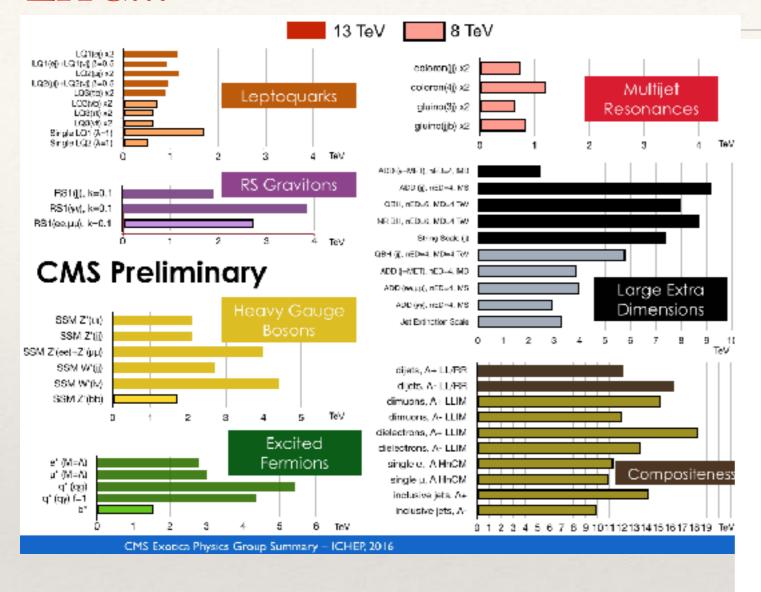


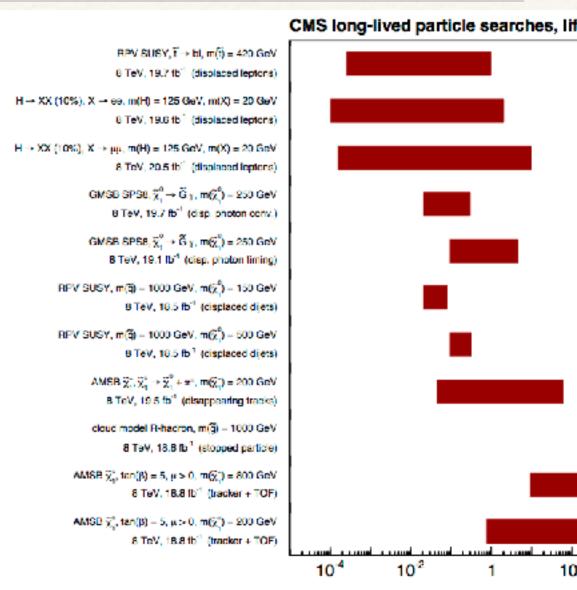
Christoph Englert

Theories killed by the LHC

YETI Durham, 10/01/2018

LHC...

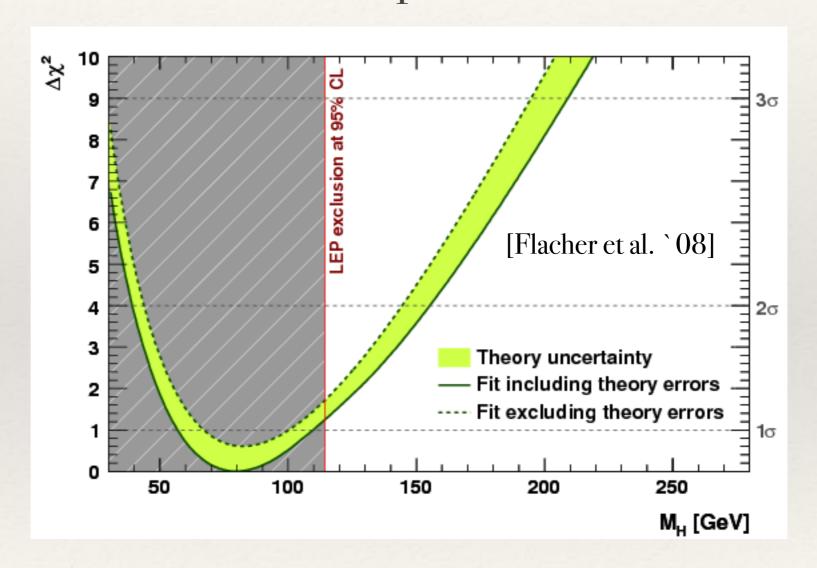




- left lots of theories constrained
- some constrained with little comeback options
- this talk: focus on concepts (and tensions) rather than quoting results
- put this in context with motivations for new experiments

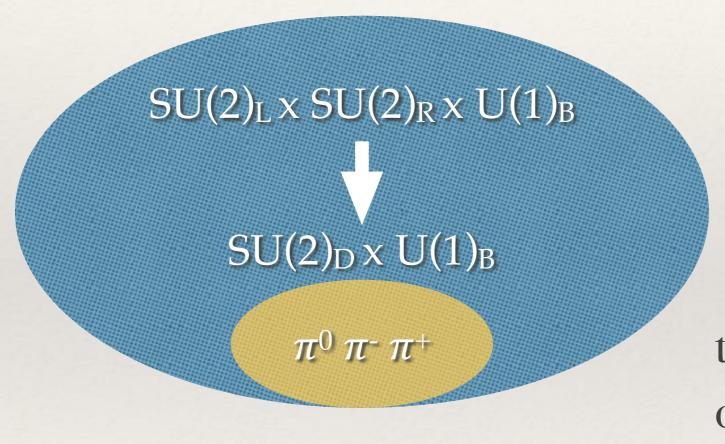
Before the LHC

- EWSB in the SM minimal yet ad-hoc $V(\Phi^{\dagger}\Phi) = \mu^2 \Phi^{\dagger}\Phi + \lambda (\Phi^{\dagger}\Phi)^2$
- the fact that we can understand the SM as a perturbative QFT allowed us to make self-consistent predictions



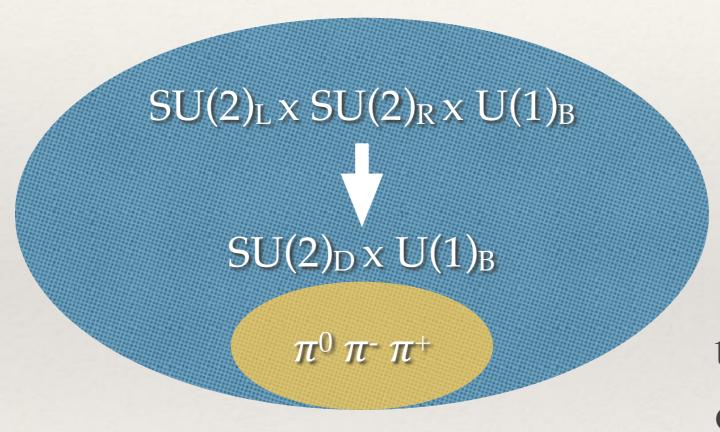
however, is perturbativity really a necessary for EWSB?

- not at all: can interpret the electroweak scale as a radiative phenomenon
- strong dynamics of QCD in chiral limit breaks global symmetries



three massless NGBs with quantum numbers of broken generators

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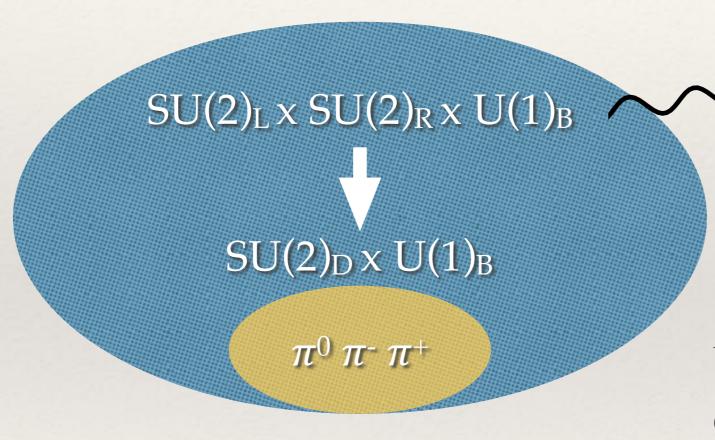


$$<0|j_{\mu}^{a5}|\pi^b>=if_{\pi}p_{\mu}\delta_{ab}$$

$$j_{\mu}^{a5} = \overline{\psi} \gamma_{\mu} \gamma^{5} \frac{\tau^{a}}{2} \psi$$
 where $\psi = (u, d)$

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 $SU(2)_L \times U(1)_Y$

gauge correlators receive QCD corrections

three massless NGBs with quantum numbers of broken generators

non-perturbative QCD effects break electroweak symmetry

$$W^{\pm} \sim \frac{\pi}{p^2 - g^2 f_{\pi}^2 / 4}$$

$$m_W = \frac{g f_{\pi}}{2} \simeq 29 \,\text{MeV}$$

• Technicolor: "scale up" QCD by adding new confining gauge interactions SU(N). The electroweak scale becomes a dimensional transmutation effect (like the QCD scale)

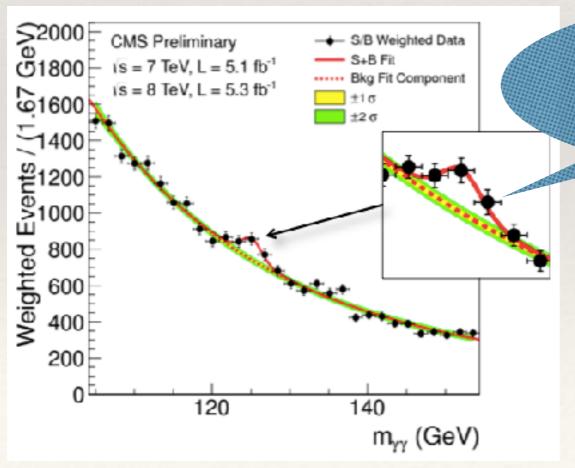
$$\mu \frac{d}{d\mu} \frac{1}{g_{TC}^2}(\mu) = -\frac{\beta_0}{8\pi^2} \implies v = M_{Pl} \exp\left(-\frac{8\pi^2}{g_{TC}^2(M_{Pl})(-\beta_0)}\right)$$

• implications:

- 1. no Higgs boson
- 2. additional strongly interacting bound states (rho-like etc.)
- 3. trouble with EWPD: $S \sim N_C/\pi$, but S < 0.3 from LEP
- 4. trouble with fermion masses: (walking) extended technicolor.

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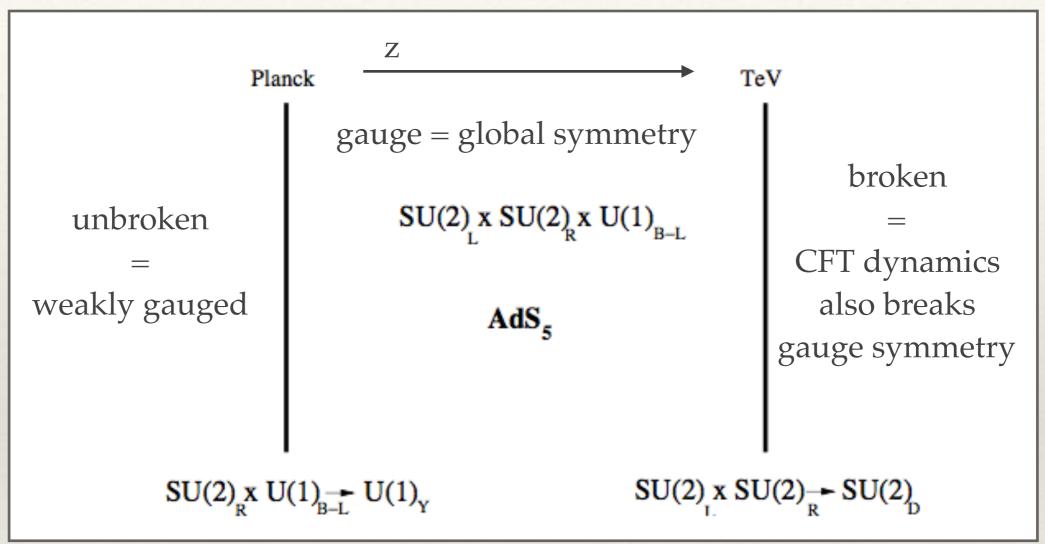


vanilla technicolor is ruled out by the LHC

Dualities

AdS/CFT dictionary

[Arkani-Hamed, Porrati, Randall `00] [Rattazzi, Zaffaroni `01]

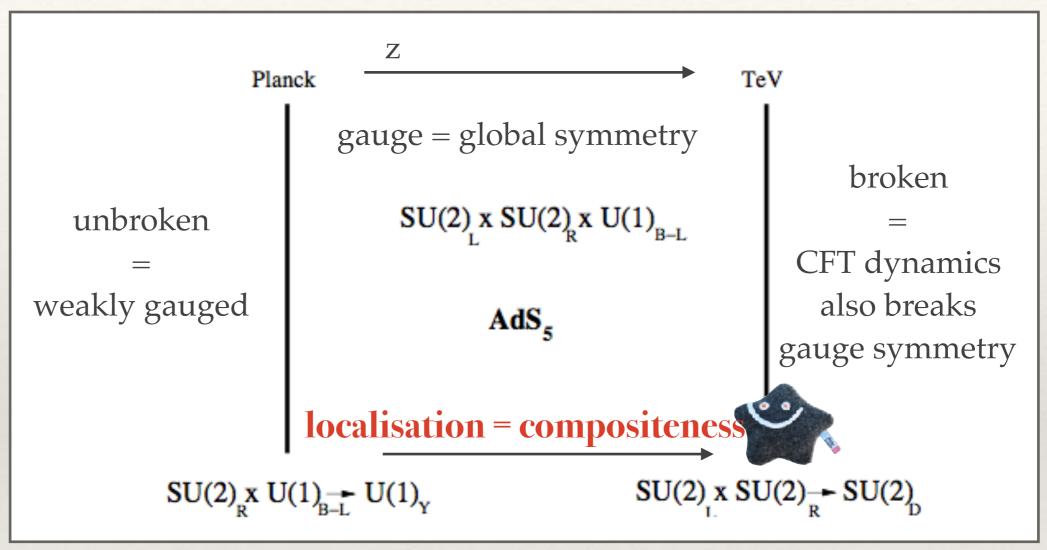


- Planck brane = UV cutoff of CFT
- bulk z = CFT energy scale, TeV brane = CFT breaks spontaneously due to strong interactions

Bringing the Higgs back

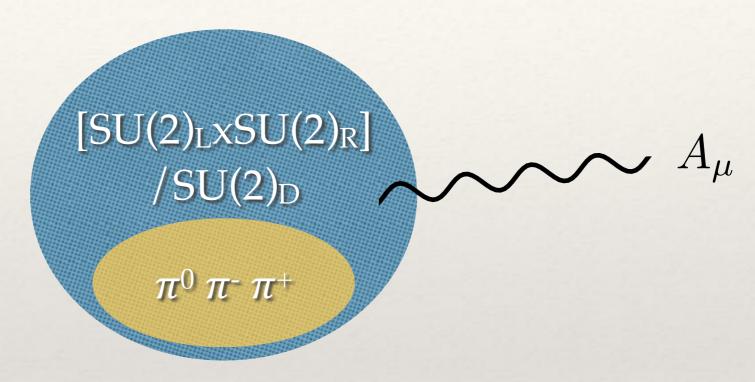
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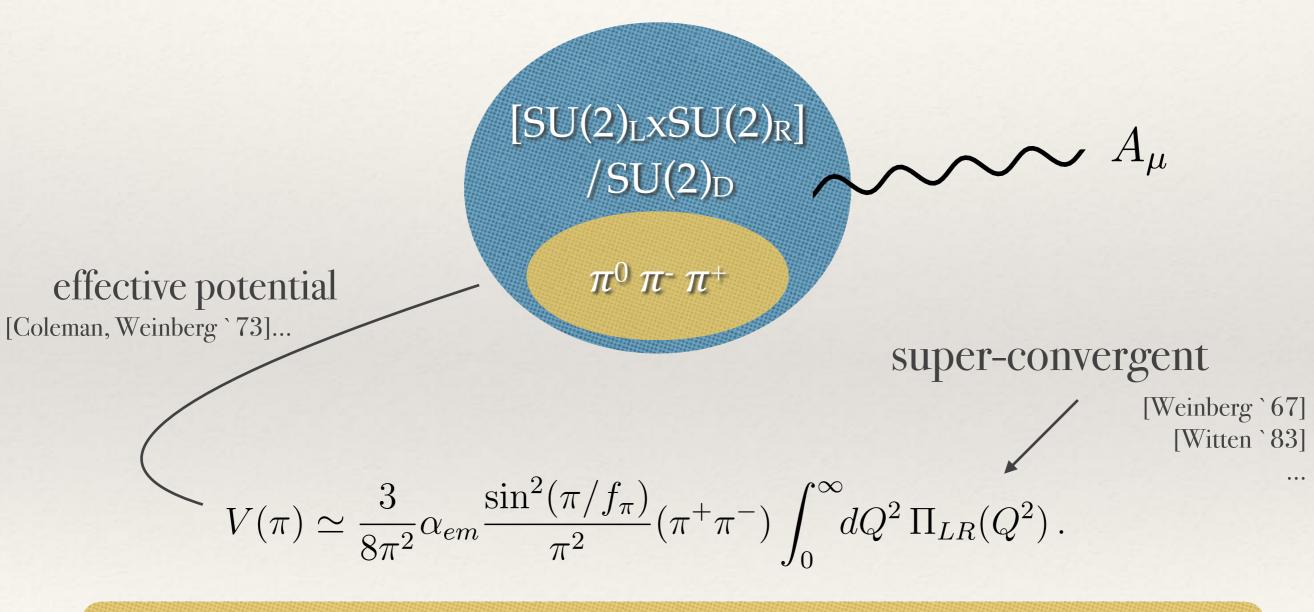


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• back to interpreting the electroweak scale as a radiative phenomenon, but this time look at the pion mass splitting



• back to interpreting the electroweak scale as a radiative phenomenon, but this time look at the pion mass splitting



$$(m_{\pi^{\pm}} - m_{\pi_0})|_{\text{TH}} \simeq 5.8 \,\text{MeV}$$
 vs $(m_{\pi^{\pm}} - m_{\pi_0})|_{\text{EXP}} \simeq 4.6 \,\text{MeV}$

• not straightforward to this adapt to the Higgs case

e.g. [Contino `10]

trigger
ELW symmetry
breaking not just
CW masses

respect global symmetries in the Higgs sector

LEP precision measurements

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LEP precision measurements

• complete vacuum mis-alignement from SU(2)_Lx U(1)_Y direction requires the presence of heavy fermions

in units of f

$$\hat{V}(\hat{h}) = \alpha \cos(2\hat{h}) - \beta \sin^2(2\hat{h})$$

gauge + fermions

fermions

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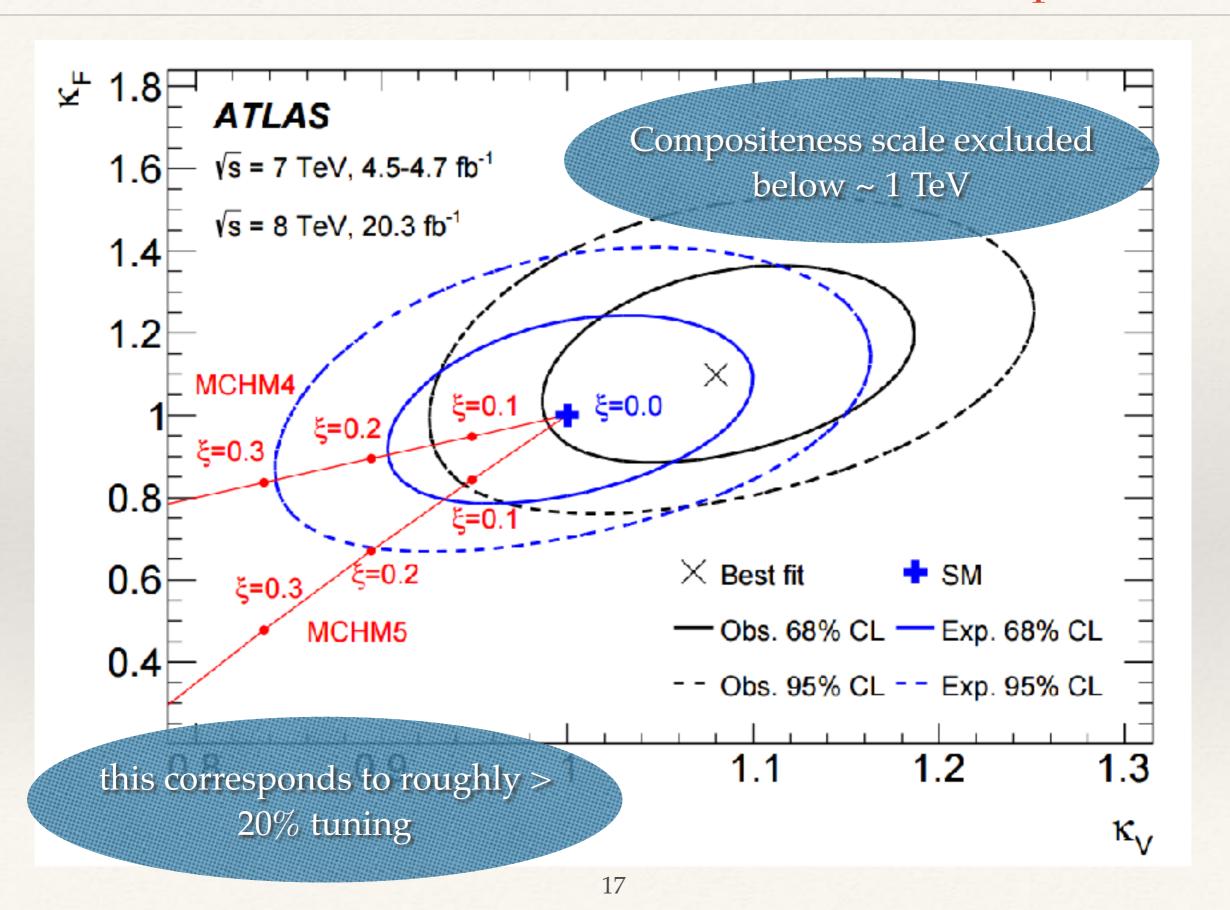
Higgs mass
$$\hat{m}_h^2 = \hat{V}''(\langle \hat{h} \rangle) = 32\beta \xi (1 - \xi) = 8\beta - 2\alpha^2/\beta$$

 \rightarrow tuning required to have $m_h \ll f$

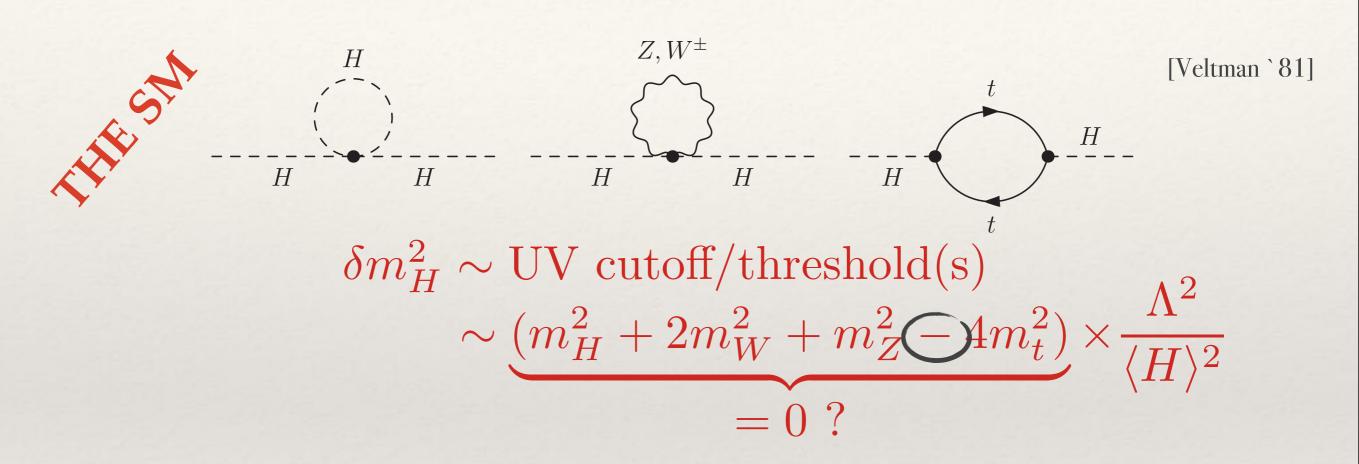
Higgs coupling modifier
$$\xi \equiv \frac{v^2}{f^2} = \sin^2(\langle \hat{h} \rangle) = \frac{\alpha + 2\beta}{4\beta}$$
.

$$g_{VVh} = \sqrt{1-\xi} g_{VVh}^{\rm SM}$$
 \rightarrow tuning not visible in couplings $\xi \sim 0$

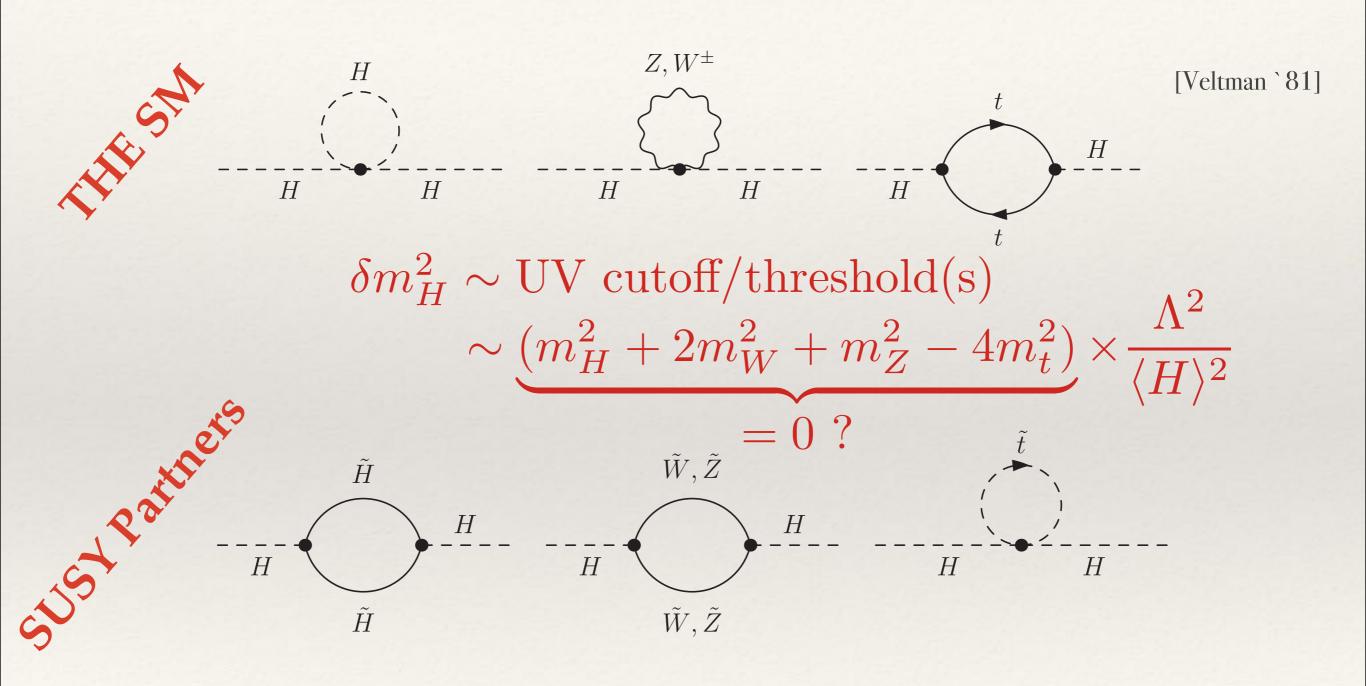
Generic hints of compositeness



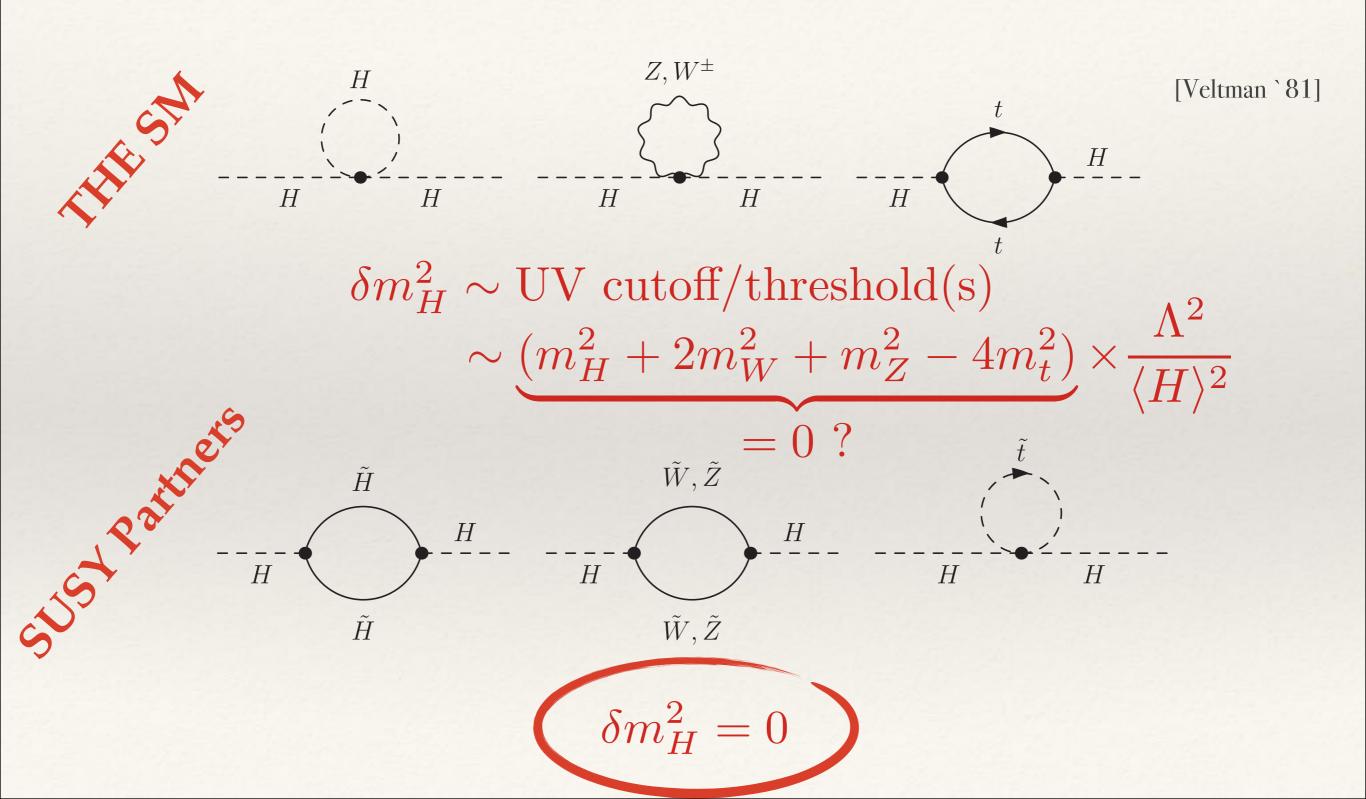
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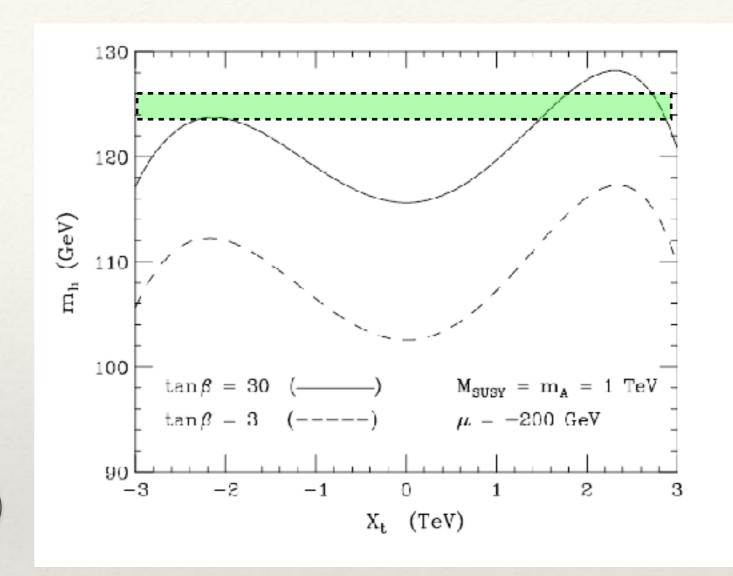


• then SUSY is not far (super well-motivated, but no clue of its scale)



- no degeneracy is observed so SUSY must be softly broken
- expressions for the CPeven Higgs mass

tree-level: $m_h \leq m_Z$

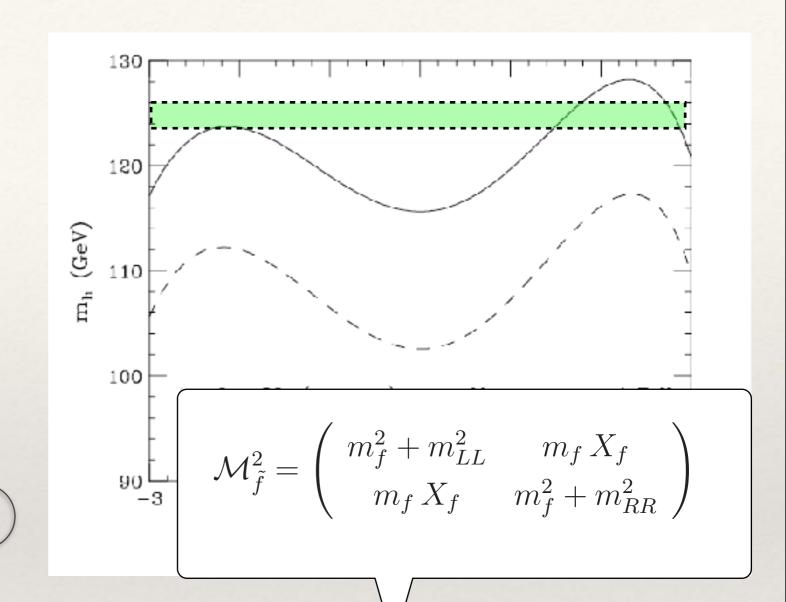


corrections:
$$m_h^2 \lesssim m_Z^2 + \frac{3g^2 m_t^4}{8\pi^2 m_W^2} \left[\ln\left(\frac{m_S^2}{m_t^2}\right) + \frac{X_t^2}{m_S^2} \left(1 - \frac{X_t^2}{12m_S^2}\right) \right]$$

 $X_t = A_t - \mu \cot \beta$ governs stop mixing and m_S^2 is average stop squared-mass

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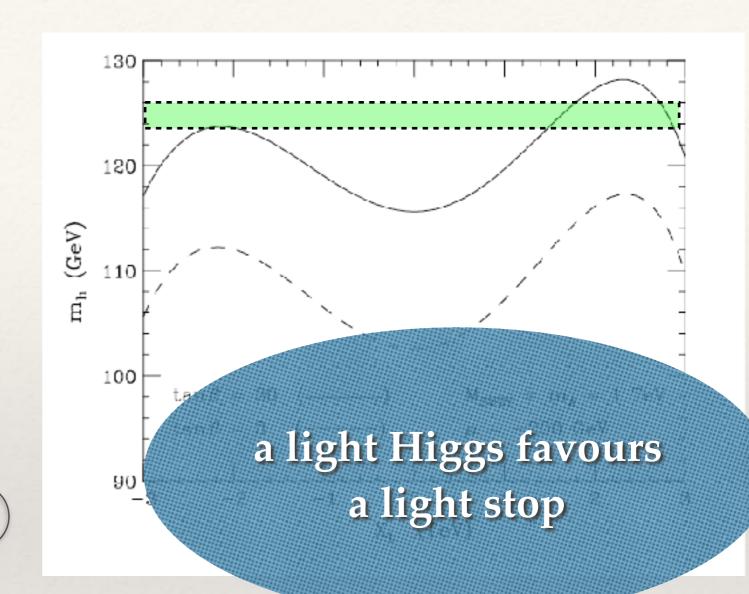


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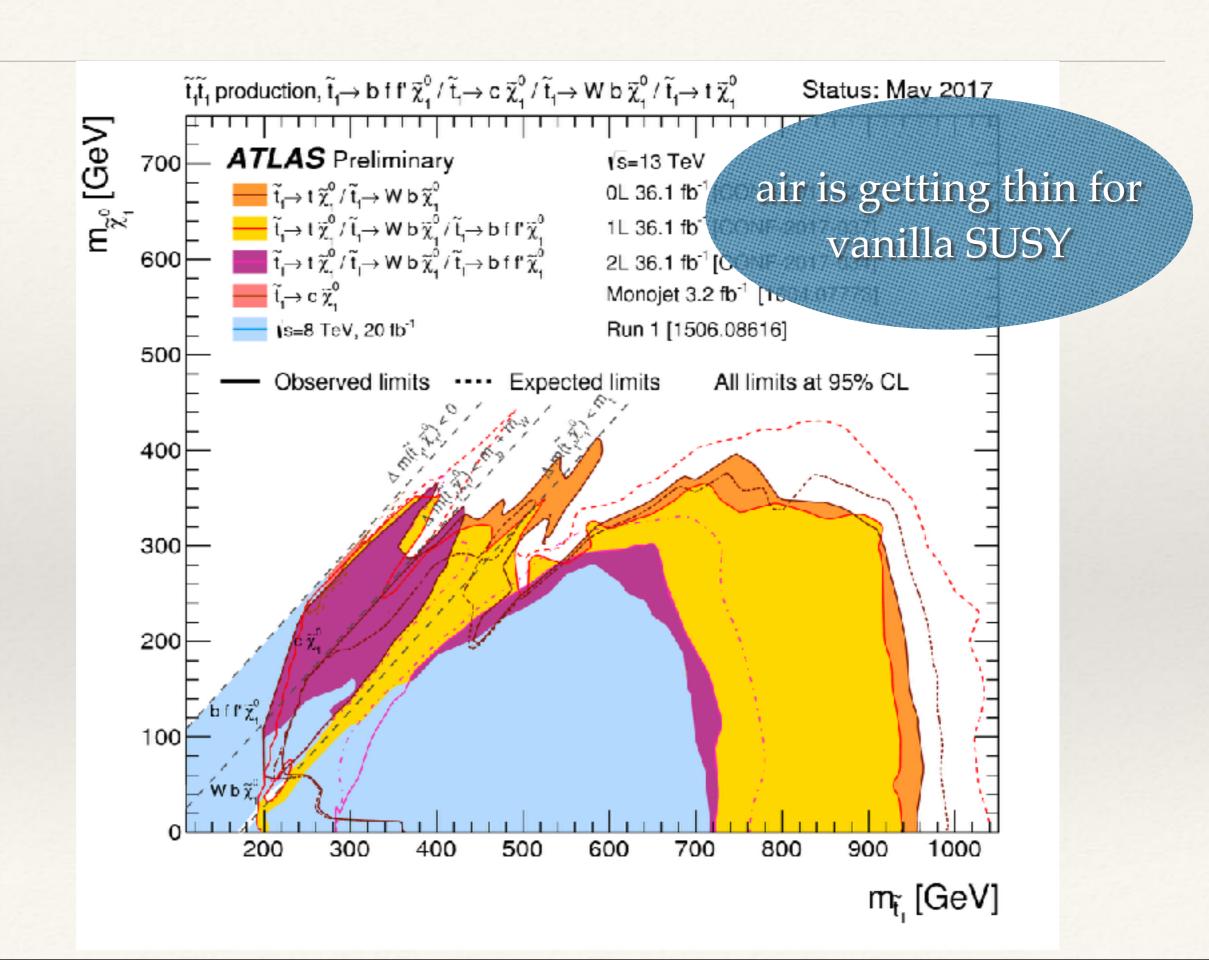
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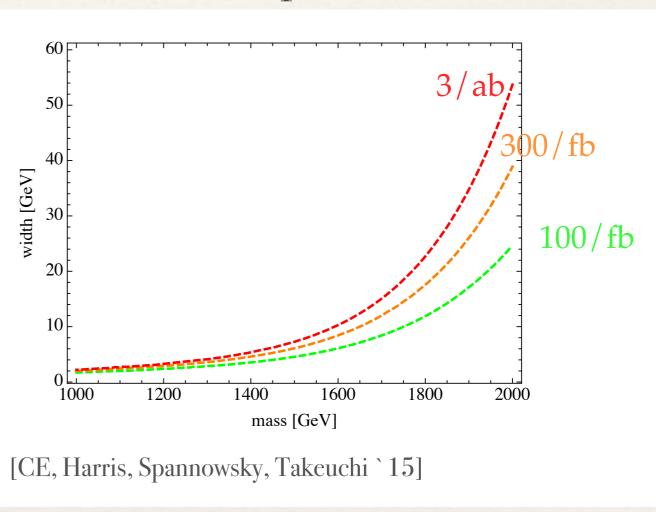
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 $X_t = A_t - \mu \cot \beta$ governs stop mixing and m_S^2 is average stop squared-mass



What are the options?

- any deviation of the SM coupling pattern induces perturbative unitarity violation $g_{VVh} = \sqrt{1 \xi} \, g_{VVh}^{\rm SM}$
- Composite scenarios have extra stuff that compensate this
 - fermiophobic = WBF
 - fermiophilic = Drell-Yan
 - LHC run 2 will zero in on those states
 - realistic spectra require lattice input

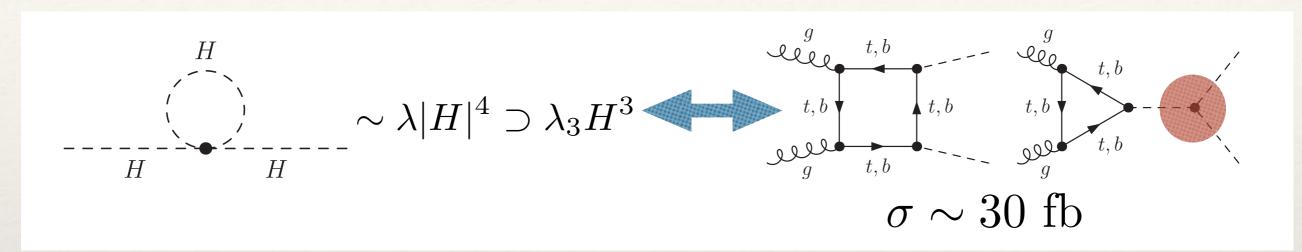


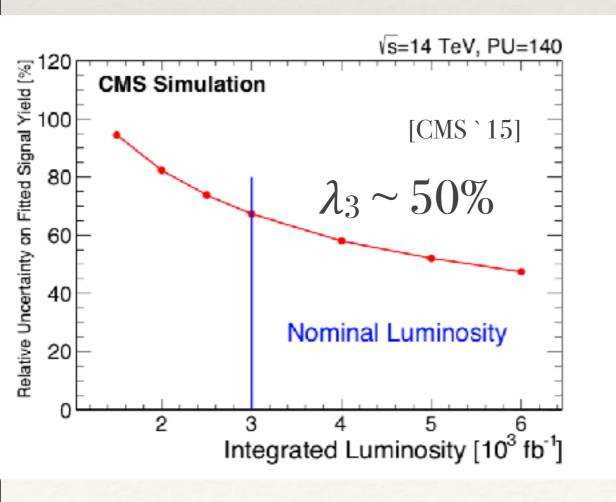
• Stops and other exotica searches scale with kinematic endpoints: 33 TeV / 100 TeV machines?

What are the options?

• enhance sensitivity to couplings that might be beyond the reach of the LHC

bkg systematics





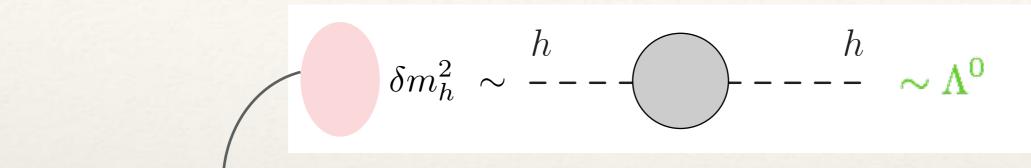
signal systematics

		$\Delta_S = 0.00$	$\Delta_S = 0.01$	$\Delta_S = 0.015$	$\Delta_S = 0.02$
0	$r_B = 0.5$	2.7%	3.4%	4.1%	4.9%
	$r_B = 1.0$	3.4%	3.9%	4.6%	5.3%
	$r_B = 1.5$	3.9%	4.4%	5.0%	5.7%
	$r_B = 2.0$	4.4%	4.8%	5.4%	6.0%
	$r_B = 3.0$	5.2%	5.6%	6.0%	6.6%

 $precision \ on \ \lambda_3$ [Mangano et al. , Physics at a 100 TeV collider `16]

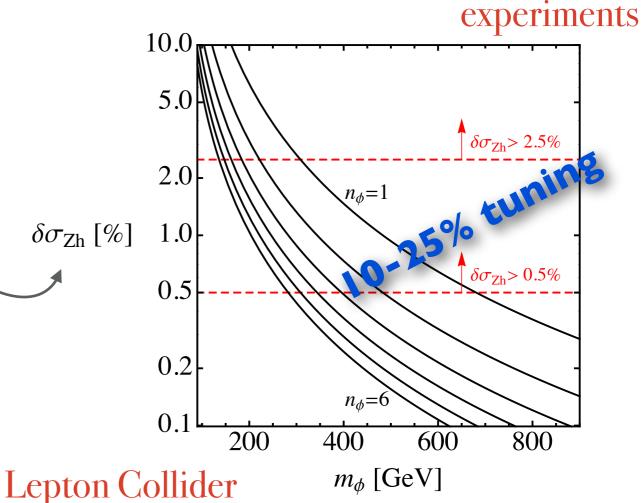
What are the options?

informed prediction about precision (model-dependent!)



$$\mathcal{L} = \mathcal{L}_{SM} + \frac{1}{2} \delta Z_h (\partial_\mu h)^2 + \dots$$

[Craig, CE, McCullough `13] [Goncalves, Han, Mukhopadhyay `17] Cancellations can't always hide at precision



- **►** We have the Higgs and it looks perturbative so far.
- **►** We have nothing else as data keeps pouring in.
- LHC: naturalness might not be such a great guiding principle after all?
- **►** future colliders: precision vs. energy?

Any BSM model gets pushed towards their decoupling limit or excluded