

Quark-mass effects in Higgs production

Jonas M. Lindert



Durham
University

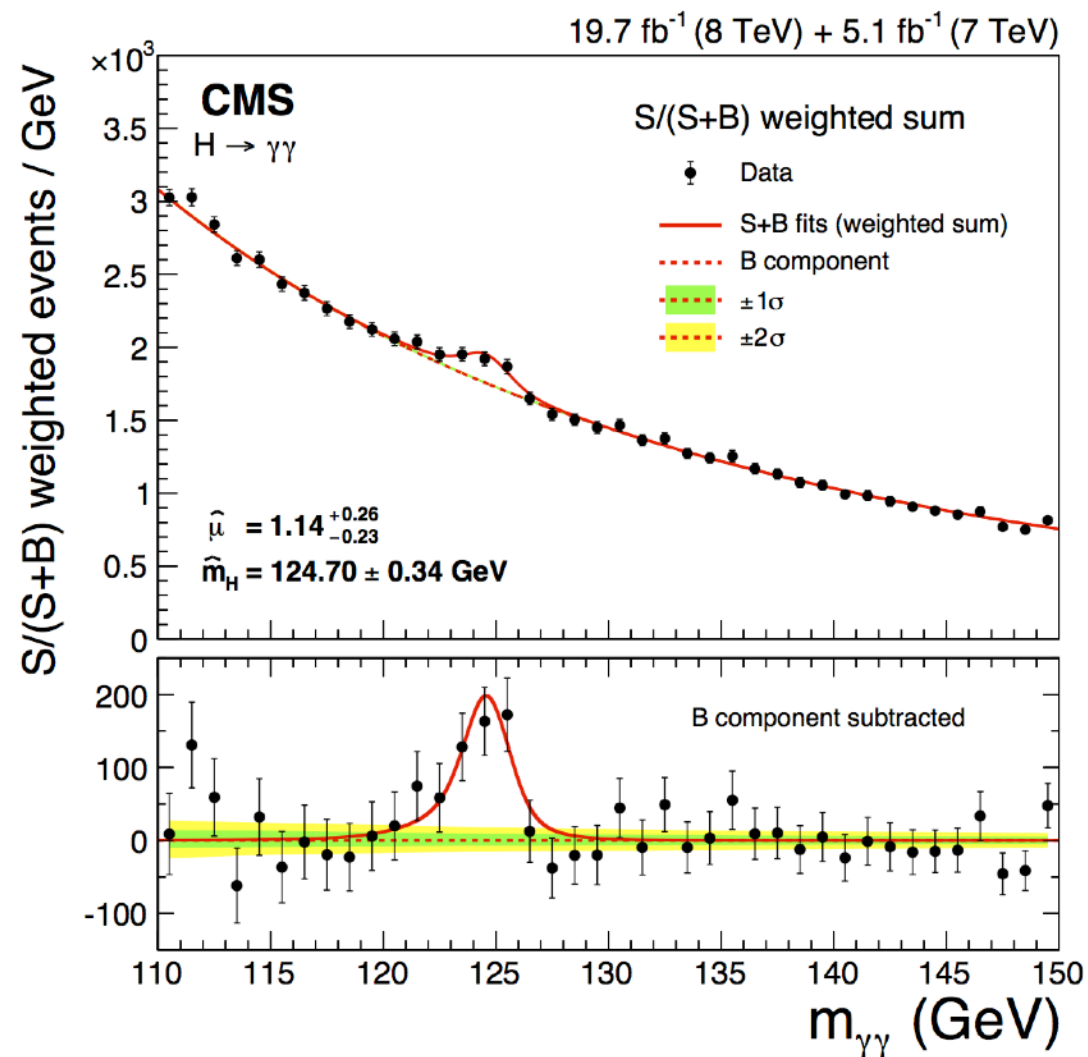


IRN Terascale
Higgs-session

IPPP, Durham, 6. August 2018

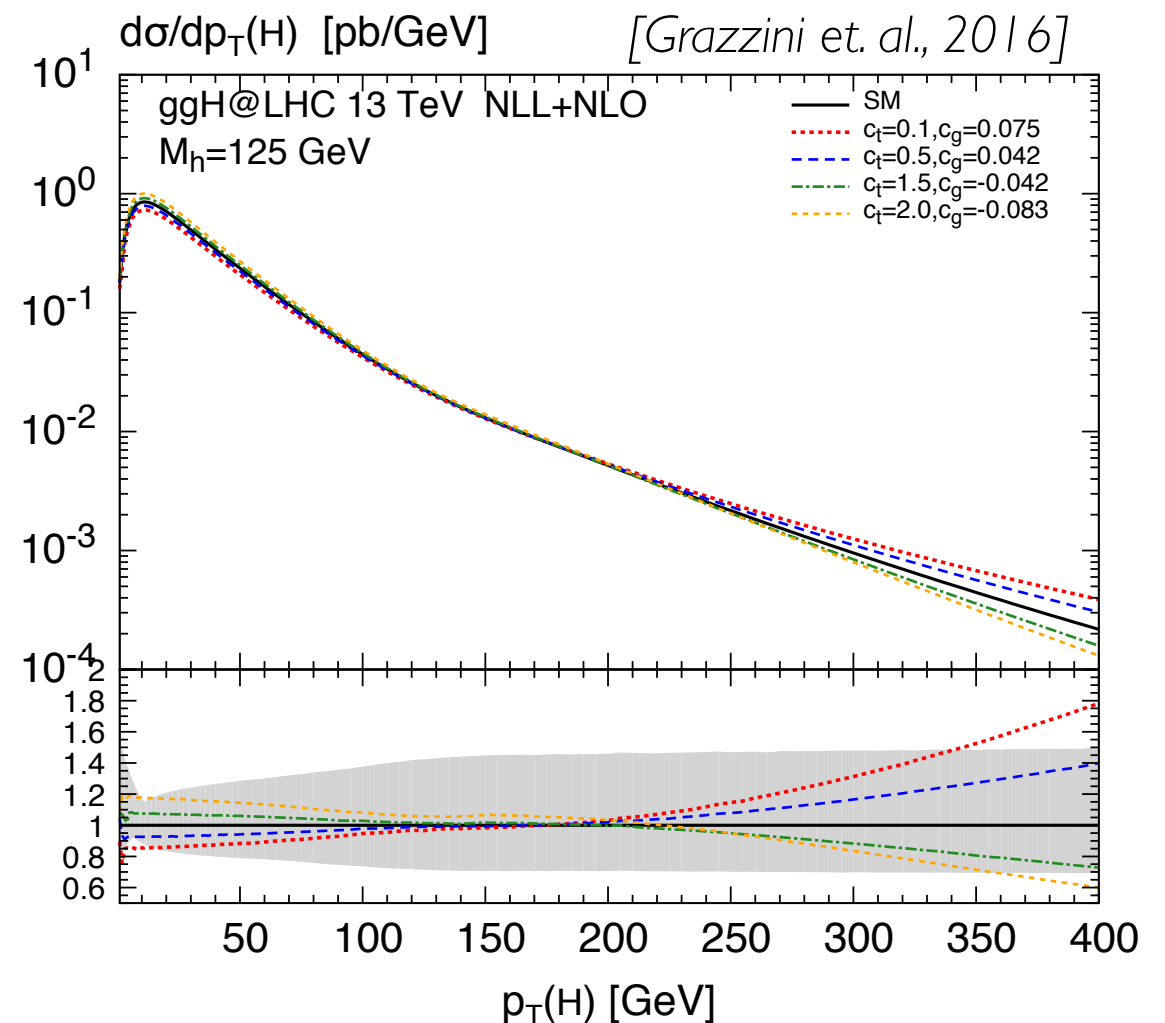
Finding the Higgs was “easy” ...

[CMS diphoton Higgs search, arXiv:1407.0558]



Bump hunting: little to no theoretical input needed.

...finding new physics might be very tough.



Look for BSM effects in small deviations from SM predictions:

- Higgs processes natural place to look at
- good control on theory necessary!

“good control” ?

Λ_{NP}

direct
bounds
 $\sim \text{TeV}$

SM $\sim \text{v.e.v.}$

Imagine to have **new physics** at a
(heavish) scale Λ_{NP}

Typical modification to observable
w.r.t. standard model prediction:

$$\delta O \sim Q^2 / \Lambda_{\text{NP}}^2$$

To gain over direct bounds:

IN THE BULK:

$Q \sim M_{\text{H}} \rightarrow \text{few percent}$

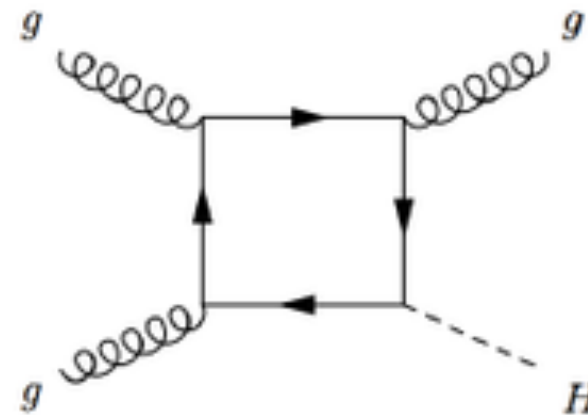
IN THE TAIL:

$Q \gtrsim 500 \text{ GeV} \rightarrow$
 $\sim 10\text{-}20\%$

[F. Caola, Moriond '18]

Outline

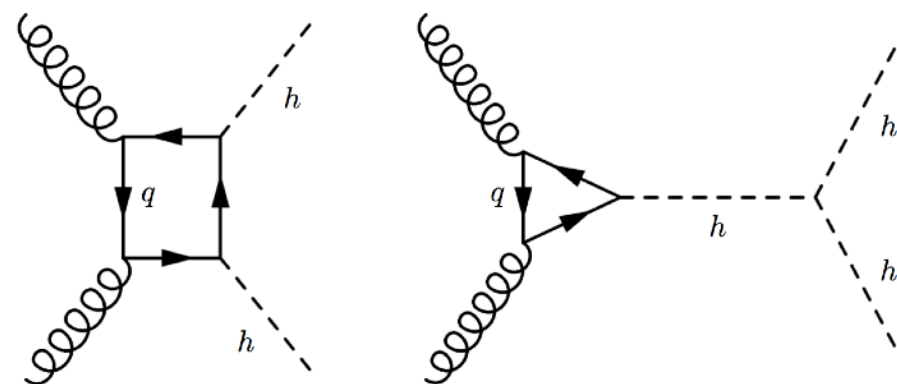
Higgs-pT



LO: Ellis et. al. '88; U. Baur et. al., '90

- $\sim 40\%$ remaining scale uncertainty
- $K^{\text{HEFT}} \sim 1.8$

Di-Higgs

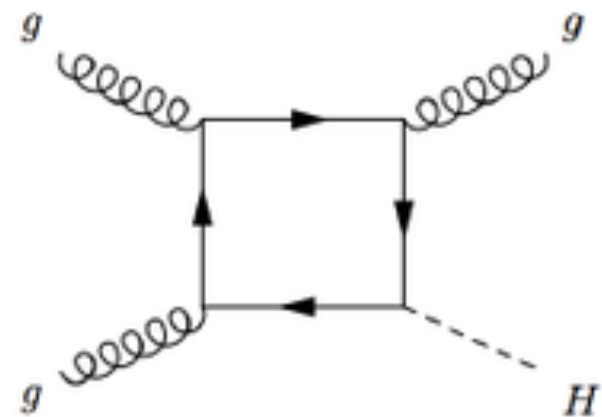


LO: Eboli et. al. '87; Glover, van der Bij '88

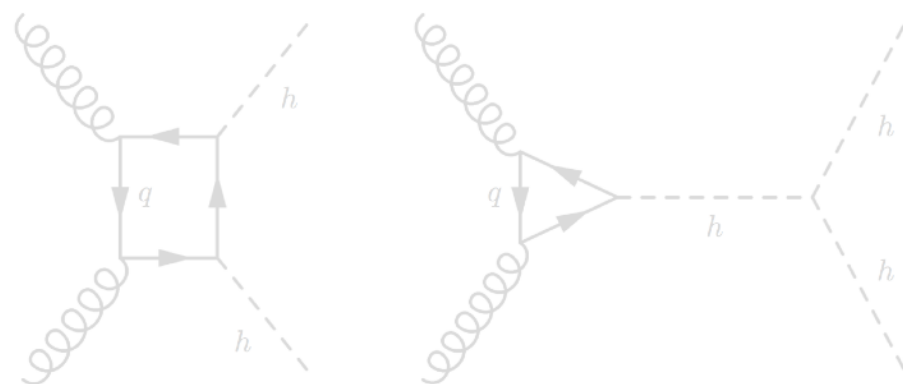
- $\sim 25\%$ remaining scale uncertainty
- $K^{\text{HEFT}} \sim 1.85$

Outline

Higgs-pT



Di-Higgs



Higgs-pT: two regimes

$$p_{\perp} \ll m_t$$

$$p_{\perp} > m_t$$

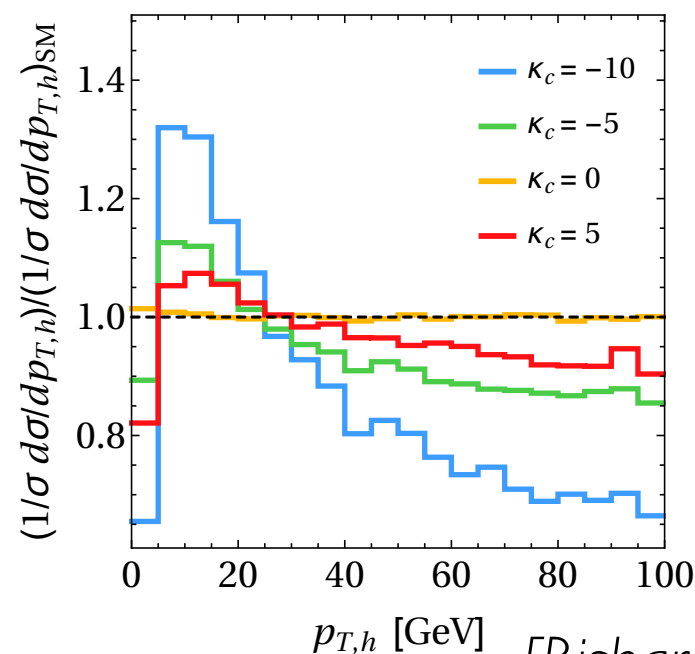
Possibility to constrain the charm-Yukawa coupling

$$d\sigma/dp_{\perp} \propto y_t^2 + y_t y_b + y_b^2 + y_t y_c + \dots$$

for $p_T \ll m_H$ $\sim 10\%$ $\sim 1\%$ $\ll 1\%$

$$A_{gg \rightarrow Hg}^Q \sim m_Q^2/m_H^2 \log^2(p_{\perp}^2/m_Q^2)$$

➔ Sudakov-like logarithmic enhancement of light-quark contribution at small pT



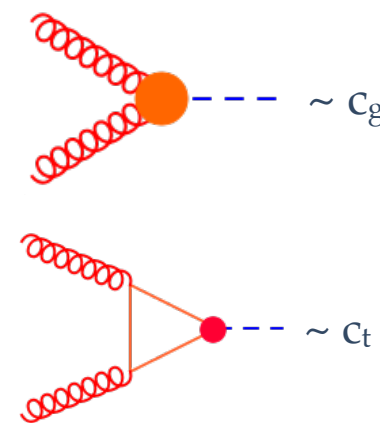
[Bishara, Haisch,
Monni, Re; '16]

Sensitive probe of New Physics

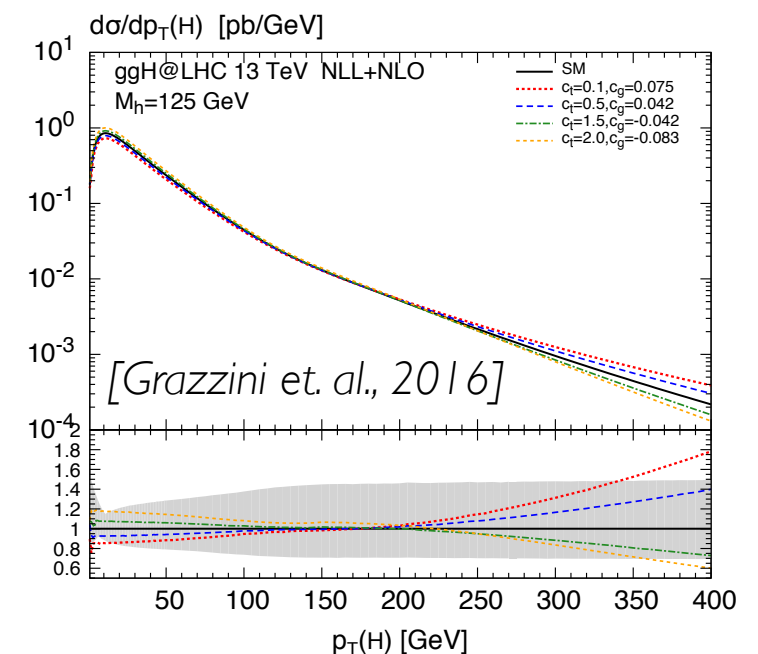
➔ In particular: disentangle c_g vs. c_t :

$$\frac{d\sigma_H}{dp_{\perp}^2} \sim \frac{\sigma_0}{p_{\perp}^2} \begin{cases} (c_g + c_t)^2, & p_{\perp}^2 < 4m_t^2, \\ \left(c_g + c_t \frac{4m_t^2}{p_{\perp}^2}\right)^2, & p_{\perp}^2 > 4m_t^2. \end{cases}$$

Note: inclusive measurements only allow to constrain $(c_g + c_t)^2$

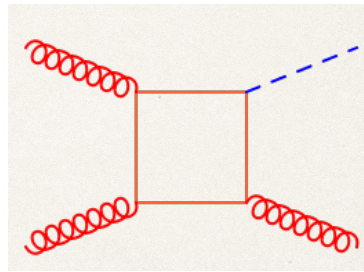


SM: $c_g=0$, $c_t=1$



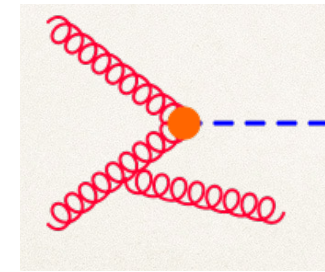
Higgs-pT: higher-order corrections

full theory: loop-induced



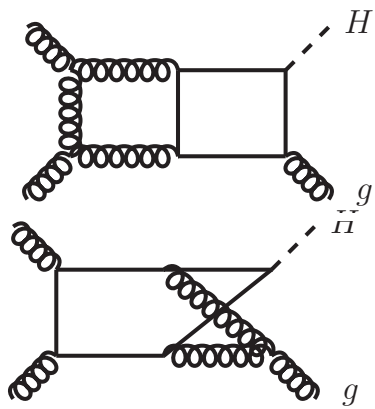
integrate-out

heavy quarks



HEFT: tree-level at LO

Bottleneck:
massive two-loop amplitudes



NLO

NNLO

[Chen et.al.; '14+'16
Boughezal et. al.; '15,
Caola et.al.; '15]

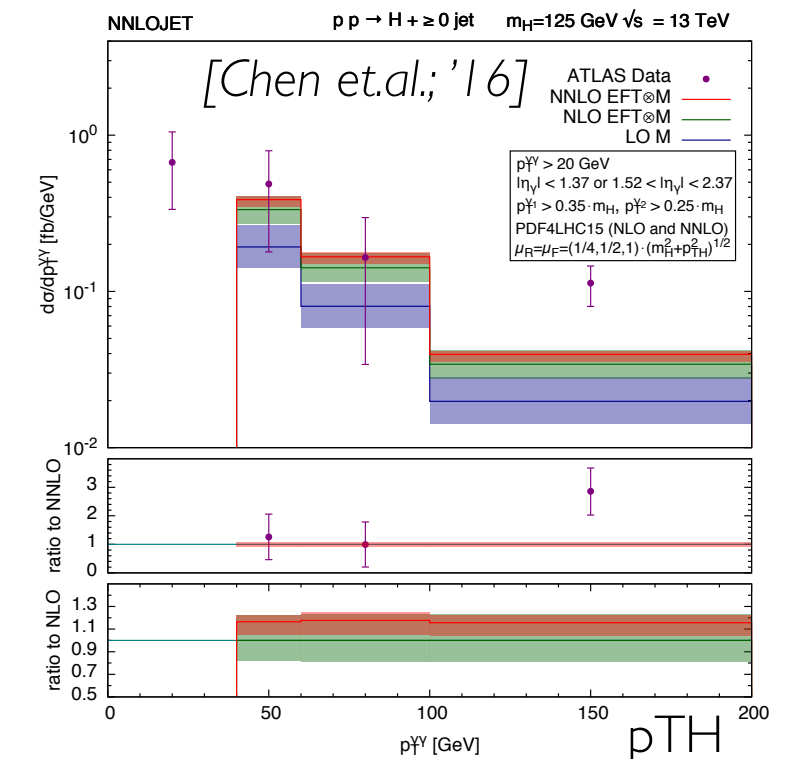
Ansätze:

- analytical: very hard, planar MI known [Bonciani et. al., '16]
- numerical: very CPU/GPU intensive [Jones et. al., '18]
- expansions: has to be performed carefully, very versatile

[Melnikov et. al., '16+'17]

Idea: QCD corrections factorize
→ apply K-factors from HEFT to lower
order predictions in full theory → check!!

Bottleneck: IR subtraction



perturb. uncertainties in HEFT
under very good control:

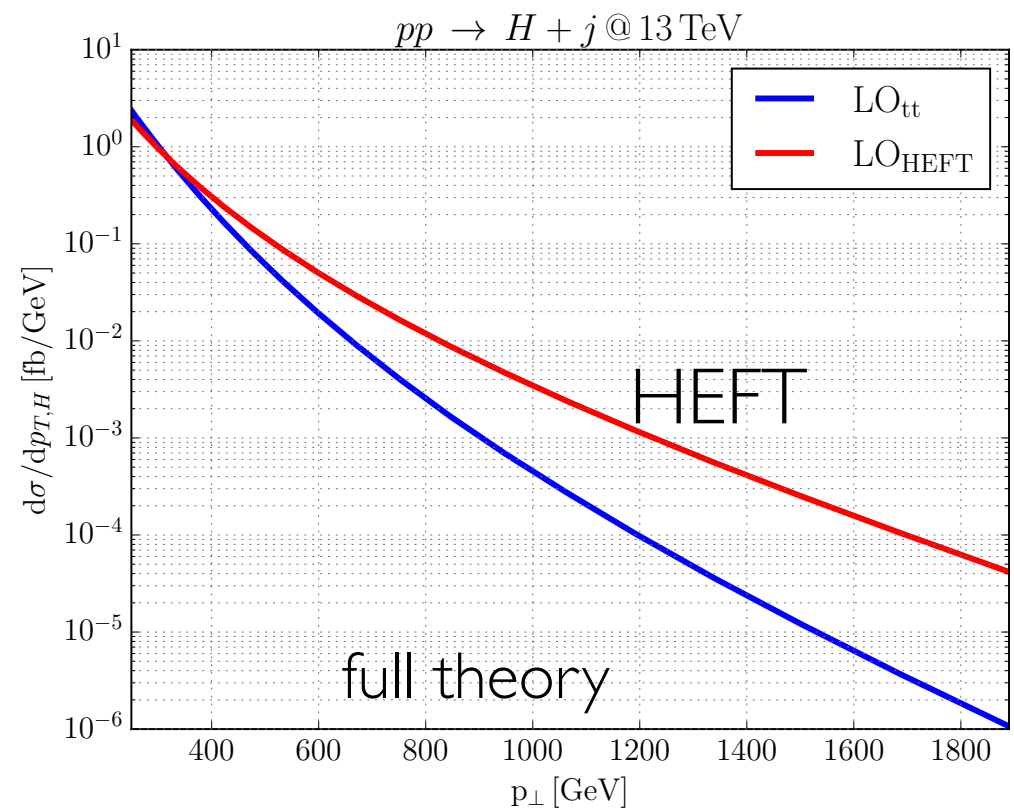
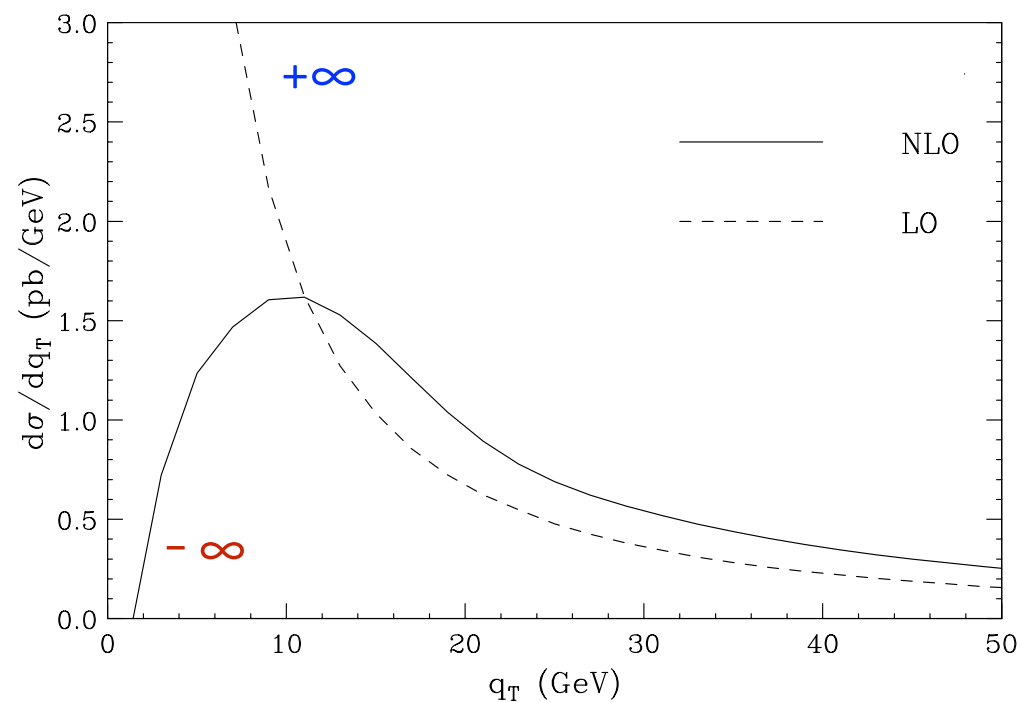
- ▶ ~10% scale variation
- ▶ stable shapes

NNLO+NXLL

Higgs-pT: two regimes

$p_{\perp} \ll m_t$

$p_{\perp} > m_t$



$$\sim p_T^{-2}$$

$$\sim p_T^{-4}$$

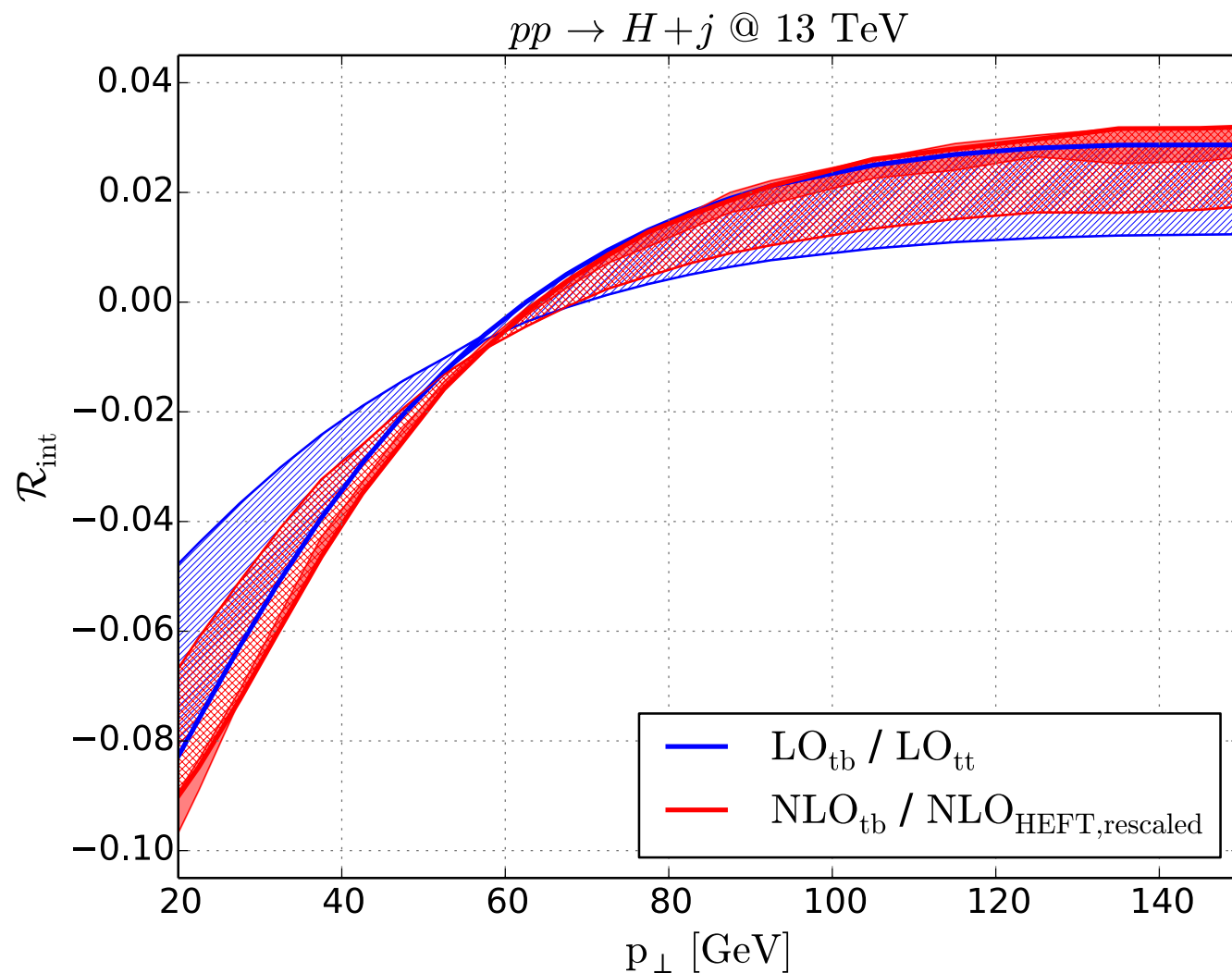
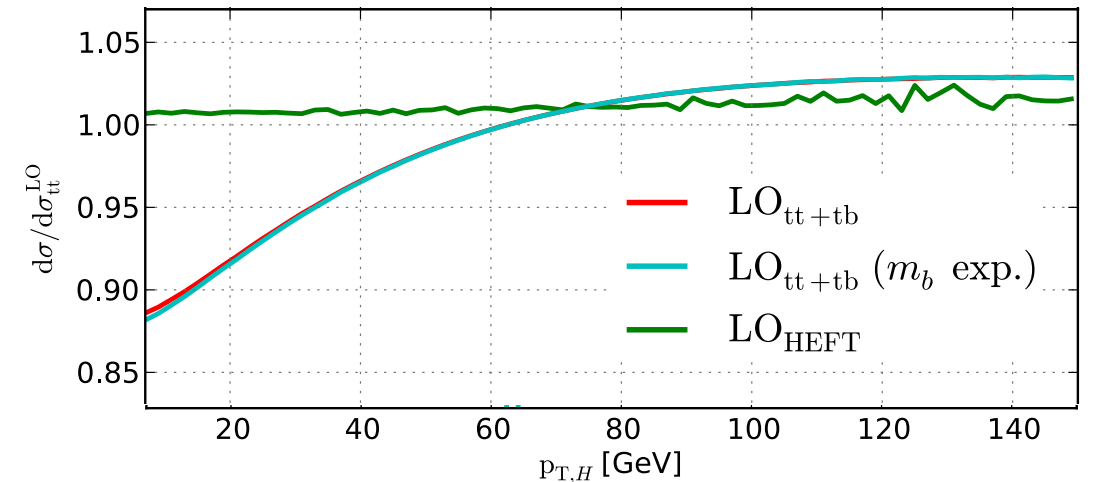
- Fixed-order breaks down at low pT

- point-like ggH (HEFT) and full theory have very different high energy behaviour.

$p_T^H \ll m_H$: bottom mass effects at NLO

[JML, Melnikov, Tancredi, Wever '17]

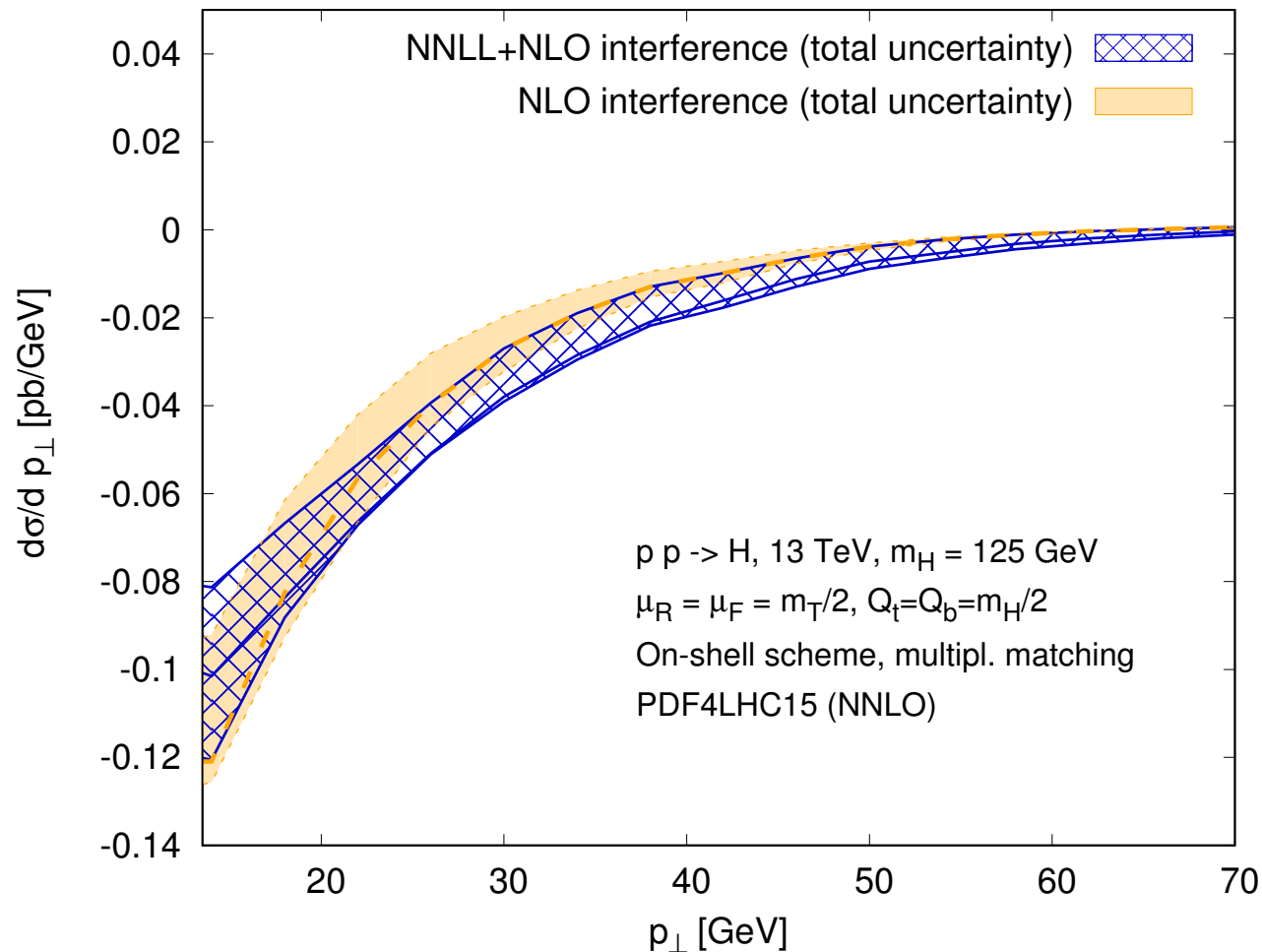
- expansion of the two-loop integrals in (m_b^2/p_T^2)
[Melnikov, Tancredi, Wever; '16+'17]
- valid at %-level down to $p_T \sim 10$ GeV
- real radiation treated exact with **OpenLoops**



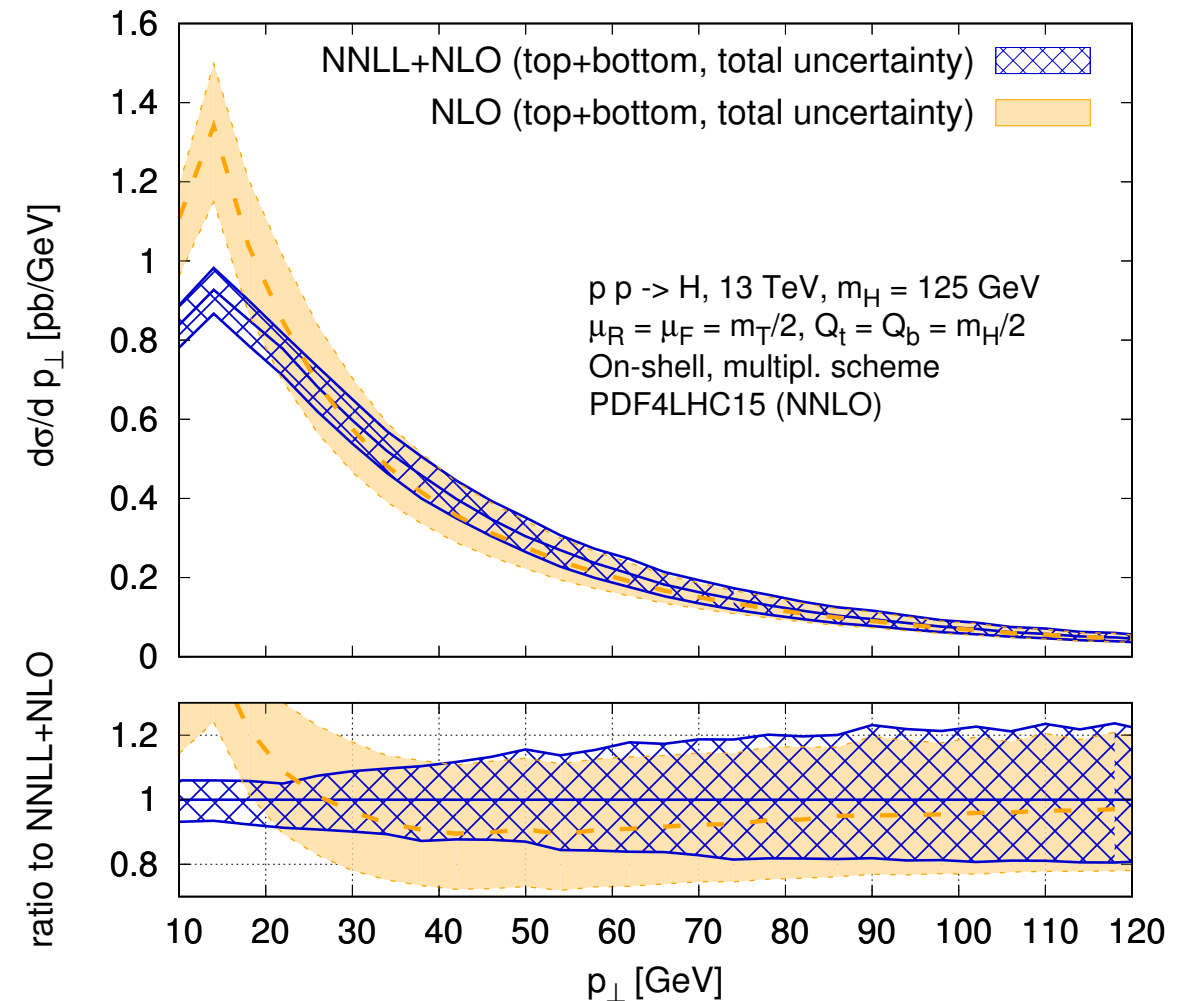
- -(5-10)% for $p_T=20-40$ GeV at LO and NLO
- Despite (large) corrections, the interference shape stable under QCD corrections
→ solid observable
- large m_b -renormalisation scheme dependence tamed at NLO

$p_{\perp} \ll m_t$: bottom mass effects at NLO+NNLL

[Caola, JML, Melnikov, Monni, Tancredi, Wever '17]



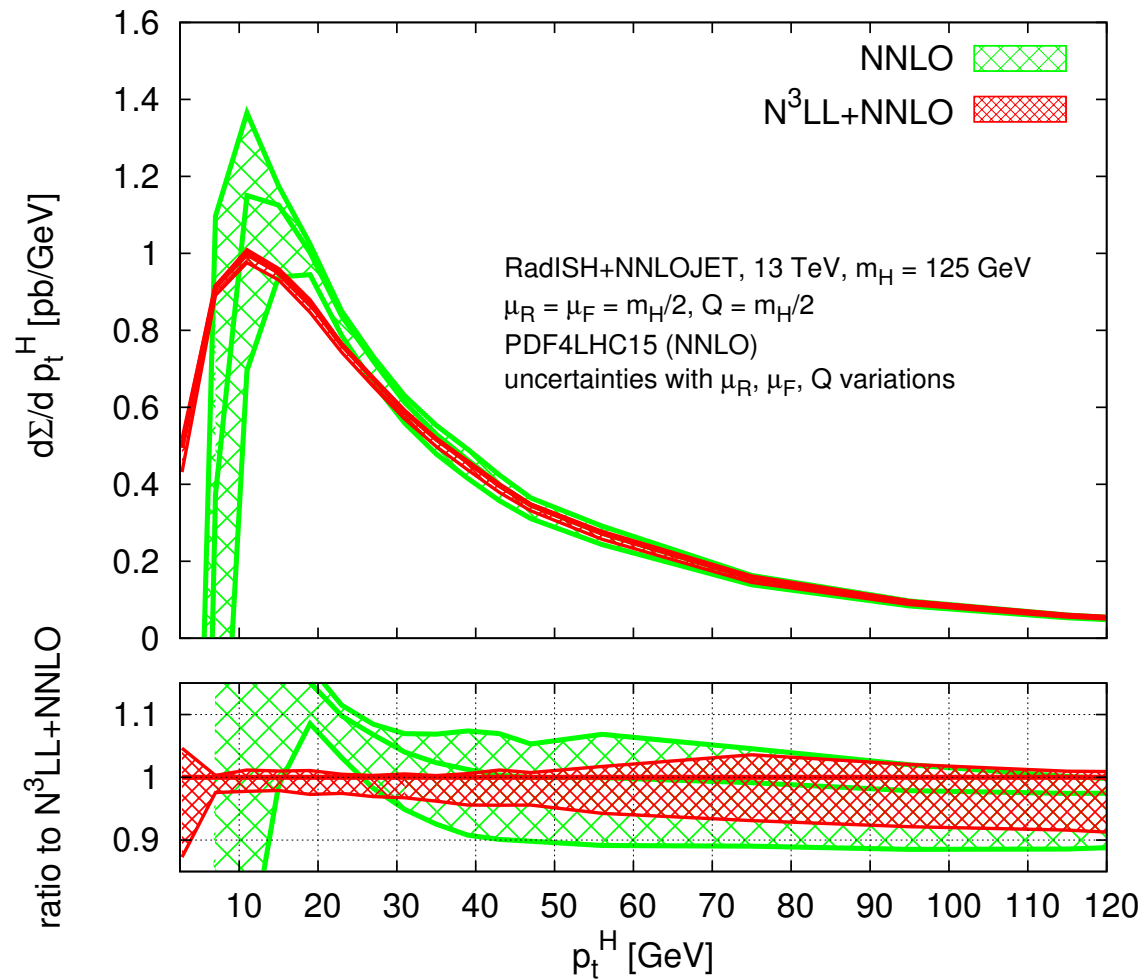
- remaining uncertainty of 20% on interference
- dominated by scheme ambiguity
- subleading uncertainties:
 - $\mu_R/\mu_F \sim 15\%$
 - $Q_b = m_H/2$ vs. $Q_b = 2m_b \sim 15\%$
- translates into ~ 1 -2% error on total spectrum



- uncertainties at the level of 5-20%
- Resummation effects relevant for $p_T \lesssim 40$ GeV
- further improvement when combined with NNLO for yt^2

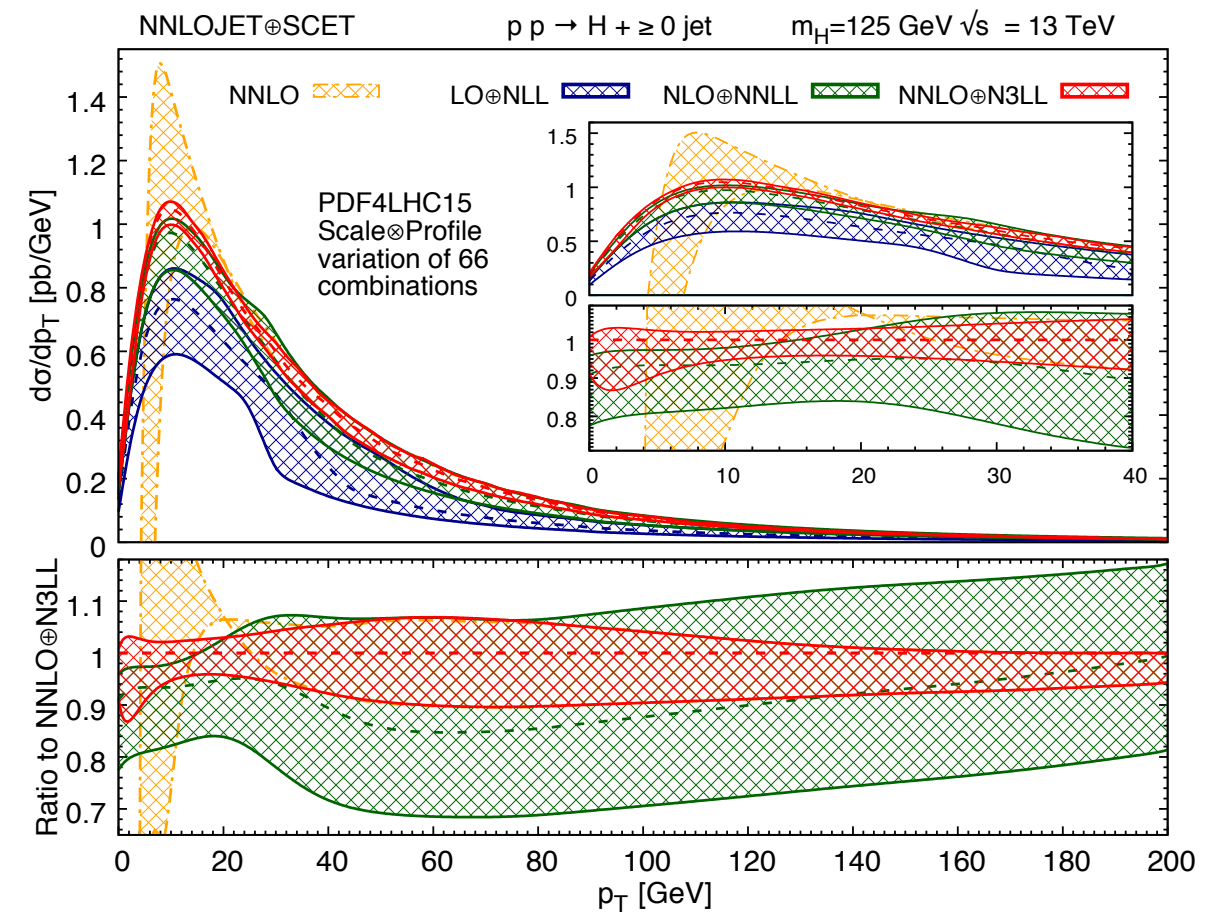
$p_{\perp} \ll m_t$: NNLO+N3LL in HEFT

[Bizon, Monni, Re, Rottoli, Torrielli+NNLOJET '17,'18]



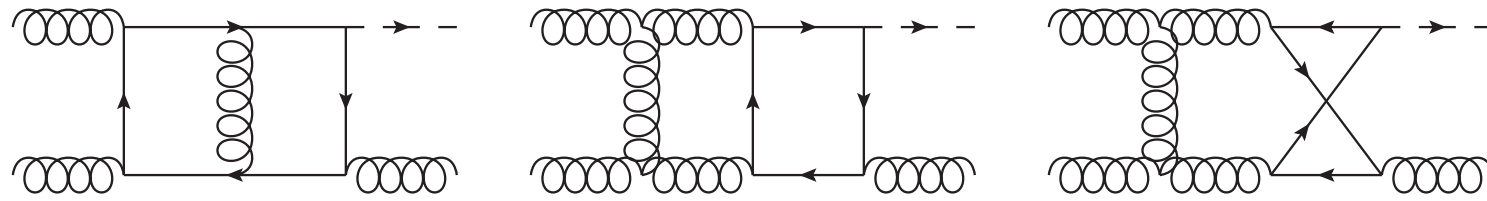
- Resummation performed in momentum space
- Multiplicative matching
- Results for fiducial phase-space available

[Li, Neill, Schulze, Stewart, Zhu+NNLOJET '18]



- Resummation performed in b-space within SCET
- Additive matching

$p_T^H \geq m_H$: virtual two-loop amplitudes

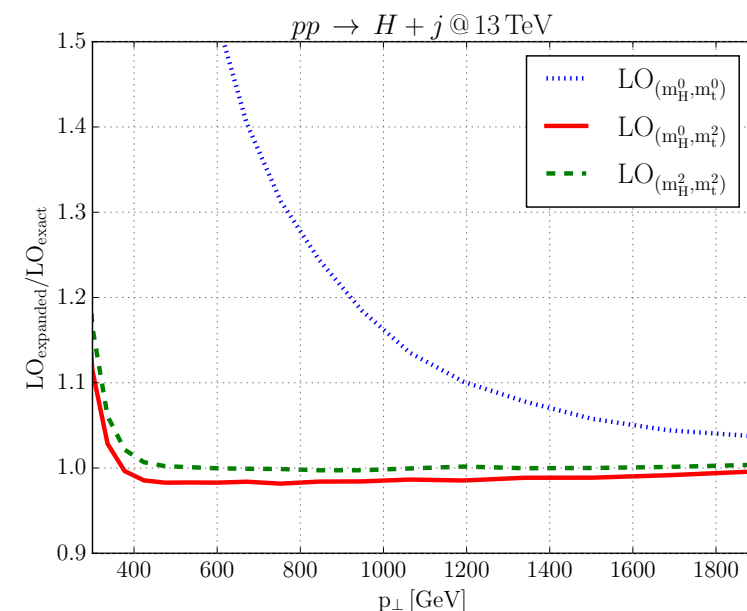


- A four-scale problem: three external (s, p_T, m_H) and one internal (m_t)
- 264 Feynman integrals, complicated reduction
- Only partial (planar topologies) analytic result with the full top mass dependence available (R. Bonciani, et al., 2016)
- Two recent approaches:

Numerical integration with SecDec
[S. P. Jones, et al., 2018]

- Sector decompose integrals with SecDec
- Numerically integrate sectors with Quasi-Monte-Carlo integration
- Accelerate with OpenCL on GPUs
- valid in all of the phase-space

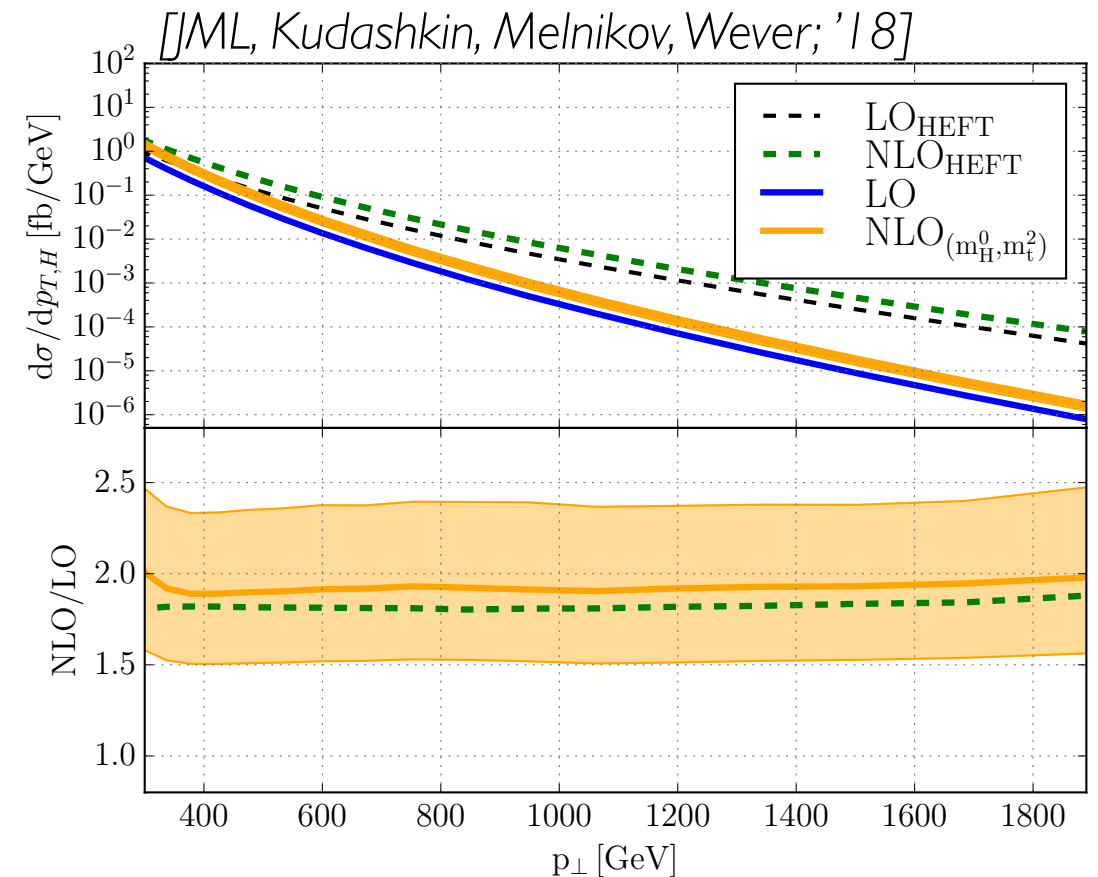
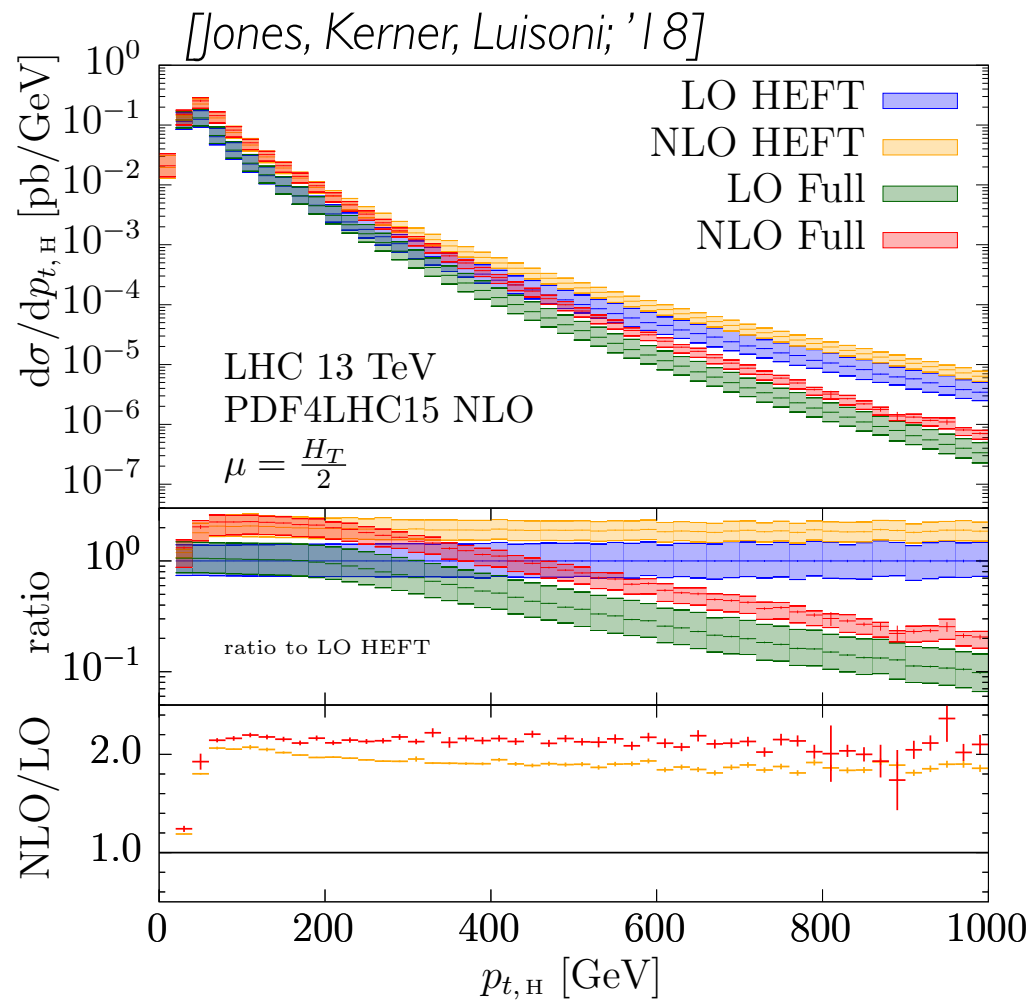
Expansion in $\eta = -\frac{m_H^2}{4m_T^2}$, $\kappa = -\frac{m_T^2}{s}$
at the level of differential equations
[Kudashkin, Melnikov, Wever; '17]



- valid at %-level for large p_T

$p_T^H \geq m_H$: top mass effects at NLO

- numerical integration of two-loop integrals based on **SecDec** [Borowka et.al.]
- valid in all of the phase-space
- expansion of the two-loop integrals up to $(m_t^2/p_T^2)^1$, $(m_H^2/p_T^2)^0$ at the level of the DE [Kudashkin, Melnikov, Wever; '17]

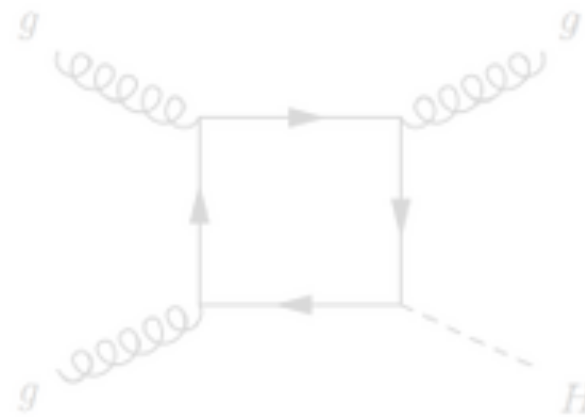


- NLO corrections very similar as in HEFT: $K \sim 2$ with remaining scale uncertainties: **~20-25%**
- hardly any shape dependence
- Outlook: combine with NNLO in HEFT

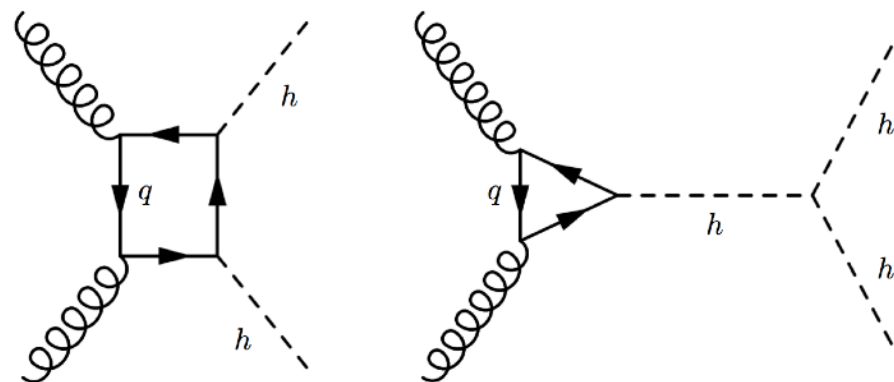
Control of the high-H- p_T tail at NLO opens the door for new physics searches in this regime!

Outline

Higgs-pT

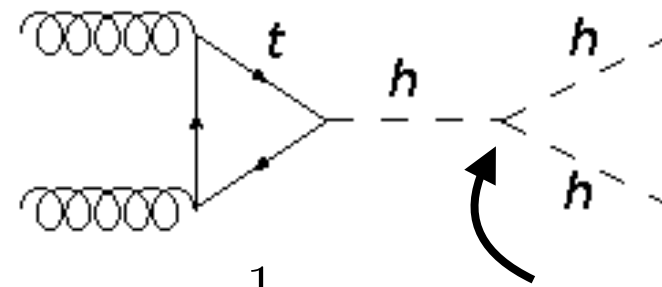
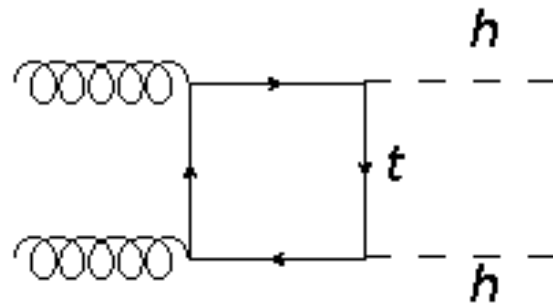


Di-Higgs

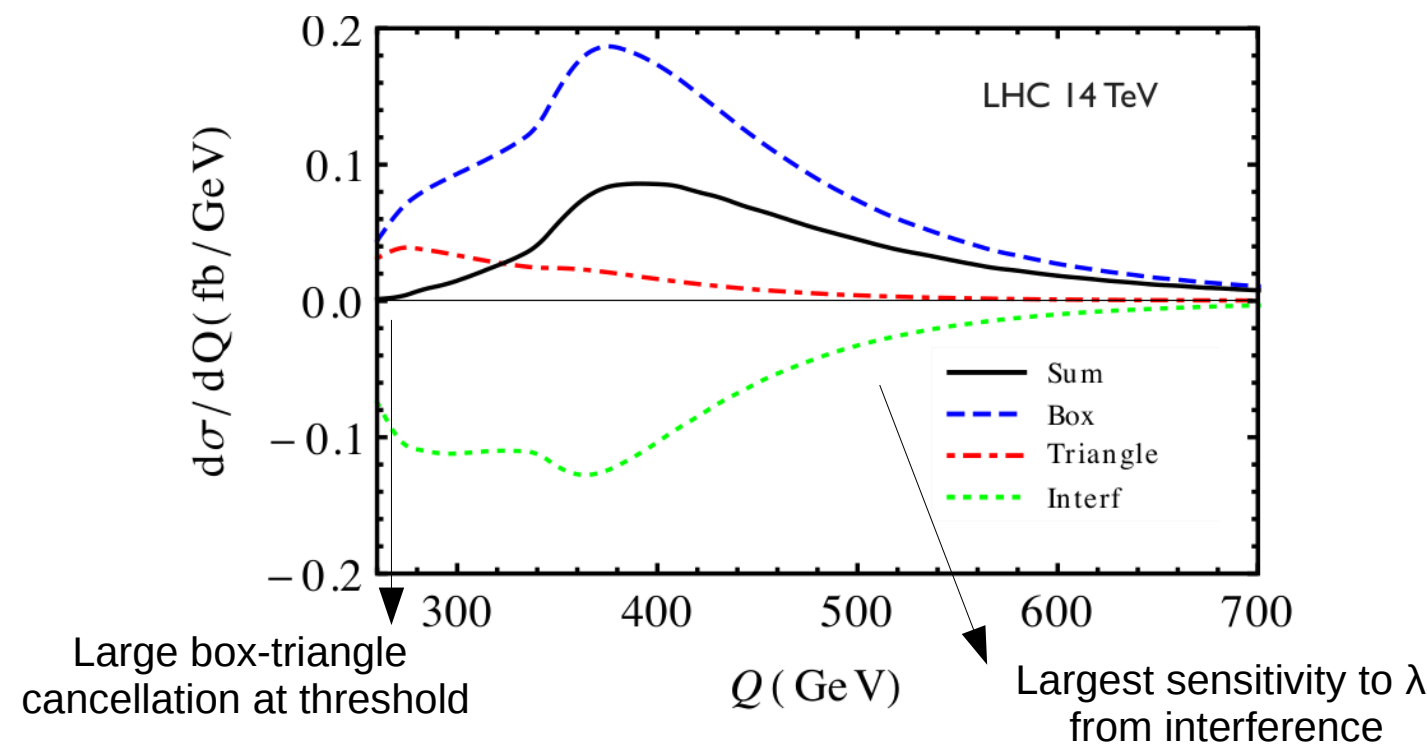
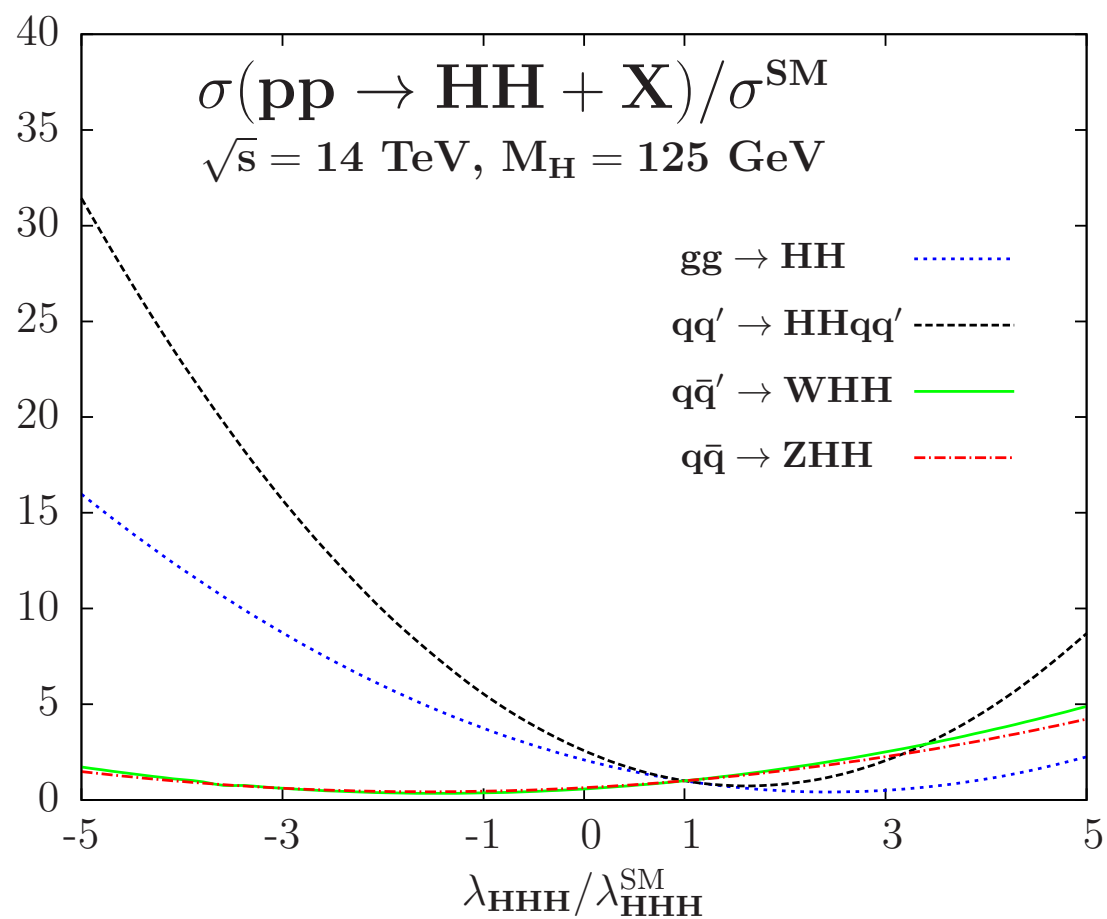


Di-Higgs

- $pp \rightarrow HH$ offers direct access to the trilinear Higgs coupling

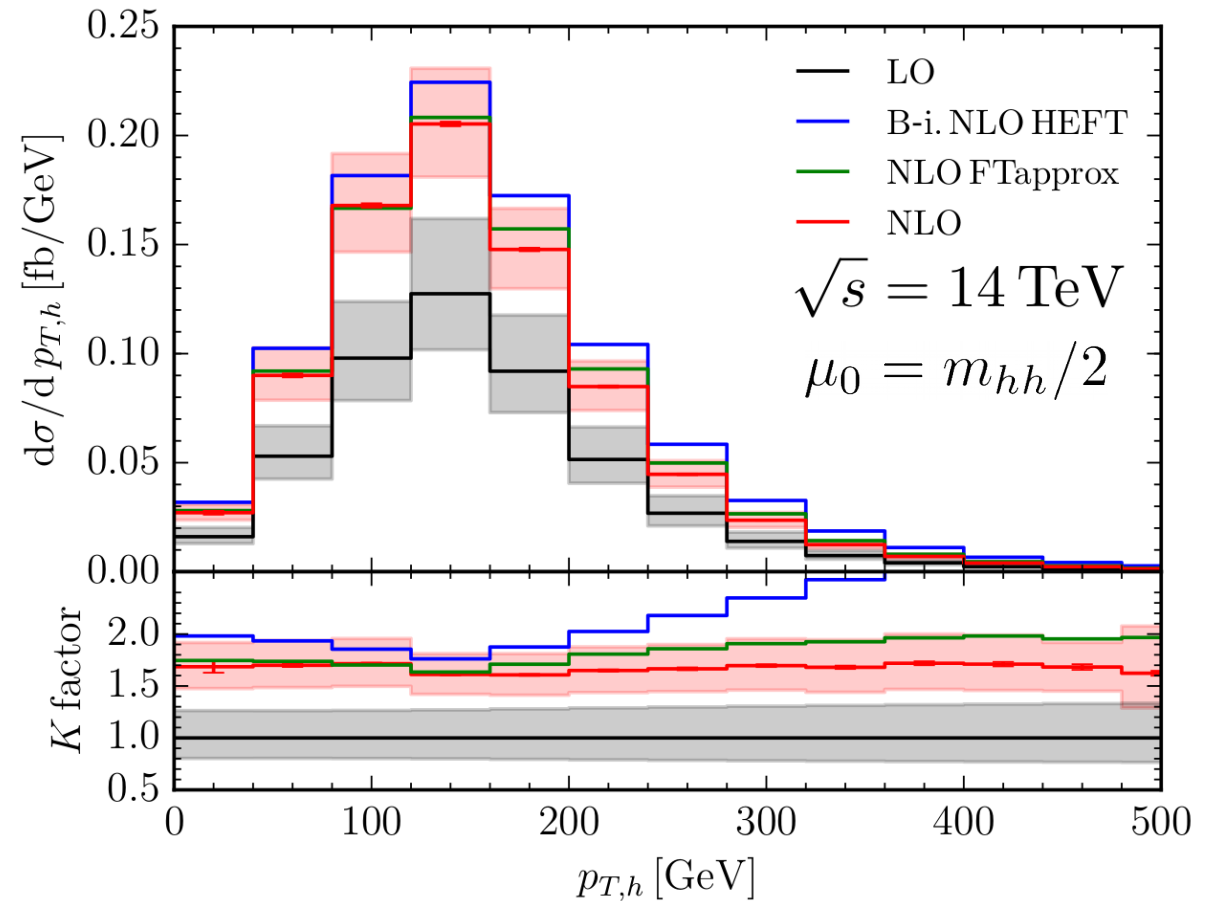
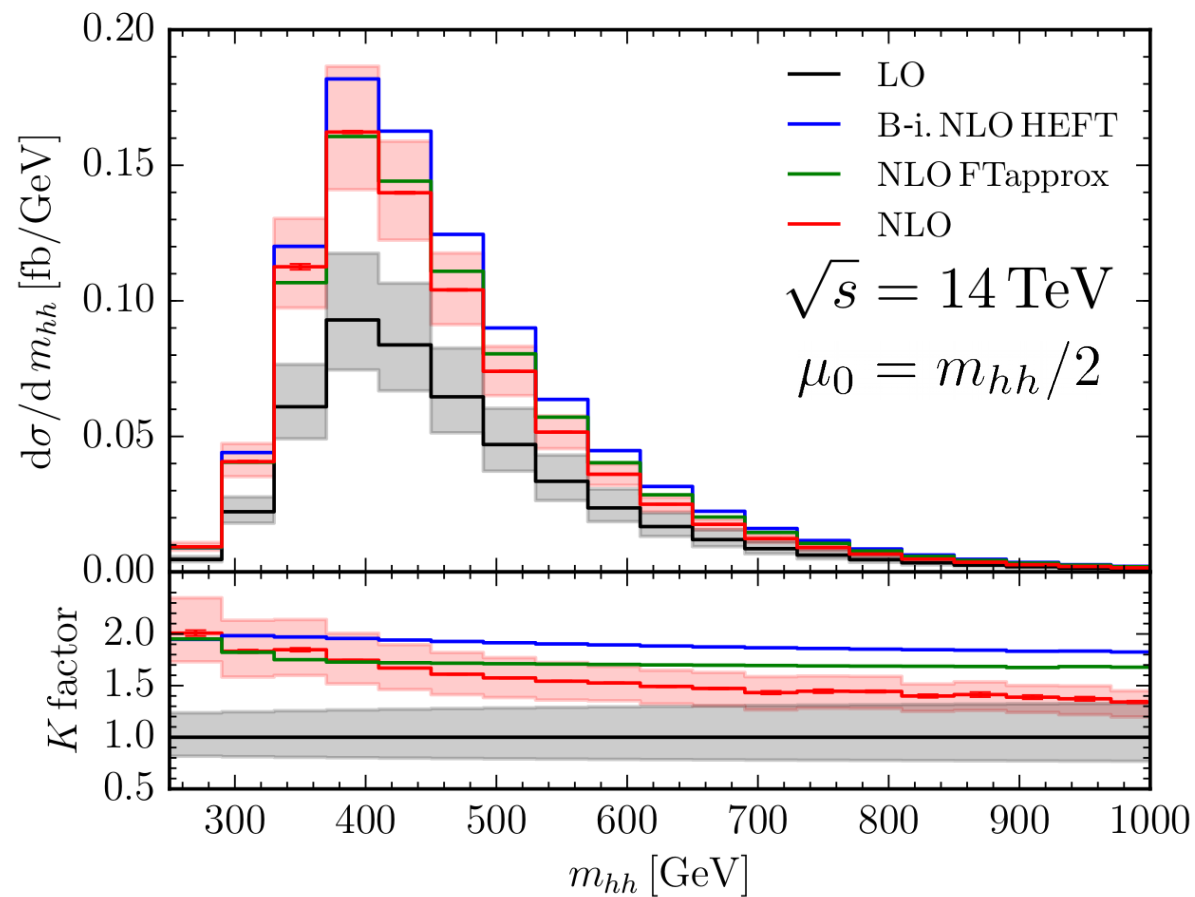


$$V(H) = \frac{1}{2}M_H^2 H^2 + \lambda v H^3 + \frac{1}{4}\lambda' H^4$$



Di-Higgs @ NLO

Again: Numerical integration with SecDec
[Borowka, et al., '16]

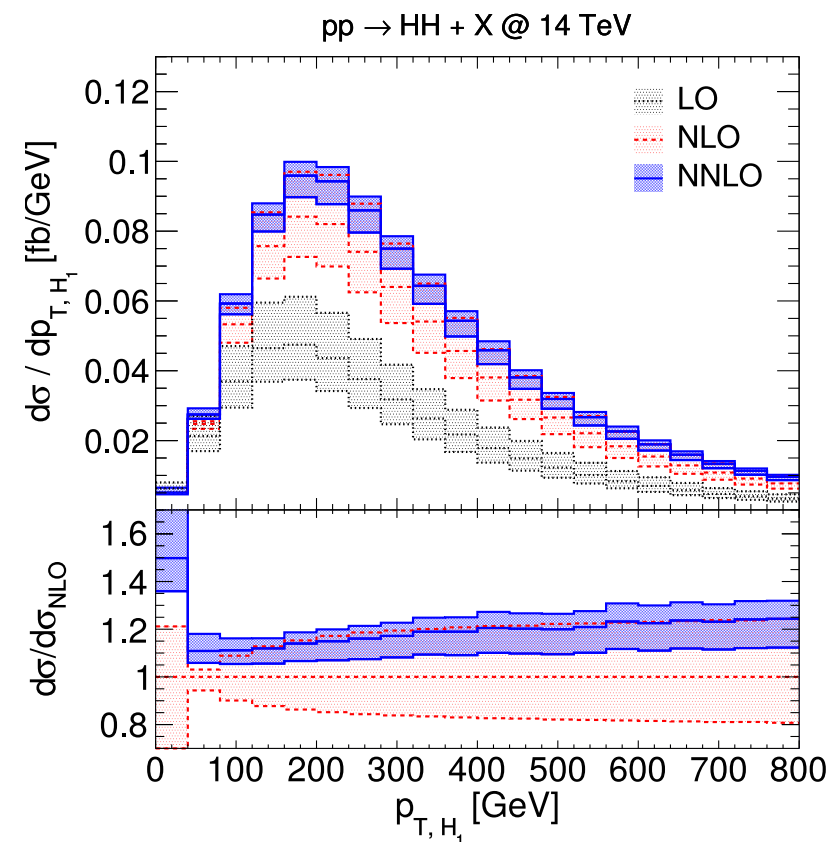
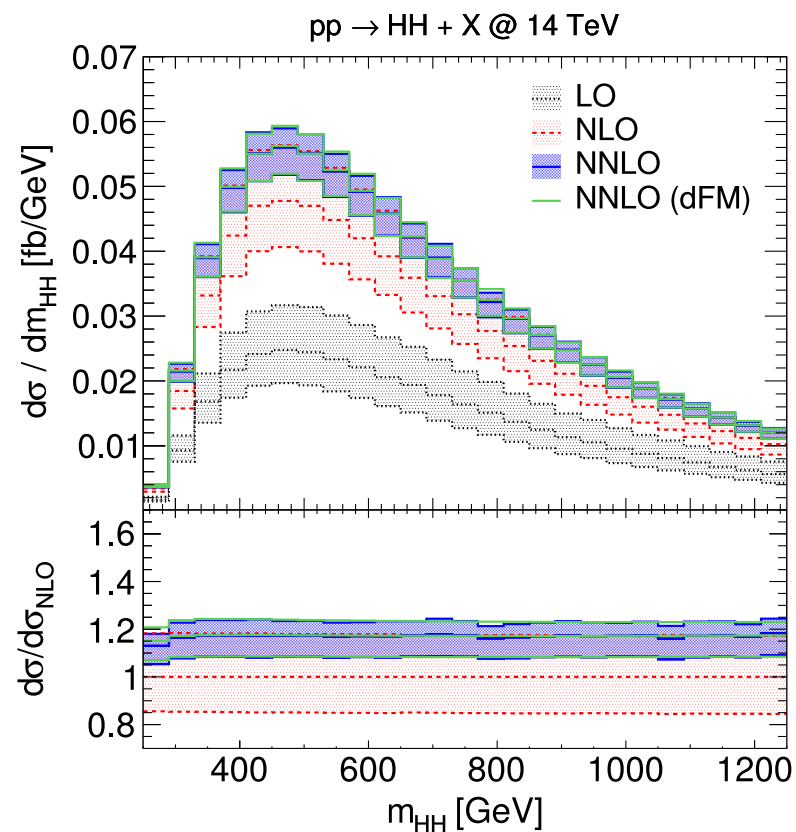


- Very large corrections NLO corrections: +66% for total XS
- Exact top-masses yield 16% smaller corrections than NLO Born-improved HEFT, and 4% smaller than NLO FTapprox (reals with exact dependence)
- In general non-trivial shape dependence of the corrections on the kinematics
- Remaining scale uncertainties: $\sim 13\%$ but no overlap with LO \rightarrow try to go beyond NLO

Di-Higgs @ NNLO in HEFT



[de Florian, JML, et.al. '16]



- ▶ NNLO corrections:
 - almost flat in m_{HH} , small shape in p_{TH1}
 - at the level of $\sim 20\%$
 - remaining scale uncertainties at the level of 10%
 - overlap with NLO uncertainty band
- ▶ How to combine with exact NLO?

Di-Higgs @ NNLO with top mass effects

Born-projected approximation:

- Reweight by LO ratio full/HEFT at the level of squared amplitudes
- Projection to Born kinematics needed for reweighting: non-unique

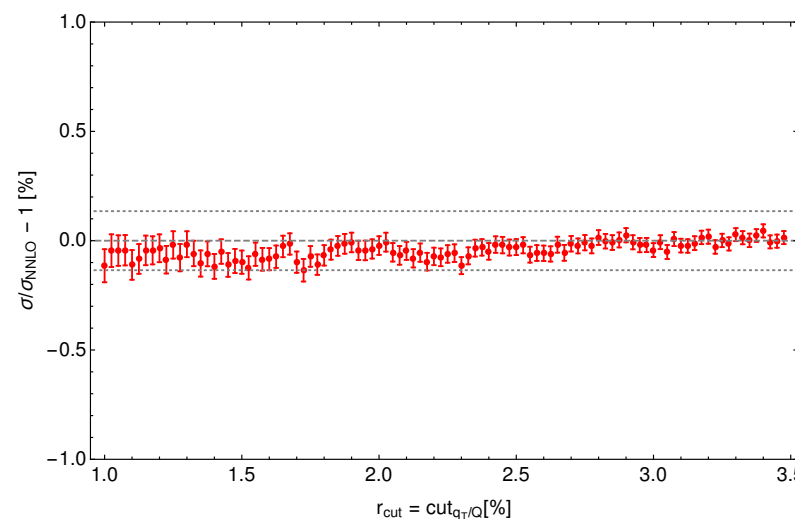
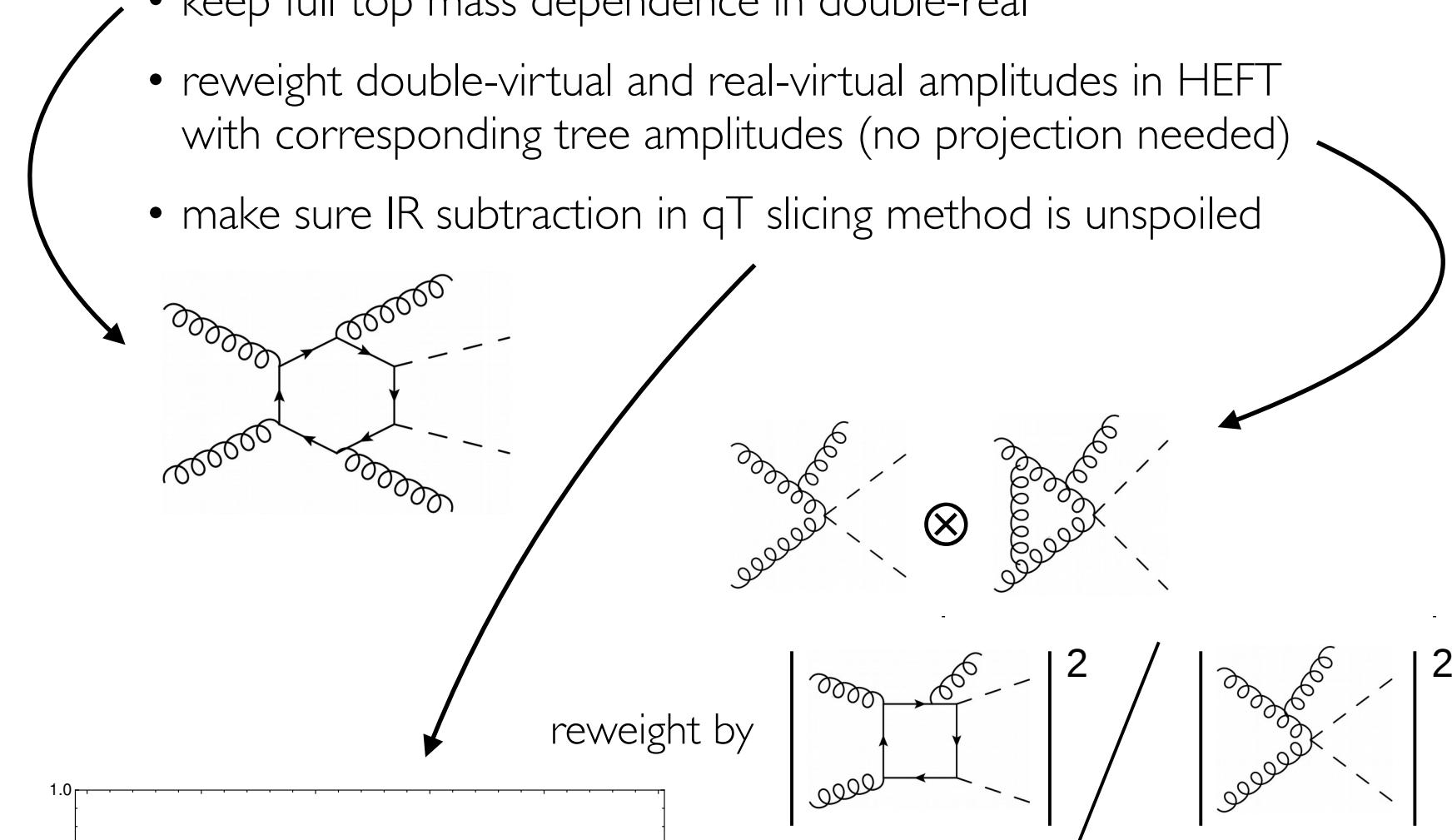
NLO-improved approximation:

- Observable level reweighting
- For each bin of each histogram do

$$d\sigma_{\text{NLO}} \times \left(\frac{d\sigma_{\text{NNLO}}}{d\sigma_{\text{NLO}}} \right)_{\text{HEFT}}$$

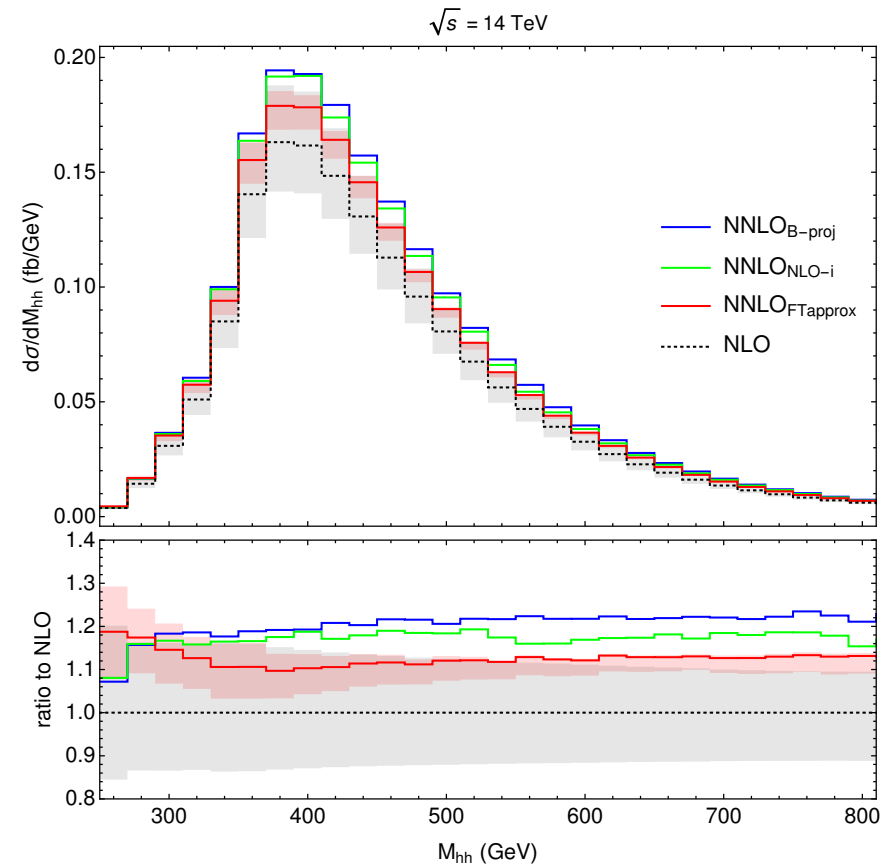
NNLO_{FTapprox}

- only approximate NNLO coefficient
- keep full top mass dependence in double-real
- reweight double-virtual and real-virtual amplitudes in HEFT with corresponding tree amplitudes (no projection needed)
- make sure IR subtraction in qT slicing method is unspoiled

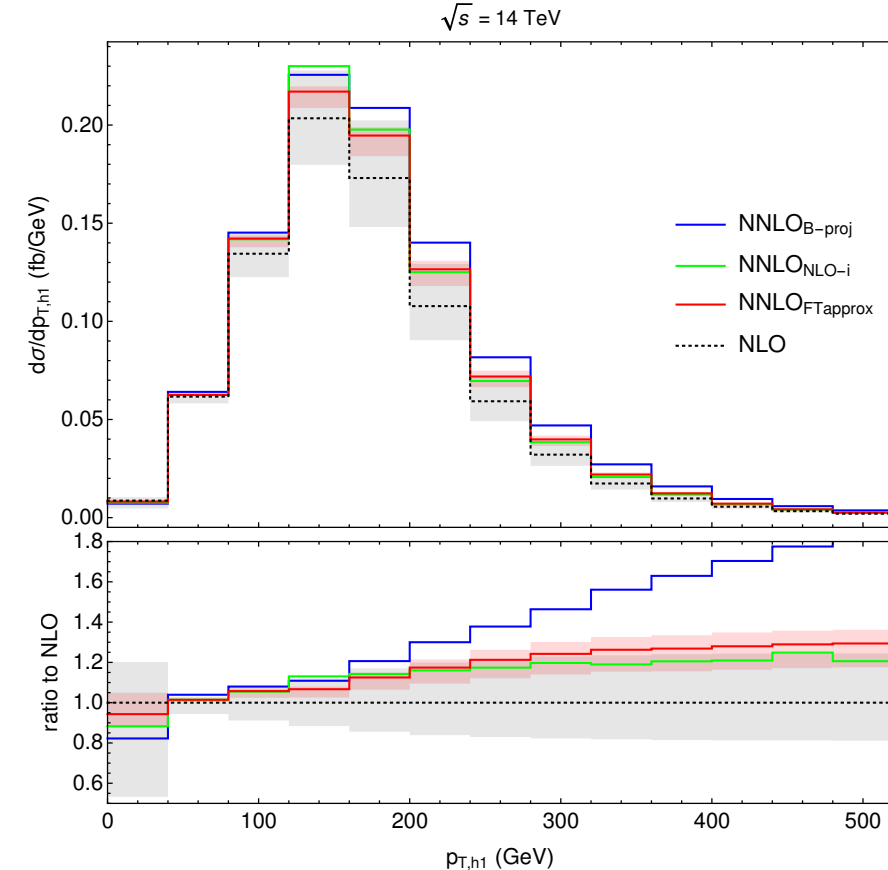


Numerical complexity in the loop-induced double-real similar to requirements for real-virtual in tree-induced $2 \rightarrow 3$ NNLO!

Di-Higgs @ NNLO with top mass effects



- FTapprox presents larger corrections at threshold, minimum corrections at $M_{hh} \sim 400\text{GeV}$, slow increase towards the tail
- Scale uncertainties in FTapprox are substantially reduced
- Overlap with the NLO band
- up to $\sim 10\%$ smaller corrections in FTapprox compared to the other approximations



- wrong scaling at large energies for $\text{NNLO}_{B\text{-proj}}$
- overlap of NLO and $\text{NNLO}_{\text{FTapprox}}$ bands
- $\text{NNLO}_{\text{FTapprox}} \sim \text{NNLO}_{\text{NLO-i}}$

\sqrt{s}	14 TeV
NLO [fb]	$32.88^{+13.5\%}_{-12.5\%}$
$\text{NLO}_{\text{FTapprox}}$ [fb]	$34.25^{+14.7\%}_{-13.2\%}$
$\text{NNLO}_{\text{NLO-i}}$ [fb]	$38.66^{+5.3\%}_{-7.7\%}$
$\text{NNLO}_{B\text{-proj}}$ [fb]	$39.58^{+1.4\%}_{-4.7\%}$
$\text{NNLO}_{\text{FTapprox}}$ [fb]	$36.69^{+2.1\%}_{-4.9\%}$
M_t unc. $\text{NNLO}_{\text{FTapprox}}$	$\pm 2.7\%$
$\text{NNLO}_{\text{FTapprox}}/\text{NLO}$	1.116

→ uncertainties on total XS largely reduced!

Conclusions

- Without a clear sign of new-physics it is crucial to perform detailed theory vs. experiment comparisons, to look for possible deviations.
 - ➡ Higgs is an obvious place to look at!
- H-pT @ NLO:
 - ▶ remaining uncertainties: $\sim 20\%$ (NNLO in HEFT: 10%)
 - ▶ corrections very similar as in HEFT
 - ▶ Outlook: NNLO with mass effects
- Di-Higgs @ NNLO_{FTapprox}
 - ▶ remaining uncertainties for inclusive cross section: $\sim 3.5\%$
 - ▶ somewhat larger in distributions due to mass effects
 - ▶ at this level: top-mass scheme ambiguity relevant