \cdot 2$$

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$$\boxed{ \begin{array}{c} \blacksquare \\ \hline \\ \blacksquare \\ \hline \\ \hline \\ \hline \\ \hline \\ \hline \\ \hline \\ \hline \\ \\ \hline \\ \\ \hline \\ \\ \hline \\ \\ \hline \\ \hline \\ \\ \hline \\ \hline \\ \hline \\ \\ \hline \\ \hline \\ \hline \\ \\ \hline \hline \\ \hline$$

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Actually, the expected ballpark of LFV effects can be predicted from $BR(B \rightarrow K \mu\mu)$ and the R_{κ} deviation alone [Glashow et al., 2015]

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• Being defined above the EWSB scale, our assumed operator

 $\bar{b}'_L \gamma^{\lambda} b'_L \bar{\tau}'_L \gamma_{\lambda} \tau'_L$

must actually be made invariant under $SU(3)_c \times SU(2)_L \times U(1)_Y$



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$$\begin{array}{c} \mathsf{SU(2)}_{\mathsf{L}} \\ \mathsf{v}^{\mathsf{T}}{}_{\mathsf{L}} \\ \mathsf{inv.} \end{array} \quad \begin{cases} \bullet \quad \bar{Q}'{}_{\mathsf{L}} \gamma^{\lambda} Q'{}_{\mathsf{L}} \quad \bar{L}'{}_{\mathsf{L}} \gamma_{\lambda} L'{}_{\mathsf{L}} \\ \bullet \quad \bar{Q}'{}_{\mathsf{L}}^{i} \gamma^{\lambda} Q'{}_{\mathsf{L}}^{j} \quad \bar{L}'{}_{\mathsf{L}}^{j} \gamma_{\lambda} L'{}_{\mathsf{L}} \\ \bullet \quad \bar{Q}'{}_{\mathsf{L}}^{i} \gamma^{\lambda} Q'{}_{\mathsf{L}}^{j} \quad \bar{L}'{}_{\mathsf{L}}^{j} \gamma_{\lambda} L'{}_{\mathsf{L}} \\ \bullet \quad \bar{Q}'{}_{\mathsf{L}}^{i} \gamma^{\lambda} Q'{}_{\mathsf{L}}^{j} \quad \bar{L}'{}_{\mathsf{L}}^{j} \gamma_{\lambda} L'{}_{\mathsf{L}} \\ \bullet \quad \mathsf{Ialso charged-current int's]} \\ \mathsf{x} U(1)_{\mathsf{v}} \end{cases}$$

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Thus, the generated structures are all of:

$$t't'\nu'_{\tau}\nu'_{\tau}, \quad b'b'\nu'_{\tau}\nu'_{\tau},$$
$$t't'\tau'\tau', \quad b'b'\tau'\tau'$$

 $\left\{ \begin{array}{c} \bullet \quad \bar{Q}'_{L} \gamma^{\lambda} Q'_{L} \quad \bar{L}'_{L} \gamma_{\lambda} L'_{L} \\ \bullet \quad \bar{Q}'_{L}^{i} \gamma^{\lambda} Q'_{L}^{j} \quad \bar{L}'_{L}^{j} \gamma_{\lambda} L'_{L}^{i} \end{array} \right.$

SU(2)_L

inv.

[neutral-current int's only]

[also charged-current int's]



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See.

and a second second second

Bhattacharya, Datta, London,

Shivashankara, PLB 15

[neutral-current int's only]



After rotation to the mass basis (unprimed), the last structure contributes to $\Gamma(b \rightarrow c \ \tau \ v)$

i.e. it can explain deviations on R(D())*

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under $SU(3)_c \times SU(2)_L \times U(1)_Y$





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• $\bar{Q}_{I}^{\prime i} \chi^{\lambda} Q_{I}^{\prime j} \bar{L}_{I}^{\prime j} \chi_{\lambda} L_{I}^{\prime i}$

[neutral-current int's only] [also charged-current int's] $t'b'\tau'\nu'$

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Bhattacharya, Datta, London,

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See

After rotation to the mass basis (unprimed), the last structure contributes to $\Gamma(b \rightarrow c \tau v)$ i.e. it can explain deviations on R(D(*))

But this coin has a flip side.

Through RGE running, one gets also LFU-breaking effects in $\tau \rightarrow \ell v v$ (tested at per mil accuracy)

SU(2)_L

inv.

Such effects "strongly disfavour an explanation of the R(D(*)) anomaly model-independently"

^{Feruglio,} Paradisi, Pattori, 2016

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And yes they are! See: [Greljo-Isidori-Marzocca] [Faroughy-Greljo-Kamenik]



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with any of the following transformation properties under the SM gauge group:

- $SU(3)_c$: 1 or 3 (\rightarrow "leptoquark")
- SU(2)₁: 1 or 2 or 3

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Recap of model-building attempts focused on models accounting for $R_{\kappa} \& R(D(*))$






















Measure more LUV ratios:
$$R_{K^*}$$
, R_{ϕ} , R_{X_s} , $R_{K_0(1430)}$, R_{f_0} Hiller, Schmaltz, JHEP 2015Interesting test:define $X_H \equiv \frac{R_H}{R_K}$, with $H = K^*$, ϕ , X_s , $K_0(1430)$, f_0

Deviations from unity in the double ratios X_{μ} can only come from RH currents

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Hiller, Schmaltz, JHEP 2015

LHCb, 1612.06764

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Zwicky, '14

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D. Guadagnoli, Flavour anomalies

Hiller, Schmaltz, JHEP 2015

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Recently, LHCb measured BR($B^{+} \rightarrow K^{+} \mu \mu$) including an accurate parameterization of the LD component in the cc region

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Measure $m_{\mu\mu}$ spectrum, including the $c\overline{c}$ resonances as a sum of BW, and fit 'em all Method:

Result: BR compatible with previous measurements, and (again) smaller than SM

What's the BR result for q^2 in [1, 6] GeV²?

Hiller, Schmaltz, JHEP 2015

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LHCb, 1612.06764













But LQCD calculation of $B \rightarrow \gamma$ f.f.'s required



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ATTENT (1111)

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 - Experiments: Results are consistent between LHCb and B factories.
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- Theory: EFT makes sense rather well of data. But hard to find convincing UV dynamics
- Timely to pursue further tests.

Examples:

- more measurements of R_{κ}
- more LUV quantities
- other observables sensitive to $C_{g} \& C_{10}$