

Probing high-scale physics using ATLAS Higgs boson measurements and SM effective field theory



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Motivation

- ► A program of precise measurements and the EFT interpretation constrains $coupling/mass^2$ factors for new particles such as dark matter. This bottom-up approach is complementary to the top-down approach of searching directly for new particles.
- Discovery in the top-down approach is through observation of specific new particles in classes of models, while discovery in the bottom-up approach is through deviations observed in precise SM measurements.

Overview

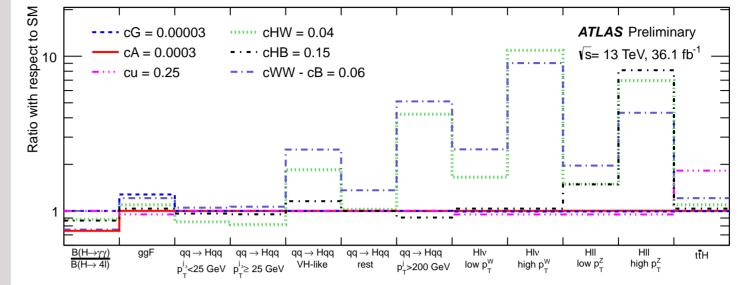
- We use a simplified set of six operators from the SM EFT to constrain new Higgs boson interactions, which are a promising source of new physics.
- \blacktriangleright The operator coefficients are fit using the combined ATLAS ${f H}
 ightarrow \gamma\gamma$ and $H \rightarrow ZZ^{(*)} \rightarrow 4I$ measurements with the 2015-2016 dataset, corresponding to an integrated luminosity of 36.1 fb^{-1} .

EFT Lagrangian

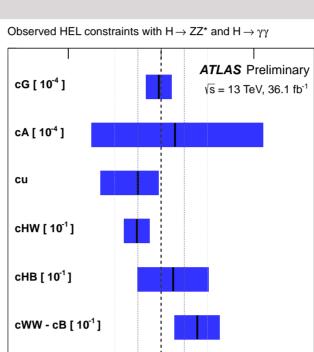
► The general form of the Lagrangian including dimension-6 operators is: $\mathcal{L} = \mathcal{L}_{\mathrm{SM}} + \mathrm{c}_{\mathrm{i}}^{(6)} \mathcal{O}_{\mathrm{i}}^{(6)} / \Lambda^2$

EFT impact plot on **STXS** measurements

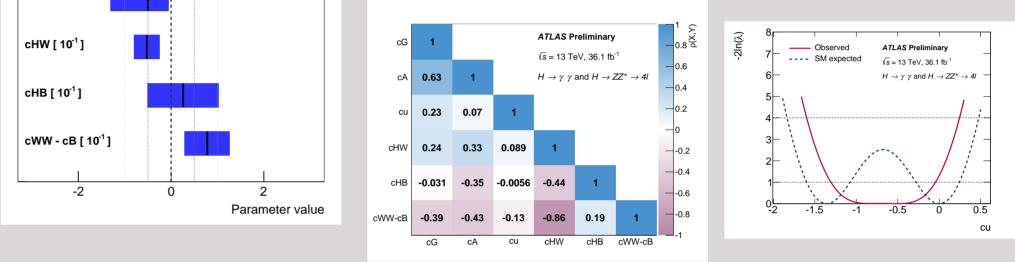
The impact plot shows how various STXS regions are affected by parameters corresponding to $+1\sigma$ expected sensitivity [2].



Results of fit to ATLAS measurements



For O(1) couplings of new particles to the SM, the constraints extend to new-physics scales of 10 TeV for ggH vertex, 5 TeV for H $\gamma\gamma$ vertex, and 500 GeV for HVV and ttH vertices. [2]



- ► To determine the relationship between STXS measurements and EFT parameters, we separate the cross section into SM, SM-BSM interference, and BSM components: $\sigma = \sigma_{\rm SM} + \sigma_{\rm int} + \sigma_{\rm BSM}$.
- ► Then the cross section dependence on the couplings can be expressed as:

$$\frac{\sigma}{\sigma_{\text{SM}}} = 1 + \sum_{i} A_i c_i + \sum_{ij} B_{ij} c_i c_j$$

EFT at LO dimension-6 and selected Wilson coefficients

The SM EFT has 59 operators. The majority of these operators do not affect Higgs physics at LO. We used the 'Higgs effective Lagrangian' model, which is

implemented in the Madgraph

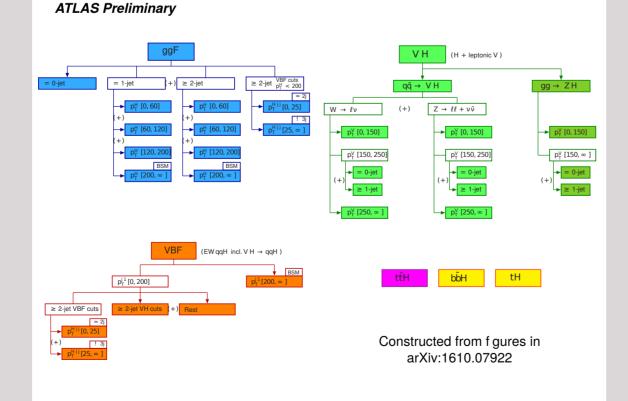
We selected the following Wilson coefficients to constrain, based on the expected sensitivity of the measurements:

| Operator | Expression | HEL coefficient | Vertices |
|----------------------|---|--|---------------------|
| \mathcal{O}_{g} | $ H ^2 G^{A}_{\mu u} G^{A\mu u}$ | $cG = \frac{m_W^2}{g_s^2} \bar{c}_g$ | Hgg |
| \mathcal{O}_γ | $ H ^2B_{\mu u}B^{\mu u}$ | $cA = \frac{m_W^2}{{g'}^2} \bar{c}_{\gamma}$ | $H\gamma\gamma,HZZ$ |
| \mathcal{O}_{u} | yu∣H ²ū₁Hu _R + h.c. | $cu = v^2 \bar{c}_u$ | Htī |
| \mathcal{O}_{HW} | $i\left(D^{\mu}H ight)^{\dagger}\sigma^{a}\left(D^{ u}H ight)W^{a}_{\mu u}$ | $cHW = \frac{m_W^2}{g} \bar{c}_{HW}$ | HWW, HZZ |
| \mathcal{O}_{HB} | $i\left(D^{\mu}H ight)^{\dagger}\left(D^{ u}H ight)B_{\mu u}$ | $cHB = \frac{m_W^2}{g'} \overline{c}_{HB}$ | HZZ |
| \mathcal{O}_{W} | ${f i}\left({f H}^{\dagger}\sigma^{{f a}}{f D}^{\mu}{f H} ight){f D}^{ u}{f W}^{{f a}}_{\mu u}$ | $cWW = \frac{m_W^2}{g} \bar{c}_W$ | HWW, HZZ |
| \mathcal{O}_{B} | ${f i}\left({f H}^{\dagger}{f D}^{\mu}{f H} ight)\partial^{ u}{f B}_{\mu u}$ | $cB = \frac{m_W^2}{g'} \overline{c}_B$ | HZZ |

Simplified template cross sections (STXS)

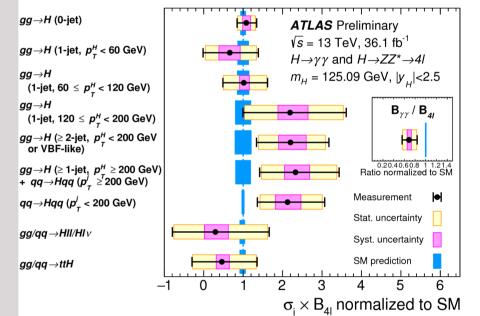
- ► The STXS are an evolution from inclusive production cross section measurements to production cross sections split into a few kinematic regions, which are most sensitive to new physics.
- Cross sections are measured in each STXS region at the generator level, with correlations.

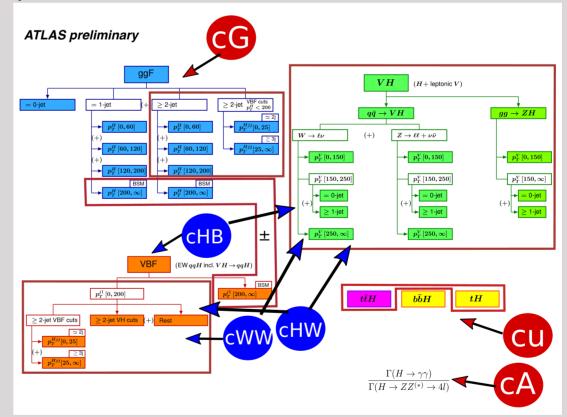
The regions are: $\sigma_{i} \times B_{4i}$

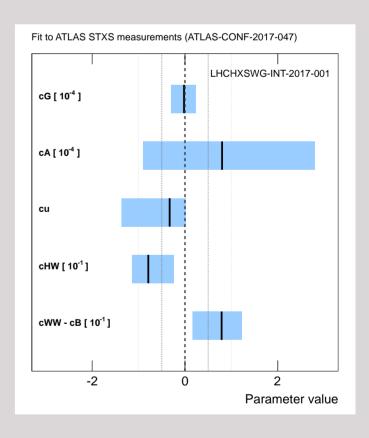


Alternative fit to a reduced set of STXS regions

- ATLAS previously used the same data to directly measure a reduced set of 11 STXS [3]. This reduced set can also be used to fit for EFT parameters to investigate the reduction in sensitivity to the parameters [4].
- The ATLAS measurements merged regions most sensitive to cWW-cB, cHW, and cHB. In the fit to these regions we remove cHB, leaving five parameters.







Summary

and $\mathbf{B_i}/\mathbf{B_{4l}}$.

generator.

The STXS measurement regions on which the fit is based are shown in the figure on the right [1] [2].

Example equations and validation

The fit uses equations expressing the STXS regions in terms of EFT parameters, for example:

 $\sigma_{\mathsf{BSM}}/\sigma_{\mathsf{SM}} = 1 + 1.546 \cdot \mathsf{cWW} - 3.631 \cdot \mathsf{cHW} - 0.2361 \cdot \mathsf{cHB} + 11.11 \cdot \mathsf{cWW} \cdot \mathsf{cWW} + 0.2245 \cdot \mathsf{cB} \cdot \mathsf{cB} + 12.54 \cdot \mathsf{cB} \cdot \mathsf{cB} + \mathsf{$ $cHW \cdot cHW + 0.2189 \cdot cHB \cdot cHB + 0.9 \cdot cWW \cdot cB + 3.43 \cdot cWW \cdot cHW + 0.43 \cdot cB \cdot cHW + 1.52 \cdot cHW \cdot cHB$

MC samples were generated with a variety of **c**_i values and a statistical precision <1%. Example validation samples show consistency with the equations (right).

| STXS region | $\sigma_{ m MG}/\sigma_{ m SM}$ | $\sigma_{ m eq}/\sigma_{ m SM}$ | δ_{MG}^{stat} |
|---|---------------------------------|---------------------------------|----------------------|
| $gg \rightarrow H \ (\geq 2 \text{ jets}, \ p_T^H \geq 200 \text{ GeV})$ | 0.859 | 0.851 | 0.006 |
| $gg \rightarrow H \ (\geq 2 \text{ jets}, \ p_T^H \geq 200 \text{ GeV})$ | 0.948 | 0.945 | 0.006 |
| ${ m q}ar{ m q} ightarrow { m HI} u \left({ m p}_{ m T}^{ m V} \ge 250~{ m GeV} ight)$ | 1.112 | 1.110 | 0.002 |
| $q\bar{q} ightarrow HI u (p_T^{\dot{V}} \ge 250 \text{ GeV})$ | 1.277 | 1.276 | 0.002 |
| $q\bar{q} \rightarrow H l \nu (p_T^{\dot{V}} \ge 250 \text{ GeV})$ | 1.420 | 1.419 | 0.002 |
| $gg/q\bar{q} \rightarrow ttH$ | 0.848 | 0.848 | 0.001 |
| ${ m gg}/{ m qar q} ightarrow { m ttH}$ | 1.653 | 1.653 | 0.002 |

- ► These results represent the first step towards a full EFT fit to Higgs boson measurements by the LHC experiments, whose results will provide important information about the possible couplings and mass scales of new physics.
- ► The six-parameter fit demonstrates the general procedure for transitioning from constraints on coefficients of dimension-4 operators to those of dimension-6 operators in a general EFT framework for combined Higgs measurements.
- ► As more measurements are included and more robust theoretical tools become available, the fits will expand to larger operator sets and more deeply probe distributions and individual couplings.

References

- [1] D. de Florian et al. Handbook of LHC Higgs Cross Sections: 4. Deciphering the Nature of the Higgs Sector. 2016.
- [2] ATLAS Collaboration. Constraints on an effective Lagrangian from the combined $H \rightarrow ZZ^* \rightarrow 4\ell$ and $H \rightarrow \gamma\gamma$ channels using 36.1 fb⁻¹ of $\sqrt{s} = 13$ TeV pp collision data collected with the ATLAS detector. ATL-PHYS-PUB-2017-018, 2017.
- [3] ATLAS Collaboration. Combined measurements of Higgs boson production and decay in the $H \rightarrow ZZ^* \rightarrow 4\ell$ and $H \rightarrow \gamma\gamma$ channels using $\sqrt{s} = 13$ TeV pp collision data collected with the ATLAS experiment. ATLAS-CONF-2017-047, 2017.
- [4] Chris Hays, Veronica Sanz, and Gabija Zemaityte. Constraining eft parameters using simplified template cross sections. LHCHXSWG-INT-2017-001, Oct 2017.