

Axions and anomalous $U(1)$'s

Quentin Bonnefoy

in collaboration with E. Dudas

Centre de Physique Théorique - École Polytechnique

IRN Terascale
IPPP, Durham University
September 6th 2018

Axions: pseudo Nambu-Goldstone bosons (pNGB)

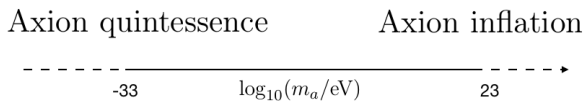
Axions: pseudo Nambu-Goldstone bosons (pNGB)

Used for: **naturalness** (strong-CP problem), **cosmology** (DM/quintessence/inflaton candidates), flavor physics, ...

Axions: pseudo Nambu-Goldstone bosons (pNGB)

Used for: **naturalness** (strong-CP problem), **cosmology** (DM/quintessence/inflaton candidates), flavor physics, ...

Phenomenology characterized by **axion mass**:



Axions: pseudo Nambu-Goldstone bosons (pNGB)

Used for: **naturalness** (strong-CP problem), **cosmology** (DM/quintessence/inflaton candidates), flavor physics, ...

Phenomenology characterized by **axion mass** and **couplings to SM fields**:

$$\begin{aligned} \mathcal{L} \supset & \frac{g_{a\gamma\gamma}}{f_a} a F_{\mu\nu} \tilde{F}^{\mu\nu} + \frac{i g_{a,\text{EDM}}}{f_a} a \bar{N} \gamma_{\mu\nu} \gamma^5 N F^{\mu\nu} \\ & + \frac{g_{aNN}}{f_a} \partial_\mu a \bar{N} \gamma^\mu \gamma^5 N + \frac{g_{aee}}{f_a} \partial_\mu a \bar{e} \gamma^\mu \gamma^5 e \end{aligned}$$

Axions: pseudo Nambu-Goldstone bosons (pNGB)

Used for: **naturalness** (strong-CP problem), **cosmology** (DM/quintessence/inflaton candidates), flavor physics, ...

Phenomenology characterized by **axion mass** and **couplings to SM fields**:

$$\mathcal{L} \supset \frac{g_{a\gamma\gamma}}{f_a} a F_{\mu\nu} \tilde{F}^{\mu\nu} + \frac{i g_{a,\text{EDM}}}{f_a} a \bar{N} \gamma_{\mu\nu} \gamma^5 N F^{\mu\nu} \\ + \frac{g_{aNN}}{f_a} \partial_\mu a \bar{N} \gamma^\mu \gamma^5 N + \frac{g_{aee}}{f_a} \partial_\mu a \bar{e} \gamma^\mu \gamma^5 e$$

Constrained by **numerous complementary experiments**

UV origin of axions:

UV origin of axions:

- **Field theory axions:** models with a spontaneously broken global symmetry

UV origin of axions:

- **Field theory axions:** models with a spontaneously broken global symmetry
⇒ **light NGB** with model-dependent couplings

Peccei-Quinn (77), Weinberg, Wilczek (78), Kim, Shifman-Vainshtein-Zakharov (79), Zhitnisky (80), Dine-Fischler-Srednicki (81), ...

UV origin of axions:

- **Field theory axions:** models with a spontaneously broken global symmetry
⇒ **light NGB** with model-dependent couplings
Perturbative or non-perturbative explicit breaking

Peccei-Quinn (77), Weinberg, Wilczek (78), Kim, Shifman-Vainshtein-

Zakharov (79), Zhitnisky (80), Dine-Fischler-Srednicki (81), ...

UV origin of axions:

- **Field theory axions:** models with a spontaneously broken global symmetry
⇒ **light NGB** with model-dependent couplings
Perturbative or non-perturbative explicit breaking

Peccei-Quinn (77), Weinberg, Wilczek (78), Kim, Shifman-Vainshtein-

Zakharov (79), Zhitnisky (80), Dine-Fischler-Srednicki (81), ...

- **String theory axions:** 10D low-energy supergravities contain p -forms

UV origin of axions:

- **Field theory axions:** models with a spontaneously broken global symmetry
 \implies **light NGB** with model-dependent couplings
 Perturbative or non-perturbative explicit breaking

Peccei-Quinn (77), Weinberg, Wilczek (78), Kim, Shifman-Vainshtein-

Zakharov (79), Zhitnisky (80), Dine-Fischler-Srednicki (81), ...

- **String theory axions:** 10D low-energy supergravities contain p -forms
 After compactification, internal d.o.f \implies **massless 4D scalars** with higher-dim. gauge protection

Witten (84), Arvanitaki-Dimopoulos-Dubovsky-Kaloper-March-Russell (10), ...

UV origin of axions:

- **Field theory axions:** models with a spontaneously broken global symmetry
 \implies **light NGB** with model-dependent couplings
 Perturbative or non-perturbative explicit breaking

Peccei-Quinn (77), Weinberg, Wilczek (78), Kim, Shifman-Vainshtein-

Zakharov (79), Zhitnisky (80), Dine-Fischler-Srednicki (81), ...

- **String theory axions:** 10D low-energy supergravities contain p -forms
 After compactification, internal d.o.f \implies **massless 4D scalars** with higher-dim. gauge protection
 Gauge or stringy instantons explicit breaking

Witten (84), Arvanitaki-Dimopoulos-Dubovsky-Kaloper-March-Russell (10), ...

Anomalous $U(1)$'s and axions:

Anomalous $U(1)$'s and axions:

Numerous effective string models have an **anomalous** $U(1)$ gauge factor

$$\mathcal{L} = -\bar{\psi}_i \not{D} \psi_i$$

Dine-Seiberg-Witten (87), Atick-Dixon-Sen (87), ...

Anomalous $U(1)$'s and axions:

Numerous effective string models have an **anomalous** $U(1)$ gauge factor with Green-Schwarz mechanism:

$$\mathcal{L} = -\bar{\psi}_i \not{D} \psi_i - \frac{1}{2} (\partial_\mu \theta - \delta_{GS} A_{X,\mu})^2 + \frac{k_G}{16\pi^2} \theta \text{Tr}(G\tilde{G})$$

Dine-Seiberg-Witten (87), Atick-Dixon-Sen (87), ...

Anomalous $U(1)$'s and axions:

Numerous effective string models have an **anomalous** $U(1)$ gauge factor with Green-Schwarz mechanism:

$$\mathcal{L} = -\bar{\psi}_i \not{D} \psi_i - \frac{1}{2} (\partial_\mu \theta - \delta_{GS} A_{X,\mu})^2 + \frac{k_G}{16\pi^2} \theta \text{Tr}(G\tilde{G})$$

\implies **anomaly-given axion with anomaly-given gauge couplings**

Dine-Seiberg-Witten (87), Atick-Dixon-Sen (87), ...

Questions:

Questions:

- which **predictions for the SM couplings?**
 - Scales/decay constant?
 - Model-independent values?
 - Hierarchies?

Questions:

- which **predictions for the SM couplings?**
 - Scales/decay constant?
 - Model-independent values?
 - Hierarchies?
- can those axions be **QCD axions?**
- ...

I. Anomalous $U(1)$ theories: physical axions and gauge couplings

$$\mathcal{L} \supset -\bar{\psi}_i \not{D} \psi_i - \frac{1}{2} (\partial_\mu \theta - \delta_{GS} A_{X,\mu})^2 + \frac{k_G}{16\pi^2} \theta \text{Tr}(G\tilde{G}) + \dots$$

I. Anomalous $U(1)$ theories: physical axions and gauge couplings

$$\mathcal{L} \supset -\bar{\psi}_i \not{D} \psi_i - \frac{1}{2} (\partial_\mu \theta - \delta_{GS} A_{X,\mu})^2 + \frac{k_G}{16\pi^2} \theta \text{Tr}(G\tilde{G}) + \dots$$

$U(1)_X$ already in Stueckelberg phase: additional charged scalar with a vev \implies physical axion. Moreover, perturbatively massless.

I. Anomalous $U(1)$ theories: physical axions and gauge couplings

$$\mathcal{L} \supset -\bar{\psi}_i \not{D} \psi_i - \frac{1}{2} (\partial_\mu \theta - \delta_{GS} A_{X,\mu})^2 + \frac{k_G}{16\pi^2} \theta \text{Tr}(G\tilde{G}) + \dots$$

$U(1)_X$ already in Stueckelberg phase: additional charged scalar with a vev \implies physical axion. Moreover, perturbatively massless.

If massive fermions, integrate out:

$$\mathcal{L} \supset \left(\underbrace{\frac{k_G}{16\pi^2} \theta - \# \theta - \sum \#_i \theta_i}_{\text{gauge-invariant physical axion}} \right) \text{Tr}(G\tilde{G}) = \frac{k_G}{16\pi^2} (\theta - \#' \theta - \#'_i \theta_i) \text{Tr}(G\tilde{G})$$

I. Anomalous $U(1)$ theories: physical axions and gauge couplings

$$\mathcal{L} \supset -\bar{\psi}_i \not{D} \psi_i - \frac{1}{2} (\partial_\mu \theta - \delta_{GS} A_{X,\mu})^2 + \frac{k_G}{16\pi^2} \theta \text{Tr}(G\tilde{G}) + \dots$$

$U(1)_X$ already in Stueckelberg phase: additional charged scalar with a vev \implies physical axion. Moreover, perturbatively massless.

If massive fermions, integrate out:

$$\mathcal{L} \supset \underbrace{\left(\frac{k_G}{16\pi^2} \theta - \# \theta - \sum \#_i \theta_i \right)}_{\text{gauge-invariant physical axion}} \text{Tr}(G\tilde{G}) = \frac{k_G}{16\pi^2} (\theta - \#' \theta - \#'_i \theta_i) \text{Tr}(G\tilde{G})$$

\implies anomaly-given gauge couplings

I. Anomalous $U(1)$ theories: physical axions and gauge couplings

If one physical axion:

$$\mathcal{L} \supset \frac{k_G}{16\pi^2} (\theta - \#'\theta - \#'_i\theta_i) \text{Tr}(G\tilde{G}) + \frac{k'_G}{16\pi^2} (\theta - \#'\theta - \#'_i\theta_i) \text{Tr}(G'\tilde{G}') + \dots$$

I. Anomalous $U(1)$ theories: physical axions and gauge couplings

If one physical axion:

$$\mathcal{L} \supset \frac{k_G}{16\pi^2} (\theta - \#'\theta - \#'_i\theta_i) \text{Tr}(G\tilde{G}) + \frac{k'_G}{16\pi^2} (\theta - \#'\theta - \#'_i\theta_i) \text{Tr}(G'\tilde{G}') + \dots$$

Ex: $SU(5)$ unification of the SM gauge group:

$$\frac{5}{3}k_Y = k_2 = k_3 \implies \frac{\mathbf{E}}{\mathbf{N}} = \frac{\mathbf{8}}{\mathbf{3}} \text{ (at the GUT scale)}$$

II. Anomalous $U(1)$ theories: heterotic string effective theories

II. Anomalous $U(1)$ theories: heterotic string effective theories

Heterotic string theory: **one** possible **anomalous $U(1)_X$** and **one (axion-dilaton) field**

$$\mathcal{L} = \mathcal{L}_{GS} + \mathcal{L}_{\text{matter}}$$

$$\mathcal{L}_{GS} = - \int d^4\theta \ln(S + \bar{S} - \delta_{GS} V_X) + \int d^2\theta \frac{k_G S}{4} \text{Tr}(W_G^\alpha W_{G,\alpha}) + h.c.$$

II. Anomalous $U(1)$ theories: heterotic string effective theories

Heterotic string theory: **one** possible **anomalous $U(1)_X$** and **one (axion-dilaton) field**

$$\mathcal{L} = \mathcal{L}_{GS} + \mathcal{L}_{\text{matter}}$$

$$\mathcal{L}_{GS} = - \int d^4\theta \ln(S + \bar{S} - \delta_{GS} V_X) + \int d^2\theta \frac{k_{GS}}{4} \text{Tr}(W_G^\alpha W_{G,\alpha}) + h.c.$$

as well as SUSY preserving matter

$$\mathcal{L}_{\text{matter}} \supset \int d^4\theta \phi^\dagger e^{-2V_X} \phi \implies V_D = \frac{g_X^2}{2} \left(\frac{\delta_{GS}}{2(S + \bar{S})} - |\phi|^2 \right)^2 \approx 0$$

II. Anomalous $U(1)$ theories: heterotic string effective theories

Heterotic string theory: **one** possible **anomalous $U(1)_X$** and **one (axion-dilaton) field**

$$\mathcal{L} = \mathcal{L}_{GS} + \mathcal{L}_{\text{matter}}$$

$$\mathcal{L}_{GS} = - \int d^4\theta \ln(S + \bar{S} - \delta_{GS} V_X) + \int d^2\theta \frac{k_{GS}}{4} \text{Tr}(W_G^\alpha W_{G,\alpha}) + h.c.$$

as well as SUSY preserving matter

$$\mathcal{L}_{\text{matter}} \supset \int d^4\theta \phi^\dagger e^{-2V_X} \phi \implies V_D = \frac{g_X^2}{2} \left(\frac{\delta_{GS}}{2(S + \bar{S})} - |\phi|^2 \right)^2 \approx 0$$

In this setup: **a** **physical axion**

II. Anomalous $U(1)$ theories: heterotic string effective theories

Heterotic string theory: **one** possible **anomalous $U(1)_X$** and **one (axion-dilaton) field**

$$\mathcal{L} = \mathcal{L}_{GS} + \mathcal{L}_{\text{matter}}$$

$$\mathcal{L}_{GS} = - \int d^4\theta \ln(S + \bar{S} - \delta_{GS} V_X) + \int d^2\theta \frac{k_{GS}}{4} \text{Tr}(W_G^\alpha W_{G,\alpha}) + h.c.$$

as well as SUSY preserving matter

$$\mathcal{L}_{\text{matter}} \supset \int d^4\theta \phi^\dagger e^{-2V_X} \phi \implies V_D = \frac{g_X^2}{2} \left(\frac{\delta_{GS}}{2(S + \bar{S})} - |\phi|^2 \right)^2 \approx 0$$

In this setup: **a/only one physical axion**

II. Anomalous $U(1)$ theories: heterotic string effective theories

Heterotic string theory: **one** possible **anomalous $U(1)_X$** and **one (axion-dilaton) field**

$$\mathcal{L} = \mathcal{L}_{GS} + \mathcal{L}_{\text{matter}}$$

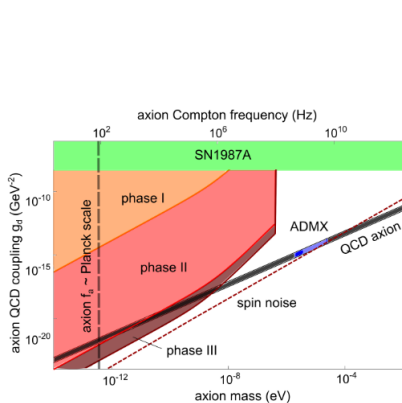
$$\mathcal{L}_{GS} = - \int d^4\theta \ln(S + \bar{S} - \delta_{GS} V_X) + \int d^2\theta \frac{k_{GS}}{4} \text{Tr}(W_G^\alpha W_{G,\alpha}) + h.c.$$

as well as SUSY preserving matter

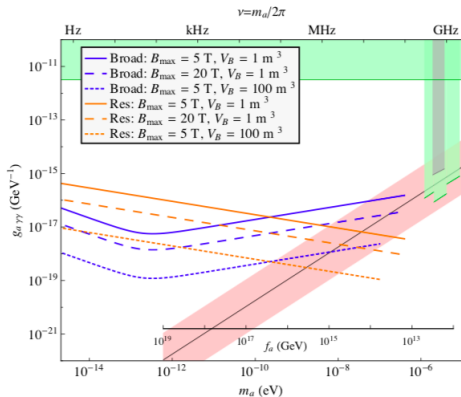
$$\mathcal{L}_{\text{matter}} \supset \int d^4\theta \phi^\dagger e^{-2V_X} \phi \implies V_D = \frac{g_X^2}{2} \left(\frac{\delta_{GS}}{2(S + \bar{S})} - |\phi|^2 \right)^2 \approx 0$$

In this setup: **a/only one physical axion, string/GUT scale decay constant**

II. Anomalous $U(1)$ theories: heterotic string effective theories



Kimball et al., arXiv:1711.08999

Kahn-Safdi-Thaler, *PRL* 117 (2016)

III. Anomalous $U(1)$ theories: Froggatt-Nielsen mechanism

Froggatt-Nielsen mechanism \implies **anomalous $U(1)$ symmetry**

III. Anomalous $U(1)$ theories: Froggatt-Nielsen mechanism

Froggatt-Nielsen mechanism \implies **anomalous $U(1)$ symmetry**, can be gauged using a Green-Schwarz mechanism

III. Anomalous $U(1)$ theories: Froggatt-Nielsen mechanism

Froggatt-Nielsen mechanism \implies **anomalous $U(1)$ symmetry**, can be gauged using a Green-Schwarz mechanism

$$\begin{aligned}\mathcal{L} \supset & -\frac{1}{2}(\partial_\mu\theta - \delta_{GS}A_{X,\mu})^2 + \sum_i \frac{k_i}{16\pi^2}\theta \operatorname{Tr}(G_i\tilde{G}_i) \\ & + \lambda_{u,ij} \left(\frac{\phi}{\Lambda}\right)^{X_{q_i}+X_{u_j}+X_{h_u}} \bar{u}_j Q_i H_u + \lambda_{d,ij} \left(\frac{\phi}{\Lambda}\right)^{X_{q_i}+X_{d_j}+X_{h_d}} \bar{d}_j Q_i H_d \\ & + \lambda_{e,ij} \left(\frac{\phi}{\Lambda}\right)^{X_{l_i}+X_{e_j}+X_{h_d}} \bar{e}_j L_i H_d + \mu \left(\frac{\phi}{\Lambda}\right)^{X_{h_u}+X_{h_d}} H_u H_d + h.c.\end{aligned}$$

III. Anomalous $U(1)$ theories: Froggatt-Nielsen mechanism

Froggatt-Nielsen mechanism \implies **anomalous $U(1)$ symmetry**, can be gauged using a Green-Schwarz mechanism

$$\begin{aligned}\mathcal{L} \supset & -\frac{1}{2}(\partial_\mu\theta - \delta_{GS}A_{X,\mu})^2 + \sum_i \frac{k_i}{16\pi^2}\theta \text{Tr}(G_i\tilde{G}_i) \\ & + \lambda_{u,ij} \left(\frac{\phi}{\Lambda}\right)^{X_{q_i}+X_{u_j}+X_{h_u}} \bar{u}_j Q_i H_u + \lambda_{d,ij} \left(\frac{\phi}{\Lambda}\right)^{X_{q_i}+X_{d_j}+X_{h_d}} \bar{d}_j Q_i H_d \\ & + \lambda_{e,ij} \left(\frac{\phi}{\Lambda}\right)^{X_{l_i}+X_{e_j}+X_{h_d}} \bar{e}_j L_i H_d + \mu \left(\frac{\phi}{\Lambda}\right)^{X_{h_u}+X_{h_d}} H_u H_d + h.c.\end{aligned}$$

\rightarrow **anomaly predictions for SM gauge field couplings**

III. Anomalous $U(1)$ theories: Froggatt-Nielsen mechanism

Froggatt-Nielsen mechanism \implies **anomalous $U(1)$ symmetry**, can be gauged using a Green-Schwarz mechanism

$$\begin{aligned} \mathcal{L} \supset & -\frac{1}{2}(\partial_\mu\theta - \delta_{GS}A_{X,\mu})^2 + \sum_i \frac{k_i}{16\pi^2}\theta \operatorname{Tr}(G_i\tilde{G}_i) \\ & + \lambda_{u,ij} \left(\frac{\phi}{\Lambda}\right)^{X_{q_i}+X_{u_j}+X_{h_u}} \bar{u}_j Q_i H_u + \lambda_{d,ij} \left(\frac{\phi}{\Lambda}\right)^{X_{q_i}+X_{d_j}+X_{h_d}} \bar{d}_j Q_i H_d \\ & + \lambda_{e,ij} \left(\frac{\phi}{\Lambda}\right)^{X_{l_i}+X_{e_j}+X_{h_d}} \bar{e}_j L_i H_d + \mu \left(\frac{\phi}{\Lambda}\right)^{X_{h_u}+X_{h_d}} H_u H_d + h.c. \end{aligned}$$

→ **anomaly predictions for SM gauge field couplings**

→ **spin couplings correlated with mass matrix entries:**

$$\frac{\partial_\mu a}{f_a} (\overline{\psi_{L,i}} U_{L,ij}^\dagger X_{L,j} \gamma^\mu U_{L,jk} \psi_{L,k} + \overline{\psi_{R,i}} U_{R,ij}^\dagger X_{R,j} \gamma^\mu U_{R,jk} \psi_{R,k})$$

III. Anomalous $U(1)$ theories: Froggatt-Nielsen mechanism

→ spin couplings correlated with mass matrix entries:

$$\frac{\partial_\mu a}{f_a} (\overline{\psi_{L,i}} U_{L,ij}^\dagger X_{L,j} \gamma^\mu U_{L,jk} \psi_{L,k} + \overline{\psi_{R,i}} U_{R,ij}^\dagger X_{R,j} \gamma^\mu U_{R,jk} \psi_{R,k})$$

⇒ possible **flavor-changing effects**

III. Anomalous $U(1)$ theories: Froggatt-Nielsen mechanism

→ spin couplings correlated with mass matrix entries:

$$\frac{\partial_\mu a}{f_a} (\overline{\psi_{L,i}} U_{L,ij}^\dagger X_{L,j} \gamma^\mu U_{L,jk} \psi_{L,k} + \overline{\psi_{R,i}} U_{R,ij}^\dagger X_{R,j} \gamma^\mu U_{R,jk} \psi_{R,k})$$

⇒ possible **flavor-changing effects**

Bounds from meson physics: $f_a \gtrsim 10^{10-11} \text{ GeV}$

Calibbi-Goertz-Redigolo-Ziegler-Zupan (2016), Ema-Hamaguchi-Moroi-Nakayama (2016)

IV. Anomalous $U(1)$ theories: QCD effects VS the rest

QCD axion if QCD mass dominates ($m_{\text{rest}} f_a \lesssim 10^{-5} m_{\pi} f_{\pi}$)

IV. Anomalous $U(1)$ theories: QCD effects VS the rest

QCD axion if QCD mass dominates ($m_{\text{rest}} f_a \lesssim 10^{-5} m_{\pi} f_{\pi}$)

SUSY/string frameworks: **saxion/modulus** need to be **non-perturbatively stabilized** \implies possible **explicit breaking of the axion shift symmetry**

IV. Anomalous $U(1)$ theories: QCD effects VS the rest

QCD axion if QCD mass dominates ($m_{\text{rest}} f_a \lesssim 10^{-5} m_\pi f_\pi$)

SUSY/string frameworks: **saxion/modulus** need to be **non-perturbatively stabilized** \implies possible **explicit breaking of the axion shift symmetry**

Tied to stabilization details (and possibly SUSY breaking, ...).

Anomalous $U(1)$ models are UV-motivated and point to axion physics.

Anomalous $U(1)$ models are UV-motivated and point to axion physics.

They make predictions about axion-gauge bosons couplings, and can be used in flavor physics.

Anomalous $U(1)$ models are UV-motivated and point to axion physics.

They make **predictions about axion-gauge bosons couplings**, and can be **used in flavor physics**.

Their associated **high-scales** may be **within reach of proposed/developed/current experiments**.

Thank you!