

# Probing Particle Physics with Gravitational Wave Detectors

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## A couple of examples

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# The last weasel that doubted that there is Physics Beyond the Standard Model



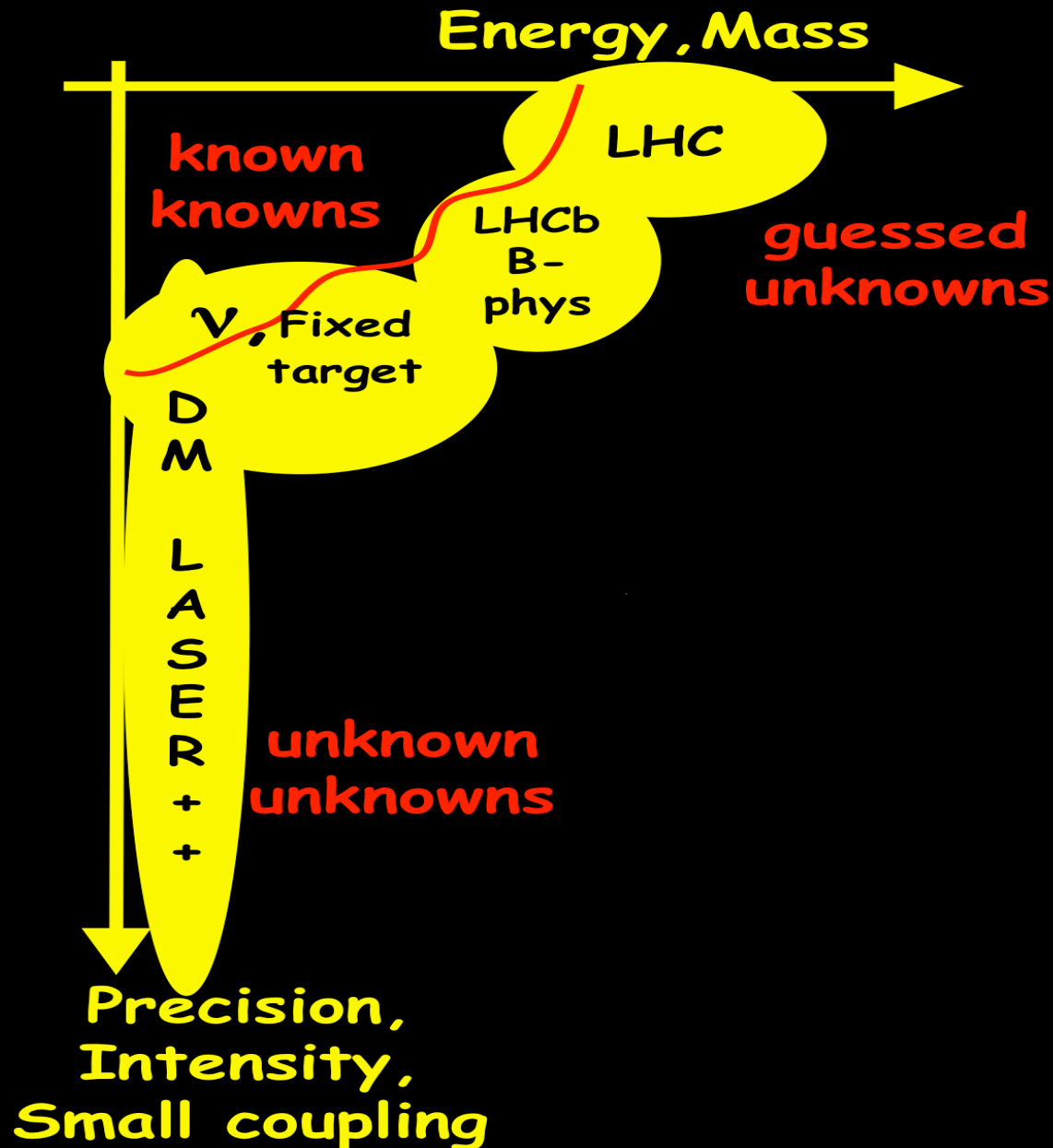
**i** The Cern stone marten, secured for inclusion in the Rotterdam Natural History Museum's Dead Animal Tales exhibition. Photograph: Kees Moeliker

From The Guardian website

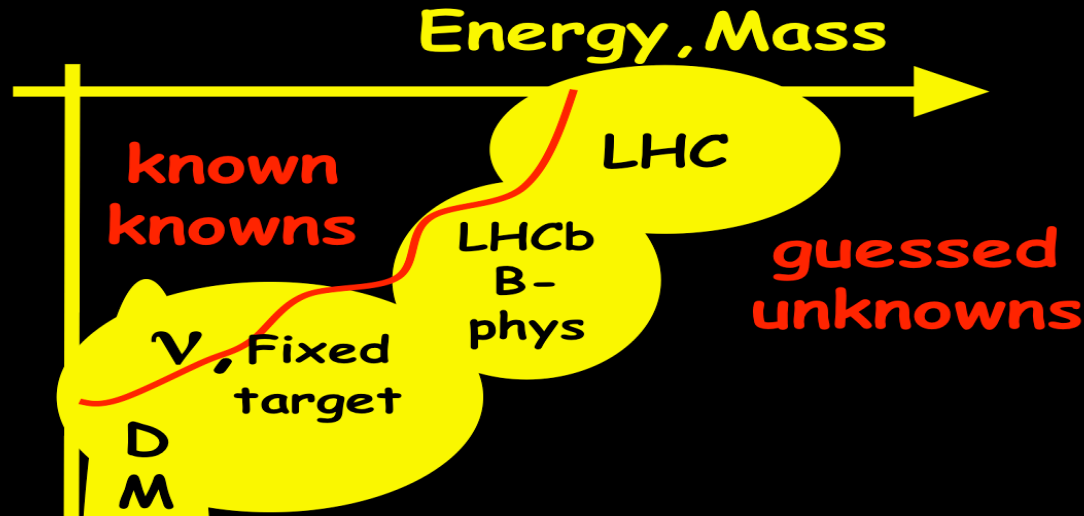


Where is the  
New Physics?

# Exploring is (at least) 2 dimensional



# Exploring is (at least) 2 dimensional



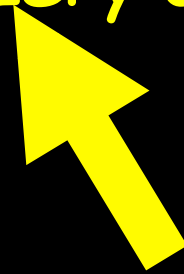
**THE DARK  
SECTOR!!!**

Precision,  
Intensity,  
Small coupling



# Running through walls

(searching for WISPy domain walls)

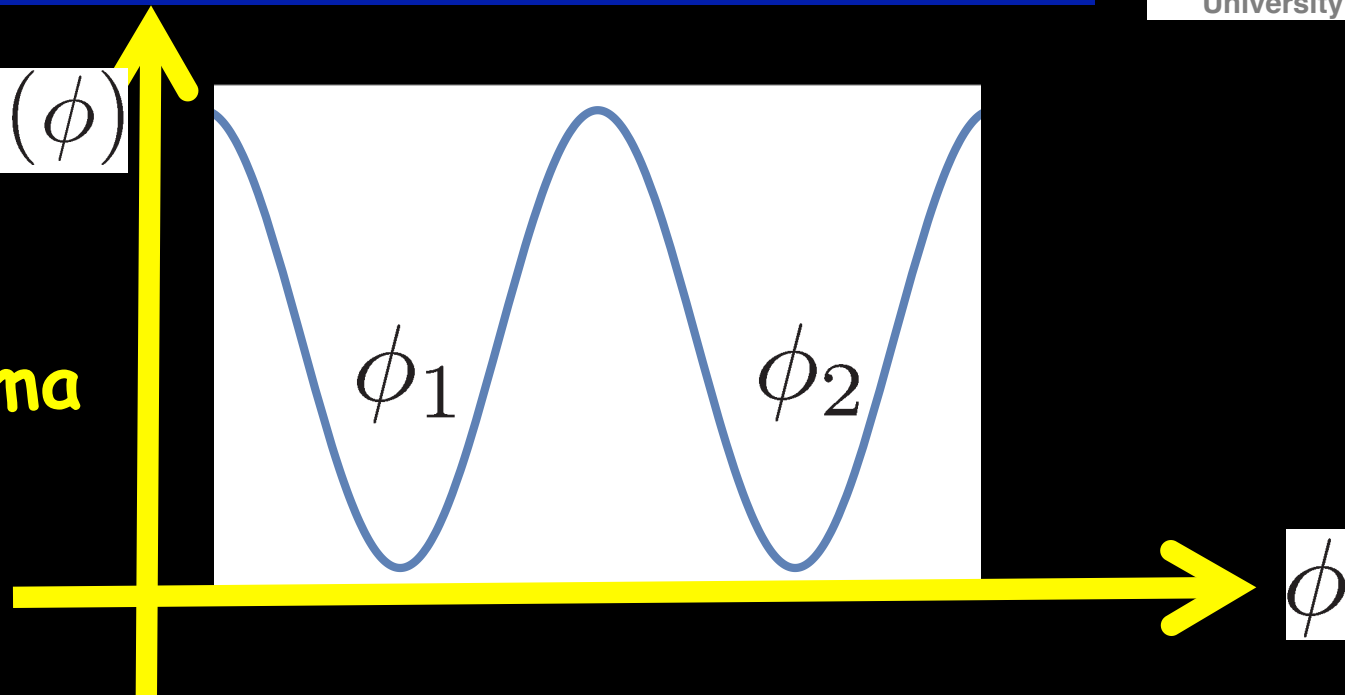


Weakly interacting sub-eV particle

# A Domain Wall

$V(\phi)$

Potential  
with two minima



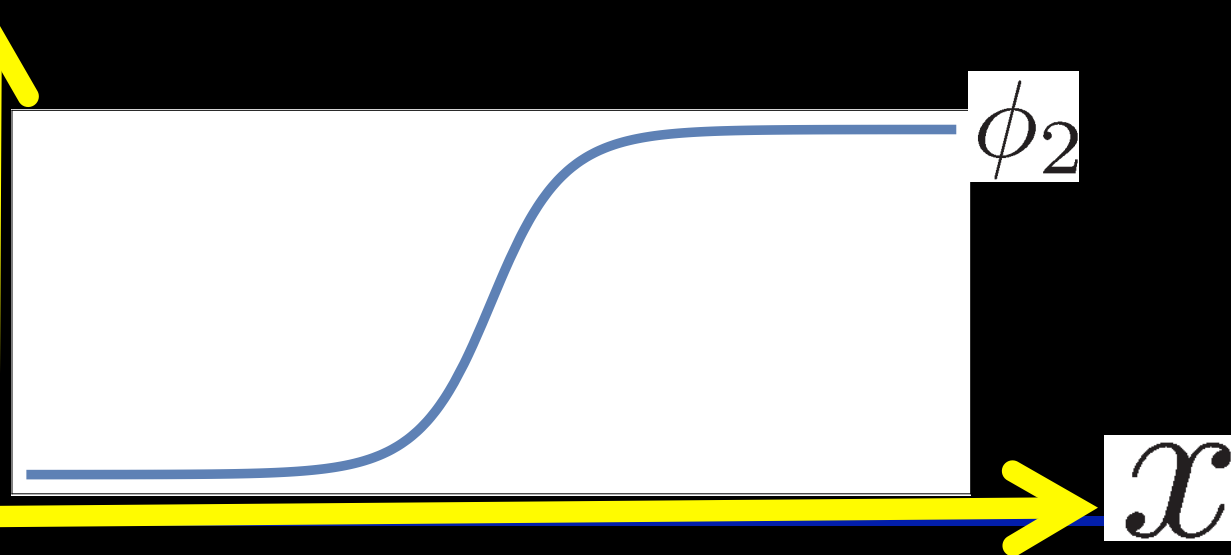
Domain wall  
from side 1  
To side 2

$\phi$

$\phi_1$

$\phi_2$

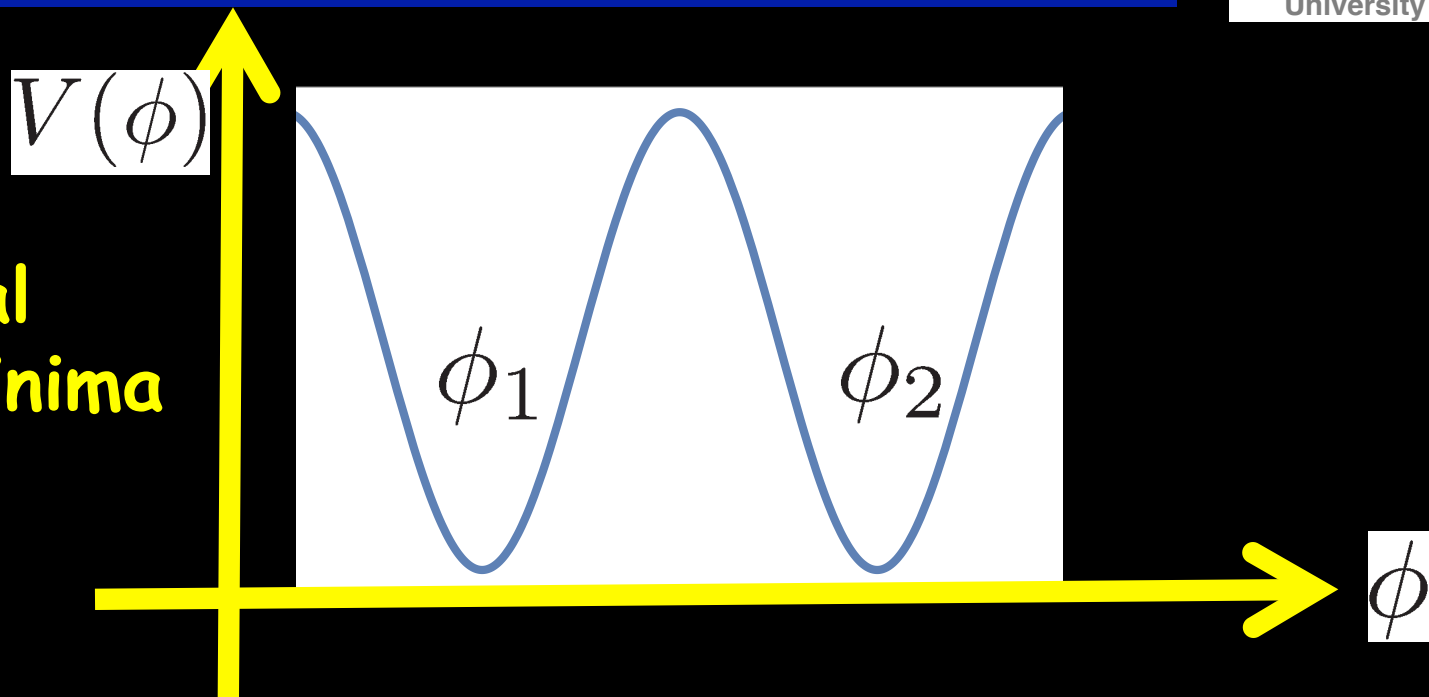
$x$



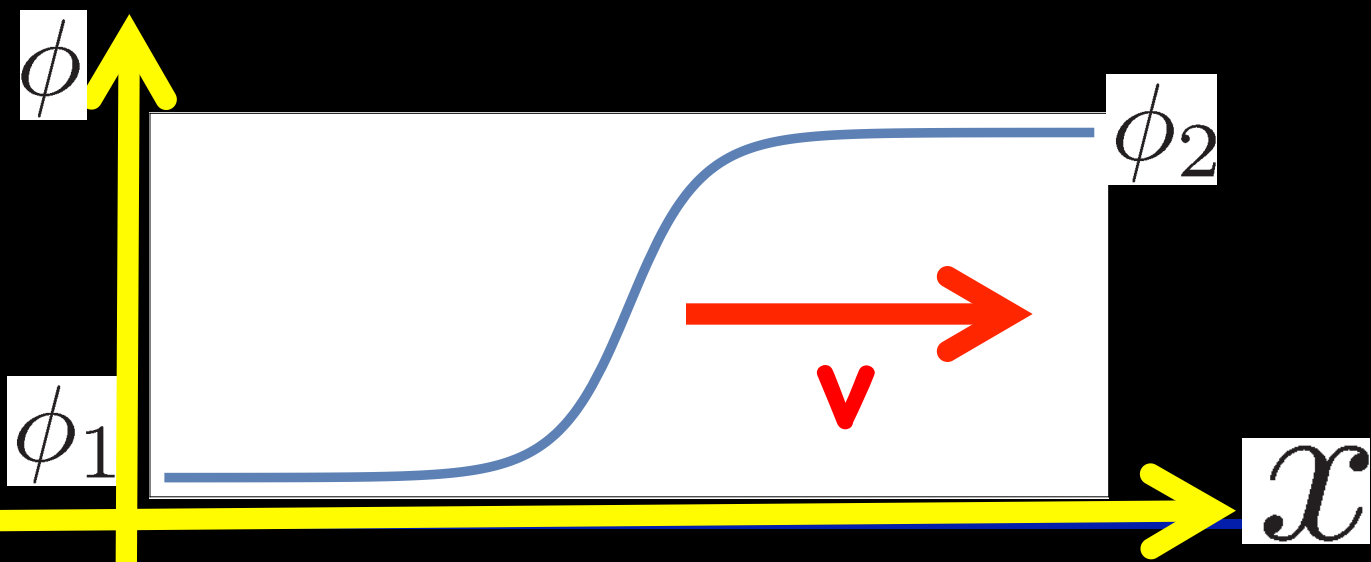


# A Domain Wall

Potential  
with two minima



Domain wall  
from side 1  
To side 2



# Adventureous assumption

- Domain walls (significantly) contribute to DM in galaxy
- This requires some pushing...

$$\frac{f}{N_\phi} \lesssim \text{TeV} \times \left( \frac{L}{10^{-2} \text{Ly}} \right)^{1/2} \left( \frac{\text{neV}}{m} \right)^{1/2} \left( \frac{\rho_{\text{DW}}}{\rho_{\text{DM}}} \right)^{1/2}$$

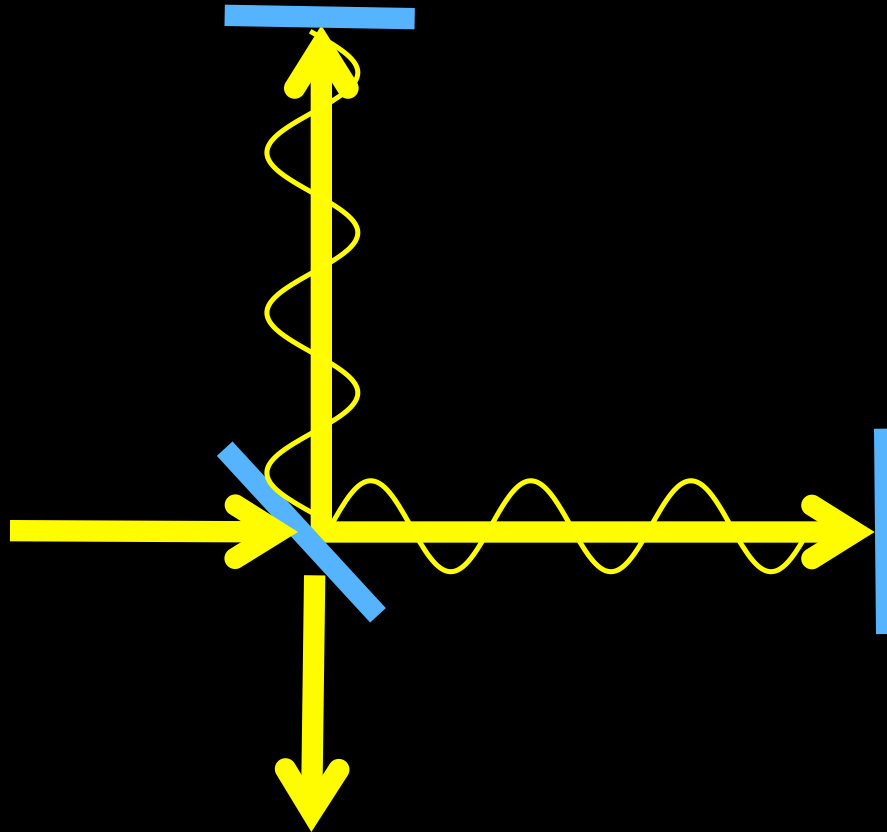
→ Event rate

$$\text{Event Rate} \sim \frac{1}{10 \text{ years}} \left( \frac{10^{-2} \text{Ly}}{L} \right) \left( \frac{v}{10^{-3}} \right)$$

# aLIGO

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- Has detected gravitational waves!!
- Is an Interferometer



Interference pattern

# Causing a phase shift

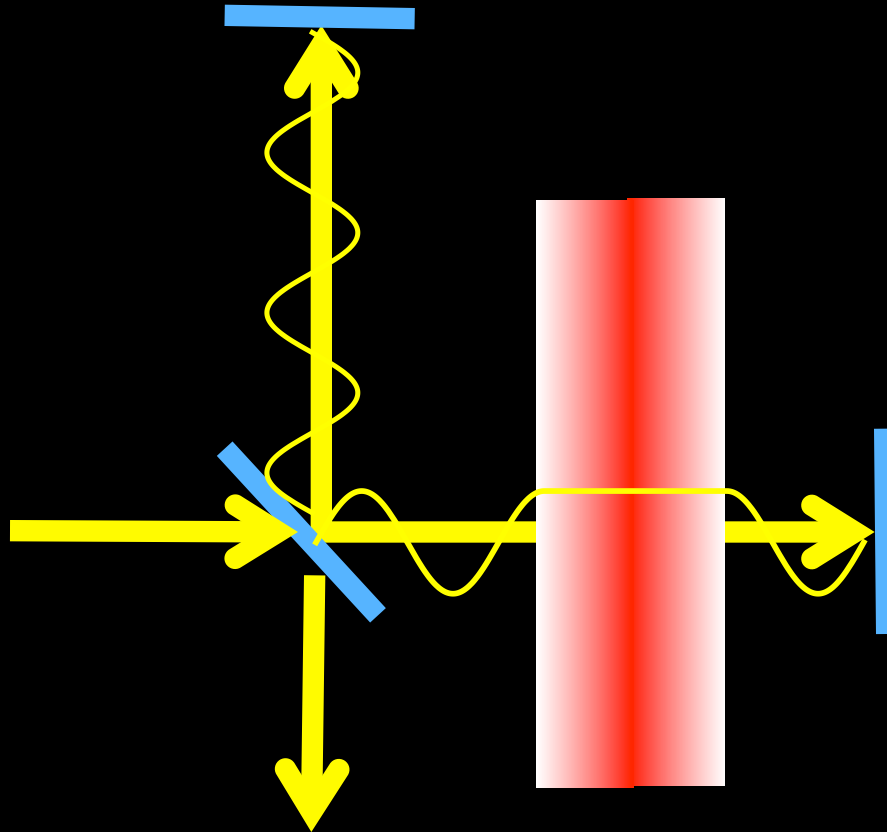
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- Interaction inside wall creates photon mass

$$\mathcal{L}_A = -\frac{1}{4}F^{\mu\nu}F_{\mu\nu} - \frac{1}{2}m_{0,\gamma}^2 \sin^2 \left( \frac{N_A \phi}{f} \right) A^\mu A_\mu$$

# aLIGO

- Has detected gravitational waves!!
- Is an Interferometer



— Interference pattern **changed** —

# Signal shapes

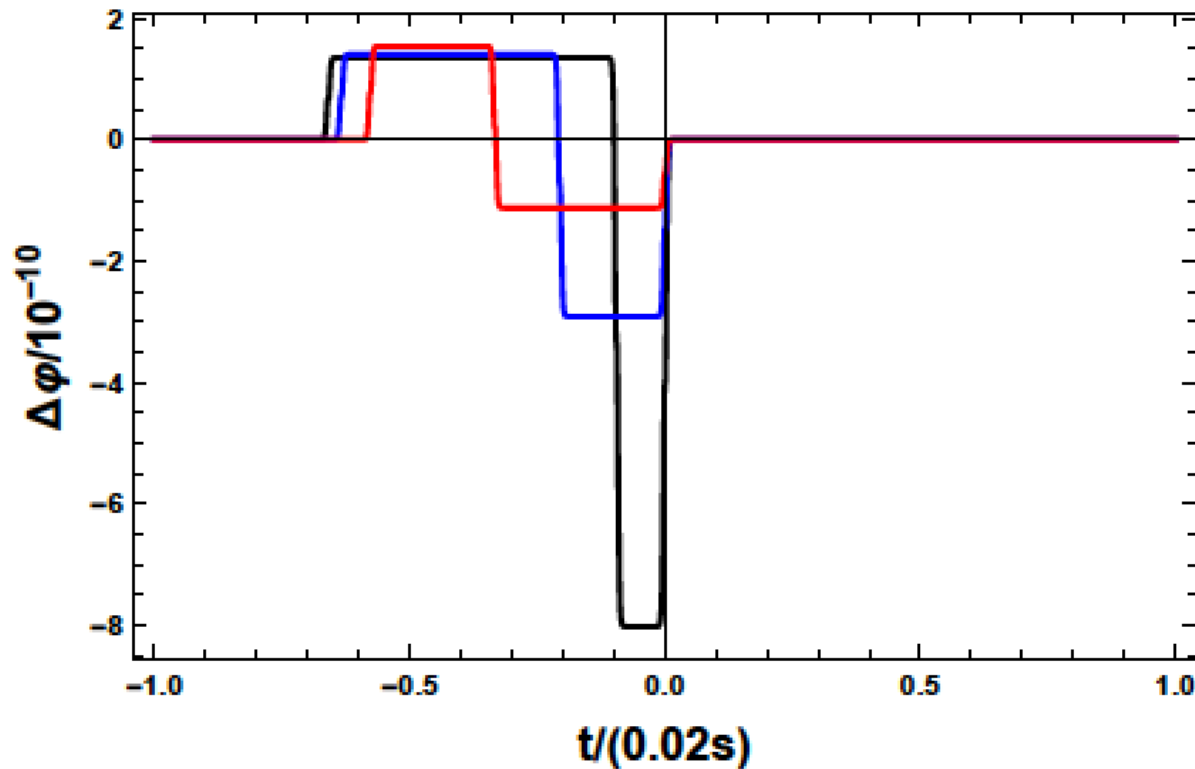


FIG. 6:  $L = 4000$  m,  $\omega \approx 1$  eV,  $m = 10$  neV,  $m_{\gamma,0} = 1$  neV,  $N_A/N_\phi = 1$ ,  $\alpha = \pi/2.2, \pi/2.5, \pi/3$  (black, blue, red),  $v$  chosen such that signal has roughly a length of  $0.02s \sim 1/(50$  Hz) this corresponds to  $v = 1 \times 10^{-3}$ .

# Signal shapes

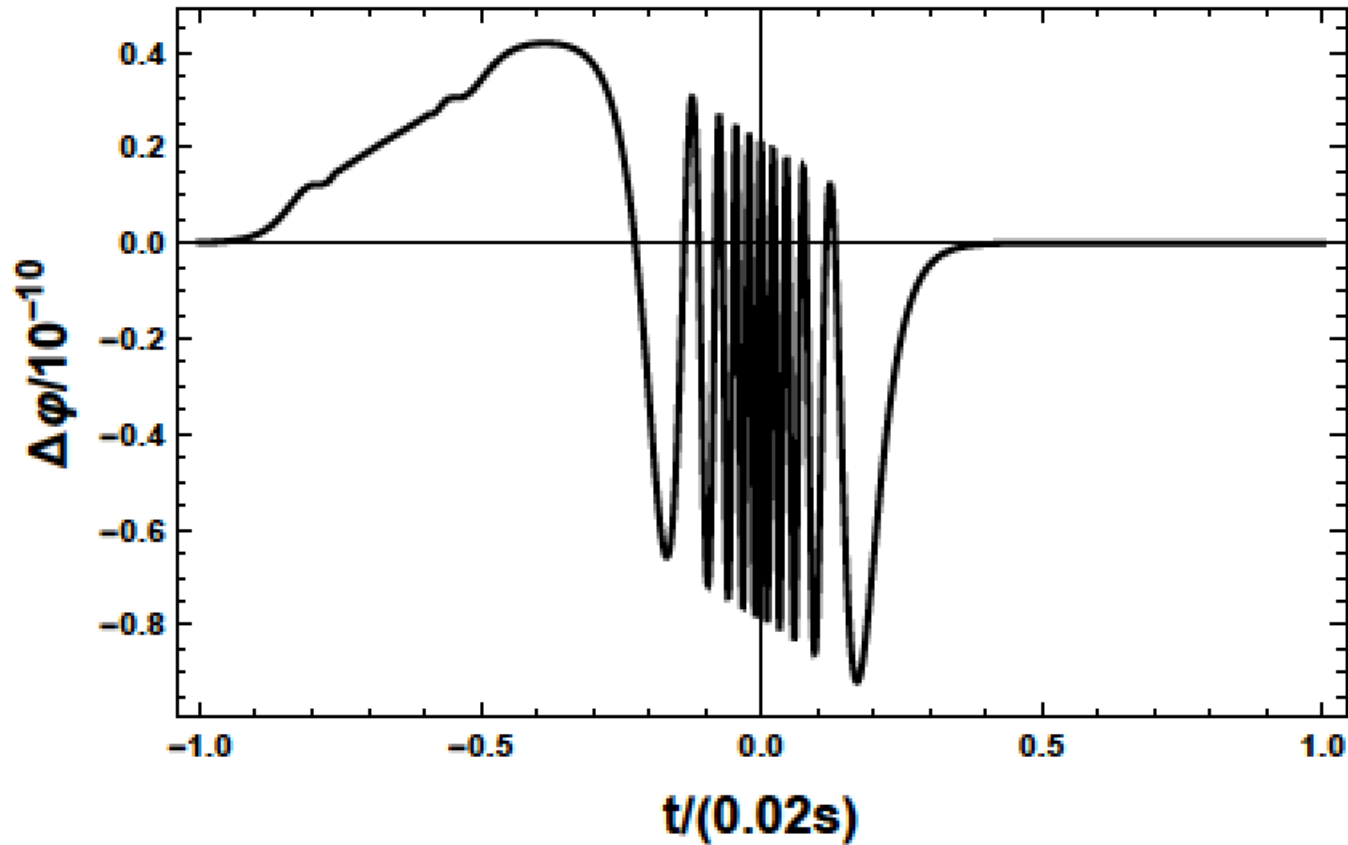


FIG. 8: As in Fig. ?? but  $m_{\gamma,0} = 0.1 \text{ neV}$ ,  $N_A/N_\phi = 5$ ,  $m = 0.5 \text{ neV}$ ,  $\alpha = \pi/2$  and  $v = 1 \times 10^{-3}$ .

# How to distinguish from grav waves?

- velocity  $\ll c$
- $v \sim 10^{-3}$

→ Time difference between two sites  
~few seconds

→ Need careful analysis strategies

Promised sensitivity:

$$m_{0,\gamma} \sim \text{neV} \left( \frac{m}{10 \text{ neV}} \right)^{1/2} \quad \text{for } m \gtrsim 0.1 \text{ neV},$$
$$\sim 0.1 \text{ neV} \quad \text{for } m \lesssim 0.1 \text{ neV}.$$



# Exploring heavier dark sectors with gravitational waves

Example I:

Scale invariant EW symmetry breaking

# Scale invariant EW symmetry breaking

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- Scale invariance
  - No explicit mass scales in the Lagrangian
- EW symmetry breaking?

$$V(\phi) = \frac{1}{4}\phi^4$$

No EW breaking at tree level

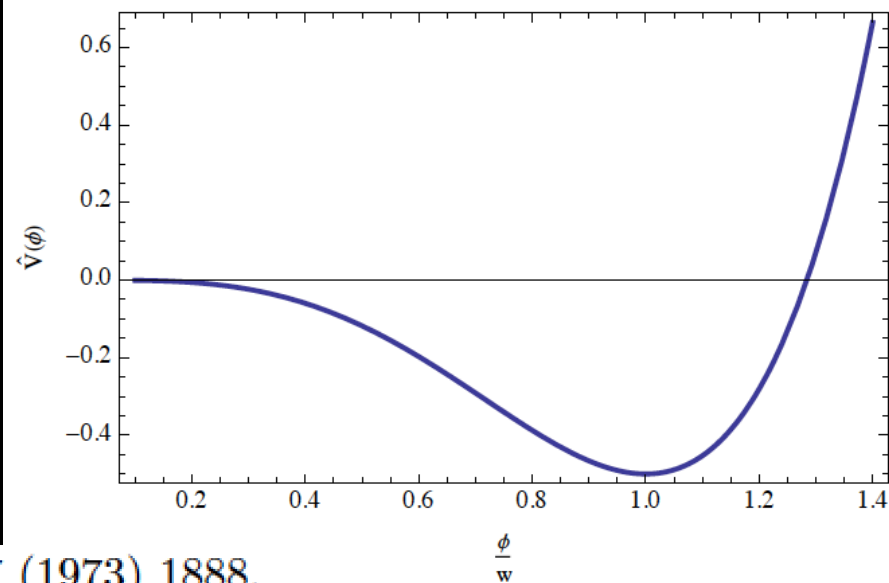
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# Scale invariant EW symmetry breaking

- Scale invariance
  - No explicit mass scales in the Lagrangian
- EW symmetry breaking?

$$V(\phi; \mu) = \frac{\lambda_\phi(\mu)}{4} \phi^4 + \frac{ng_D(\mu)^4}{64\pi^2} \phi^4 \left( \log \left( \frac{\phi^2}{\mu^2} \right) - \frac{25}{6} \right)$$

**Works  
at 1-loop level!**



# Need dark=hidden sector

- In the CW-SM Higgs would be too light
- top mass too high → unstable

→ Remedy:

Generate mass scale in Dark sector

$$V_0(h, \phi) = \frac{\lambda_\phi}{4} \phi^4 + \frac{\lambda_H}{4} h^4 - \frac{\lambda_P}{4} h^2 \phi^2$$

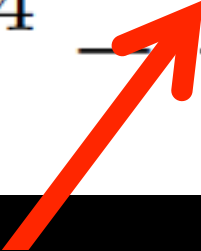
$\langle \phi \rangle \neq 0$  generates negative mass for h  
→ EW symm breaking

# Need dark=hidden sector

- In the CW-SM Higgs would be too light
- top mass too high → unstable

→ Remedy:

Generate mass scale in Dark sector

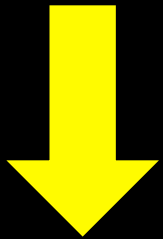
$$V_0(h, \phi) = \frac{\lambda_\phi}{4} \phi^4 + \frac{\lambda_H}{4} h^4 - \frac{\lambda_P}{4} h^2 \phi^2$$


Weak coupling to dark sector

# Features

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- $\lambda_P \ll 1$



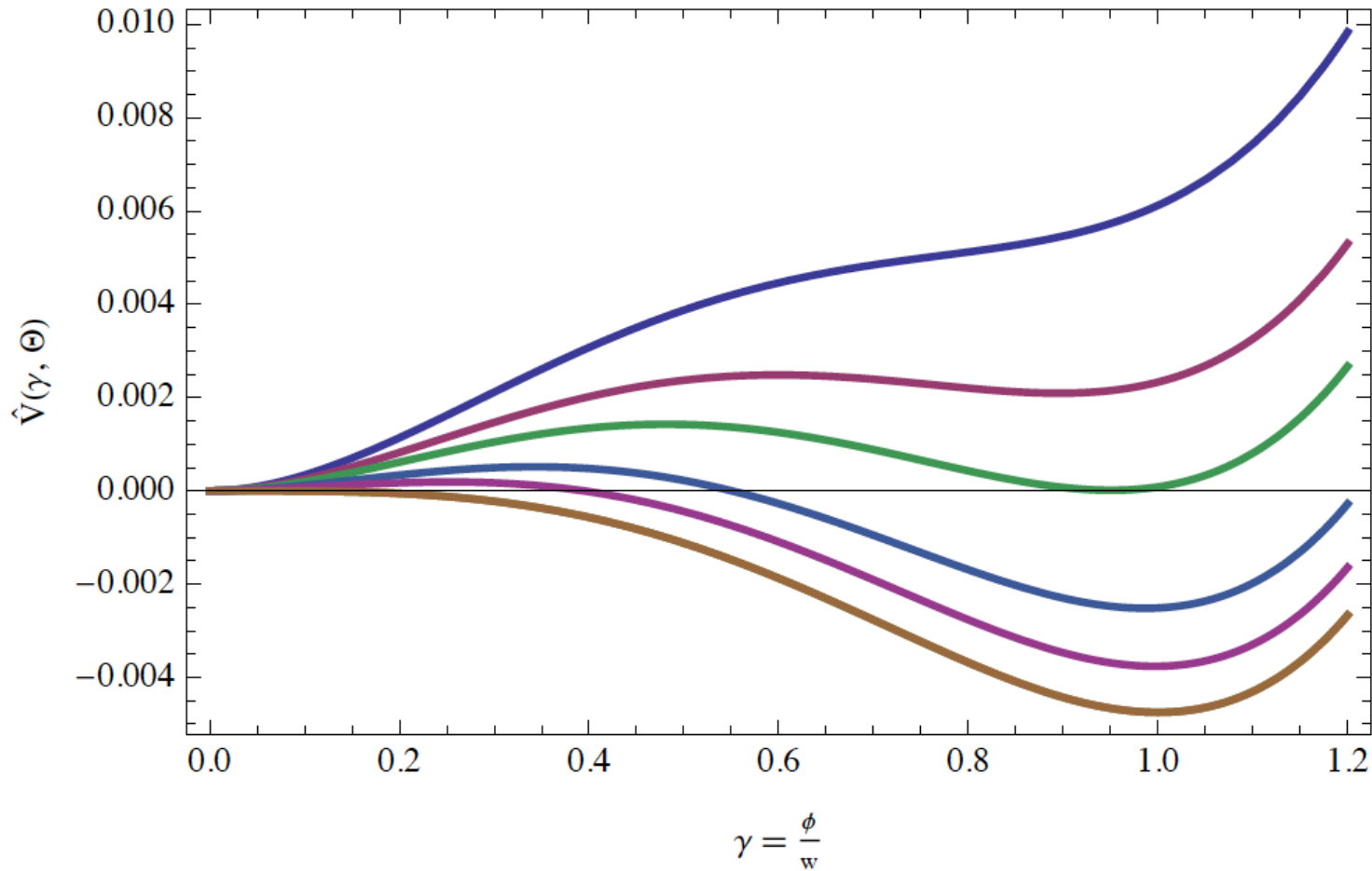
- $\langle \phi \rangle \gg v_{EW}$

- $m_\phi$  typically sizeable often  $\sim v$  or bigger

Hard to test weakly coupled and heavy dark sector

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# Temperature dependent potential



→ First order phase transition

→ Bubble formation

→ Gravitational waves

# Detectable gravitational waves

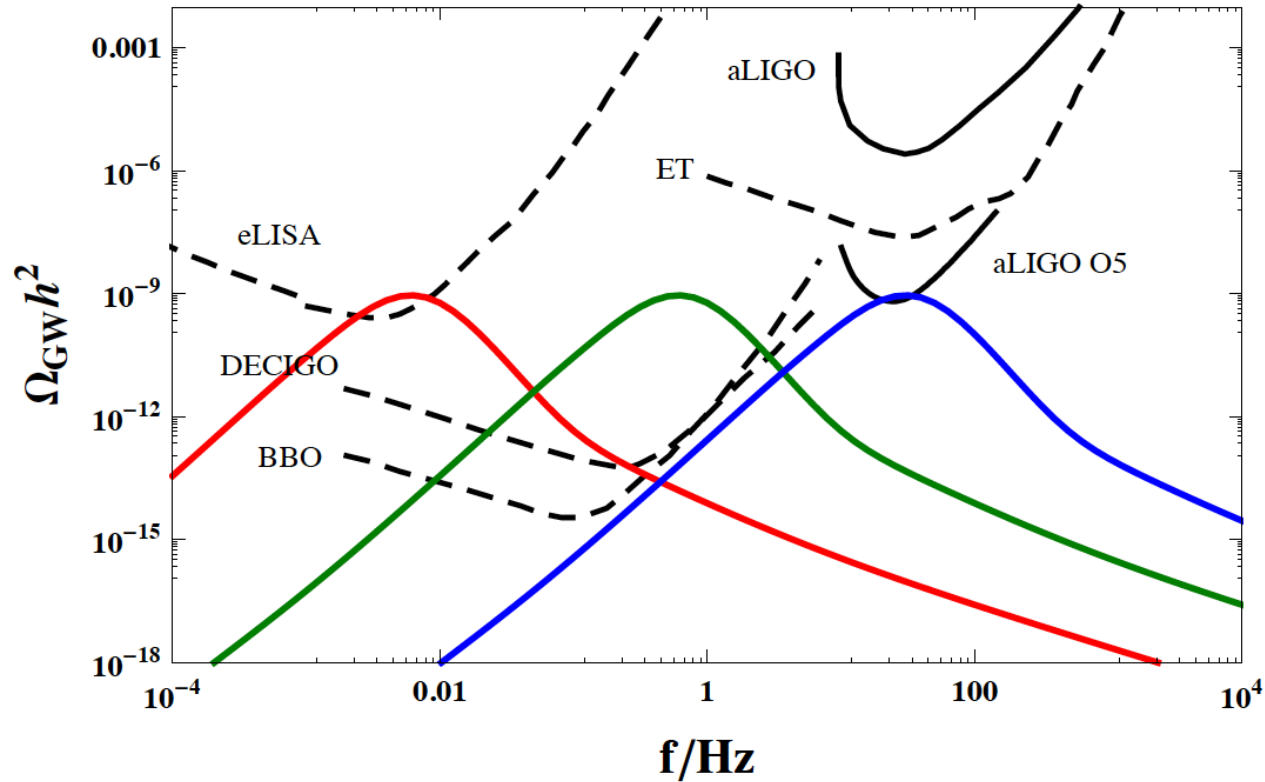


FIG. 6: Reach of gravitational wave detectors: We show aLIGO together with the fifth phase of aLIGO (both solid black), and the proposed detectors BBO, DECIGO, ET and eLISA [dashed black] (the sensitivities are taken from the gravitational wave plotter <http://rhcole.com/apps/GWplotter/> [29]). For the curves of the CW phase transition – going from left to right – we choose  $v_w = 1$  throughout, and respectively  $(\kappa = 1.0, g_D = 0.6, T_* = 100 \text{ GeV})$  [in red],  $(\kappa = 1.0, g_D = 0.6, T_* = 10 \text{ TeV})$  [green] and  $(\kappa = 1.0, g_D = 0.6, T_* = 500 \text{ TeV})$  [in blue].



# Detectable gravitational waves

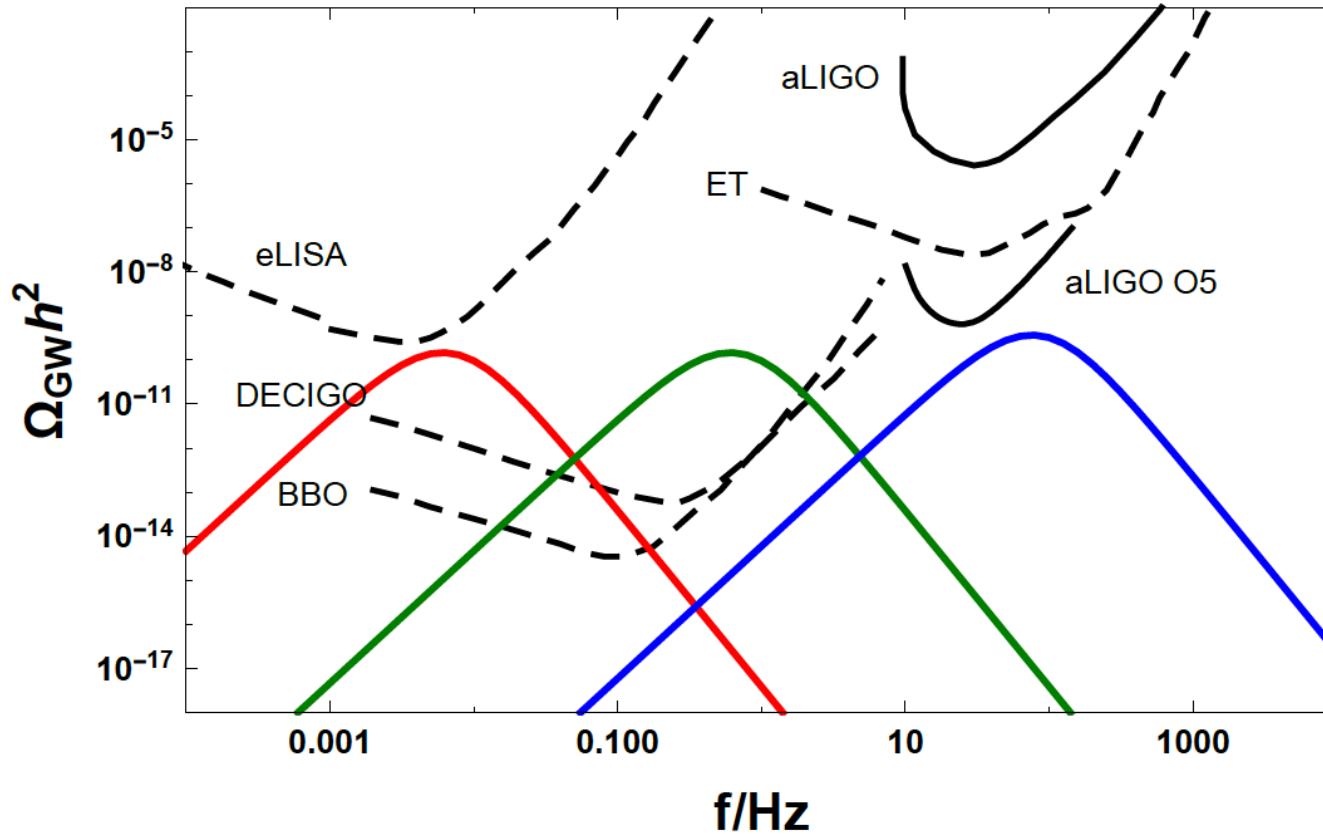


FIG. 7: Reach of gravitational wave detectors for a more conservative scenario  $\kappa_{sw} = 0.4$  (all other parameters as in Fig. 6).

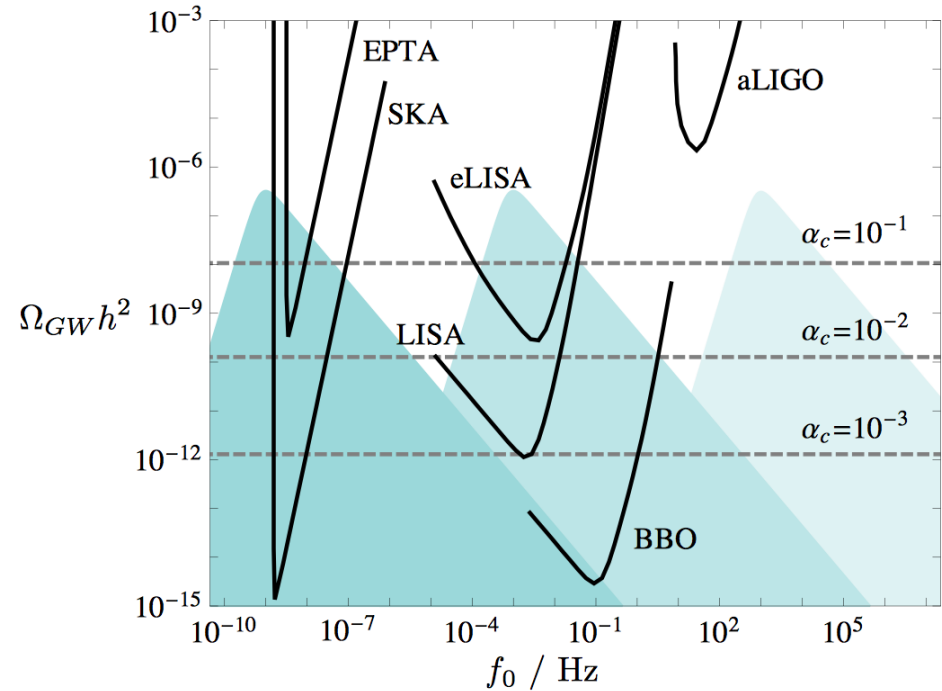
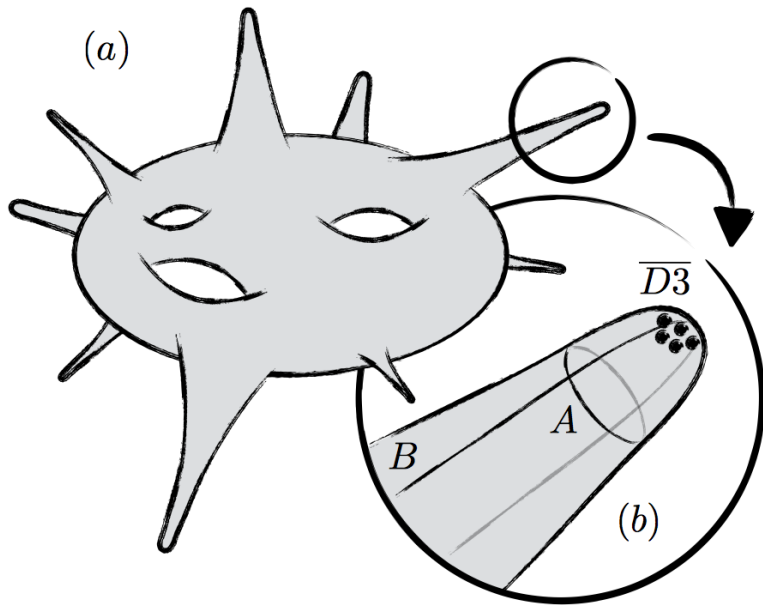
# Example I b: Deep Throats

The String Soundscape at Gravitational Wave Detectors

Isabel García García,<sup>1,\*</sup> Sven Krippendorf,<sup>1,†</sup> and John March-Russell<sup>1,‡</sup>

<sup>1</sup>*Rudolf Peierls Centre for Theoretical Physics, University of Oxford, 1 Keble Road, Oxford OX1 3NP, UK*

# Other models also produce GWs



## The String Soundscape at Gravitational Wave Detectors

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# Exploring heavier dark sectors with gravitational waves

Example II:  
GW Waves from Monodromy Inflation

# Monodromy potential

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$$V(\phi) = \Lambda^4 \cos\left(\frac{\phi}{f} + \gamma\right)$$

Monodromy add-on



"Axion" potential  
(pseudo-Goldstone pot.)

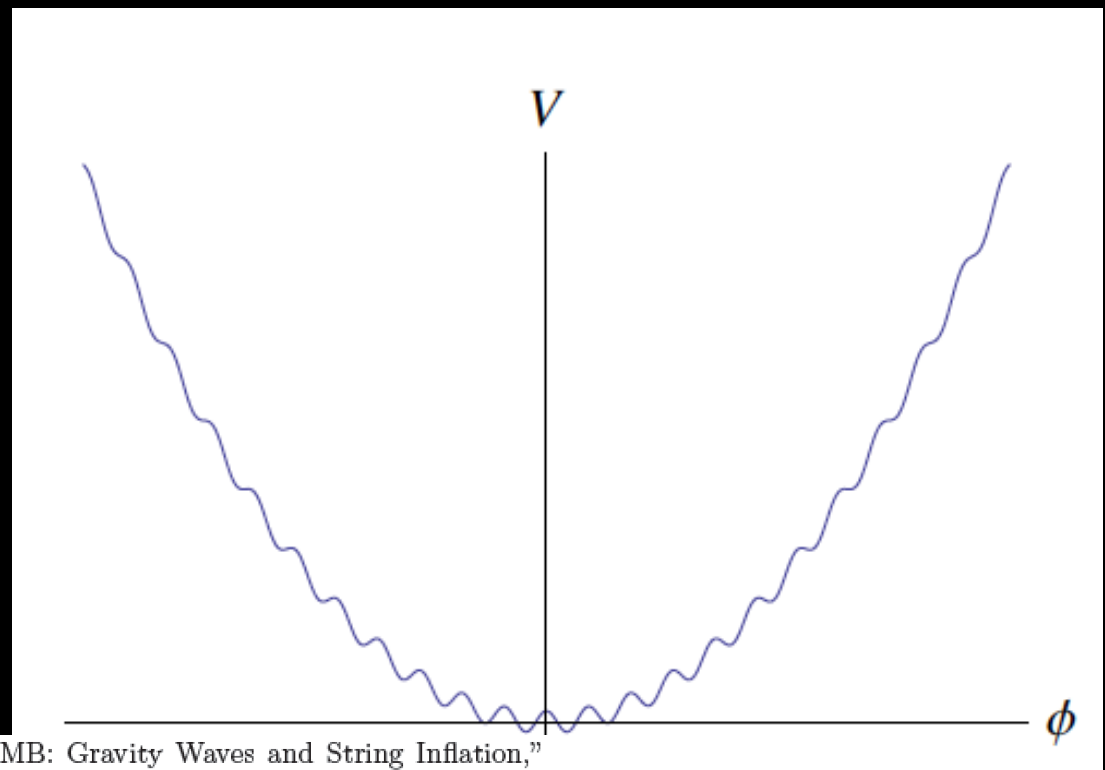


# Monodromy potential

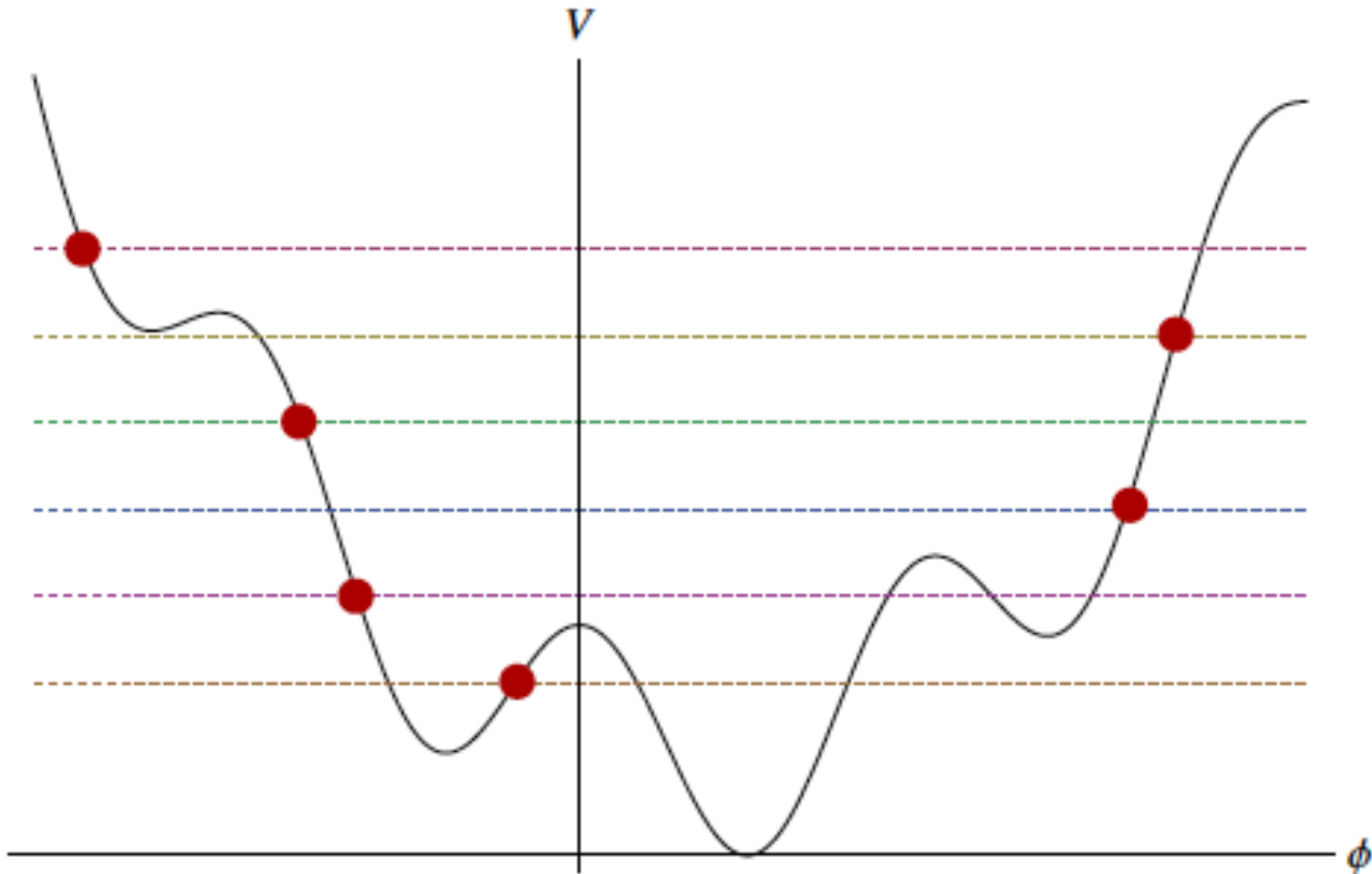
$$V(\phi) = \frac{1}{2}m^2\phi^2 + \Lambda^4 \cos\left(\frac{\phi}{f} + \gamma\right)$$

Funny potential

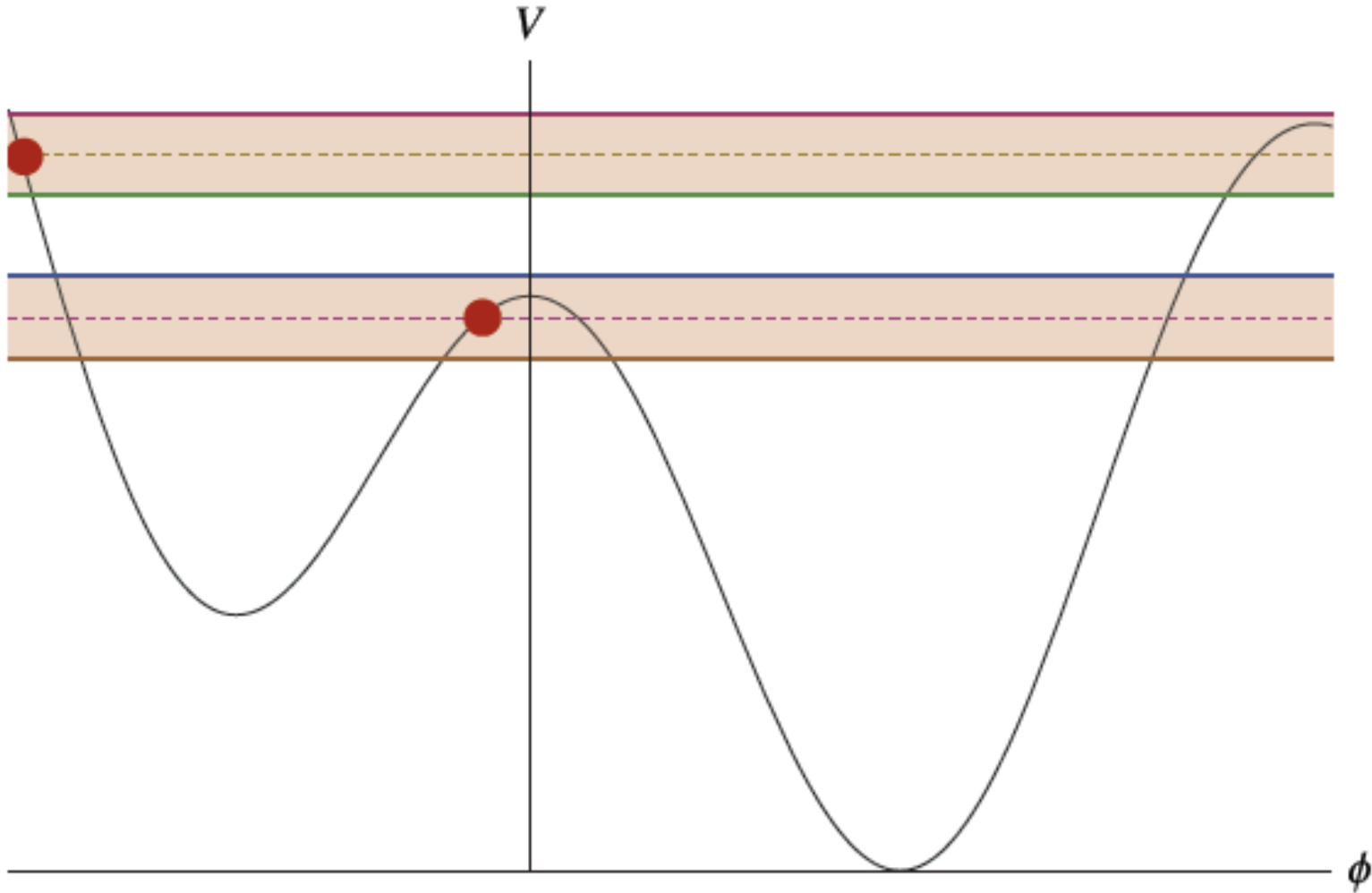
+ enlarged  
field range



# Field oscillations

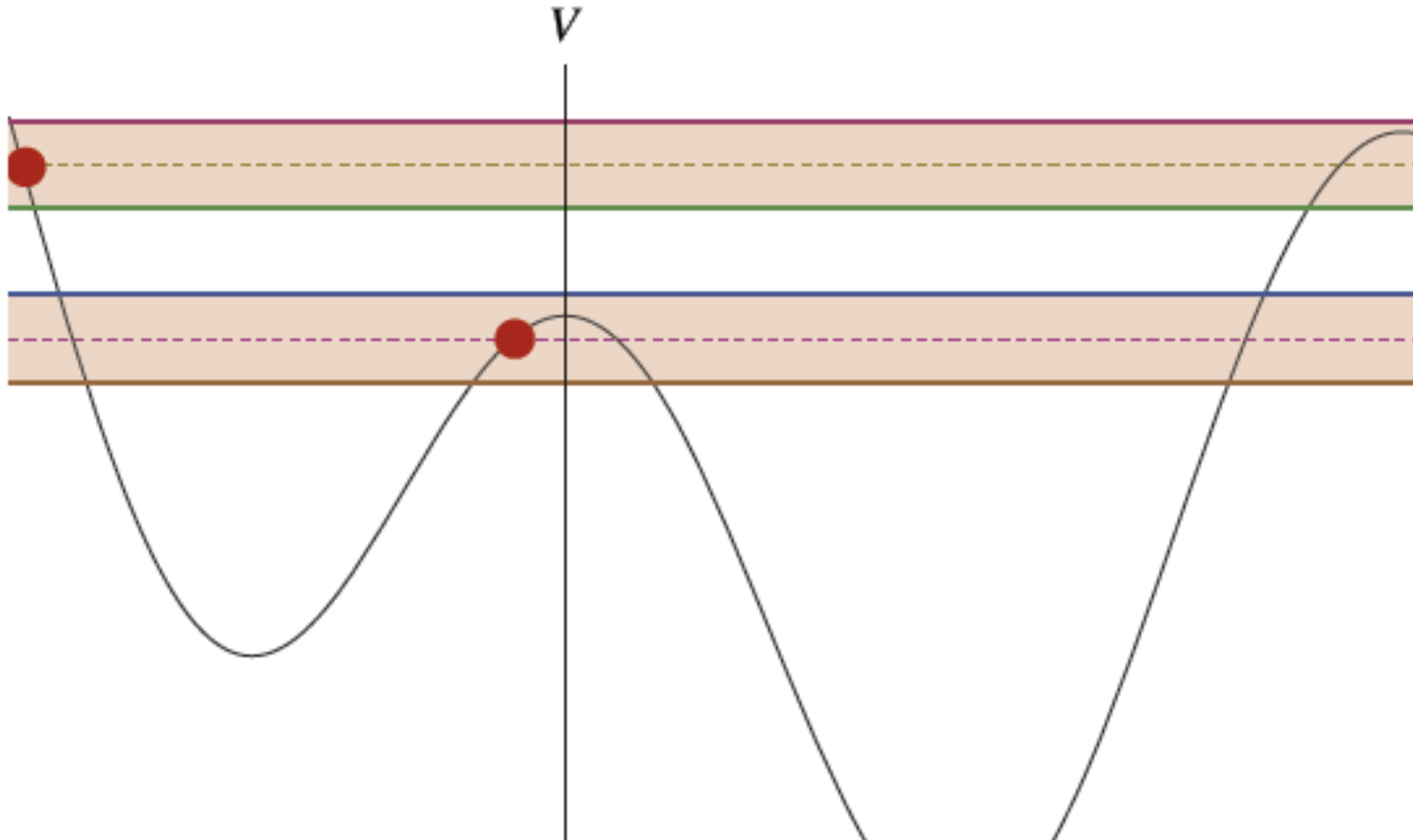


# Including fluctuations





# Including fluctuations



Different parts of Universe end in different minima  
→ Bubbles form  
→ GW are produced

# Gravitational wave spectra

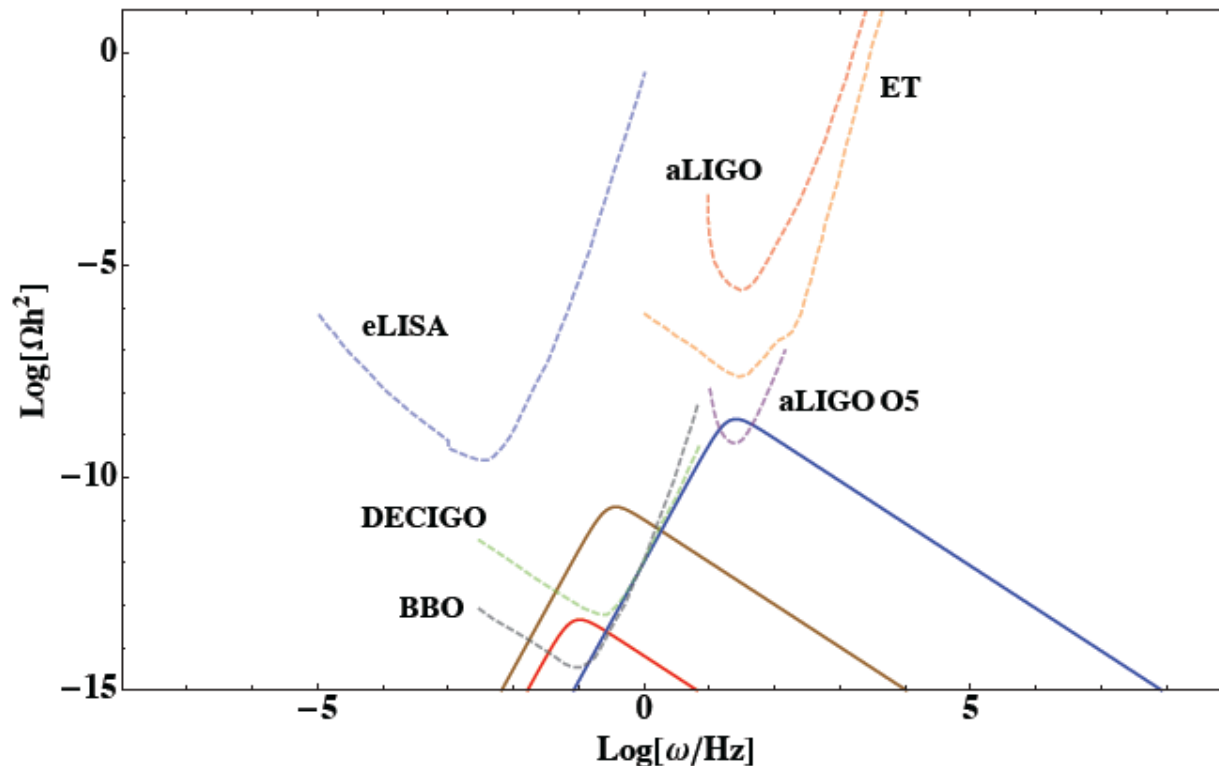
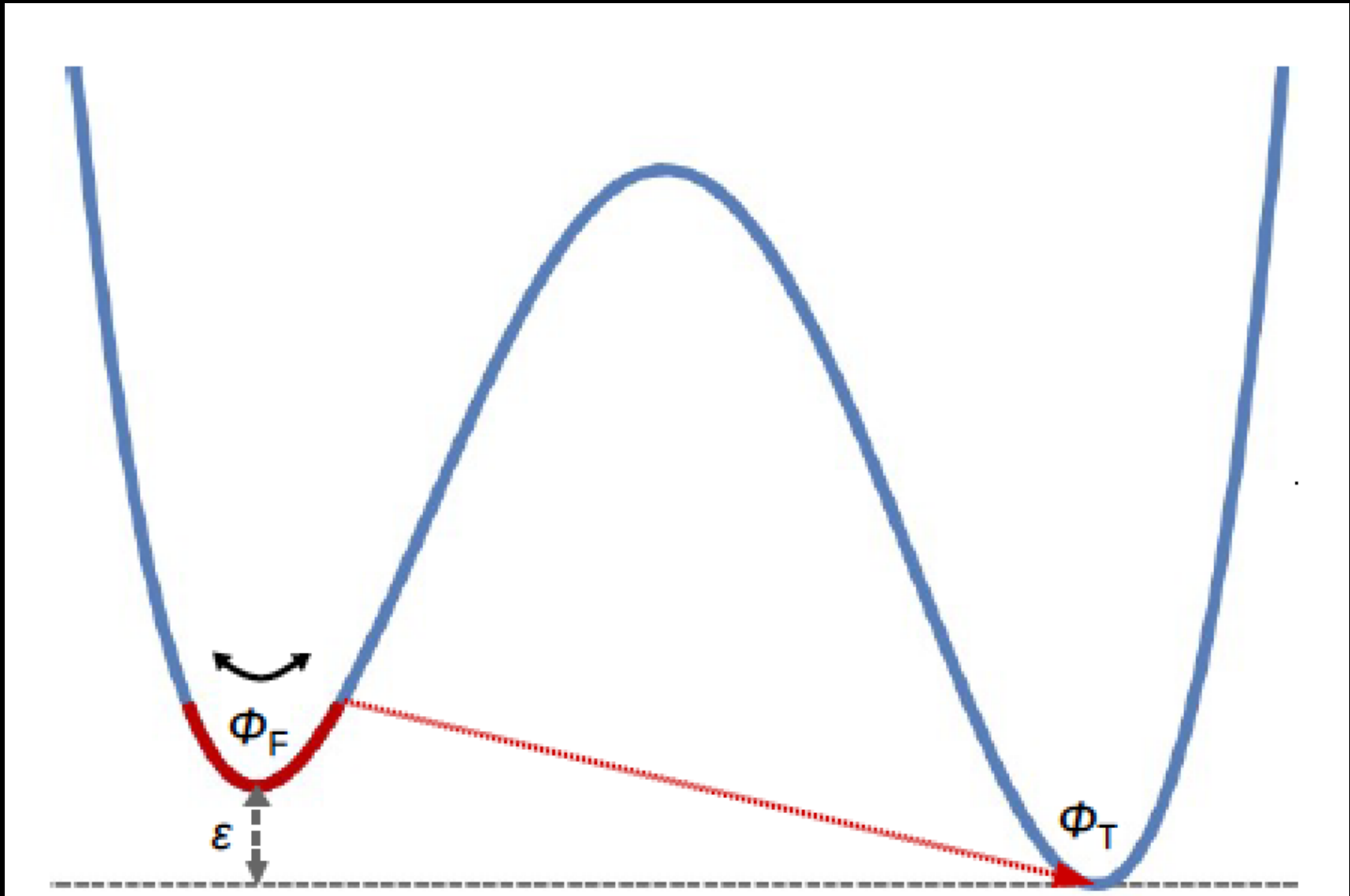


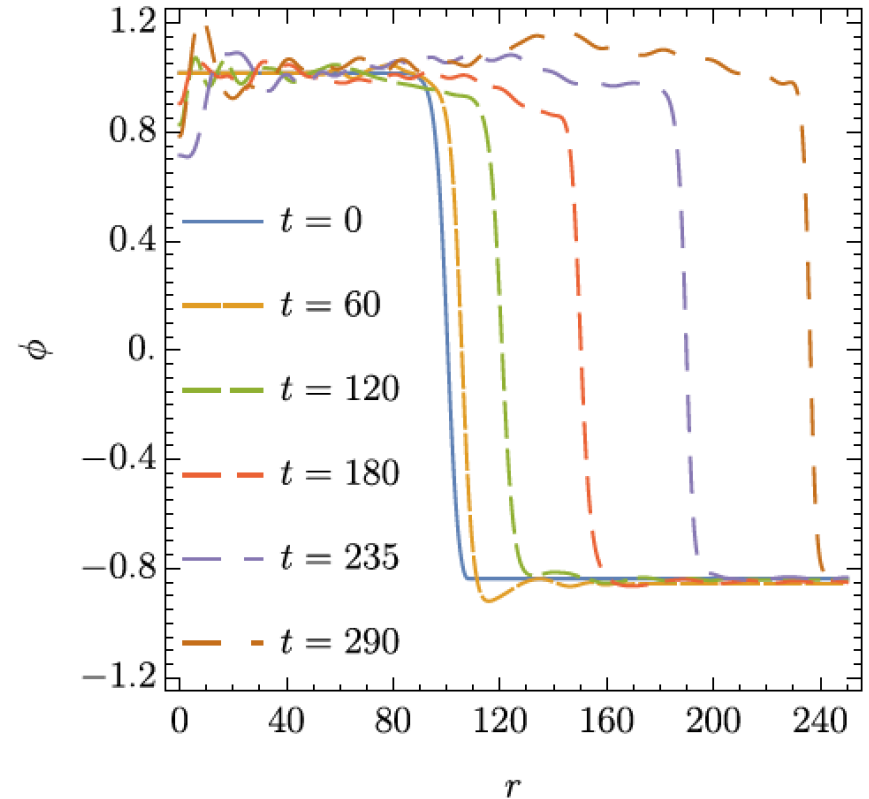
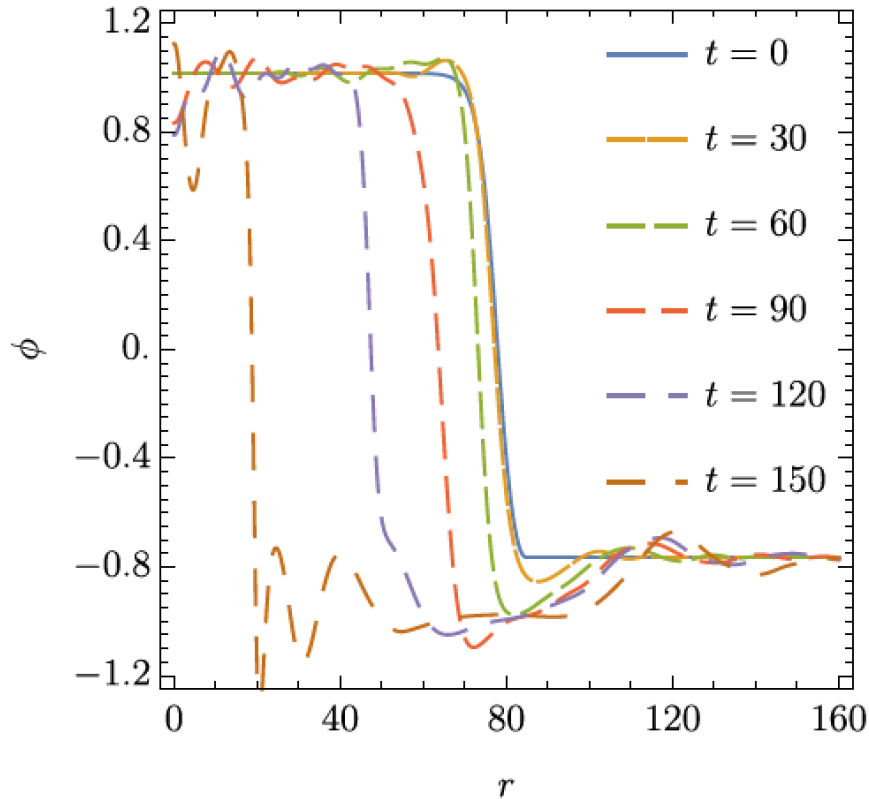
Figure 12: Gravitational wave spectra as in (5.22) with  $w = 1/3$ . The inflaton mass is fixed to  $m \sim 10^{-5}M_p$ . Spectra are shown as solid lines for different values of  $\kappa$ ,  $f$  and  $T_{RH}$ : the blue curve is obtained for  $\kappa = 5$ ,  $f = 0.1M_p$ ,  $T_{RH} \sim 10^{12}$  GeV; the brown curve for  $\kappa = 10$ ,  $f = 0.01M_p$ ,  $T_{RH} \sim 10^{11}$  GeV; the red one for  $\kappa = 70$ ,  $f = 0.001M_p$ ,  $T_{RH} \sim 10^{11}$  GeV. We have also taken  $w = 1/3$ ,  $\theta_0 = 10^{-2}$ ,  $\sigma = 10^{-1}$  in (5.22). For the values of the reheating temperature considered here, we have  $g_*(T_{RH}) \sim 10^2$ . Sensitivity curves of some ground- and space-based interferometers are shown for comparison as dashed curves (data taken from [74]).

Technical note:  
Tunneling from  
an oscillating vacuum...

# Non-trivial initial state



# Bubbles can collapse



**Figure 8:** Field profiles showing lattice bubble evolution for oscillation reaching  $1.2\sqrt{2\alpha_{tw}}$  (left panel) and  $0.8\sqrt{2\alpha_{tw}}$  (right panel). The values defining the potential were set to  $g = 1/10$ ,  $b = 1/300$  and  $c = 1$ .

# Another really cool example

## Discovering the QCD Axion with Black Holes and Gravitational Waves

Asimina Arvanitaki<sup>\*</sup>

*Perimeter Institute for Theoretical Physics, Waterloo, Ontario, N2L 2Y5, Canada*

Masha Baryakhtar<sup>†</sup> and Xinlu Huang<sup>‡</sup>

*Stanford Institute for Theoretical Physics, Department of Physics,  
Stanford University, Stanford, CA 94305, USA*

### Black Hole Mergers and the QCD Axion at Advanced LIGO

Asimina Arvanitaki,<sup>1,\*</sup> Masha Baryakhtar,<sup>1,†</sup> Savas Dimopoulos,<sup>2,‡</sup> Sergei Dubovsky,<sup>3,§</sup> and Robert Lasenby<sup>1,¶</sup>

<sup>1</sup>*Perimeter Institute for Theoretical Physics, Waterloo, Ontario N2L 2Y5, Canada*

<sup>2</sup>*Stanford Institute for Theoretical Physics, Stanford University, Stanford, California 94305, USA*

<sup>3</sup>*Center for Cosmology and Particle Physics, New York University New York, NY, 10003, USA*

# GWs probe axion detectors

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## Discovering the QCD Axion with Black Holes and Gravitational Waves

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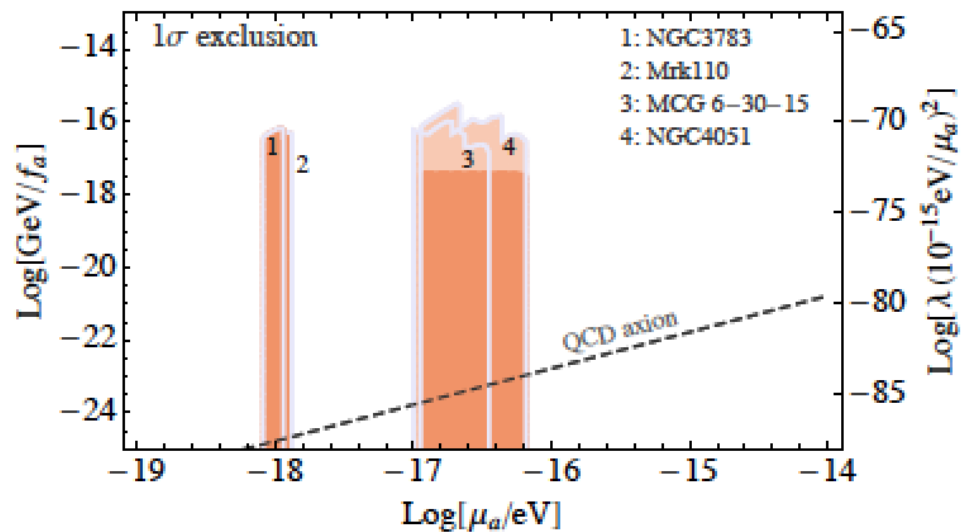
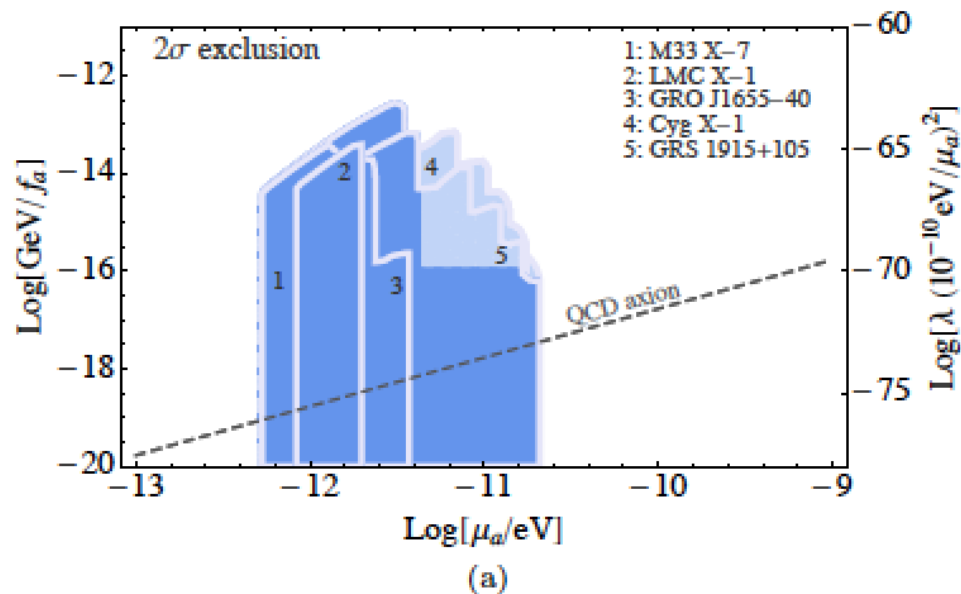
Masha Baryakhtar<sup>†</sup> and Xinlu Huang<sup>‡</sup>

*Stanford Institute for Theoretical Physics, Department of Physics,  
Stanford University, Stanford, CA 94305, USA*

## If light bosons

- superradiance for Spinning Black holes
  - efficient if  $R_{\text{BH}} \sim 1/m_{\text{Boson}}$
  - spin is quickly “used up” to produce bosons
  - Measure spin of black holes (aLIGO, LISA)
-

# Sensitivity to "zero" couplings





# Do statistics with aLIGO

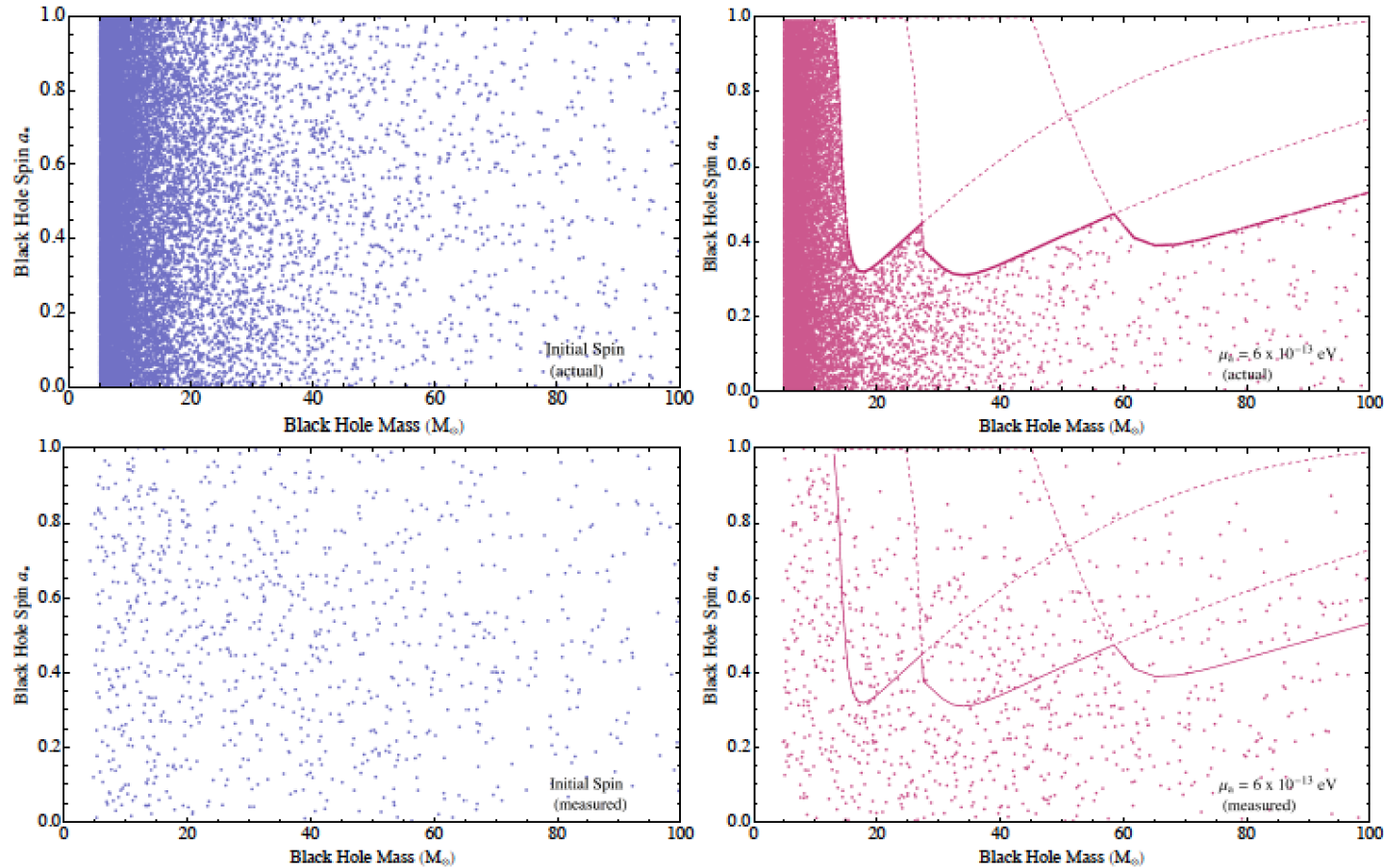


FIG. 2. Expected distribution of intrinsic (top) and measured (bottom) spins and masses of merging BHs in the absence (left) and the presence (right) of an axion of mass  $6 \times 10^{-13}$  eV, normalized to 1000 events detected at aLIGO. We assume  $\sigma_M/M \sim 10\%$  measurement error in the mass and  $\sigma_{\alpha_*} \sim 0.25$  error in the spin [30, 31]. We have assumed that all BBHs formed at a distance such that they take  $10^{10}$  years to merge. The theoretical curves shown are boundaries of the regions where SR had at most  $10^{10}$  years to spin down the BHs, and the effect of the companion BH does not significantly affect the SR rate.

## Black Hole Mergers and the QCD Axion at Advanced LIGO

Asimina Arvanitaki,<sup>1,\*</sup> Masha Baryakhtar,<sup>1,†</sup> Savas Dimopoulos,<sup>2,‡</sup> Sergei Dubovsky,<sup>3,§</sup> and Robert Lasenby<sup>1</sup>

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<sup>3</sup>Center for Cosmology and Particle Physics, New York University New York, NY, 10003, USA

# More signals...

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Bose enhanced level transitions

→ Monofrequency signals

Annihilation of bosons into gravitons

→ Monofrequency signal  $\omega = 2m_a$

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Conclusions

- Gravitational wave detectors are amazing tools
  - ➔ use to learn about dark sectors
- Use to directly search for new particles (Example WISPy domain walls and ultralight bosons such as axions)
- Use gravitational wave signatures to probe dark sectors populated by heavier particles that are hard to detect otherwise
  - ➔ Challenge: Need to develop specificity to distinguish different scenarios!

Hidden sector



# What is PBC?

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**Study group mandated by CERN management to prepare for the next European HEP strategy update**

(coordinators Mike Lamont, Claude Vallee, JJ)

**“Explore the opportunities offered by the CERN accelerator complex to address some of today's outstanding questions in particle physics through experiments complementary to high-energy colliders and other initiatives in the world” (Excerpt from the mandate)**

**Time scale ~ 2040**

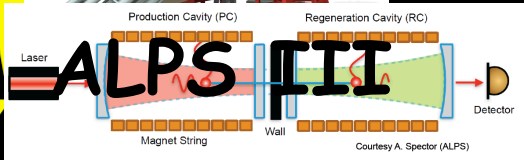
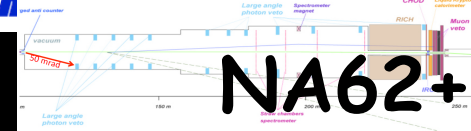
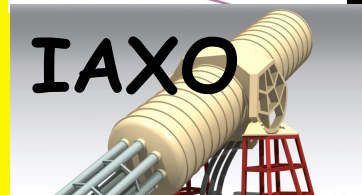
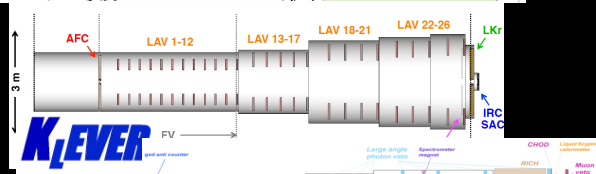
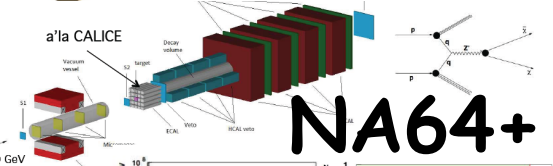
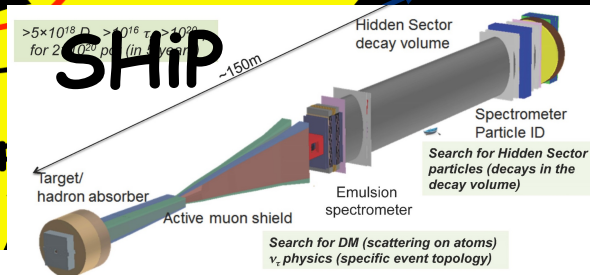
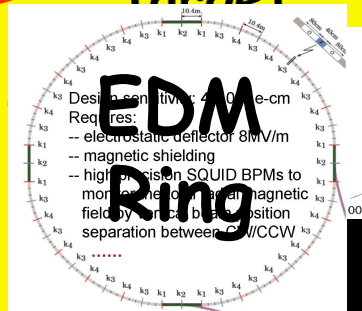
**[pbc.web.cern.ch](http://pbc.web.cern.ch)**

# Here we want to go in the Future...

## Energy, Mass

known  
knowns

$\nu$ , Fixed  
target



Precision,  
Intensity,  
Small coupling

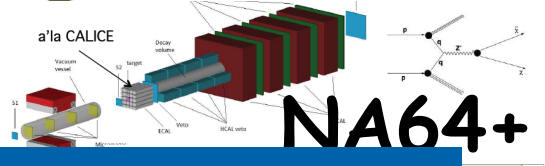
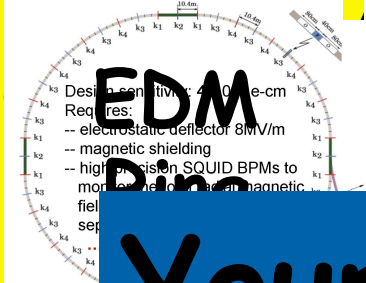
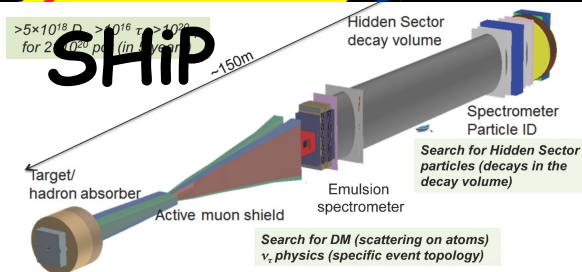
# Here we want to go in the Future...

## Energy, Mass

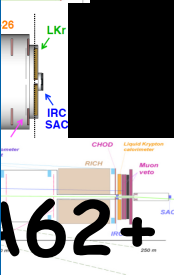
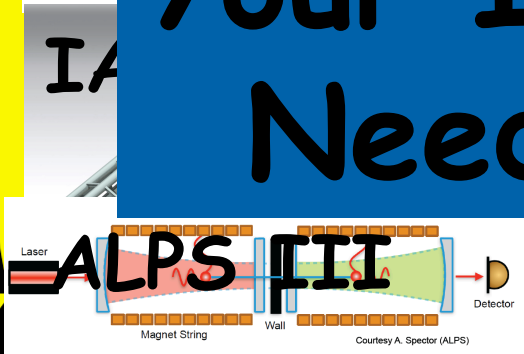
known  
knowns

LHC

$\nu$ , Fixed  
target



Your Ideas  
Needed



Precision,  
Intensity,  
Small coupling