Supersymmetry in Dark Matter allowed regions

Alexander Belyaev



OUTLINE

- The status of the Standard Model: problems and solutions
- Supersymmetry as one of the best candidate for underlying theory
 - status of the Supersymmetry: theory versus experiment
 - dark Matter motivated regions and collider phenomenology
 - complementarity of the ILC and Dark matter search experiments
 - motivations for non-minimal models: beyond mSUGRA and beyond MSSM
- Conclusions

The present status of the SM

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The present status of the SM

- Based on SU(3)xSU(2)_LxU(1)_Y gauge symmetry spontaneously broken down to SU(3)xU(1)_e:
- Matter: 3 generations of quarks and leptons
- One of the central role is played by Higgs field
 - one higgs doublet, interacts with all fields
 - develops condensate
 - W,Z bosons, lepton and quarks and Higgs field itself acquires mass



Higgs boson is the most wanted particle! The present Higgs mass limit is M_H>114.4 GeV from LEP2

SM describes perfectly almost all data ...



- Experimental problems
 - Dark Matter & Dark Energy problem



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- Theoretical problems
 - the problem of large quantum corrections: fine-tuning problem

 $- \frac{f}{H} \frac{f}{f} - \frac{f}{H}$

 $\mathsf{SM}:\Delta M_{H}^{2}\sim\Lambda_{UV}^{2}$

 $M_{H}^{2} = M_{H^{0}}^{2} - \Delta M_{H}^{2},$ (100 GeV)² = (10¹⁶ GeV)² - (10¹⁶ GeV)² the cancellation is at the 28th digit for $\Lambda_{UV} \sim 10^{16}$ GeV

SM describes perfectly almost all ♂ data ... but has serious problems ♀

Experimental problems

- Dark Matter & Dark Energy problem
- matter anti-matter asymmetry: baryogenesis problem
- the origin of EWSB is still unknown.
 Higgs boson is not found yet ...
- Theoretical problems
 - the problem of large quantum corrections: fine-tuning problem
 - at very high energy forces start to behave similar log₁₀Q due to effect of different 'running' of coupling constants for abelian and non-abelian fields. But unification is not exact!



- boson-fermion symmetry aimed to unify all forces in nature $Q|BOSON\rangle = |FERMION\rangle, \quad Q|FERMION\rangle = |BOSON\rangle$
- extends Poincare algebra to Super-Poincare Algebra: the most general set of space-time symmetries! (1971-74)



Golfand and Likhtman'71; Ramond'71; Neveu,Schwarz'71; Volkov and Akulov'73; Wess and Zumino'74

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MSSM Higgs sector: two Higgs doublets

provide masses for up- and down-type fermions, cancellation of anomalies

→ 5 Higgs bosons h,H,A,H^{+/-:} M_A , $tan\beta = v_u/v_d$ define Higgs sector at tree-level

SUSY invented more then 30 years ago has 'little' problem

SUSY invented more then 30 years ago has 'little' problem it has not been found yet! Why it is still so attractive?



Consequences of SUSY

- Provides good DM candidate LSP
- CP violation can be incorporated baryogenesis via leptogenesis
- Radiative EWSB
- Solves fine-tuning problem
- Provides gauge coupling unification
- local supersymmetry requires spin 2 boson – graviton!
- allows to introduce fermions into string theories

 $\frac{h}{h_{t}} \frac{(TOP)}{h_{t}} \frac{h}{h_{t}} \frac{(STOP)}{h_{t}} \frac{h}{h_{t}}$ $\frac{h}{h_{t}^{2}} \frac{h}{h_{t}^{2}}$ $\Delta M_{H}^{2} \sim M_{SUSY}^{2} \log(\Lambda/M_{SUSY})$



SUSY was not deliberately designed to solve the SM problems!

SUSY is not observed, it must be broken



Gravity mediation Gauge mediation Anomaly mediation Gaugino mediation

$$\mathcal{L}_{soft}^{MSSM} = \underbrace{\sum_{i,j} B_{ij} \mu_{ij} S_i S_j}_{bilinear \ terms} + \underbrace{\sum_{ij} m_{ij}^2 S_i S_j^{\dagger}}_{scalar \ mass \ terms} + \underbrace{\sum_{i,j,k} A_{ijk} f_{ijk} S_i S_j S_k}_{trilinear \ scalar \ interactions} + \underbrace{\sum_{A,\alpha} M_{A\alpha} \overline{\lambda}_{A\alpha} \lambda_{A\alpha}}_{gaugino \ mass \ terms}$$

Minimal Supergravity Model (mSUGRA)

- visible-Hidden sectors interact with each other via gravity
- weak scale model constructed via RGE evolution, assuming:



Crucial constraint from Cosmology: DM candidate should be heavy, neutral, stable, non-baryonic Dark Matter candidate



SUSY has a perfect DM candidate, but this is only a beginning of the story ...

Evolution of neutralino relic density



relic density depends crucially on $\langle \sigma_A v \rangle$ thermal equilibrium stage: $T > m_{\chi}, \quad \chi \chi \leftrightarrow f \bar{f}$ universe cools: $T \leq m_{\chi}, \quad \chi \chi \not\leftrightarrow f \bar{f}$, $n = n_{eq} \sim e^{-m/T}$ neutralinos "freeze-out" at $T_F \sim m/25$

ISARED code: complete set of processes Baer, A.B., Balazs '02 exact tree-level calculations using CompHEP



Neutralino relic density in mSUGRA

most of the parameter space is ruled out! $\Omega h^2 \gg 1$ special regions with high σ_A are required to get $0.094 < \Omega h^2 < 0.129$



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Collider signatures in DM allowed regions

 DM allowed regions are difficult for the observation at the colliders: stau(stop) co-annihilation , FP region: small visible energy release



Why FP region is important

- small value of |µ|-parameter: mixed higgsino-bino LSP
- Light mass spectum of chargino and neutralinos
- low value of |μ|-parameter was advocated as "fine-tuning" measure Chan, Chattopadhyay,Nath '97; Feng, Matchev, Moroi '99; Baer, Chen,Paige,Tata '95
- DM motivated mSUGRA region with 'natural' neutralino mass ~100 GeV !
- ILC connection: the signal observation at the LHC is crucial for the fate of ILC





A. Belyaev Supersymmetry in Dark Matter Allowed Regions

IPPP, Durham, November 23

Recent Studies in FP region



'Far' FP analysis at the LHC

A.B, Genest, Leroy, Mehdiyev'07

- 'far' FP region dominated by EW chargino-neutralino production requires special cuts/analysis
- the signal observation in the 'far' FP region could be crucial for the fate of ILC



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Relative contributions of SUSY subprocesses (before cuts)

	$[3500,600] { m GeV}$	$[4670,975] { m GeV}$
Produced sparticles	Fraction of SUSY $\operatorname{events}(\%)$	Fraction of SUSY $\operatorname{events}(\%)$
$\tilde{W}_1 + \tilde{W}_1$	16.42	15.78
$\tilde{W}_2 + \tilde{W}_2$	5.88	4.46
$\tilde{W}_1 + \tilde{W}_2$	0.68	0.22
$\tilde{Z}_1 + \tilde{W}_1$	8.48	8.66
$\tilde{Z}_1 + \tilde{W}_2$	0.02	0.04
$\tilde{Z}_2 + \tilde{W}_1$	21.36	25.88
$\tilde{Z}_2 + \tilde{W}_2$	0.56	0.20
$\tilde{Z}_3 + \tilde{W}_1$	20.10	22.48
$\tilde{Z}_3 + \tilde{W}_2$	0.56	0.16
$\tilde{Z}_4 + \tilde{W}_2$	10.34	6.98
$\tilde{Z}_4 + \tilde{W}_1$	0.46	0.26
$\tilde{Z}_1 + \tilde{Z}_1$	0.02	0.02
$\tilde{Z}_1 + \tilde{Z}_2$	< 0.02	4.46
$\tilde{Z}_1 + \tilde{Z}_3$	3.72	< 0.02
$\tilde{Z}_2 + \tilde{Z}_3$	8.72	10.20
$\tilde{Z}_2 + \tilde{Z}_4$	< 0.02	0.04
$\tilde{Z}_3 + \tilde{Z}_4$	0.34	0.02
$\tilde{g} + \tilde{g}$	2.12	0.06

Signal and Backgrounds

signature $1\ell + jets + \not\!\!E_T$ signal $[m_0,m_{1/2}] = [3500,600] \longrightarrow ~240 \text{ fb}$ $t\bar{t}$ background $\longrightarrow ~20.7 \text{ pb}$ W+jets background $\longrightarrow ~366 \text{ pb}$

- $p_T^e > 20 \ GeV, \ p_T^\mu > 10 \ GeV$
- $p_T^J > 40 \ GeV \ within \ |\eta| < 3.0$
- Number of jets to be ≥ 4
- Number of leptons = 1
- $p_T^{J_1} \ge 500 \ GeV$
- $p_T^{J_2} \ge 300 \ GeV$
- $\Delta \phi(p_T^{lep}, E_T) \ge 20^\circ$



W+jets is dominant: PYTHIA W+jets underestimates BG by factor>3 as compared to Madgraph W+4jets which is used in our study

Improved strategy: softer preselection + new kinematical cuts



Improved strategy: softer preselection + new kinematical cuts



5

Further analysis of kinematical variables and correlations



Significance optimization



For SUSY datapoint $[m_0,m_{1/2}]=[3500,600]$ GeV produced in ISAJET v7.72, the statistical significance of the signal observation is shown as a function of the cut values for i) maximum R (with preselection cuts only), ii) maximum M_T (for preselection cuts only). The arrows represent the chosen cut values.

Signal and background efficiencies

	Pre-cuts	$p_T^{lep} < 200 \mathrm{GeV}$	$M_T \ge 160$	$R \le 1.5 { m ~GeV}$	All cuts
$[3500,\!600]$	2.65	97.01	39.21	91.14	0.92
v7.72					
[4000,700]	1.19	94.39	34.41	93.93	0.36
v7.72					
$t\bar{t}$	0.075	95.13	0.027	66.67	$1.3 \mathrm{x} 10^{-5}$
W+jets	0.09	85.01	0.27	20.0	$4.0 \mathrm{x} 10^{-5}$

Relative contributions of SUSY subprocesses (before/after cuts)

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$\tilde{W}_2 + \tilde{W}_2$	13.59	19.60
$\tilde{W}_1 + \tilde{W}_2$	< 0.49	0.35
$\tilde{Z}_1 + \tilde{W}_1$	2.43	4.90
$\tilde{Z}_1 + \tilde{W}_2$	< 0.49	< 0.35
$\tilde{Z}_2 + \tilde{W}_1$	6.31	14.00
$\tilde{Z}_2 + \tilde{W}_2$	< 0.49	0.30
$\tilde{Z}_3 + \tilde{W}_1$	7.77	12.90
$\tilde{Z}_3 + \tilde{W}_2$	0.97	0.35
$\tilde{Z}_4 + \tilde{W}_2$	26.21	31.50
$\tilde{Z}_4 + \tilde{W}_1$	1.94	0.70
$\tilde{Z}_1 + \tilde{Z}_1$	< 0.49	< 0.35
$\tilde{Z}_1 + \tilde{Z}_2$	< 0.49	< 0.35
$\tilde{Z}_1 + \tilde{Z}_3$	0.49	< 0.35
$\tilde{Z}_2 + \tilde{Z}_3$	0.49	0.70
$\tilde{Z}_2 + \tilde{Z}_4$	< 0.49	0.35
$\tilde{Z}_3 + \tilde{Z}_3$	< 0.49	< 0.35
$\tilde{g} + \tilde{g}$	29.61	1.40

Extended LHC reach



Extended LHC reach



Complementarity of Direct and Indirect DM search





LEP2 constraints

Light Higgs mass and LEP2 constraints: $M_H^{SM} > 114$ GeV pushes SUSY scale to 1TeV $M_h^2 = \frac{1}{2} \left[m_A^2 + M_Z^2 - \sqrt{(M_A^2 + M_Z^2)^2 - 4m_A^2 M_Z^2 \cos^2 2\beta} \right] \Rightarrow M_h \simeq M_Z |\cos 2\beta| \text{ for } M_a \gg M_Z$ Top-stop Radiative corrections to the light Higgs mass drive its mass up! $\delta M_h = \frac{3g^2 m_t^4}{8\pi^2 m_W^2} \left[\ln\left(\frac{M_S^2}{m_t^2}\right) + x_t^2 \left(1 - \frac{x_t^2}{12}\right) \right] \qquad \frac{h}{h_t} \begin{pmatrix} \text{TOP} & h & (\text{STOP}) \\ h_t & h & (\text{STOP}) \\ h$ $M_h \leq 135$ GeV for $M_S \sim 1$ TeV, for $x_t = \sqrt{6}$ (max mixing) Top-quark mass and EW fit: $m_t : 170.9 \rightarrow 178.0 \text{ GeV} \Rightarrow M_H : 76 \rightarrow 117.0 \text{ GeV}$ LEP2 SUSY particle search

• pair slepton production: $e^+e^- \rightarrow \tilde{\ell}^+_{L,R}\tilde{\ell}^-_{L,R} \rightarrow \ell^+\tilde{Z}_1\ell^-\tilde{Z}_1$ $\Rightarrow m_{\tilde{e}} > 99.6 \,\text{GeV}, \, m_{\tilde{\mu}} > 94.6 \,\text{GeV}, \, m_{\tilde{\tau}} > 85.9 \,\,\text{GeV}$

• pair chargino production: $e^+e^- \rightarrow \widetilde{W}_1^+ \widetilde{W}_1^-$, $\widetilde{W}_1 \rightarrow \widetilde{Z}_1 \ell \nu(\widetilde{Z}_1 q q')$, $\Rightarrow m_{\widetilde{W}_1} \gtrsim 100 \text{GeV}$

amplitude for H-mediated decay grows as $tan\beta^3$ (!) \Rightarrow relevant to high $tan\beta$ scenario [Babu,Kolda; Dedes,Dreiner,Nierste; Arnowitt,Dutta,Tanaka; Mizukoshi,Tata,Wang]

mSUGRA: combined constraints



Baer, A.B., Krupovnickas, Mustafayev hep-ph/0403214

mSUGRA: $\chi^2 = \chi^2_{\delta a_{\mu}} + \chi^2_{\Omega h^2} + \chi^2_{b \to s \gamma}$ analysis



Baer, A.B., Krupovnickas, Mustafayev hep-ph/0403214

Global CMSSM fit

68% (dotted) and 95% (solid) confidence level regions



O. Buchmueller, R. Cavanaugh, A. De Roeck, S. Heinemeyer, G. Isidori, P. Paradisi, F. Ronga, A. Weber, G. W. '07



NMH: SUSY spectra and LHC signatures



Scenario with non-universal Higgs masses (NUHM)

- universality of m₀ is motivated by the need to suppress unwanted flavor changing processes (generation blind mech for matter scalars in SUSY GUTs)
- ▶ this does not apply to soft breaking Higgs masses. In SO(10) SUSY GUTs: $(10 + \overline{5} + \overline{\nu}) \in \hat{\psi}(16), (5_H, \overline{5}_H) \in \hat{\phi}(10)$, different repres \Rightarrow SUSY breaking scalar mass terms for $\hat{\psi}(16)$ and $\hat{\phi}(10)$ are not expected to be the same

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the minimal non-universal Higgs extension of mSUGRA \Rightarrow NUHM1: $m_0, m_{\phi}, m_{1/2}, A_0, \tan\beta$ and $sign(\mu)$ $m_{\phi} = sign(m_{H_u,d}^2) \cdot \sqrt{|m_{H_{u,d}}^2|}$ $m_{H_u,d}^2$ are allowed to be negative

- μ becomes small for $m_{\phi} > m_0$ \Rightarrow FP! can be reached even for low m_0 and $m_{1/2}$!
- $\begin{tabular}{ll} \hline M_A \ decrease \ down \ to \ 2m_{\widetilde{Z}_1} \ for \\ m_\phi \ going \ down \Rightarrow \ Funnel! \ Even \\ for \ low \ tan \ \beta! \ Requires \ m_\phi^2 < 0. \end{tabular}$

Baer, A.B., Mustafayev, Profumo, Tata '05

 $m_0 = 300 \text{GeV}, m_{1/2} = 300 \text{GeV}, \tan\beta = 10, A_0 = 0, \mu > 0, m_r = 178 \text{GeV}$



Collider signatures for NUMH1 scenario



Tevatron: 3ℓ from $p\bar{p} \to W_1Z_2X$ followed by $\widetilde{W}_1 \to \ell \nu_\ell \widetilde{Z}_1$ and $\widetilde{Z}_2 \to \ell \overline{\ell} \widetilde{Z}_1$. When $m_\phi > m_0$ and $|\mu|$ is small \Rightarrow improved prospects for clean 3ℓ (no dominant events with tau leptons)

LHC: similar reach (in terms of $m_{\tilde{q}}$ and $m_{\tilde{g}}$ parameters) in the mSUGRA and NUHM1 models, but detailed gluino and squark cascade decays will change! H and AHiggs could be much lighter \Rightarrow direct production followed by $H, A \rightarrow \tau \bar{\tau}$.

ILC: In addition to "standard" mSUGRA FP signatures, H^0Z^0 , A^0h (possibly a good determ of $\tan \beta$), H^+H^- become accessible to study; $\tilde{Z}_1\tilde{Z}_3$, $\tilde{Z}_1\tilde{Z}_4$, $\tilde{Z}_2\tilde{Z}_2$, $\tilde{Z}_2\tilde{Z}_3$, $\tilde{Z}_2\tilde{Z}_4$ and even $\tilde{Z}_3\tilde{Z}_4$ as well as $\tilde{W}_1^{\pm}\tilde{W}_2^{\mp}$ are kinematically accessible.

SUSY spectroscopy would become a reality!

Conclusions

- SUSY is very compelling theory
- CDM constraints are crucial
- LHC: covers funnel region and stau-coannihilation region, just small portion of FP/HB is covered
- ILC: greatly extends LHC reach in FP/HB
- Extention of LHC reach in FP region could be crucial for ILC fate,
 - 2.4 TeV gluino mass is (indirectly) accessible with new analysis!
- direct/indirect DM search experiments: high degree of complementarity to LHC/ILC
- combined constraints: mSUGRA is practically excluded!
- one step beyond the universality solves many problems: NMH, NUMH, non-universal gauginos; motivated by SUSY GUTS

Present constraints/data, especially CDM one give a good idea how SUSY should look like at the upcoming experiments aimed to finally hunt down EW scale Supersymmetry!

Appendix

Sparticle reach of LHC for 100 fb⁻¹



Sparticle reach of LHC various luminosities





FP Region

HB/FP region for $m_{1/2} = 225$ GeV, $\tan \beta = 30$, $A_0 = 0$, $\mu > 0$: $\sqrt{s} = 500$ GeV



Simulations

ATLFAST $\Delta \eta \times \Delta \phi = 0.025 \times 0.025$ • Leptons $-2.5 < \eta < 2.5$ $0.1/\sqrt{E}(GeV) \bigoplus 0.007$ $E_T > 5 \text{ GeV}$ $\Delta \eta \times \Delta \phi = 0.1 \times 0.1$ • Jets $-2.5 < \eta < 2.5$ $0.5/\sqrt{E(GeV)} \oplus 0.03$ $E_T > 20 \text{ GeV}$ $\Delta R = 0.4$

SUSY GUTs

Gauge couplings unifi cation in the MSSM is the compelling hint for SUSY GUTs

- **SU(5)**[Georgi, Glashow(1974)] : { $Q = (u d), e^c, u^c$ } $\in \mathbf{10} \{d^c \ L = (v e)\} \in \mathbf{5}$ Higgs doublets have color triplet SU(5) partners: $(H_u T), (H_d T) \in \mathbf{5}_{\mathbf{H}}, \mathbf{\overline{5}}_{\mathbf{H}}$
- **SO**(10)[Georgi,Glashow;Fritzsch,Minkowski(1974)] : gauge and family AND two Higgs multiplet unifi cation: $(10 + \overline{5} + \overline{v}) \in 16$, $(5_H, \overline{5}_H) \in 10_H$

SO(10) SUSYGUT models are particularly intriguing:

- unify all matter of a single generation into the 16-d spinorial multiplet of SO(10)
- The 16 of SO(10) contains a gauge singlet v_R convenient for giving neutrinos mass (sea-saw: $m_{v_{\tau}} \simeq 0.03 \text{ eV} \Rightarrow M_N \sim 10^{14} \text{ GeV}$)
- SO(10) explains the cancellation of triangle anomalies
- Neutrino sector of SO(10) models lends itself to a theory of baryogenesis via leptogenesis
- Minimal SO(10) SUSYGUT: SM Higgs doublets are both 10d Higgs multiplet \Rightarrow Yukawa coupling unification: $f_t = f_b = f_{\tau}$, $W \ni f\hat{\psi}(\mathbf{16})^T\hat{\psi}(\mathbf{16})\hat{\phi}(\mathbf{10}) + \cdots$

However, 4D SUSY GUTs models have problems: large Higgs reps – cumbersome; spectrum of SM matter fi elds; rapid proton decay and doublet-triplet splitting problem.

SUSY GUTs

Recent progress in constructing SUSY GUT in 5+ space-time: GUT symmetry can be broken by compactifi cation of the extra dimensions on an appropriate topological manifold, such as an S₁/(Z₂ × Z'₂) orbifold [Kawamura; Hall,Nomura; Altarelli,Feruglio; Kobakhidze; Hebecker,March-Russel; Asaka,Buchmuller,Covi; Dermisek,Mafi ; Hall,Nomura,Okui,Smith]

Maintain positive features of 4D SUSYGUTS and solve many problems

Reduction of the gauge group upon breaking $SO(10) \Rightarrow D$ -term For $SO(10) \rightarrow SU(5) \times U(1)_X \rightarrow SU(3)_c \times SU(2)_L \times U(1)_Y$ one has: $m_Q^2 = m_E^2 = m_U^2 = m_{16}^2 + M_D^2$, $m_D^2 = m_L^2 = m_{16}^2 - 3M_D^2$, $m_N^2 = m_{16}^2 + 5M_D^2$, $m_{H_{u,d}}^2 = m_{10}^2 \mp 2M_D^2$ Parameters: m_{16} , m_{10} , M_D^2 , $m_{1/2}$, A_0 , $sign(\mu)$

results for SO(10) Yukawa unifi ed models: [Auto,Baer,Balazs,AB,Ferrandis,Tata; Blazek,Dermisek,Raby; Tobe,Wells] very specifi c param space and strong correlations

- $\tan \beta \sim 50$, $m_{16} \gtrsim 5$ TeV, $m_{1/2} \lesssim 100 150$ GeV light charginos and neutralinos may be accessible at Tevatron!
- $m_{10} \simeq \sqrt{2}m_{16}$, $A_0 \simeq -2m_{16}$, $M_D/M16 \simeq 0.33$, radiatively driven inverted scalar mass hierarchy regime [Bagger et al.]

Results for SO(10) model



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SO(10): DT (D-term) and HS (Higgs-split) models



