

# Predicting Right-Handed Neutrino Masses from the Littlest Seesaw and Leptogenesis

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YTF10, 12 Jan 2018

ongoing work with S. Molina-Sedgwick and Prof. S.F. King

- Background
  - ▶ Seesaw Models: Type-I
  - ▶ The Littlest Seesaw
  - ▶ Leptogenesis
- This Work: *Predictions for RH neutrinos*
- Preliminary Results
- Outlook

# Why Seesaw Models?

- Standard Model cannot explain neutrino masses and oscillations or the observed BAU
- Need extension theory in order to be consistent with data
- Possible solution: models that utilise a Type-I Seesaw mechanism

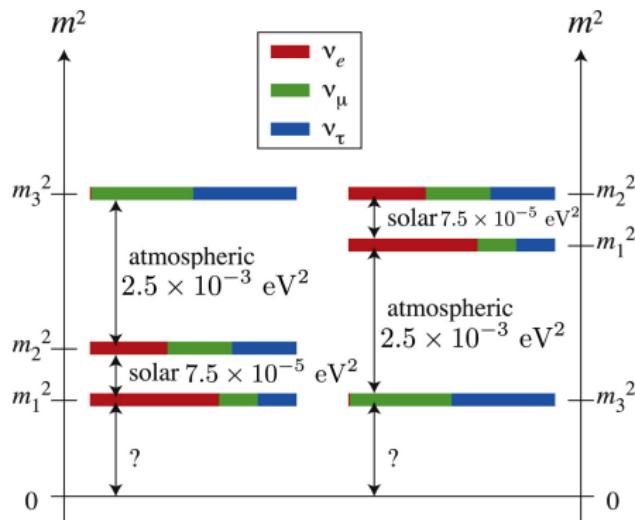


Figure: [King:1701.04413]

# The Type-I Seesaw

- Extend the SM by a number of right-handed neutrinos -  $\nu_R$
- Give rise to additional terms in the Lagrangian:

$$\mathcal{L}_m^D = -Y_\nu \bar{\ell}_L \tilde{H} \nu_R + h.c. \rightarrow -m_D \bar{\nu}_L \nu_R + h.c.$$

$$\mathcal{L}_m^M = -\frac{1}{2} M_R \bar{\nu}_R^c \nu_R + h.c.$$

- Collect terms together  $\implies$  neutrino mass matrix:

$$\begin{pmatrix} \bar{\nu}_L & \bar{\nu}_R^c \end{pmatrix} \begin{pmatrix} 0 & m_D \\ m_D^T & M_R \end{pmatrix} \begin{pmatrix} \nu_L^c \\ \nu_R \end{pmatrix}$$

- In the limit  $M_R \gg m_D$ , approximately diagonalise to obtain light neutrino masses:

$$m_\nu = m_D M_R^{-1} m_D^T = v^2 Y_\nu M_R^{-1} Y_\nu^T$$

Extend the SM by **two** RH  $\nu$  singlets:  $\nu_R = \begin{pmatrix} \nu_R^{atm} \\ \nu_R^{sol} \end{pmatrix}$

$$-\mathcal{L}_{LS} = -\mathcal{L}_{SM} + \left( Y_\nu \bar{\ell}_L \tilde{H} \nu_R + \frac{1}{2} M_R \overline{\nu_R^c} \nu_R + h.c. \right)$$

N.B.  $Y_\nu$  is a  $2 \times 3$  matrix and  $M_R$  is  $2 \times 2$

*Constrained Sequential Dominance*  $\implies$  heaviest RH $\nu$  gives dominant contribution to the heaviest LH neutrino mass

Two RH $\nu$ s  $\implies$  lightest LH neutrino is **massless**,  $m_1 = 0$

# The Littlest Seesaw

Fixes the absolute scale of neutrino masses

Constraints on  $\Delta m_{12}$  and  $\Delta m_{13} \rightarrow$  constraints on  $m_2$  and  $m_3$

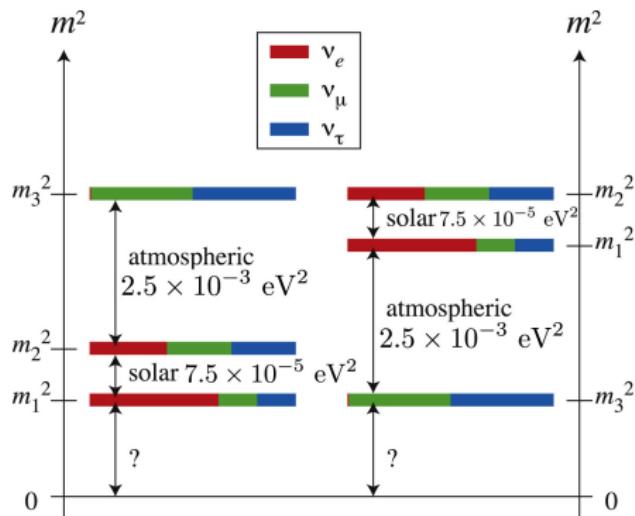


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Four distinct cases of the LS, labelled A through D.

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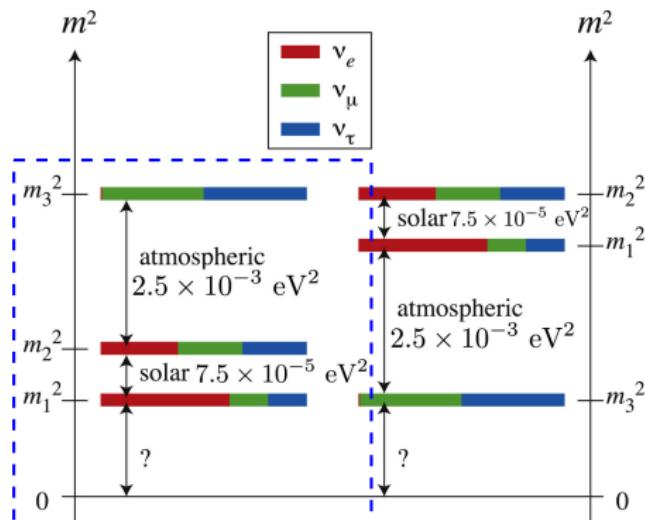


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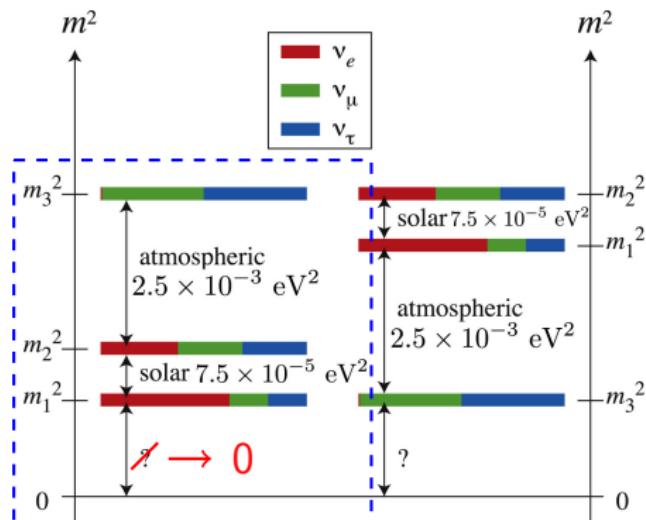


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Four distinct cases of the LS, labelled A through D.

# The Littlest Seesaw: A and B

- Case A:

$$Y_\nu^A = \begin{pmatrix} 0 & be^{i\eta/2} \\ a & nbe^{i\eta/2} \\ a & (n-2)be^{i\eta/2} \end{pmatrix}, \quad M_R^A = \begin{pmatrix} M_{atm} & 0 \\ 0 & M_{sol} \end{pmatrix}$$

- Case B:

$$Y_\nu^B = \begin{pmatrix} 0 & be^{i\eta/2} \\ a & (n-2)be^{i\eta/2} \\ a & nbe^{i\eta/2} \end{pmatrix}, \quad M_R^B = \begin{pmatrix} M_{atm} & 0 \\ 0 & M_{sol} \end{pmatrix}$$

$\eta$  is a complex phase and  $n$  is the order of CSD.

# The Littlest Seesaw: C and D

- Case C:

$$Y_\nu^C = \begin{pmatrix} be^{i\eta/2} & 0 \\ nbe^{i\eta/2} & a \\ (n-2)be^{i\eta/2} & a \end{pmatrix}, \quad M_R^C = \begin{pmatrix} M_{sol} & 0 \\ 0 & M_{atm} \end{pmatrix}$$

- Case D:

$$Y_\nu^D = \begin{pmatrix} be^{i\eta/2} & 0 \\ (n-2)be^{i\eta/2} & a \\ nbe^{i\eta/2} & a \end{pmatrix}, \quad M_R^D = \begin{pmatrix} M_{sol} & 0 \\ 0 & M_{atm} \end{pmatrix}$$

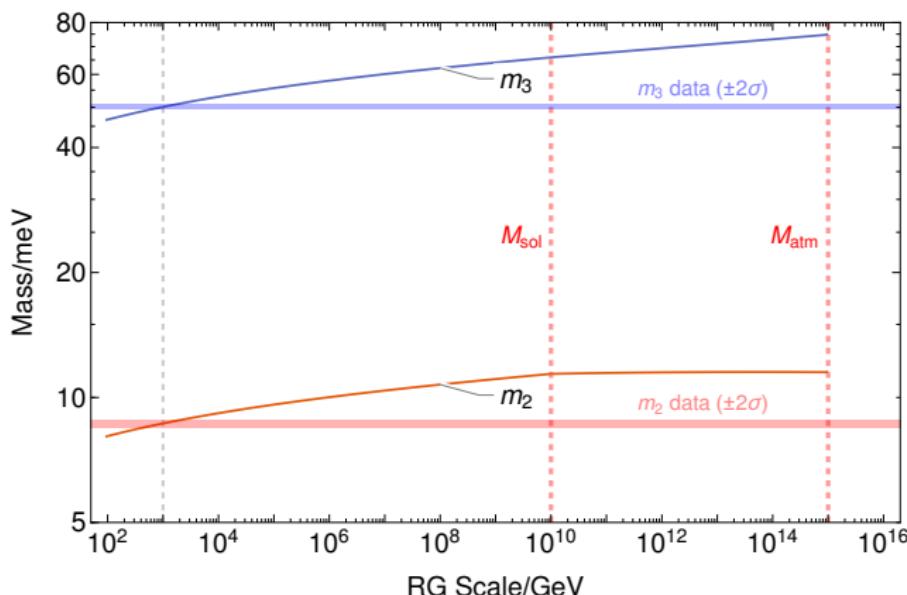
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- Theory defined at  $\mu = \Lambda_{GUT}$ , data available at low energies
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# RG Evolution

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- Evolve observables to low scales through **RG running**

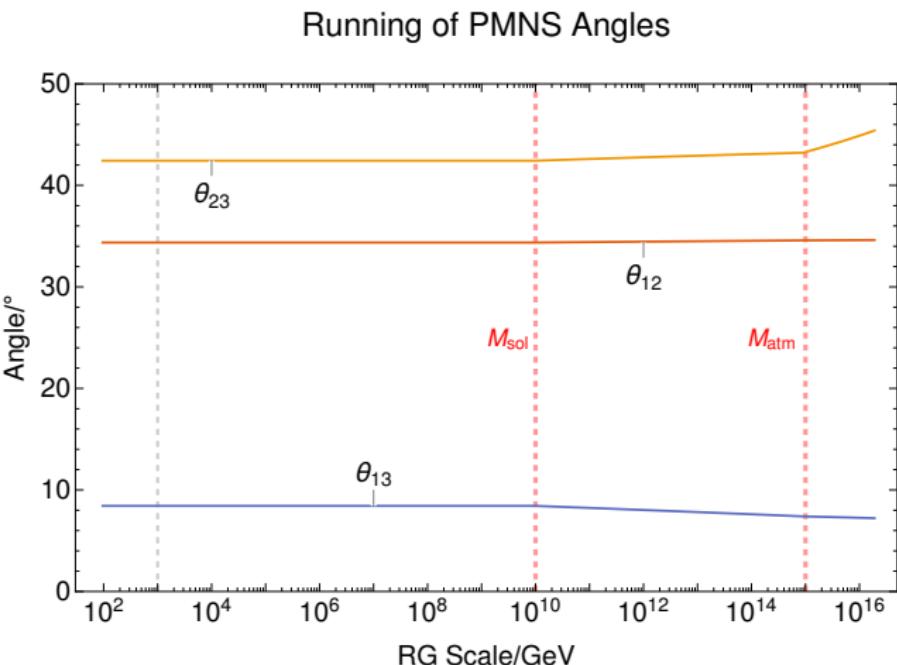
Running of Light  $v$  Mass Eigenstates



Case D:

$$\begin{aligned}M_{atm} &= 1.0 \times 10^{15} \text{ GeV} \\M_{sol} &= 1.0 \times 10^{10} \text{ GeV} \\a &= 1.24691 \\b &= 0.00114306 \\n &= 3 \\\eta &= -2\pi/3\end{aligned}$$

# RG Evolution (2)



Case D:

$$M_{atm} = 1.0 \times 10^{15} \text{ GeV}$$

$$M_{sol} = 1.0 \times 10^{10} \text{ GeV}$$

$$a = 1.24691$$

$$b = 0.00114306$$

$$n = 3$$

$$\eta = -2\pi/3$$

- Observed asymmetry normalised to the entropy density of the universe:

$$Y_B \simeq (0.87 \pm 0.03) \times 10^{-10}$$

- Leptogenesis: **mechanism for generating a BAU**
- Type-I seesaw models are able to generate a lepton asymmetry through the decay of the lightest RH $\nu$
- Lepton asymmetry converted to baryon imbalance through sphaleron processes in the SM

- Can we accurately predict the masses of the RH neutrinos in the Littlest Seesaw?

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- Can we accurately predict the masses of the RH neutrinos in the Littlest Seesaw?
- Use precision neutrino data

<i>Observable</i>	<i>Measured Value</i>
$\sin^2(\theta_{12})$	$0.306^{+0.012}_{-0.012}$
$\sin^2(\theta_{13})$	$0.02166^{+0.00075}_{-0.00075}$
$\sin^2(\theta_{23})$	$0.441^{+0.027}_{-0.021}$
$\delta$	$-99^\circ {}^{+51^\circ}_{-59^\circ}$
$\Delta m_{12}^2$	$7.50^{+0.19}_{-0.17} \times 10^{-5} \text{ eV}^2$
$\Delta m_{13}^2$	$2.524^{+0.039}_{-0.040} \times 10^{-3} \text{ eV}^2$

[Esteban et. al.(NuFit 3.0):  
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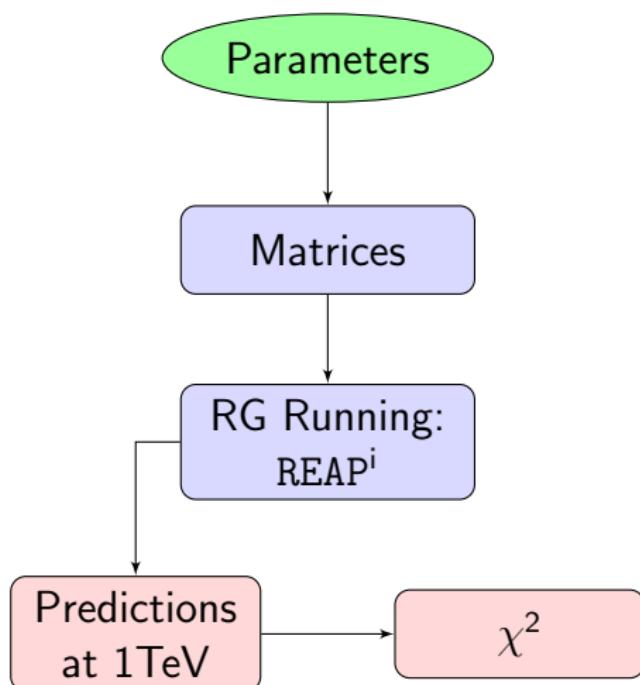
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- Use precision neutrino data and BAU from Leptogenesis

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[Spergel et. al.:  
[astro-ph/0603349](https://arxiv.org/abs/astro-ph/0603349)]



Scan over RH neutrino masses:

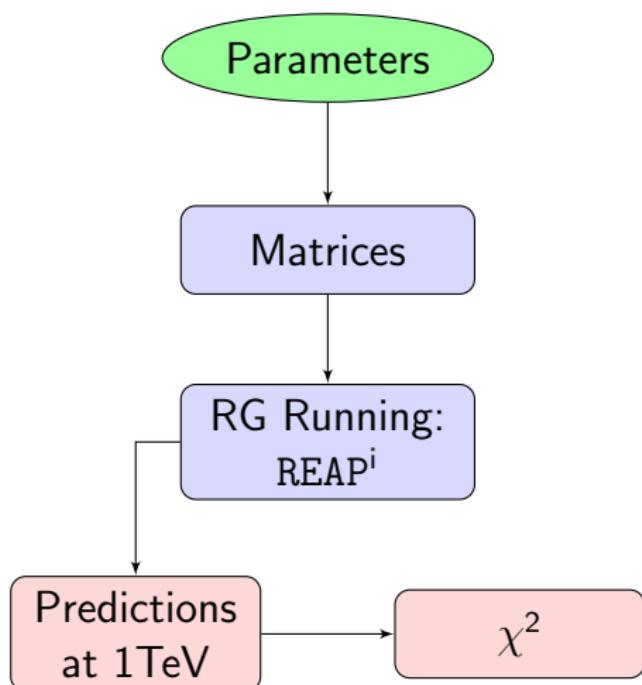
$$1.0 \times 10^9 \leq M_{atm} \leq 5.0 \times 10^{15}$$

$$1.0 \times 10^9 \leq M_{sol} \leq 5.0 \times 10^{15}$$

$a$  and  $b$  are left as free parameters.

$n$  fixed to be 3 and  $\eta$  fixed to be  $\pm 2\pi/3$

<sup>i</sup>Antusch et. al.:hep-ph/0501272.



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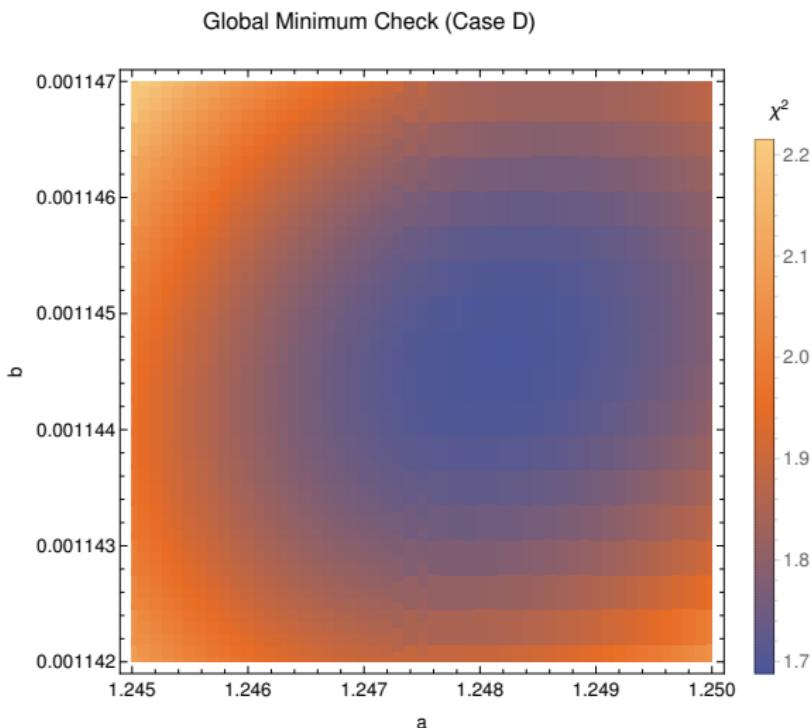
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Grid over parameter space, determine the best fit point for each case by minimising  $\chi^2$

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# Checking 2D Parameter Space

- Minimising  $\chi^2$  only useful if there is a stable global minimum in the parameter space
- Checked 2D parameter space (fix  $M_{atm}$  and  $M_{sol}$ ) to ensure global minimum is stable



# Preliminary Best Fit Points

Set of current best-fit points after a single coarse scan  
(no leptogenesis)

	Parameters				$\chi^2$
	$M_{atm}/\text{GeV}$	$M_{sol}/\text{GeV}$	$a$	$b$	
Case A	$1.0 \times 10^{11}$	$1.0 \times 10^{15}$	0.0113	0.370534	7.31496
Case B	$1.0 \times 10^{10}$	$1.0 \times 10^{15}$	0.00354	0.370534	4.93358
Case C	$1.0 \times 10^{15}$	$1.0 \times 10^{10}$	1.24691	0.00114	3.46989
Case D	$1.0 \times 10^{15}$	$1.0 \times 10^{10}$	1.24691	0.00114	<b>1.78567</b>

# Preliminary Best Fit Points: Case D

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Concentrate on Case D, look at variation of  $\chi^2$  with RH $\nu$  masses.

Varying  $M_{atm}$ :

Parameters				$\chi^2$
$M_{atm}/\text{GeV}$	$M_{sol}/\text{GeV}$	$a$	$b$	
$5.0 \times 10^{12}$	$1.0 \times 10^{10}$	0.0782107	0.00108097	10.0119
$1.0 \times 10^{13}$	$1.0 \times 10^{10}$	0.110607	0.00108097	13.0470
$5.0 \times 10^{13}$	$1.0 \times 10^{10}$	0.251836	0.00109801	8.56636
$1.0 \times 10^{14}$	$1.0 \times 10^{10}$	0.35615	0.00109801	11.0418
$5.0 \times 10^{14}$	$1.0 \times 10^{10}$	0.839184	0.00111505	3.50215
$1.0 \times 10^{15}$	$1.0 \times 10^{10}$	0.836253	0.00114306	1.78567

# Preliminary Best Fit Points: Case D

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Concentrate on Case D, look at variation of  $\chi^2$  with RH $\nu$  masses.

Varying  $M_{sol}$ :

Parameters				$\chi^2$
$M_{atm}/\text{GeV}$	$M_{sol}/\text{GeV}$	$a$	$b$	
$1.0 \times 10^{15}$	$1.0 \times 10^{10}$	0.836253	0.00114306	1.78567
$1.0 \times 10^{15}$	$5.0 \times 10^{10}$	0.836253	0.00257733	1.79572
$1.0 \times 10^{15}$	$1.0 \times 10^{11}$	1.24691	0.00367512	2.17086
$1.0 \times 10^{15}$	$5.0 \times 10^{11}$	1.24691	0.00828538	2.03957
$1.0 \times 10^{15}$	$1.0 \times 10^{12}$	1.24691	0.0117173	1.78793
$1.0 \times 10^{15}$	$5.0 \times 10^{12}$	1.24691	0.0264144	1.98070

- Neutrino data seems to be fixing heaviest RH neutrino mass

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- Include leptogenesis observables/data into our  $\chi^2$  fitting  $\implies$  constraints on lighter RH neutrino
- Perform more rigorous fitting by ‘homing in’ on desirable areas of the parameter space
- Possible use of Markov Chain Monte Carlo (MCMC) methods to find best fit in desirable region

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**Thank you for your attention**