

s and c quark mass determination from lattice QCD

Dan Hatton

University of Glasgow

Supervisor: Prof. Christine Davies

HPQCD

A. T. Lytle, G. P. Lepage, C. Sturm



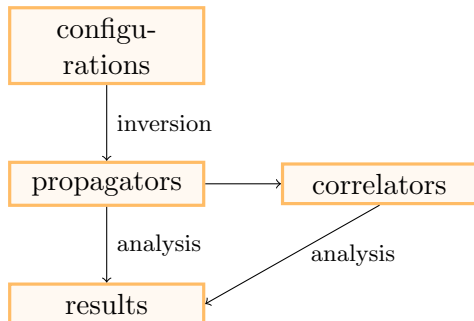
Introduction and Motivation

- ▶ Quark masses are difficult to determine precisely due to confinement
- ▶ Precise determination desirable for some precision Higgs calculations
- ▶ PDG strange and charm masses:
 - ▶ $\overline{m}_s(2 \text{ GeV}) = 96 \pm_4^8 \text{ MeV}$
 - ▶ $\overline{m}_c(\overline{m}_c) = 1.28 \pm 0.03 \text{ GeV}$
- ▶ Lattice results typically considerably more precise



Lattice Gauge Configurations

- ▶ Spacetime is discretised
 - ▶ Allows numerical calculations
 - ▶ Regulates divergences
- ▶ Gauge configuration is a group of $SU(3)$ matrices; one for each link between lattice sites
- ▶ Gauge configurations chosen by MC to follow e^{-S} distribution
- ▶ Correlation functions are evaluated configuration by configuration and then averaged



Propagators

- ▶ Grassmann variables in path integral can be performed analytically
 - ▶ $\int \mathcal{D}\bar{\psi}\mathcal{D}\psi\bar{\psi}\psi e^{-\bar{\psi}M\psi} = M^{-1}\det(M)$
 - ▶ $\det(M)$ included in gauge field generation
- ▶ The propagator is the inverse of the Dirac matrix $M \equiv \not{D} - m$
- ▶ M is a very large matrix that is inverted numerically
- ▶ Correlation functions are constructed from the propagators
- ▶ Calculation time diverges as $am \rightarrow 0$



<https://www.hpc.cam.ac.uk/images/darwin.jpg/view>



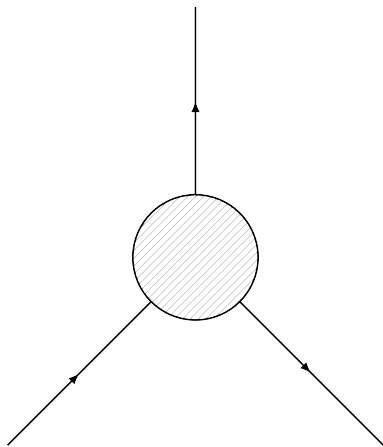
Tuning Quark Masses

- ▶ Quark masses are input into simulations and tuned to produce correct meson mass
 - ▶ Input as am
- ▶ am_s and am_c are tuned on each lattice to give correct η_s and η_c meson masses
- ▶ Lattice QCD is a regularisation and the bare input masses therefore need to be renormalised
- ▶ Renormalisation is done nonperturbatively in a momentum subtraction scheme and then perturbatively matched to $\overline{\text{MS}}$ in the continuum



RI-SMOM: Setup

- ▶ To renormalise a current, consider the vertex function of (G_Γ) the relevant operator $(\bar{\psi}\Gamma\psi)$
- ▶ Conditions are placed on the amputated vertex functions in the limit of zero quark mass
- ▶ Kinematic setup is chosen to satisfy $p_1^2 = p_2^2 = q^2$
- ▶ Momentum space Green's functions are constructed from Fourier transformed propagators



RI-SMOM: Conditions

- ▶ Propagator condition gives $Z_q = 1$ in the free theory
- ▶ Other conditions satisfy Ward identities
- ▶ $\Lambda_\Gamma = S^{-1}(p_2)G_\Gamma S^{-1}(p_1)$

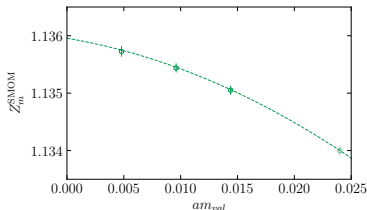
$$\begin{aligned}\frac{1}{12p^2} \text{Tr}[S^{-1}(p)\not{p}] &= -Z_q \\ \frac{1}{12} \frac{Z_S}{Z_q} \text{Tr}[\Lambda_S(p_1, p_2)] &= 1 \\ Z_m &= \frac{1}{Z_S}\end{aligned}$$

- ▶ Only propagator and scalar operator are required for mass renormalisation



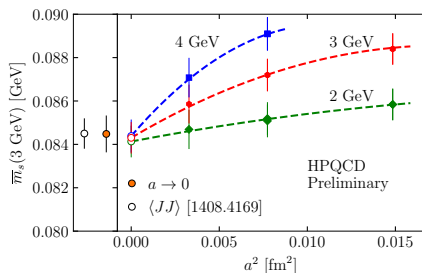
$\overline{\text{MS}}$ Matching

- ▶ To obtain masses in the $\overline{\text{MS}}$ scheme a matching calculation must be performed in the continuum
- ▶ Done perturbatively
- ▶ Perturbation theory does not include condensate effects that are present in the lattice calculation
- ▶ Chiral condensates are removed by $am_{\text{val}} \rightarrow 0$
- ▶ Other condensate is allowed for in continuum extrapolation



Continuum Extrapolation

- ▶ Calculate m_s and m_c on three different lattice spacings:
 $\sim \{0.12, 0.09, 0.06\}$ fm
- ▶ Extrapolate these results to $a = 0$
- ▶ Take account of $(a\mu)^2$, $(a\mu)^4$, $(a\Lambda)^2$ and $(a\Lambda)^4$ discretisation effects
 - ▶ HISQ action has disc. errors removed to this order
- ▶ Allow for condensate term



- ▶ Largest error from bare mass tuning
- ▶ May reduce extrapolation error by using smaller lattice spacing
 - ▶ Will also allow for higher μ points
- ▶ Look into Z_V ; relevant for $g - 2$ calculation

