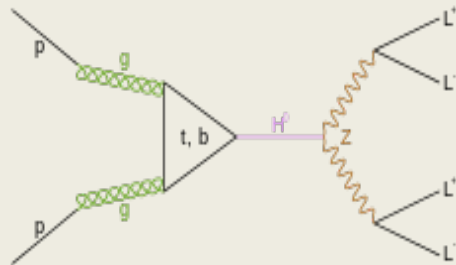
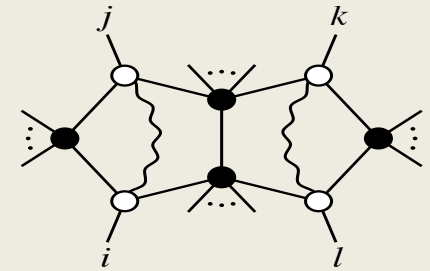
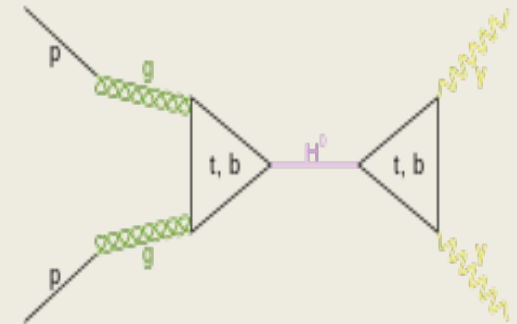


# Cluster Algebras, Landau Singularities and Scattering Amplitudes



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# Planar $N=4$ Yang-Mills

- $N=4$  Yang-Mills is a quantum field theory which describes “real world” particle interactions. Many techniques we develop for YM have been applied to QCD and Standard Model.
- $N=4$  Yang-Mills is a string theory (via AdS/CFT), so learning about  $N=4$  Yang-Mills teaches us about quantum gravity.
- $N=4$  Yang-Mills is solvable (integrable) which is a non-trivial fact for four-dimensional quantum field theory.

# N=4 Yang-Mills Amplitudes

The goal is to explore the hidden structures of amplitudes and to exploit these structures as much as possible to make previously impossible computations possible.

- Twistor string description
- Dual conformal and Yangian symmetries
- Wilson loops/amplitudes/correlation functions triality
- Color/kinematic duality
- Amplituhedron, positivity, on-shell diagrams
- Simplifying limits: fishnet, multi-Regge

Bern, Dixon, Kosower; Witten; Drummond, Henn, Korchemsky, Sokachev, Brandhuber, Heslop, Spence, Travaglini, Arkani-Hamed, Cheung, Cachazo, Kaplan, Trnka, Bourjaily, Bern, Carrasco, Johansson, Caron-Huot, He, Beisert, Eden, Staudacher, Mason, Skinner, Alday, Maldacena, Basso, Sever, Viera, Roiban, Spradlin, many others

# Plan

- Introduction
- N=4 Yang-Mills Amplitudes Bootstrap Program:  
Input = Symbol Alphabet

How Symbol Alphabet be determined?

- Symbol Alphabet and Landau Singularities
- Symbol Alphabet and Cluster Algebra
- Conclusions



# N=4 Planar Yang-Mills Amplitudes

L-loop  $n$ -points  $N^k$ MHV

- 0- and 1-loops: all  $n$ ,  $k$
- 2-loop: all  $n$ -point MHV, 6, 7 NMHV
- 3, 4-loops: 6, 7-point MHV & NMHV
- 5-loops: 6-point MHV & NMHV
- 6-loops: 6-point MHV
- L-loops: 4, 5-point MHV

[Bern, Dixon, Kosower ][Caron-Huot, Dixon, Drummond, Duhr, Foster, Gurdogan, Harrington, Henn, McLeod, Papathanasiou, Spradlin, von Hippel 2015-present]

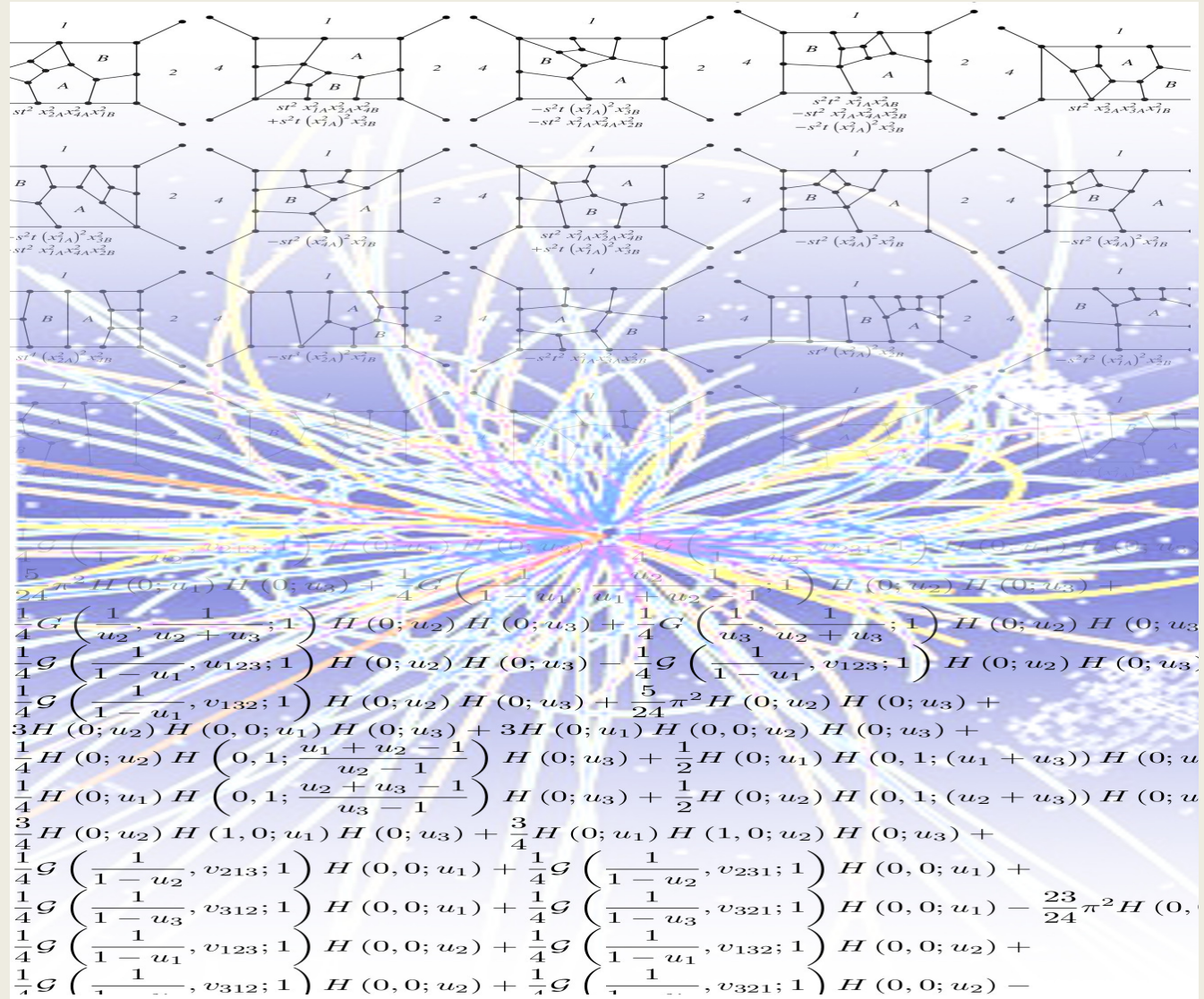
# Quantum Field Theory Textbooks:

$$\text{Amplitude} = \int d(\text{loop momenta}) (\text{Integrand})$$

Integrand

hard

Amplitude



# N=4 Yang-Mills Amplitudes

Integrands are known for all  $L$ ,  $k$  and  $n$ , but as we saw on the previous slide the resulting integral is only known for a few cases....

- We need new tools to make these calculations possible.
- Having more data is crucial for identifying hidden structures of these amplitudes.
- Modern approaches to computing amplitudes avoid knowledge of integrands completely....

Amplitudes Bootstrap

# Amplitude Bootstrap : Old vs New



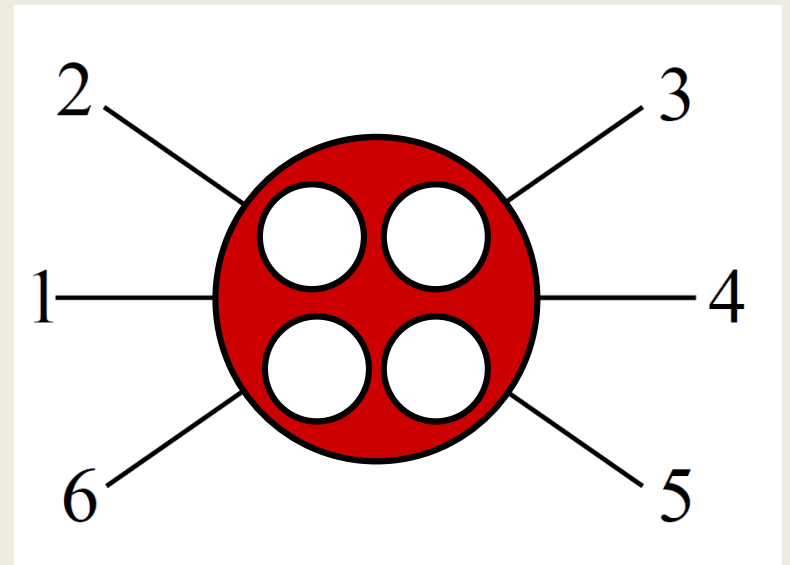
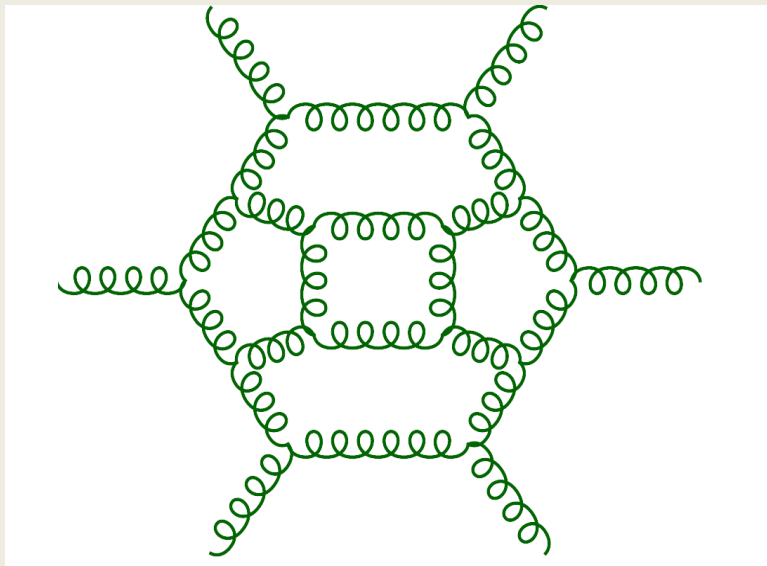
It has long been a goal of the S-matrix program to be able to construct scattering amplitudes based on a few physical principles and a thorough understanding of their analytic structure.

Important  
New  
Ingredients

1. Momentum twistors [Penrose, Hodges]
2. Symbol of Polylogarithms [GSVV]
3. Amplituhedron [Arkani-Hamed, Trnka]

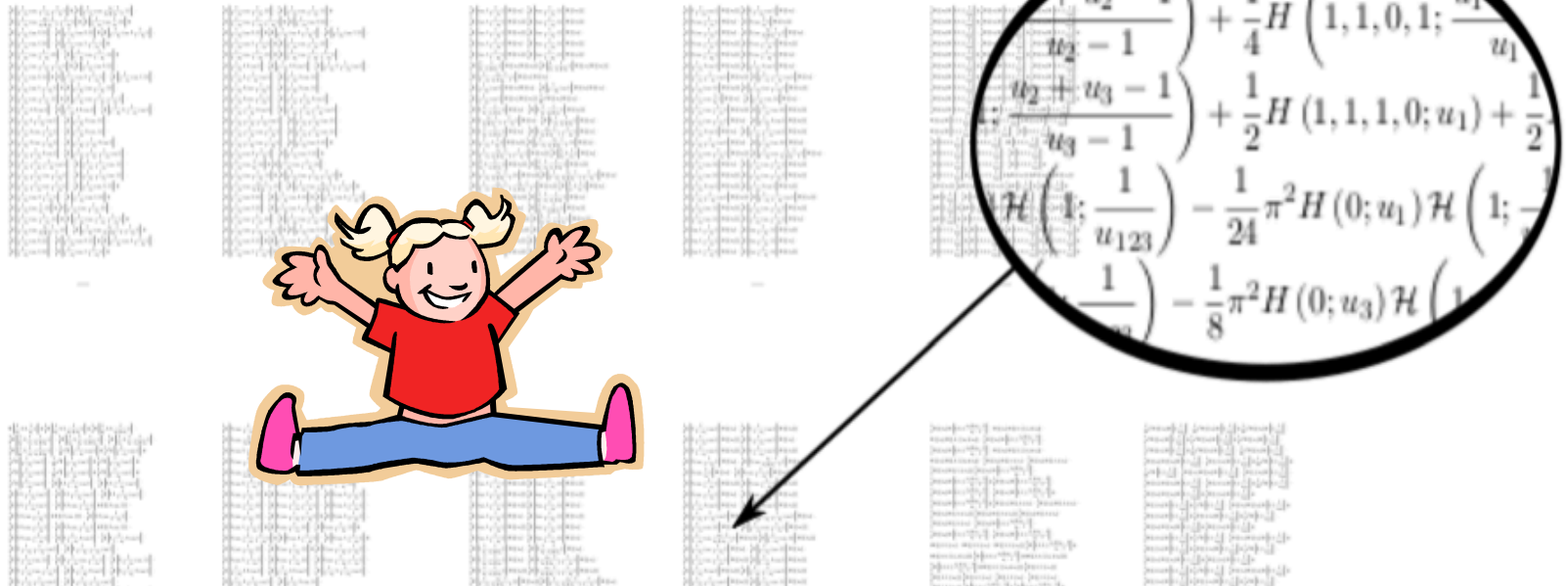
# Amplitude Bootstrap

Scattering Amplitudes:  
functions of the geometry  
of scattering configurations



Use mathematical and physical properties  
to determine these functions directly.

# Example: 2-loop 6-point MHV

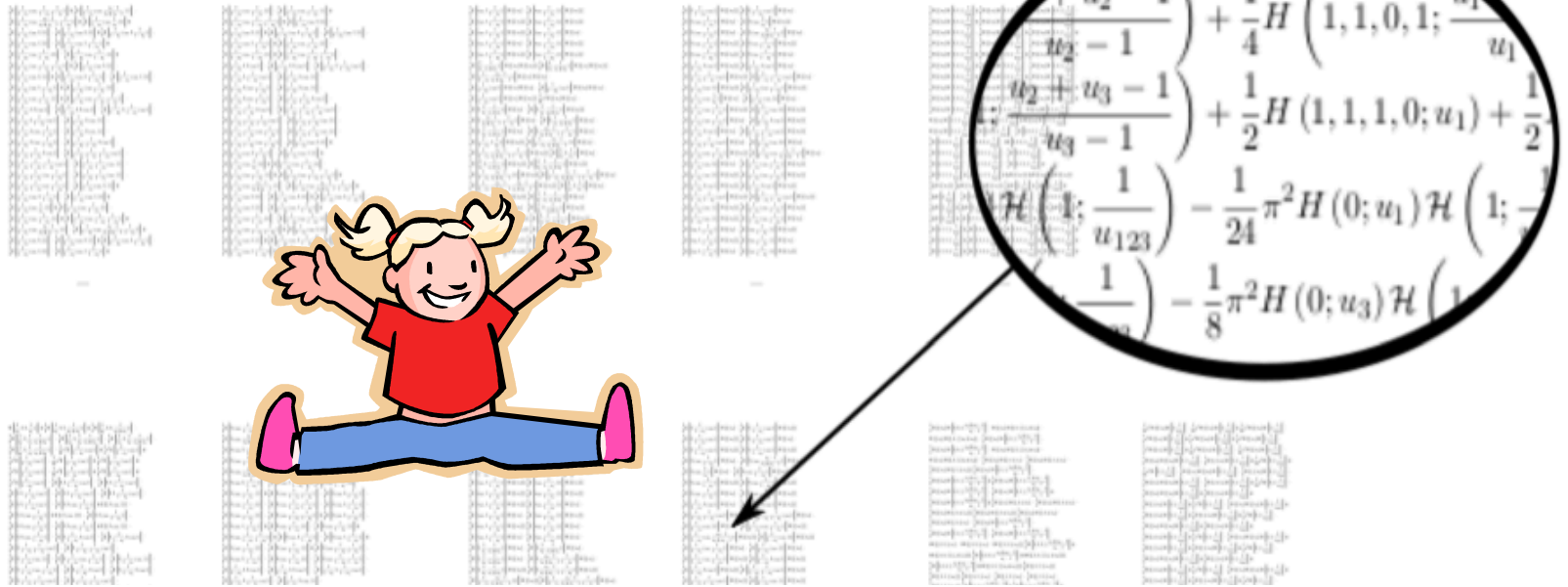


$$R_6^{(2)} = \sum_{\text{cyclic}} \text{Li}_4 \left( -\frac{\langle 1234 \rangle \langle 2356 \rangle}{\langle 1236 \rangle \langle 2345 \rangle} \right) - \frac{1}{4} \text{Li}_4 \left( -\frac{\langle 1246 \rangle \langle 1345 \rangle}{\langle 1234 \rangle \langle 1456 \rangle} \right) \\ + \text{products of } \text{Li}_k(-x) \text{ functions of lower weight}$$





# Example: 2-loop 6-point MHV



$$R_6^{(2)} = \sum_{\text{cyclic}} \text{Li}_4 \left( -\frac{\langle 1234 \rangle \langle 2356 \rangle}{\langle 1236 \rangle \langle 2345 \rangle} \right) - \frac{1}{4} \text{Li}_4 \left( -\frac{\langle 1246 \rangle \langle 1345 \rangle}{\langle 1234 \rangle \langle 1456 \rangle} \right) \\ + \text{products of } \text{Li}_k(-x) \text{ functions of lower weight}$$

Function and its Argument

GSVV

# Polylogarithm Functions

- Uniform Transcendentality  $2L$  @  $L$ -loops

- Logarithm Functions  $\log^{2L} x$

- Polylogarithm Functions

$$Li_{2L}(x) = \int_0^x \frac{dt}{t} Li_{2L-1}(t) \quad Li_1(x) = -\log(1-x)$$

- Generalized Polylogarithm Functions

$$G(a_1, \dots, a_{2L}; x) = \int_0^x \frac{dt}{t - a_1} G(a_2, \dots, a_{2L}; t)$$

- Elliptic Polylogarithms (n=10 L=2 k=3)

$$\tilde{\Gamma}\left(\begin{smallmatrix} n_1 & \dots & n_k \\ z_1 & \dots & z_k \end{smallmatrix}; z, \tau\right) = \int_0^z dz' g^{(n_1)}(z' - z_1, \tau) \tilde{\Gamma}\left(\begin{smallmatrix} n_2 & \dots & n_k \\ z_2 & \dots & z_k \end{smallmatrix}; z', \tau\right)$$



# Arguments of the Polylogarithms

Null vectors in 4d Minkowski space subject to momentum conservation can be represented by momentum twistors

$$\left( \begin{array}{c|c|c|c} | & | & \cdots & | \\ Z_1 & Z_2 & \cdots & Z_n \\ | & | & \cdots & | \end{array} \right), \quad Z_i \in \mathbb{P}^3$$

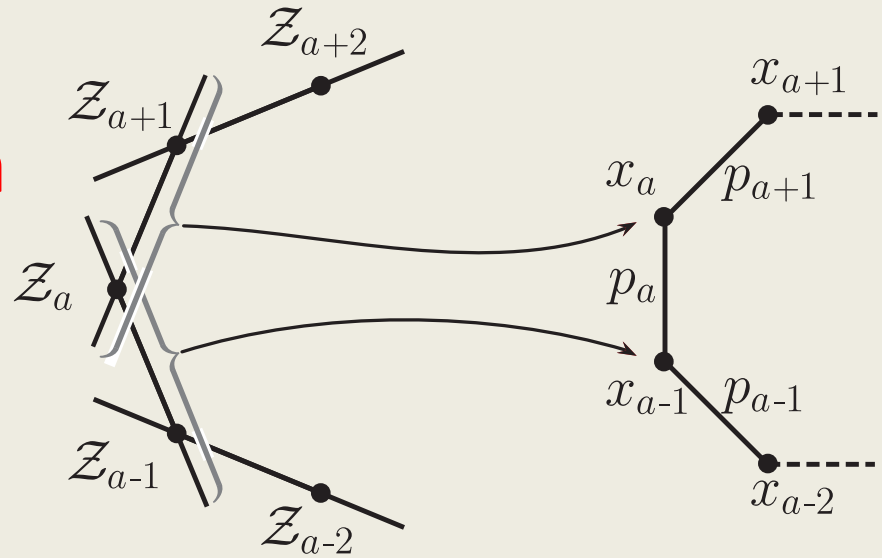
Hodges

$$\langle ijkl \rangle = \det(Z_i Z_j Z_k Z_l)$$

$n=6$

$$\begin{aligned} v_1 &= \frac{\langle 1246 \rangle \langle 1345 \rangle}{\langle 1234 \rangle \langle 1456 \rangle}, & v_2 &= \frac{\langle 1235 \rangle \langle 2456 \rangle}{\langle 1256 \rangle \langle 2345 \rangle}, & v_3 &= \frac{\langle 1356 \rangle \langle 2346 \rangle}{\langle 1236 \rangle \langle 3456 \rangle}, \\ x_1^+ &= \frac{\langle 1456 \rangle \langle 2356 \rangle}{\langle 1256 \rangle \langle 3456 \rangle}, & x_2^+ &= \frac{\langle 1346 \rangle \langle 2345 \rangle}{\langle 1234 \rangle \langle 3456 \rangle}, & x_3^+ &= \frac{\langle 1236 \rangle \langle 1245 \rangle}{\langle 1234 \rangle \langle 1256 \rangle}, \\ x_1^- &= \frac{\langle 1234 \rangle \langle 2356 \rangle}{\langle 1236 \rangle \langle 2345 \rangle}, & x_2^- &= \frac{\langle 1256 \rangle \langle 1346 \rangle}{\langle 1236 \rangle \langle 1456 \rangle}, & x_3^- &= \frac{\langle 1245 \rangle \langle 3456 \rangle}{\langle 1456 \rangle \langle 2345 \rangle}, \end{aligned}$$

## Relation to momentum



$$\langle 1\,4\,5\,6 \rangle \sim \frac{(p_1 + p_2)^2 (p_3 + p_4)^2}{(p_2 + p_3 + p_4)^2} \frac{u_1 + u_2 + u_3 - 1 + \sqrt{\Delta}}{2u_1 u_2 u_3},$$

$$\Delta = (1 - u_1 - u_2 - u_3)^2 - 4u_1 u_2 u_3$$

$$u_1 = \frac{(p_1 + p_2)^2 (p_4 + p_5)^2}{(p_1 + p_2 + p_3)^2 (p_4 + p_5 + p_6)^2}, \quad u_2, u_3 = \text{cyclic}$$

# Symbol of Polylogarithm

17-pages can be simplified using a very useful tool from modern mathematics

Symbol

$$\log(R) \rightarrow R$$

$$Li_2(R) \rightarrow -(1 - R) \otimes R$$

$$Li_2(x) + Li_2(-x) = \frac{1}{2} Li_2(x^2)$$

$$-(1-x) \otimes x - (1+x) \otimes (-x) = -(1-x^2) \otimes x = -\frac{1}{2}(1-x^2) \otimes x^2$$



QCD computations

# Symbol Alphabet: n=6

- Amplitude = Function of Cross-Ratios

$$R_6^{(2)} = \sum_{\text{cyclic}} \text{Li}_4 \left( -\frac{\langle 1234 \rangle \langle 2356 \rangle}{\langle 1236 \rangle \langle 2345 \rangle} \right) - \frac{1}{4} \text{Li}_4 \left( -\frac{\langle 1246 \rangle \langle 1345 \rangle}{\langle 1234 \rangle \langle 1456 \rangle} \right) \\ + \text{products of } \text{Li}_k(-x) \text{ functions of lower weight}$$

- Full Symbol

$\langle 1256 \rangle \otimes \langle 1346 \rangle \otimes \langle 1246 \rangle \otimes \langle 1456 \rangle + \dots$  7272 terms

- Symbol Alphabet: 9 letters

$\langle 1235 \rangle, \langle 2345 \rangle, \langle 1345 \rangle, \langle 2456 \rangle, \langle 1356 \rangle, \langle 1246 \rangle, \langle 1245 \rangle, \langle 2356 \rangle, \langle 1346 \rangle$

# Amplitudes Bootstrap

start with symbol alphabet, impose constraints,  
find unique solution

Constraint	$L = 1$	$L = 2$	$L = 3$	$L = 4$	$L = 5$
0. Functions	(10,10)	(82,88)	(639,761)	(5153,6916)	(???,???)
1. Steinmann	(7,7)	(37,39)	(174,190)	(758,839)	(3105,3434)
2. Symmetry	(3,5)	(11,24)	(44,106)	(174,451)	(???,???)
3. Final-entry	(2,2)	(5,5)	(19,12)	(72,32)	(272,83)
4. Collinear	(0,0)	(0,0)	(1,1)	(3,5)	(9,15)
5. Regge	(0,0)	(0,0)	(0,0)	(0,0)	(0,0)

[Caron-Huot, Dixon, Drummond, McLeod, Papathanasiou, Spradlin, Von Hippel]

Bootstrap in QCD amplitudes: Almelid, Duhr,  
Gardi, McLeod, White; Chicherin, Henn, Mitev

# How Do We Determine Symbol Alphabet?

1. Symbol Alphabet and Cluster Algebra
2. Symbol Alphabet and Landau Singularities

# Symbol Alphabet and Cluster Algebra

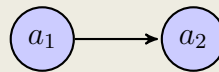
Symbol Alphabet is given by  
a subset of cluster coordinates of  
Grassmannian Cluster Algebra

$$Gr(4, n)$$

Golden, Goncharov, Spradlin, Vergu, AV

# $A_2$ Cluster Algebra

Initial Quiver

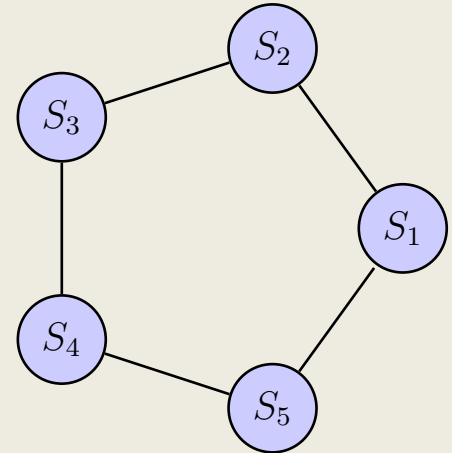


Mutation Rule

$$a_{n+1} = \frac{1 + a_n}{a_{n-1}}$$

Cluster Coordinates

$$a_1, a_2, a_3 = \frac{1 + a_2}{a_1}, a_4 = \frac{1 + a_1 + a_2}{a_1 a_2}, a_5 = \frac{1 + a_1}{a_2}$$

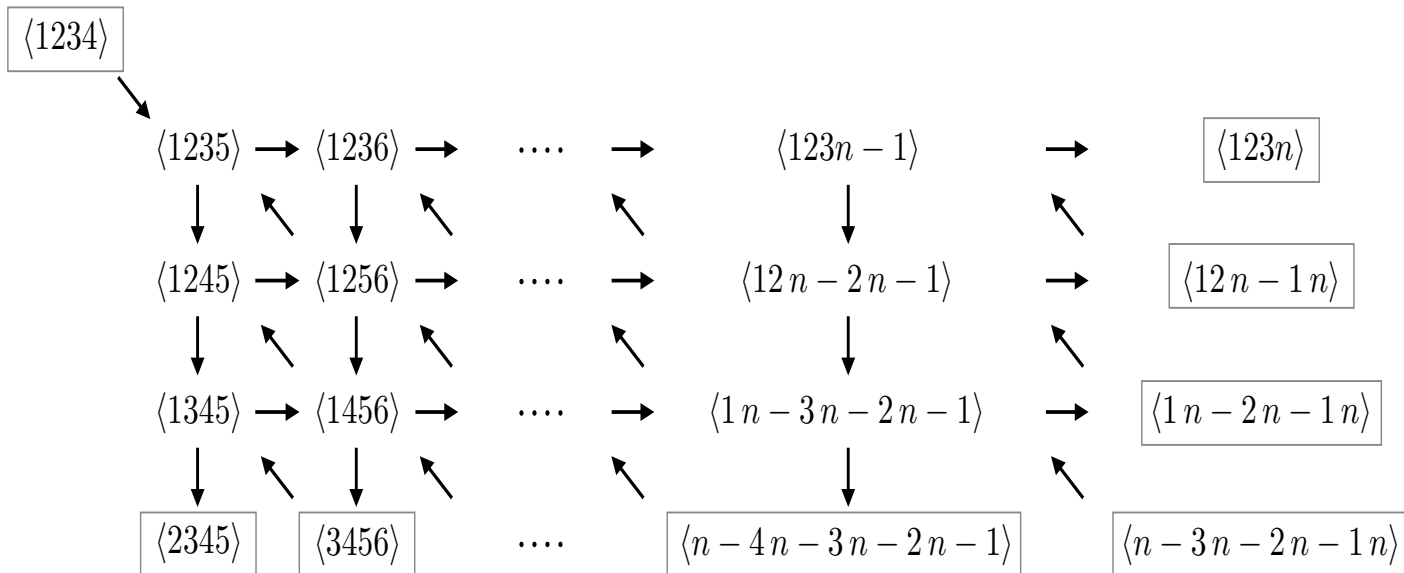




# Grassmannian Cluster Algebra

$$Gr(4, n)$$

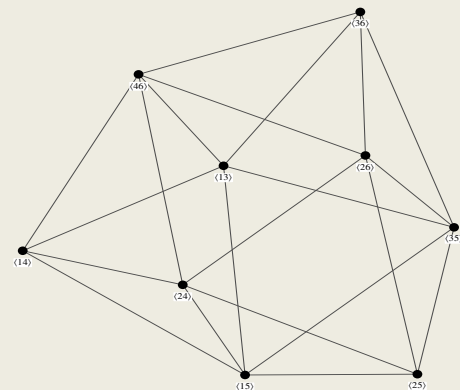
$3(n - 5)$  Quiver + Mutation Rule



Fomin, Zelevinsky; Scott; Gekhtman, Shapiro, Vainshtein

# 3-loop 6-point Symbol

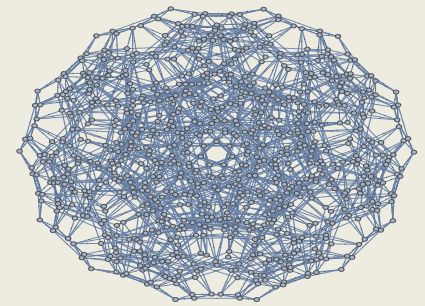
Symbol Alphabet: 9 letters



$\langle 1235 \rangle, \langle 2345 \rangle, \langle 1345 \rangle, \langle 2456 \rangle, \langle 1356 \rangle, \langle 1246 \rangle, \langle 1245 \rangle, \langle 2356 \rangle, \langle 1346 \rangle$

```
-12 T[x2, x1, x1, x1, x1, x1] - 12 T[x2, x1, x1, x1, x1, x6] +
T[x2, x1, x1, x1, x2, x1] + T[x2, x1, x1, x1, x2, x3] + T[x2, x1, x1, x1, x2, x4] +
T[x2, x1, x1, x1, x2, x6] - 12 T[x2, x1, x1, x1, x6, x1] -
12 T[x2, x1, x1, x1, x6, x6] + T[x2, x1, x1, x1, x8, x1] +
T[x2, x1, x1, x1, x8, x5] + T[x2, x1, x1, x1, x8, x6] + T[x2, x1, x1, x1, x8, x9] +
T[x2, x1, x1, x2, x1, x1] + T[x2, x1, x1, x2, x1, x3] + ... 5376 ... +
T[x8, x9, x9, x8, x9, x6] + T[x8, x9, x9, x8, x9, x9] - 12 T[x8, x9, x9, x9, x5, x5] -
12 T[x8, x9, x9, x9, x5, x9] + T[x8, x9, x9, x9, x7, x3] + T[x8, x9, x9, x9, x7, x4] +
T[x8, x9, x9, x9, x7, x5] + T[x8, x9, x9, x9, x7, x9] + T[x8, x9, x9, x9, x8, x1] +
T[x8, x9, x9, x9, x8, x5] + T[x8, x9, x9, x9, x8, x6] + T[x8, x9, x9, x9, x8, x9] -
12 T[x8, x9, x9, x9, x9, x5] - 12 T[x8, x9, x9, x9, x9, x9]
```

# 3-loop 7-point Symbol



Symbol Alphabet: 42 letters

```
-48 T[a11, a11, a11, a11, a11, a21] - 48 T[a11, a11, a11, a11, a11, a22] +  
48 T[a11, a11, a11, a11, a11, a25] - 48 T[a11, a11, a11, a11, a11, a31] +  
48 T[a11, a11, a11, a11, a11, a34] - 48 T[a11, a11, a11, a11, a11, a37] +  
12 T[a11, a11, a11, a11, a14, a21] - 12 T[a11, a11, a11, a11, a14, a24] -  
12 T[a11, a11, a11, a11, a14, a25] - 12 T[a11, a11, a11, a11, a14, a33] -  
12 T[a11, a11, a11, a11, a14, a34] + ... 703 674 ... + 12 T[a17, a66, a66, a66, a37, a37] -  
12 T[a17, a66, a66, a66, a45, a23] - 12 T[a17, a66, a66, a66, a45, a37] -  
12 T[a17, a66, a66, a66, a47, a25] - 12 T[a17, a66, a66, a66, a47, a32] -  
12 T[a17, a66, a66, a66, a55, a21] - 12 T[a17, a66, a66, a66, a55, a32] -  
12 T[a17, a66, a66, a66, a57, a23] - 12 T[a17, a66, a66, a66, a57, a34]
```

# Cluster Adjacency

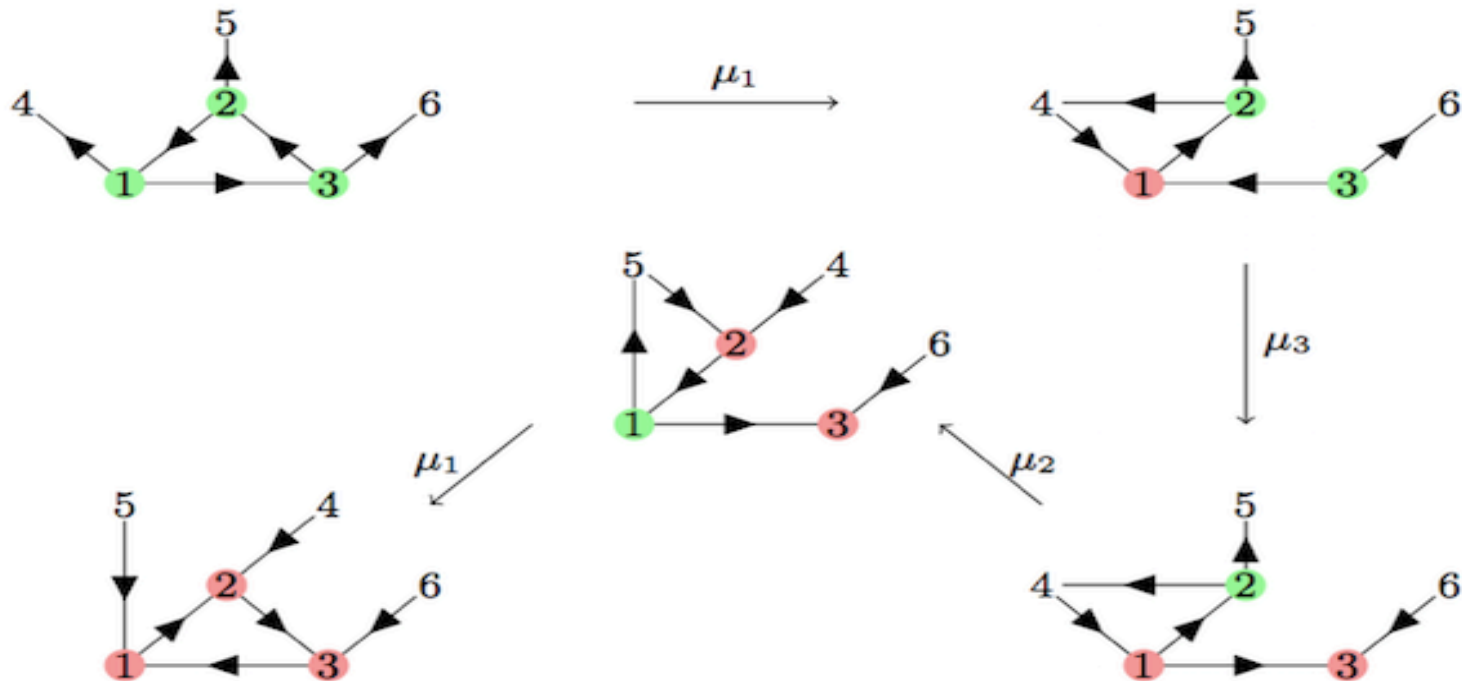
*Two distinct  $A$ -coordinates can appear consecutively in a symbol only if there exists a cluster where they both appear.*

Drummond, Foster, Gurdogan

- Works for all known examples
- It is not known whether there is a physical or mathematical principle which requires this to be true, nor how general it is...

# In Progress

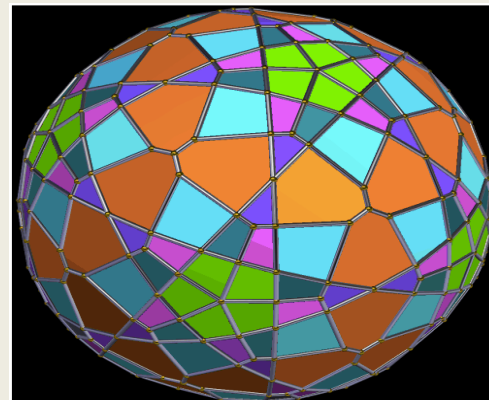
- Write symbol in terms of preferred clusters
- Look at triples, quads, etc



# Beyond Cluster Algebra

Arkani-Hamed, Lam, Spradlin

- For  $n > 7$ : cluster algebra is infinite dimensional and amplitudes involve square roots.
- Natural object:  $Gr(4, 8)/T^7$
- It is a finite polytope with 360 faces.
- There is a natural role for square roots to play in the variables associated to this polytope.



# Other Interesting Applications of Cluster Algebras to Amplitudes

- Coproduct: only cluster X-coordinates appear which constraints the class of functions to cluster polylogarithm functions [GPSV]
- Special kinematics: [Duhr et al]
- Applications to other theories: [Henn et al]
- Cluster algebras also show up in the integrand: [Arkani-Hamed et al]

# Symbol Alphabet and Landau Singularities



# Symbol and Landau Singularities

- Symbol captures the information about the singularity structure of the polylogarithm functions.
- Given the integrand, the location of its singularities is described by Landau equations.
- There should be a close connection between solutions of Landau Equations (integrand) and Symbol Alphabet (integral).

Maldacena, Simons-Duffin, Zhiboedov

Abreu, Britto, Duhr, Gardi

# Landau Singularities

Landau 1959, ELOP

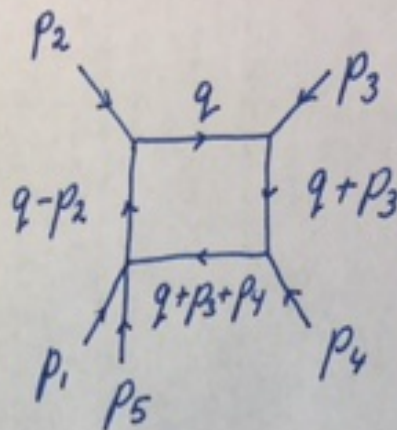
Landau Equations for a given Feynman integral are a set of kinematic constraints that are necessary for the appearance of a pole or a branch point in the integrated function:

$$\sum_{i \in \text{loop}} \alpha_i q_i^\mu = 0 \quad \forall \text{ loops},$$
$$\alpha_i (q_i^2 - m_i^2) = 0 \quad \forall i.$$

**Landau Singularities** occur for external momenta such that the Landau equations have solutions.

# Example

$$\int \frac{d^4 q}{q^2 (q-p_2)^2 (q+p_3)^2 (q+p_3+p_4)^2}$$



$$\begin{cases} (q-p_2)^2 = q^2 = (q+p_3)^2 = (q+p_3+p_4)^2 = 0 \\ \alpha_1 (q-p_2) + \alpha_2 q + \alpha_3 (q+p_3) + \alpha_4 (q+p_3+p_4) = 0 \end{cases} \quad \Leftrightarrow (p_2+p_3)^2 (p_3+p_4)^2 = 0$$

$$\langle 1234 \rangle \langle 2345 \rangle = 0$$

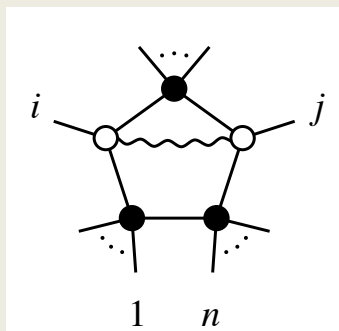
# One-loop n-point MHV

Bern, Dixon, Dunbar, Kosower

## Chiral pentagon integrals

Arkani-Hamed, Bourjaily,  
Cachazo, Trnka

$$\frac{\mathcal{A}_{\text{MHV}}^{1\text{-loop}}}{\mathcal{A}_{\text{MHV}}^{\text{tree}}} = \int_{AB} \sum_{1 < i < j < n}$$



## Landau Singularities

Dennen, Spradlin, AV

$$\langle i \bar{j} \rangle \quad \langle i-1 \ i \ j-1 \ j \rangle$$

$$\langle j(j-1 \ j+1)(i \ i+1)(k \ k+1) \rangle$$

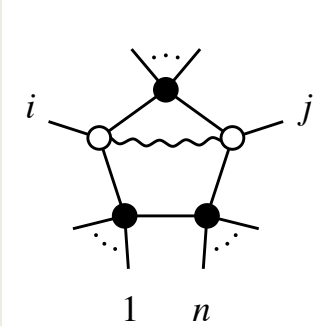
$$\langle (abc) \cap (def) gh \rangle = \langle abcg \rangle \langle defh \rangle - \langle abch \rangle \langle defg \rangle$$

$$\langle (abc) \cap (dec) gh \rangle \equiv \langle c(ab)(de)(gh) \rangle \quad \bar{a} = (a-1, a, a+1)$$

# Symbol: one-loop n-point MHV

Bern, Dixon, Dunbar, Kosower

Arkani-Hamed, Bourjaily, Cachazo, Trnka

$\int_{AB}$ 


$$= \text{Li}_2(1 - u_{n,i-1,i,j}) - \text{Li}_2(1 - u_{j,n,i,j-1}) - \text{Li}_2(1 - u_{i,j-1,n,i-1})$$

$$- \text{Li}_2(1 - u_{i,j-1,n,i-1}) + \text{Li}_2(1 - u_{i,j-1,j,i-1})$$

$$+ \log(u_{j,n,i-1,j-1}) \log(u_{n,i-1,i,j})$$

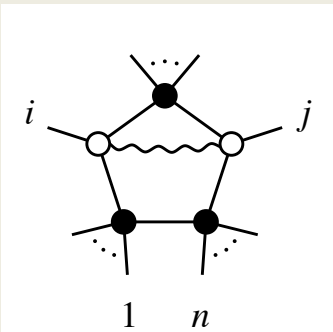
- One-loop n-point MHV Symbol Alphabet

$$\langle i-1 \ i \ j-1 \ j \rangle \quad \langle i \ \bar{j} \rangle ,$$

# Symbol vs LS: one-loop n-point MHV

Bern, Dixon, Dunbar, Kosower

Arkani-Hamed, Bourjaily, Cachazo, Trnka

$$\int_{AB} \text{Diagram} = \text{Li}_2(1 - u_{n,i-1,i,j}) - \text{Li}_2(1 - u_{j,n,i,j-1}) - \text{Li}_2(1 - u_{i,j-1,n,i-1}) \\ - \text{Li}_2(1 - u_{i,j-1,n,i-1}) + \text{Li}_2(1 - u_{i,j-1,j,i-1}) \\ + \log(u_{j,n,i-1,j-1}) \log(u_{n,i-1,i,j})$$


- One-loop n-point MHV Symbol Alphabet

$$\langle i-1 \ i \ j-1 \ j \rangle \quad \langle i \ \bar{j} \rangle ,$$

- Landau Singularities

$$\langle i-1 \ i \ j-1 \ j \rangle \quad \langle i \ \bar{j} \rangle , \quad \langle j \ (j-1 \ j+1) (i \ i+1) (k \ k+1) \rangle$$

Spurious Landau Singularity

# Landau Singularities vs Symbol Alphabet

**All** symbol entries are Landau Singularities.

There are spurious Landau Singularities  
which are **not** in Symbol Alphabet.

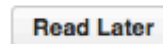
[not surprising: numerators, sum of integrals, etc]

Q: How to eliminate spurious Landau Singularities?

A: Use Amplituhedron.

BEST INVENTIONS

# The 25 Best Inventions of the Year 2013

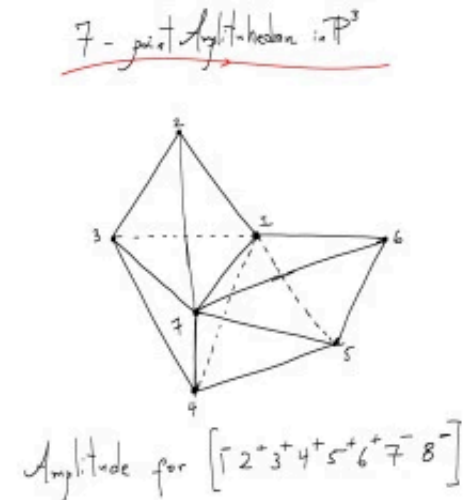


Interesting

## The Amplituhedron

By TIME Staff | Nov. 13, 2013

Physicists at the Institute for Advanced Study in Princeton, N.J., recently found a major shortcut for predicting subatomic-particle collisions. The new method **represents probabilities as pyramid-like structures**, then combines the pyramids into one elegant gemstone-like structure called an amplituhedron, thereby massively simplifying the task of calculating particle interactions. Ultimately the amplituhedron could lead to the long-sought quantum theory of gravity.



The integrand of MHV amplitude is a canonical form defined by having logarithmic singularities only on the boundary of the Amplituhedron.





# Amplituhedron

Arkani-Hamed, Trnka

The integrand of MHV amplitude is a canonical form defined by having logarithmic singularities only on the boundary of the Amplituhedron.

$$\mathcal{L}_{\alpha}^{(\ell)I} = \sum_{i=1}^n D_{\alpha i}^{(\ell)} Z_i^I$$

Diagram illustrating the components of the equation:

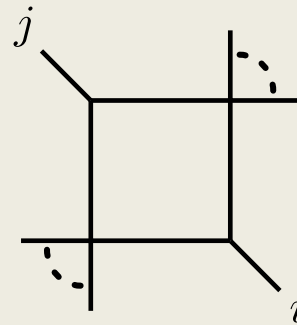
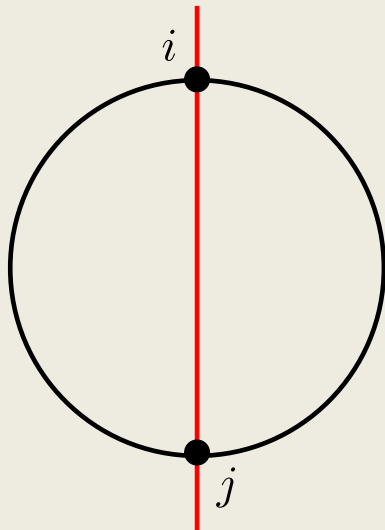
- $\mathcal{L}_{\alpha}^{(\ell)I}$  is associated with **Loop Momenta**.
- $D_{\alpha i}^{(\ell)}$  is associated with **D-matrices** and  $G_+(2, n)$ .
- $Z_i^I$  is associated with **External Momenta** and  $G_+(4, n)$ .

**Inside Amplituhedron**

$$\langle \mathcal{L}^{(\ell)} i j \rangle > 0 \text{ for } i < j \text{ and all } \ell$$

$$\langle \mathcal{L}^{(\ell_1)} \mathcal{L}^{(\ell_2)} \rangle > 0 \text{ for all } \ell_1, \ell_2.$$

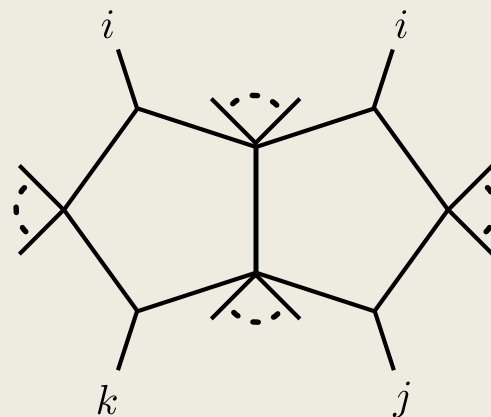
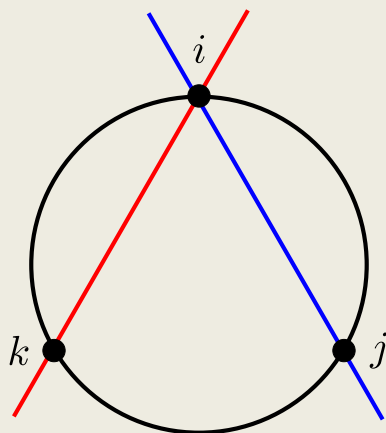
# One-loop MHV Amplitudes



$$\langle i-1 \ i \ j-1 \ j \rangle$$

$$\langle i \ \bar{j} \rangle$$

# Two-Loop MHV Amplitudes



$$\langle a \bar{b} \rangle$$

$$\langle a b c c+1 \rangle$$

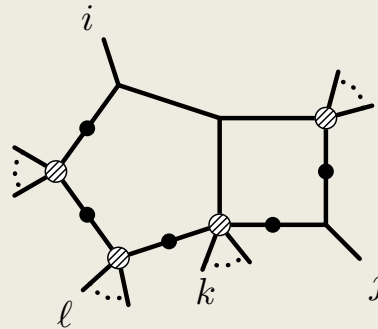
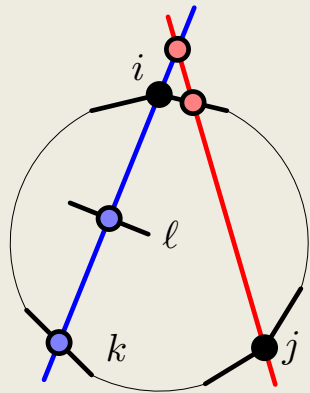
$$\langle a a+1 \bar{b} \cap \bar{c} \rangle$$

$$\langle a (a-1 a+1) (b b+1) (c c+1) \rangle$$

# Two-loop NMHV and Three-loop MHV

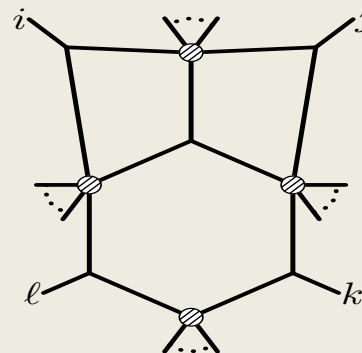
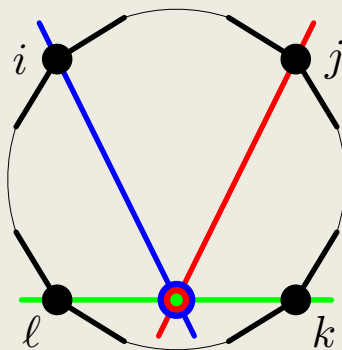
Prlina, Spradlin, Stankowicz, Stanojevic

- Two-loop NMHV: allow for sign flips



Arkani-Hamed, Thomas, Trnka

- Three-loop MHV



$$\langle a a + 1 b b + 1 \rangle$$

$$\langle a b c c + 1 \rangle$$

$$\langle a (b, b + 1) (c, c + 1) (d, d + 1) \rangle$$

$$\langle a (a - 1, a + 1) (b, b + 1) (c, c + 1) \rangle$$

$$\langle a (a - 1, d) (b, b + 1) (c, c + 1) \rangle$$

$$\langle \bar{b} \cap (d, c, c + 1) (a, a + 1) \rangle$$

$$\langle \bar{b} \cap (a a + 1 b), \bar{d} \cap (d, c, c + 1) \rangle$$

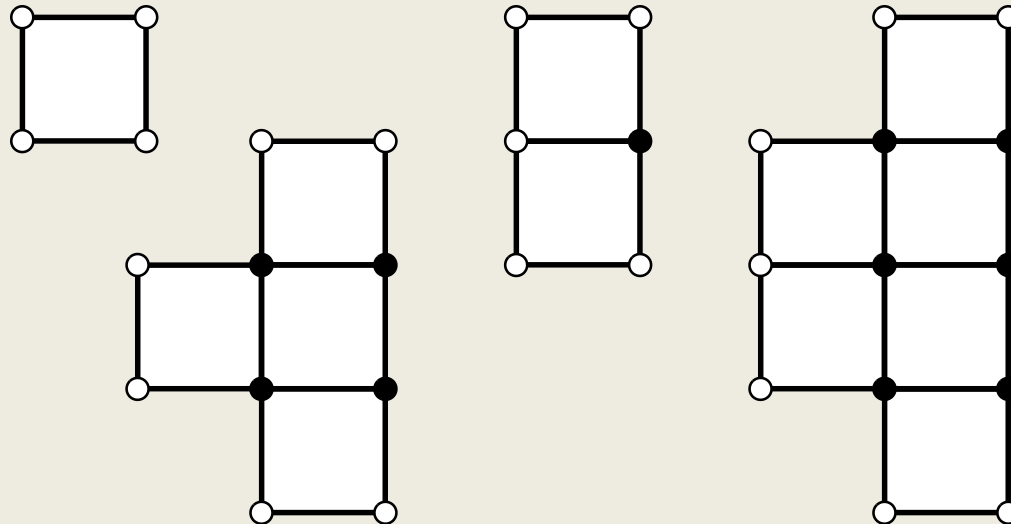
$$\langle b (b - 1, b + 1) (c, c + 1) (a, d) \rangle,$$



# Ziggurat Graphs

Prlina, Spradlin, Stanojevic

- Landau singularities of any  $n$ -particle amplitude in any massless, planar theory are a subset of those of a special type of "ziggurat graph."
- For  $n=6$  consistent with symbol alphabet.



# Conclusion

- N=4 Yang-Mills Amplitudes Bootstrap Program
- Symbol Alphabet and Cluster Algebra
- Symbol Alphabet and Landau Singularities
- Many Open Questions Remain

Thank you!