

The g-2 experiment at Fermilab



Becky Chislett

**UK HEP Forum – The spice of flavour
28th November 2018**



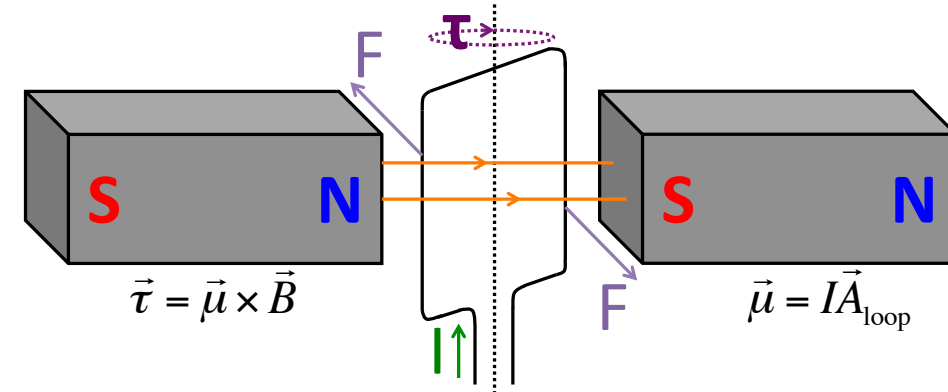
The Magnetic Moment

The magnetic moment determines how something interacts with a magnetic field

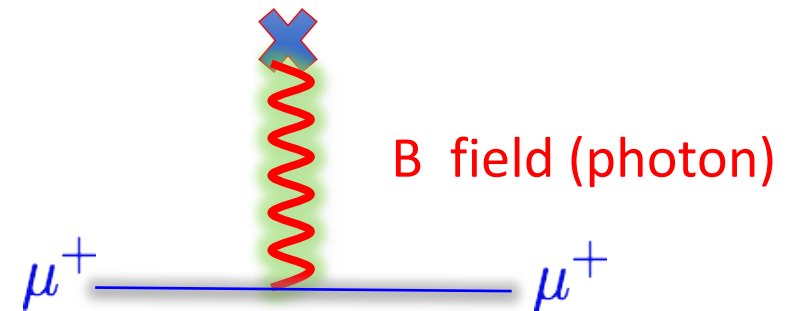
A magnetic moment will experience a force when placed in a magnetic field :

For particles the magnetic moment is related to the spin through the g factor

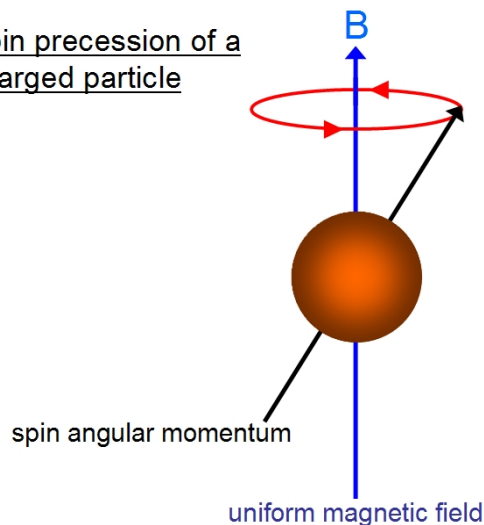
$$\vec{\mu} = g \frac{q}{2m} \vec{s}$$



Dirac predicted that $g = 2$:



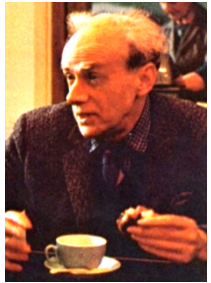
Spin precession of a charged particle



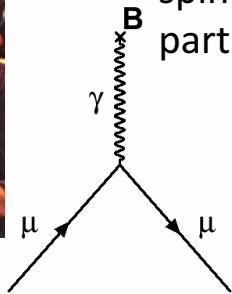
When placed in a magnetic field this causes the spin to precess

The Magnetic Moment

Dirac



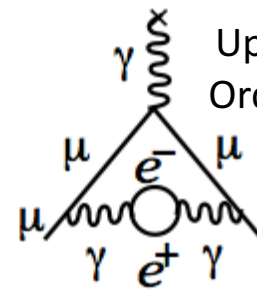
Charged,
spin 1/2
particle



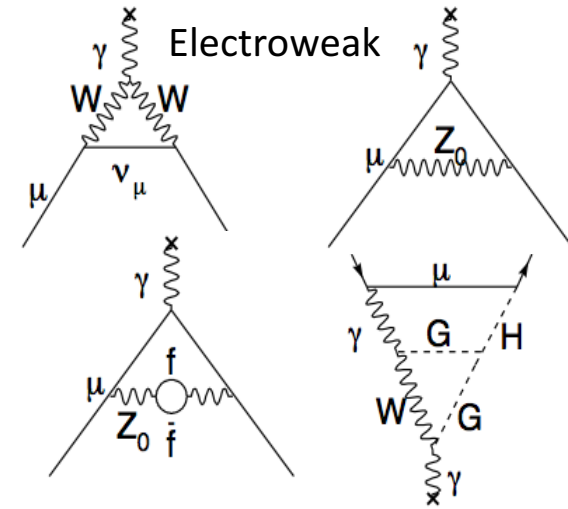
12672
diagrams



Kinoshita

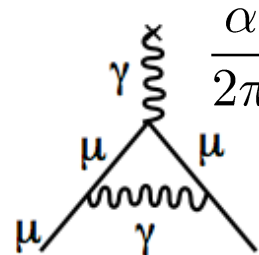
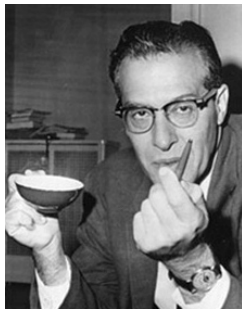


Up to 10th
Order QED



$$g_\mu = 2.002\,331\,841\,78(126)$$

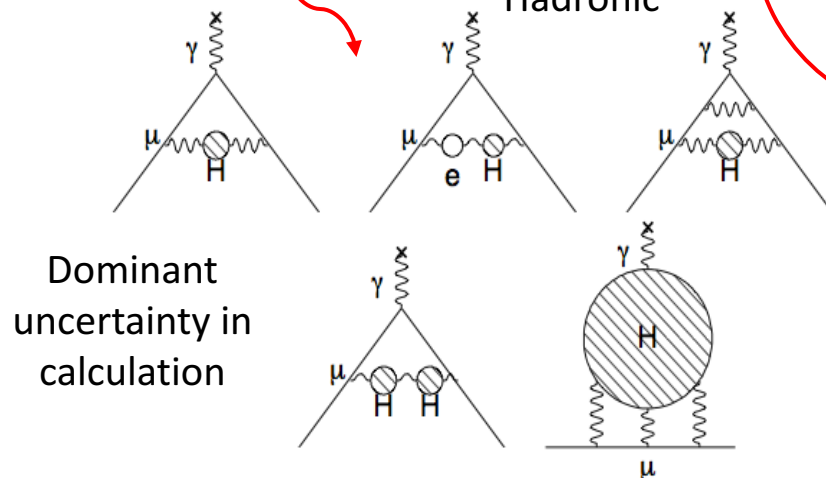
Schwinger



1st Order QED

$$\frac{\alpha}{2\pi} = 0.00232$$

Hadronic



Dominant
uncertainty in
calculation

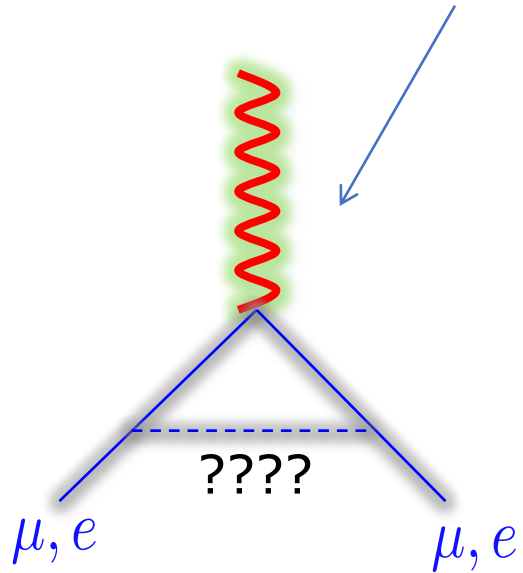
?

Motivation

The previous measurement at BNL differs from the theoretical prediction by $\sim 3.5\sigma$.

Is this :

- A mistake in the theory
- A mistake / statistical fluctuation in the experiment
- A sign of new physics



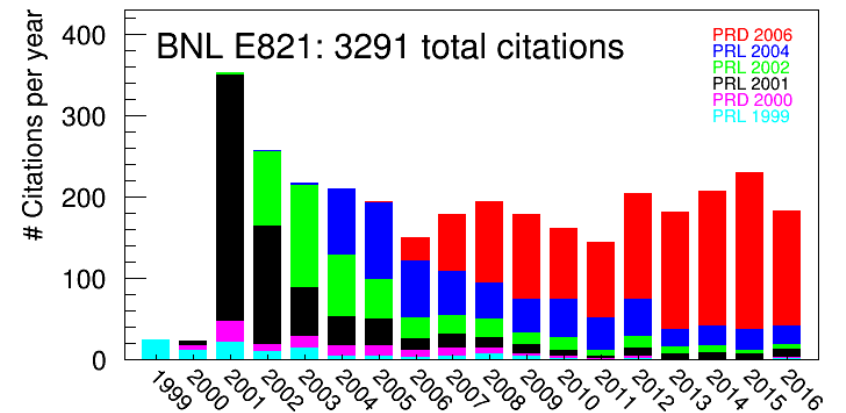
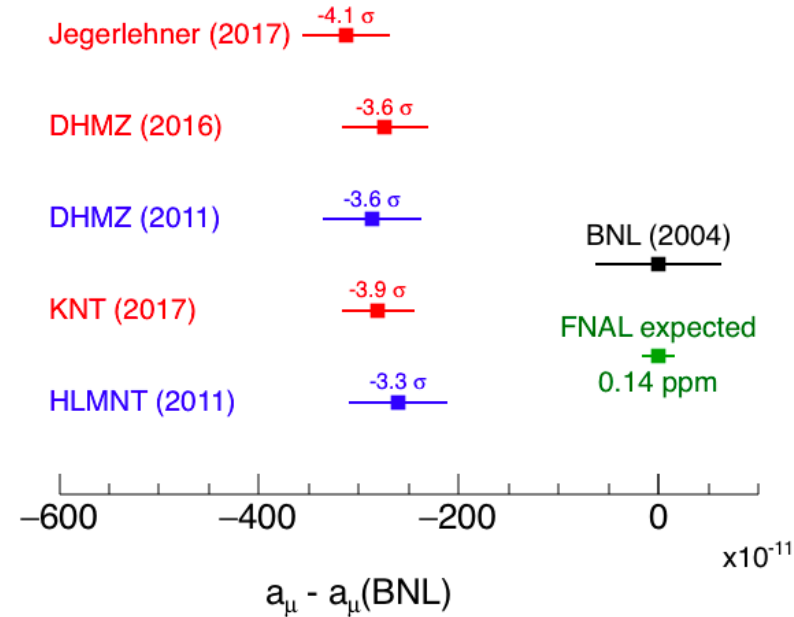
Contributes as $\left(\frac{m_\ell}{M_{\text{NEW}}}\right)^2$

The muon has a mass advantage

$$\left(\frac{m_\mu}{m_e}\right)^2 \approx 44,000$$

Muon g-2 is sensitive to new physics from MeV to TeV scales

Comparison of SM & BNL Measurement



Experimental Technique

The anomalous magnetic moment causes the spin to precess faster than the momentum vector as the muon moves around the ring

$$\omega_s = \frac{eB}{\gamma m} + \frac{eB}{m} \left(\frac{g-2}{2} \right)$$

$$\omega_c = \frac{eB}{\gamma m}$$

$$\omega_a = \omega_s - \omega_c = \left(\frac{g-2}{2} \right) \frac{eB}{m} = a \frac{eB}{m}$$

Measure the magnetic field in the ring

Measure the spin precession from the positron decays

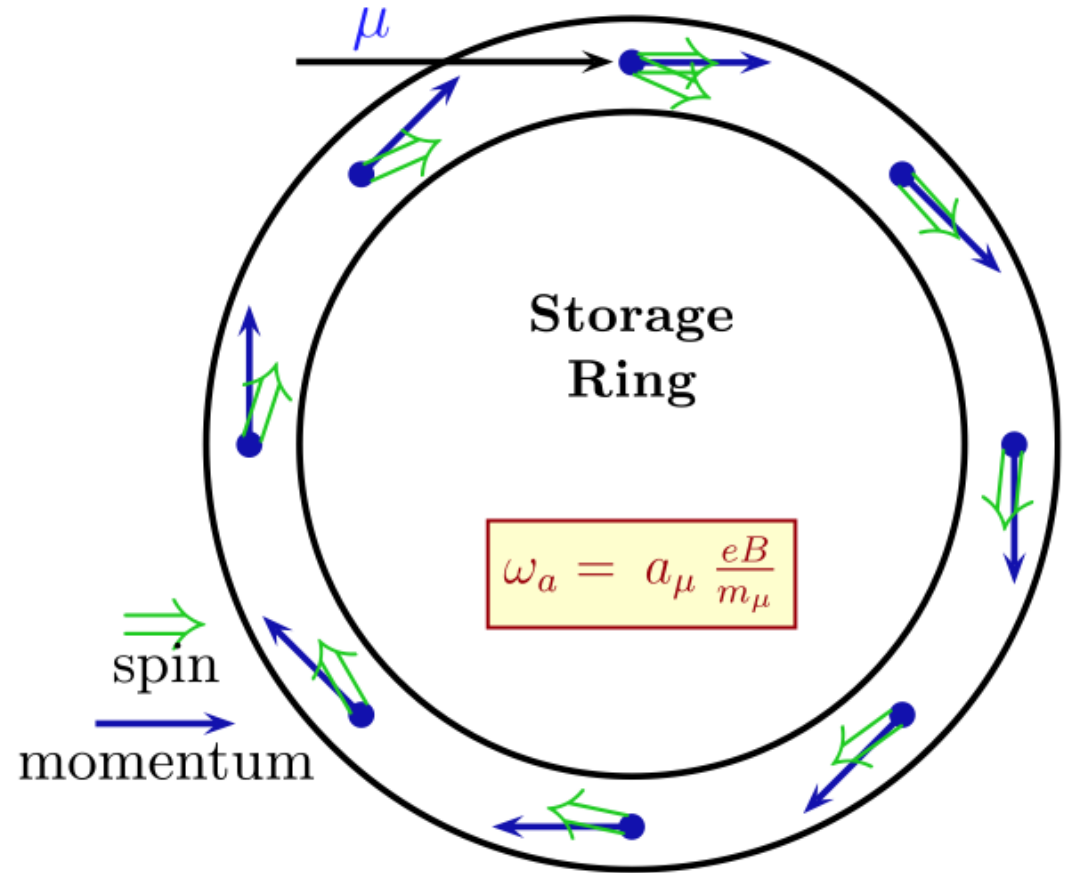
We actually measure 2 frequencies :

$$a_\mu = \frac{\omega_a}{\tilde{\omega}_p} \frac{\mu_p}{\mu_e} \frac{m_\mu}{m_e} \frac{g_e}{2}$$

3ppb

22ppb

0.0003ppb

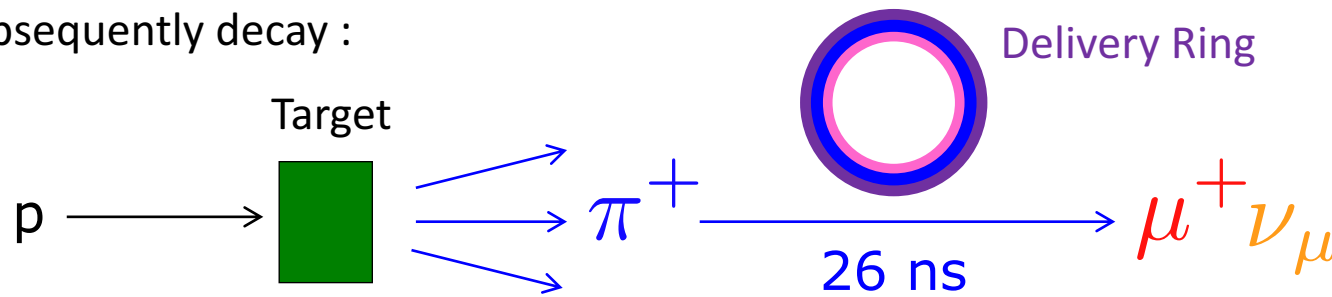


In a 1.5 T magnetic field the spin rotates in 144ns and the momentum in 149ns

Muon Production

The Fermilab accelerator provides a purer muon beam giving more muons at a lower instantaneous rate

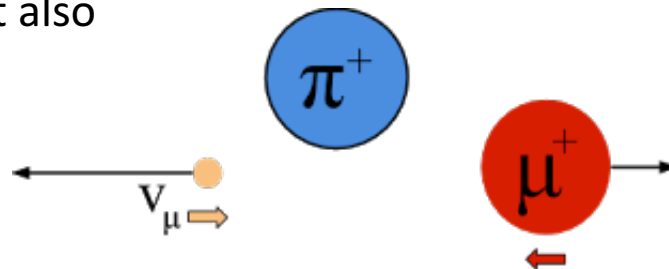
Protons hit a pion target to produce pions which subsequently decay :



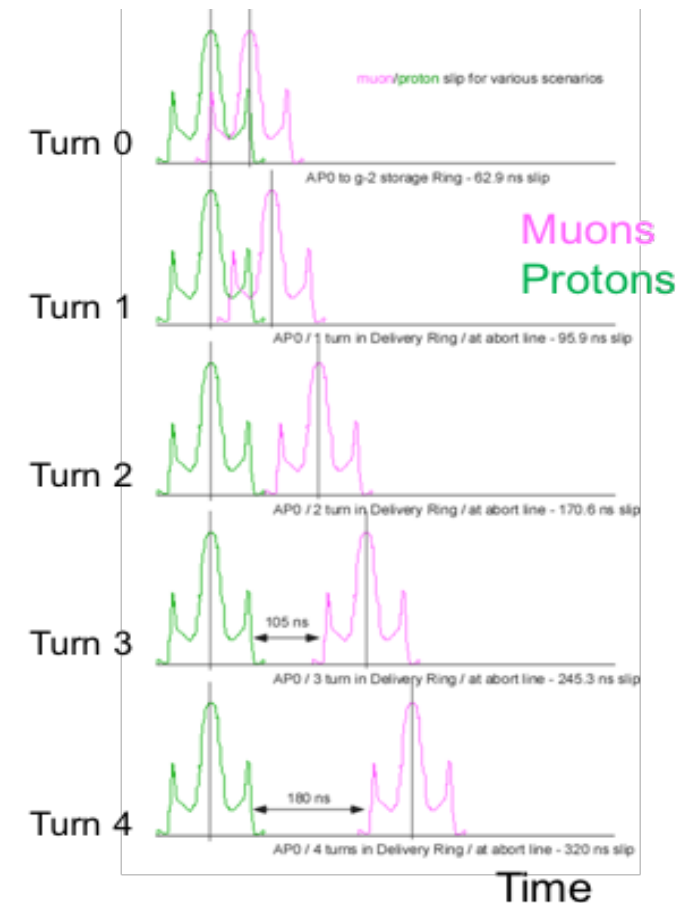
3.11 GeV pions selected using a lithium lens

In the pion decay the neutrino must have spin opposite to the momentum

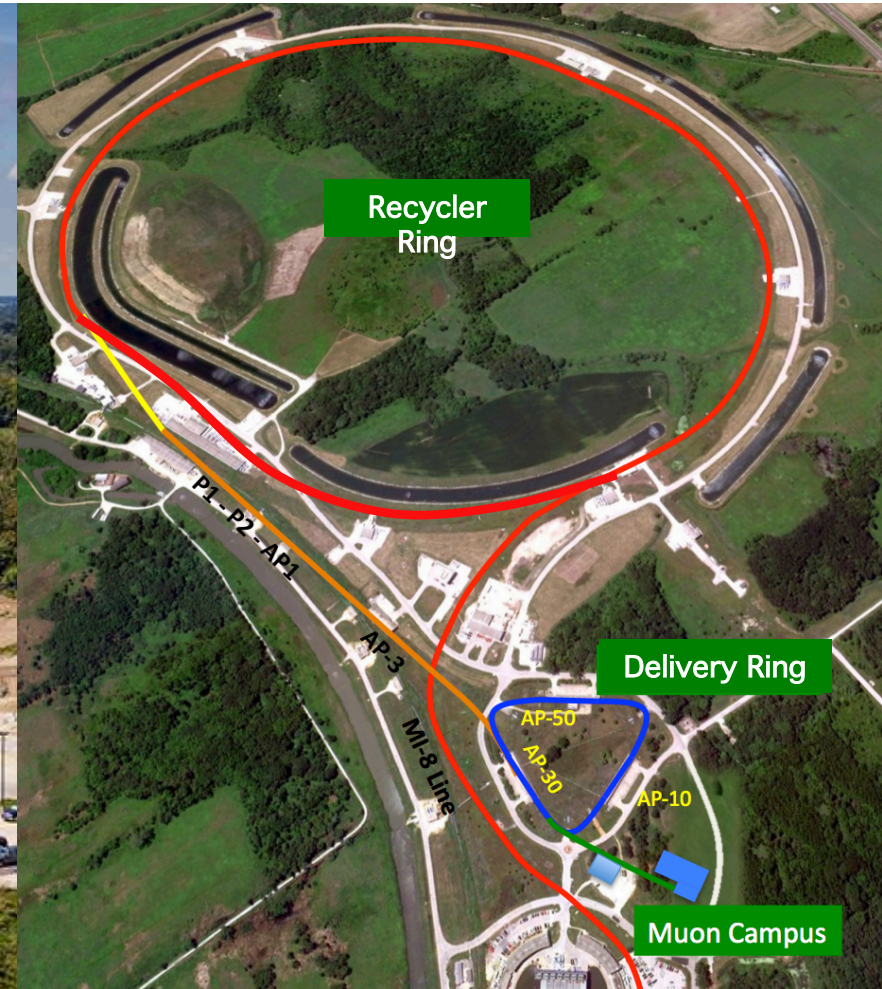
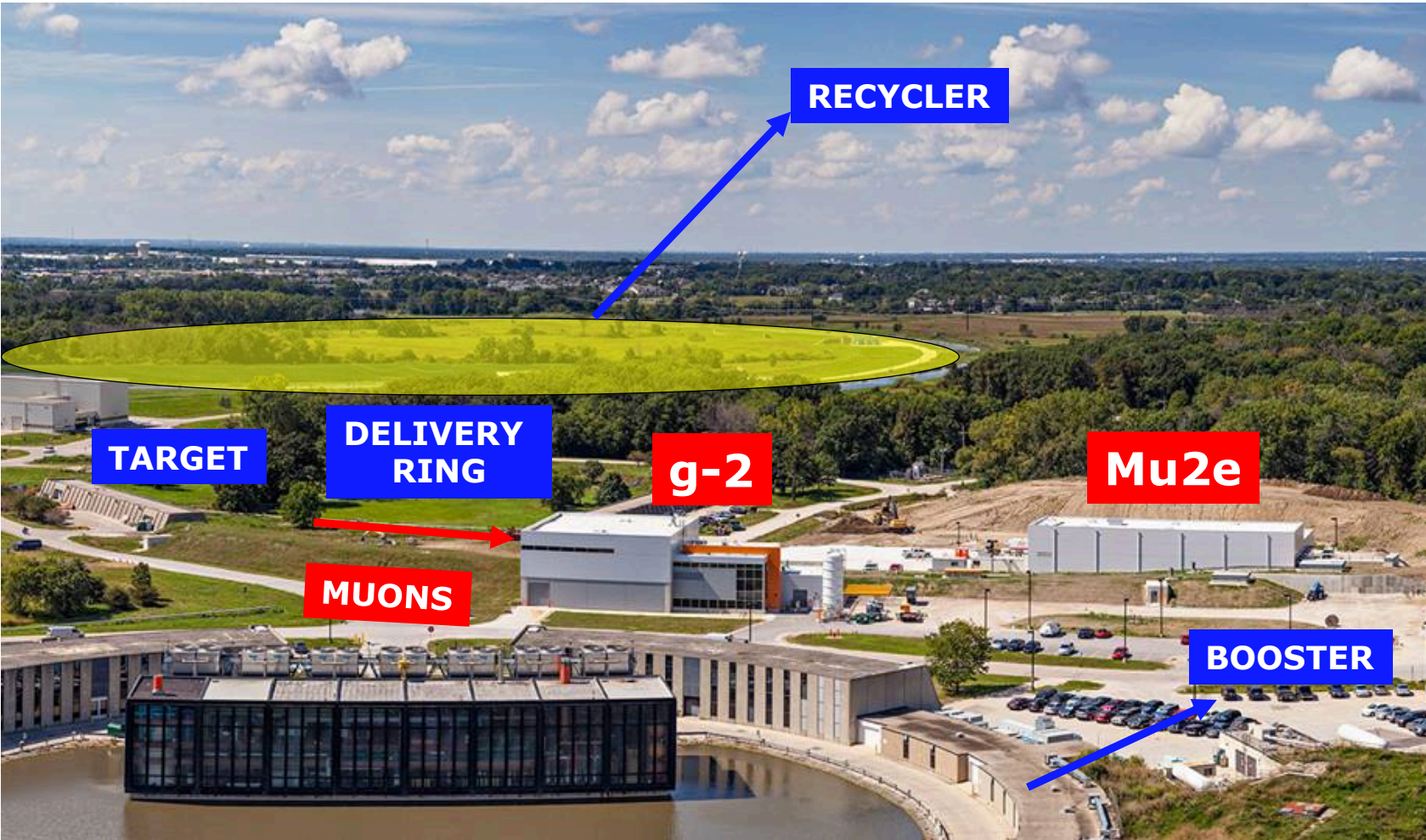
→ To conserve spin the muon spin must also be opposite to the momentum



The muons and protons separate as they go around the delivery ring



Accelerator Complex



Injection into the ring

The Fermilab accelerator provides a purer muon beam giving more muons at a lower instantaneous rate

- Inflector provides a field free region for the muons to enter the ring
- The Kickers kick the beam onto the right orbit (~11mrad)
- The electrostatic quadrupole magnets provide vertical focusing
But this field looks like a magnetic field to a moving particle :

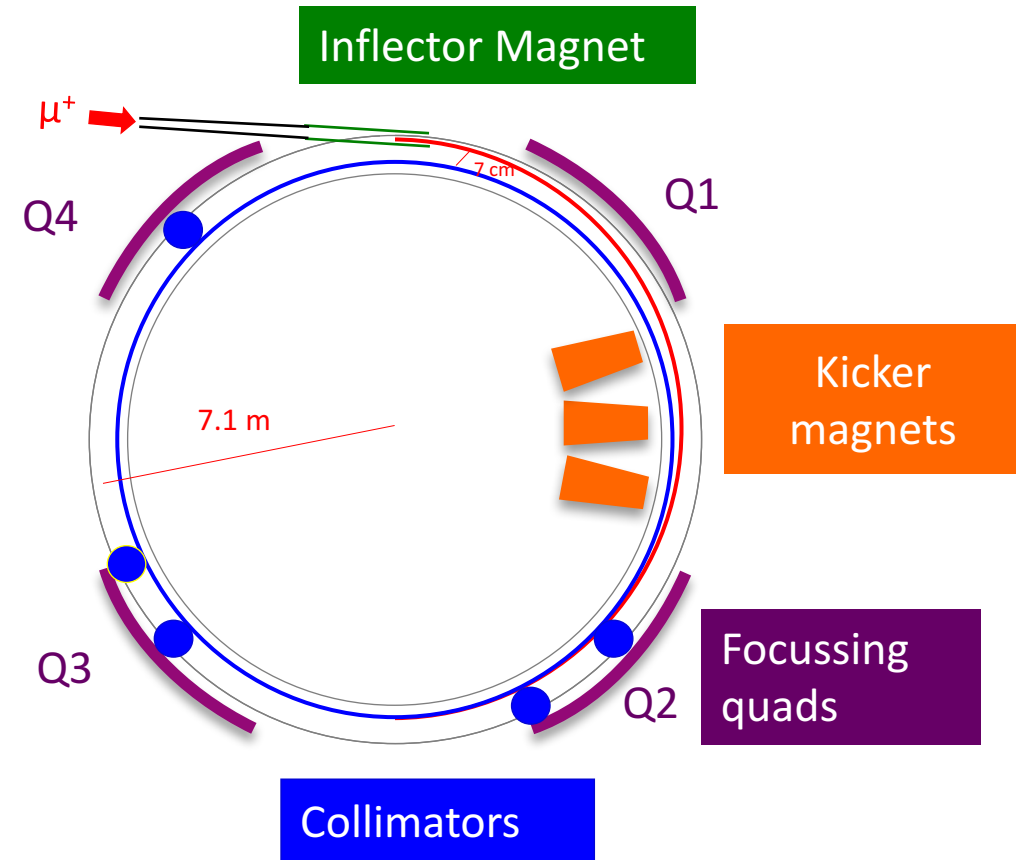
$$\vec{\omega}_a = \frac{e}{mc} \left[a_\mu \vec{B} - \left(a_\mu - \frac{1}{\gamma^2 - 1} \right) \vec{\beta} \times \vec{E} - a_\mu \left(\frac{\gamma}{\gamma + 1} \right) (\vec{\beta} \cdot \vec{B}) \vec{\beta} \right]$$

This term is cancelled by running at the magic momentum, $p = 3.094 \text{ GeV}$

Even so there are small effects :

- The muons aren't exactly at the magic momentum
- There is a small degree of vertical motion of the muons

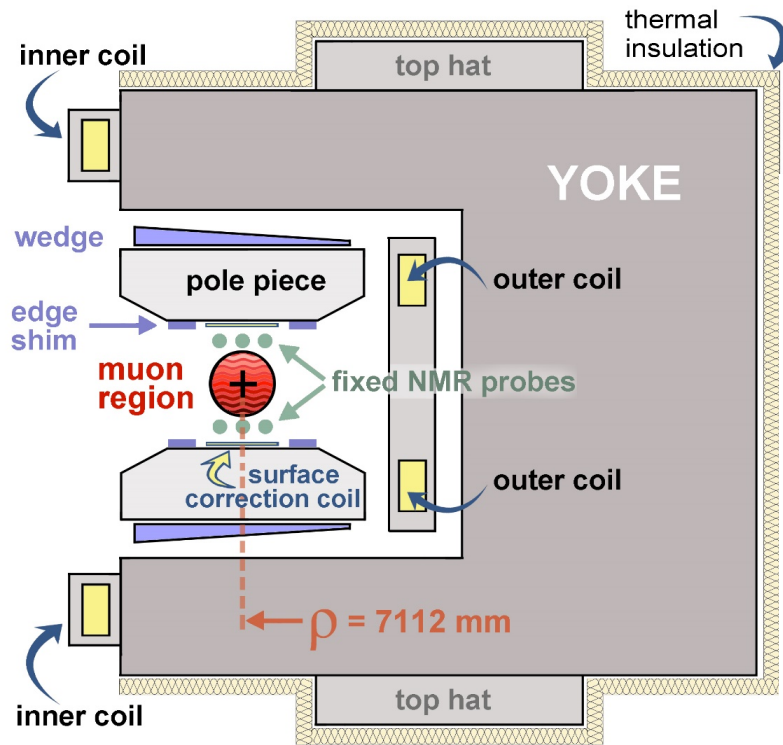
These small corrections can be calculated using the tracker and beam dynamics models



The magnetic field

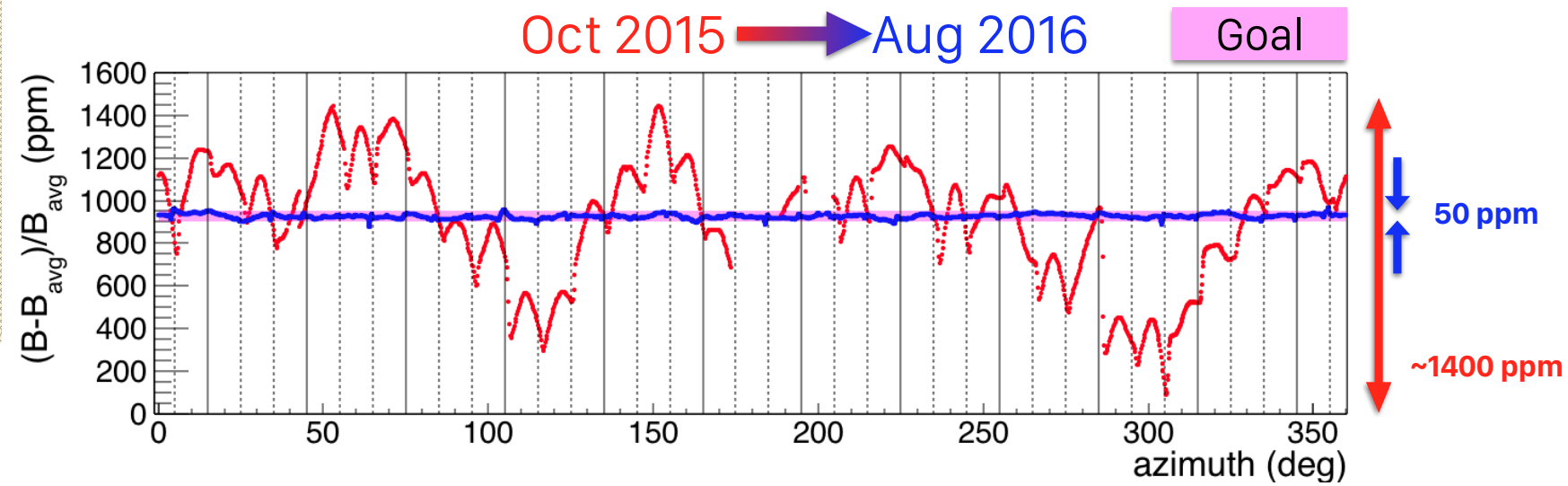
There are many knobs on the magnet that can be tuned to achieve excellent field uniformity

- 864 wedge shims, 48 top hats, 144 edge shims, 8424 laser cut iron foils
- 200 surface coils where the current can be adjusted



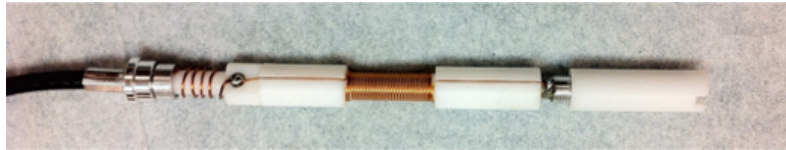
g-2 Magnet in Cross Section

Rough Shimming Results

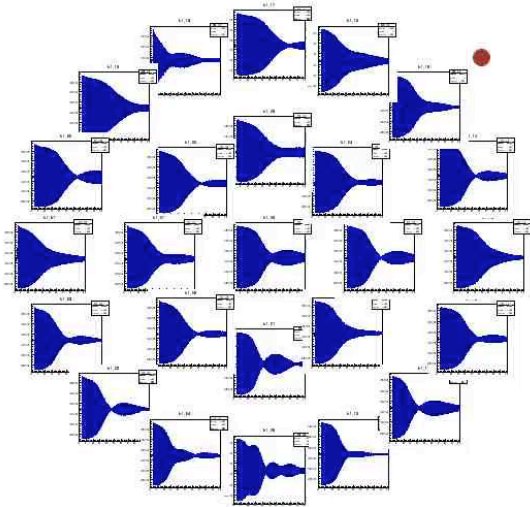
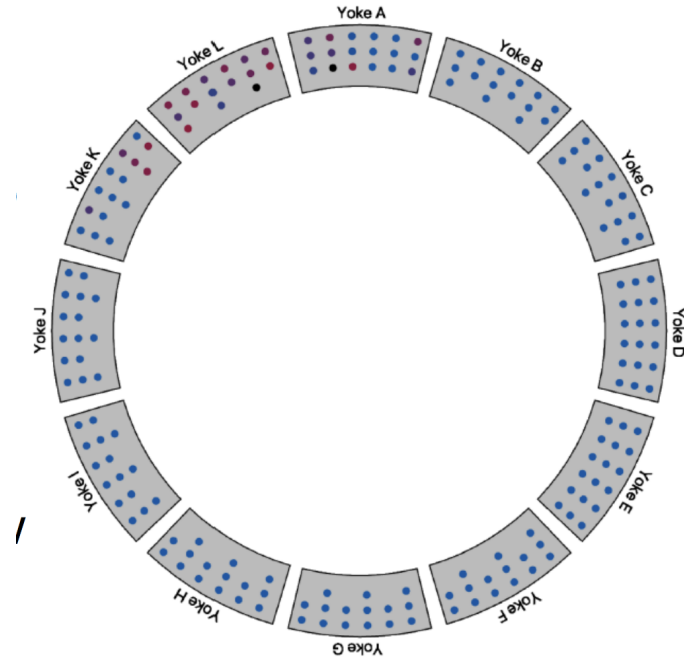
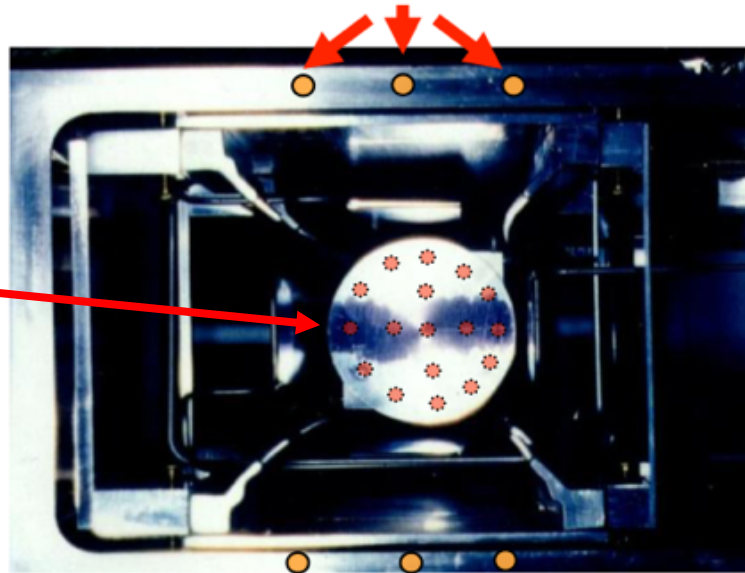


Measuring the Magnetic Field

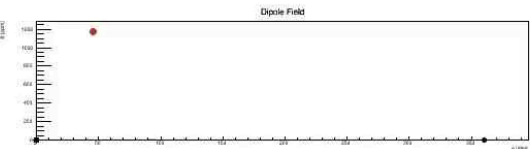
The magnetic field is measured using pulsed NMR probes (375 fixed probes, 17 probes on a trolley, plunging probe)



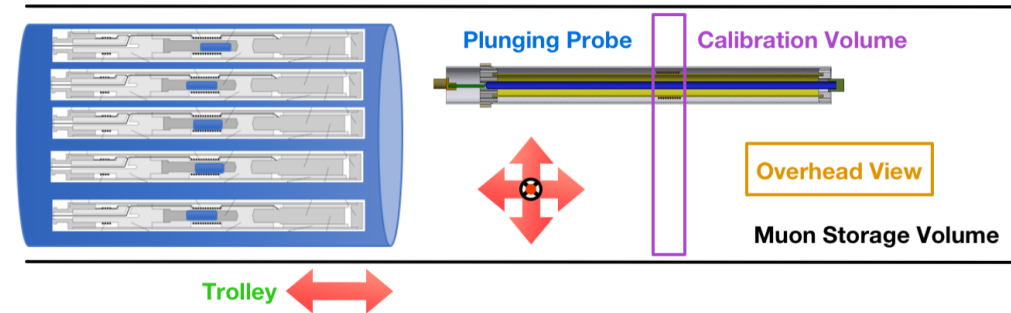
Fixed probes measure the magnetic field all the time outside the storage region



Trolley probes measure the magnetic field in the storage region during special trolley runs

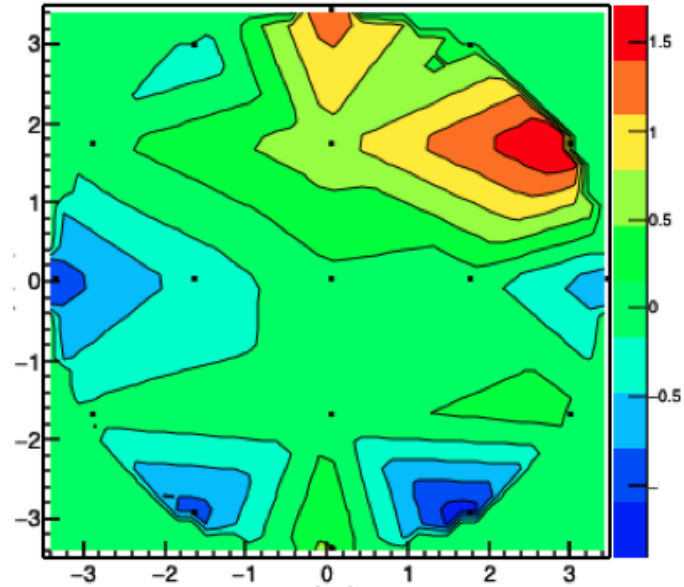
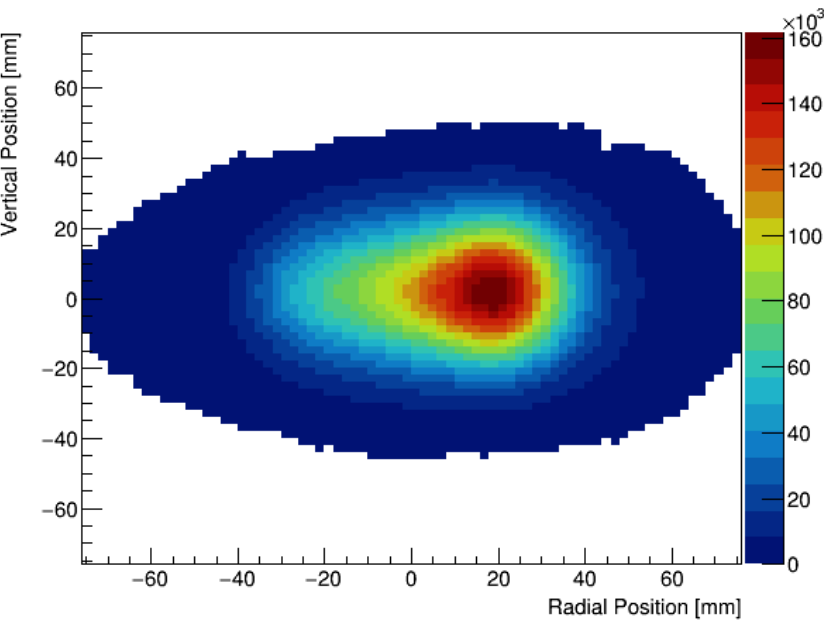


A plunging probe is used for calibration



Magnetic field measurement

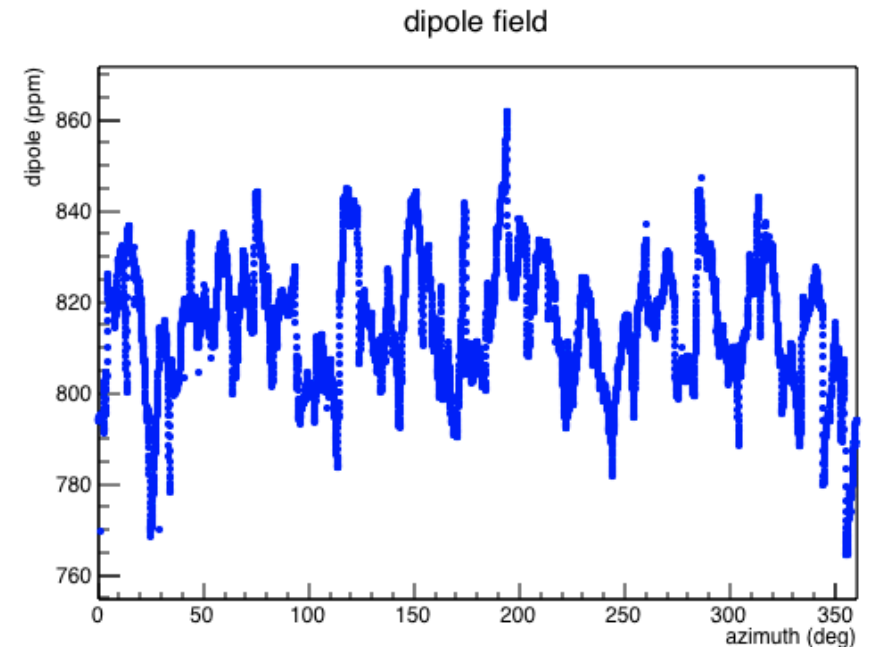
The magnetic field must be convoluted with the muon distribution to calculate the final result



Need to know the field that the muon has experienced at the point of decay

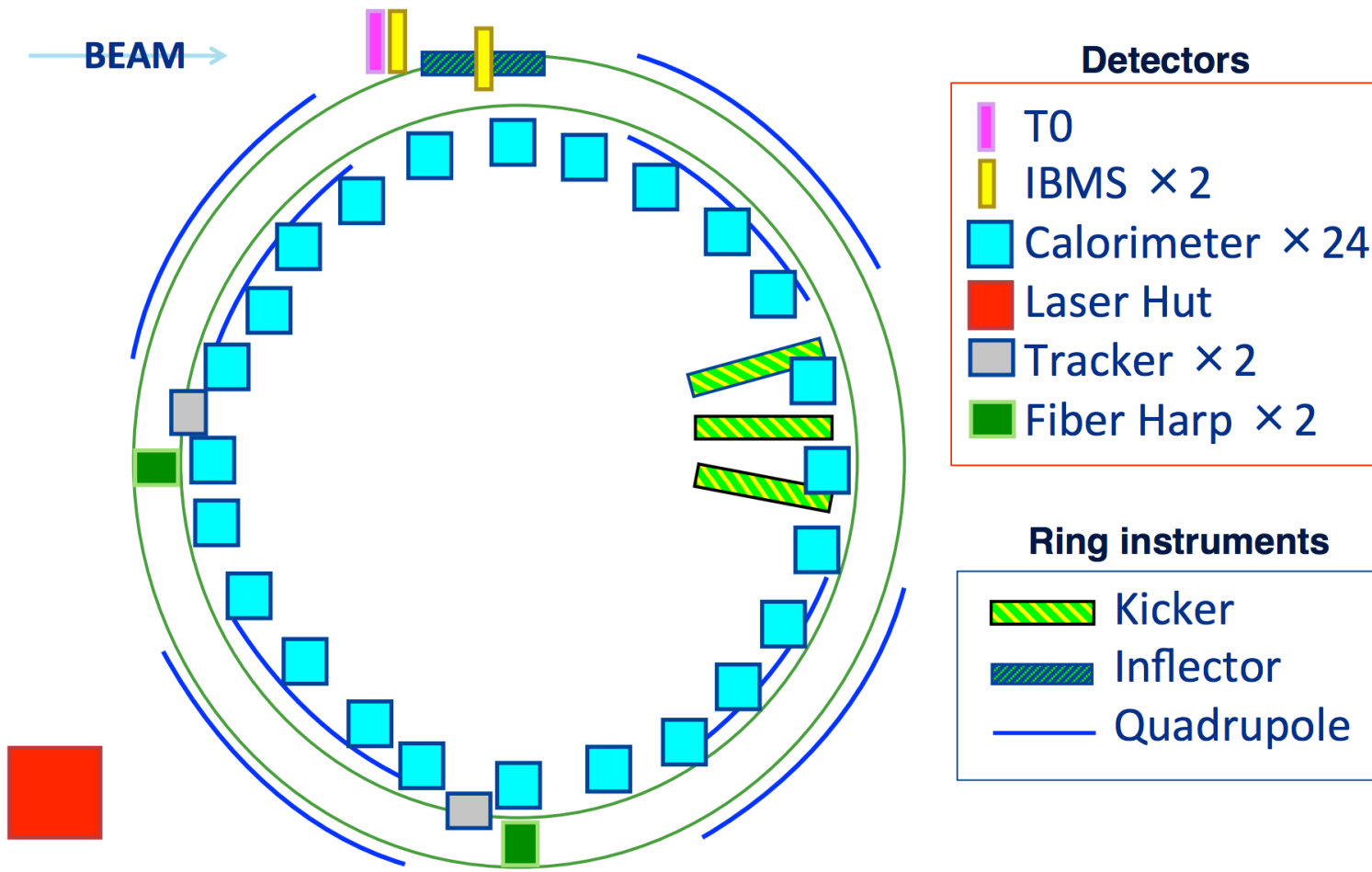
Over the past year of running trolley runs were conducted \sim every 3 days at varying times of the day

The surface coils have been adjusted to minimise the multipole moments

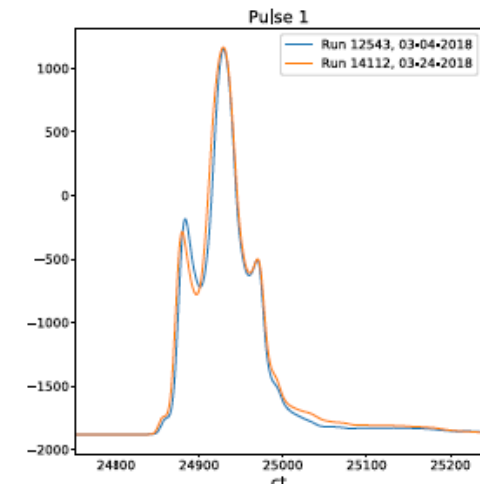


G-2 detector systems

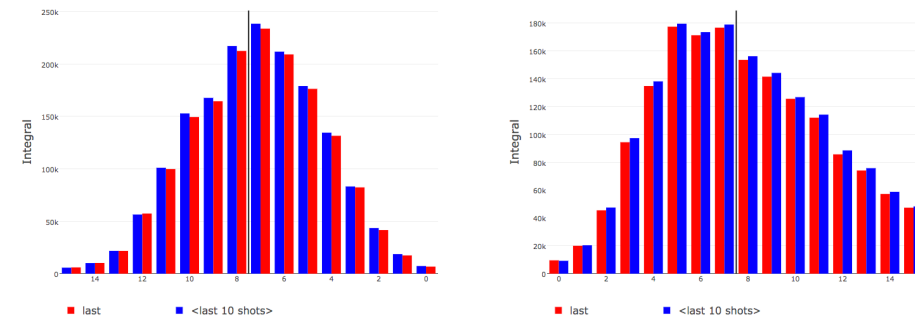
The different detector systems measure the precession frequency and monitor the beam distribution



T0 detector measures the beam arrival time and the temporal distribution :

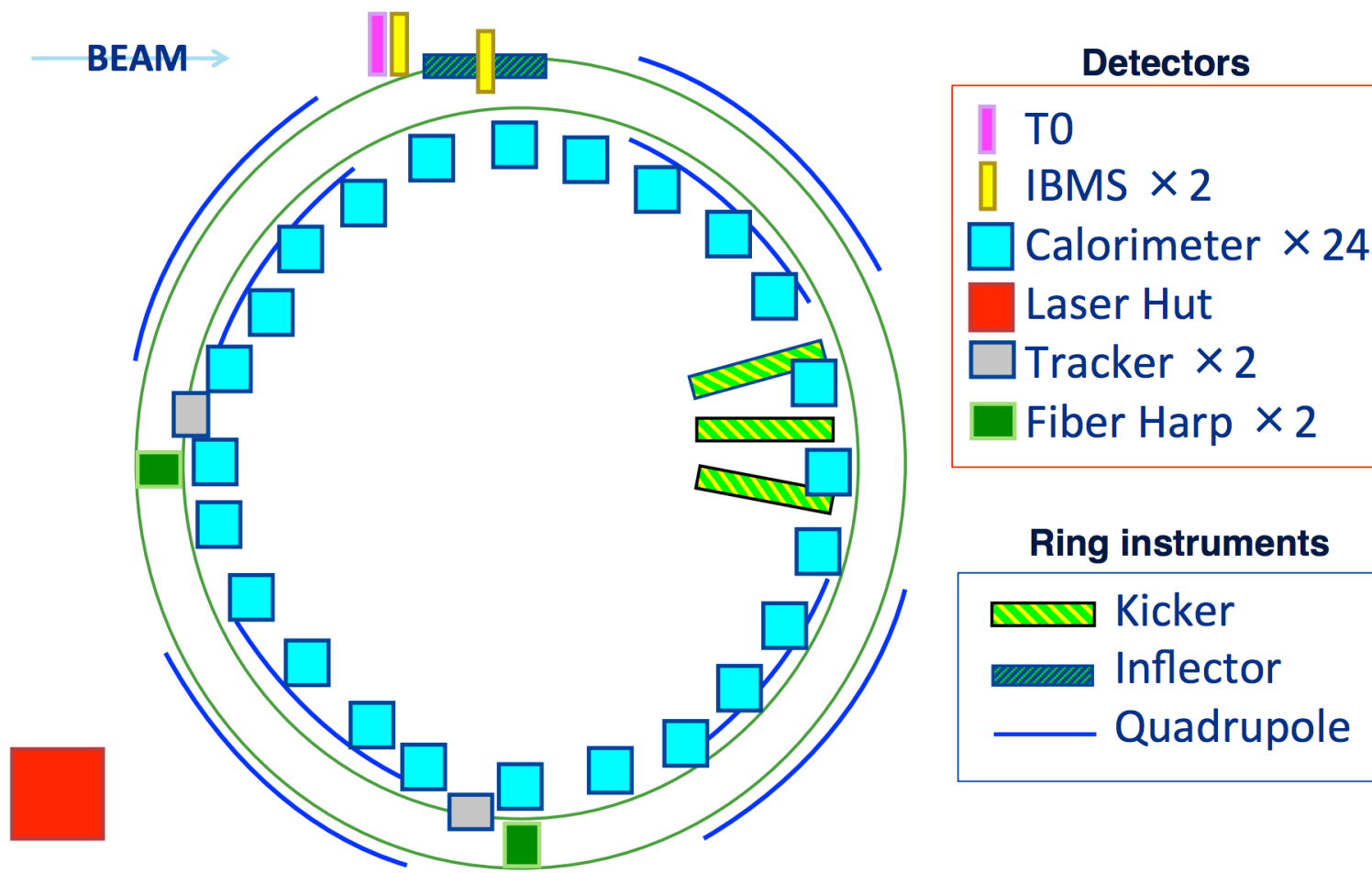


The IBMS measures the horizontal and vertical distributions on entry :

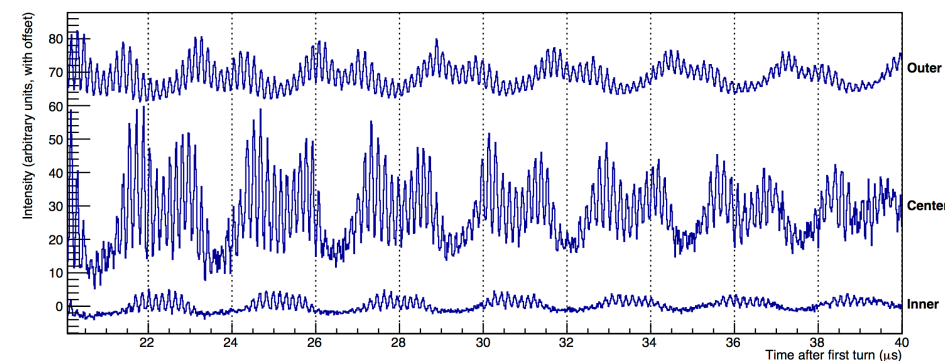
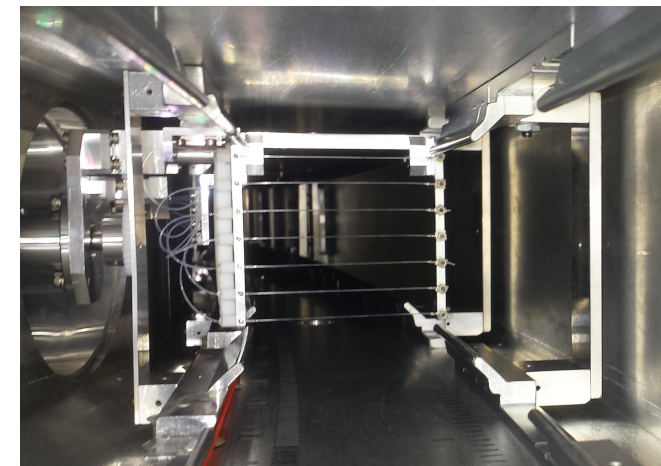


G-2 detector systems

The different detector systems measure the precession frequency and monitor the beam distribution

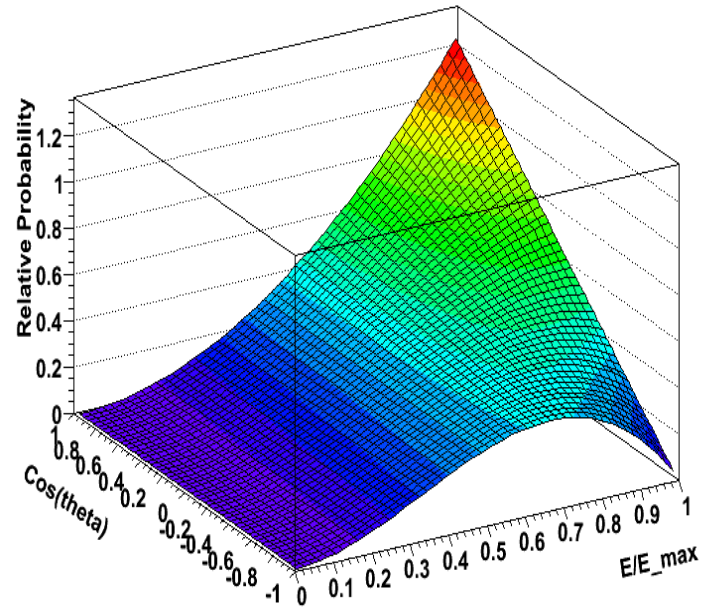


The fibre harps slide in to the beam to make a destructive measurement of the beam profile

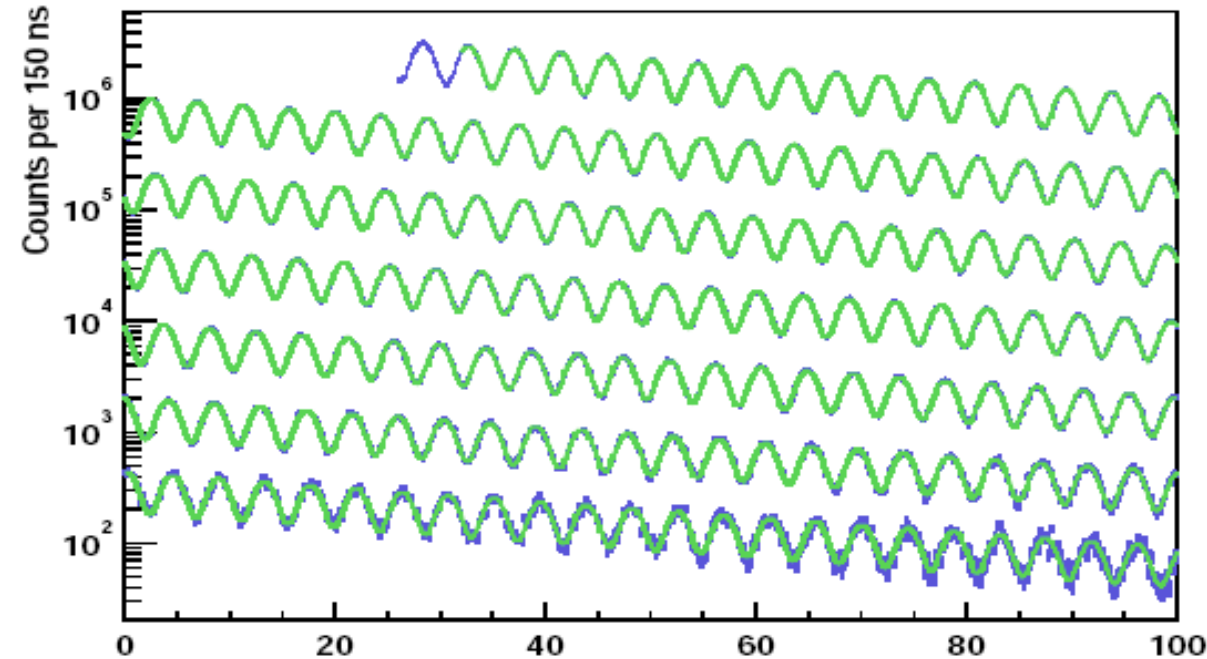


Measuring the precession frequency

The highest energy positrons from the muon decay are preferentially released along the direction of the muon spin



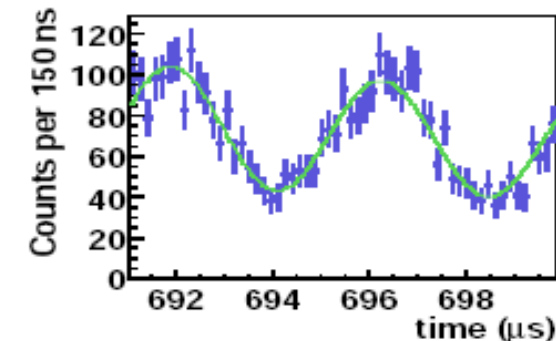
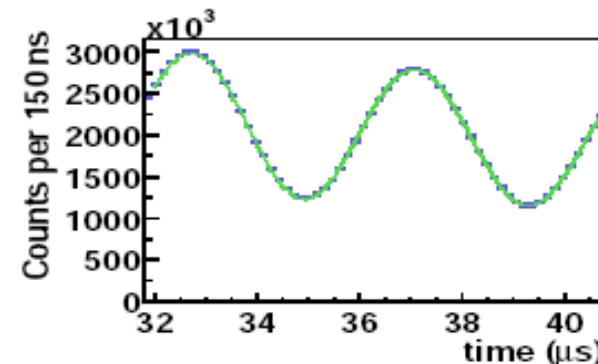
Count the number of positrons hitting the calorimeter above 1.8 GeV as a function of time



- Number oscillation due to spin precession
- Exponential decay as the number of stored muons decreases

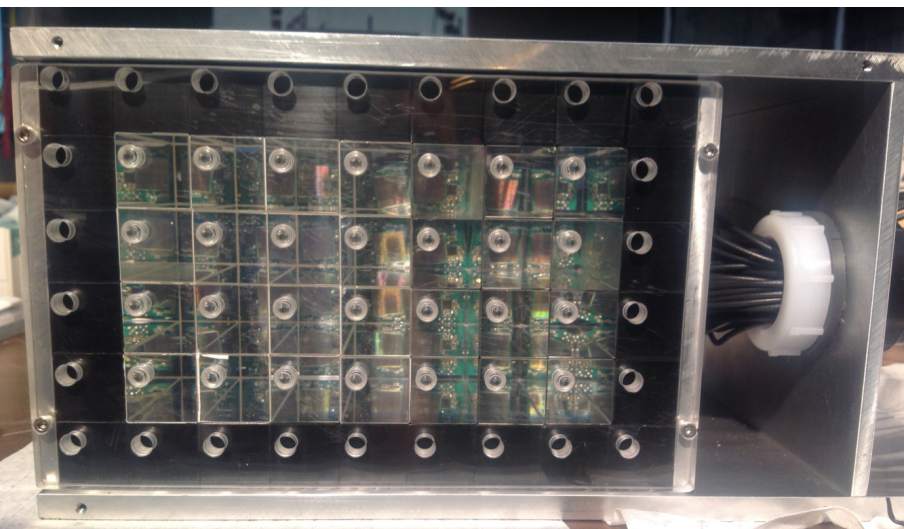
The precession frequency is extracted from a fit to the data:

$$N_e(t) \simeq N_0 e^{-\frac{t}{\gamma\tau}} [1 - A \cos(\omega_a t + \phi_a)]$$

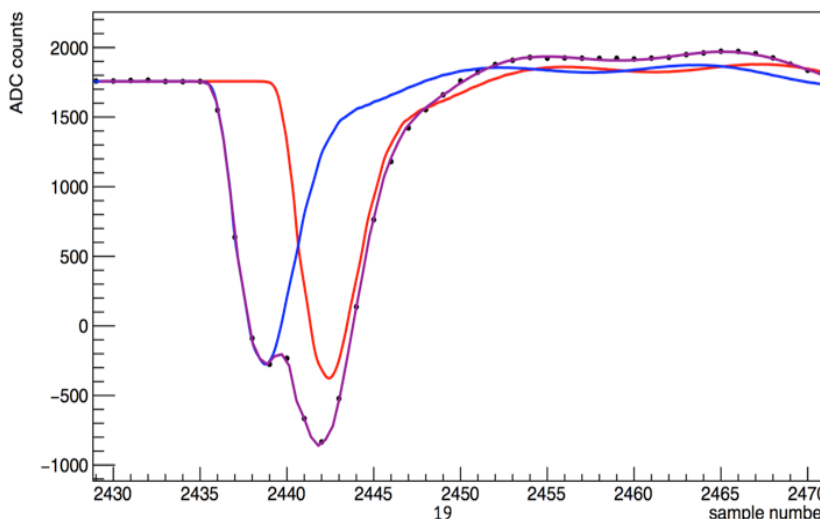
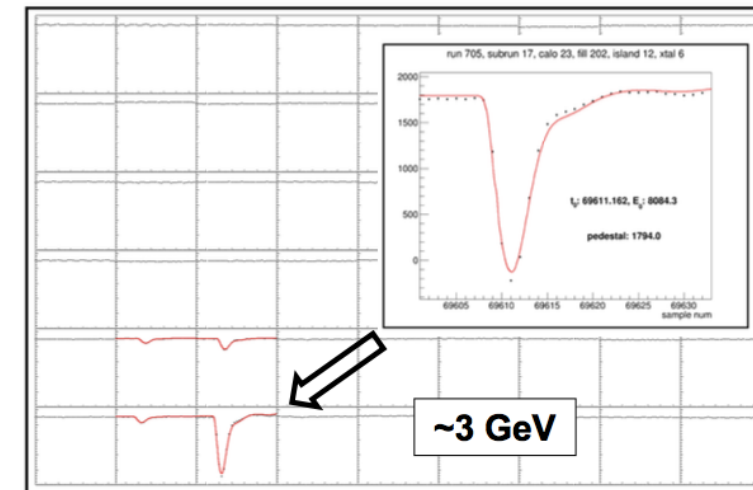


Calorimeters

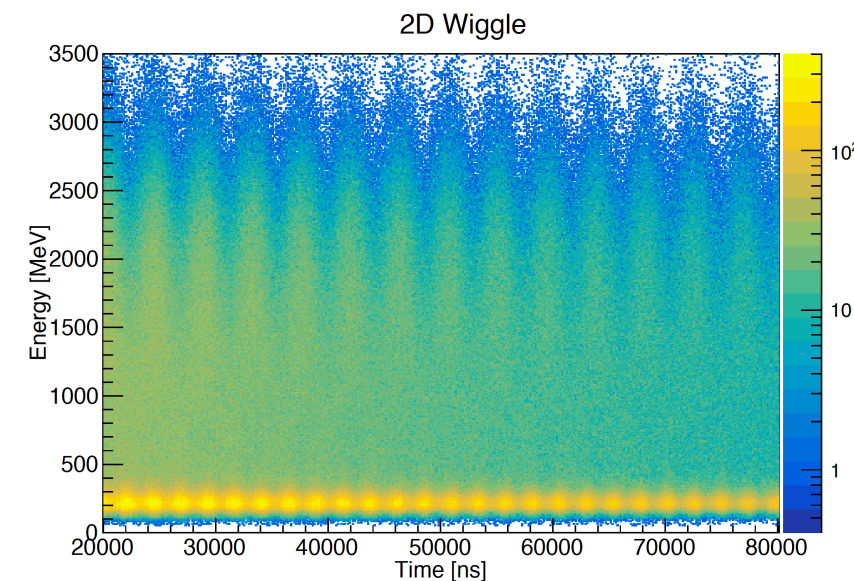
There are 24 calorimeters around the centre of the ring to measure the positrons from the muon decays



- 6 x 9 lead fluoride crystals
- Requirements :
 - Resolve pulses separated by more than 5ns
 - Better than 5% energy resolution
 - Time accurate to 100ps
 - Stable gain during a fill

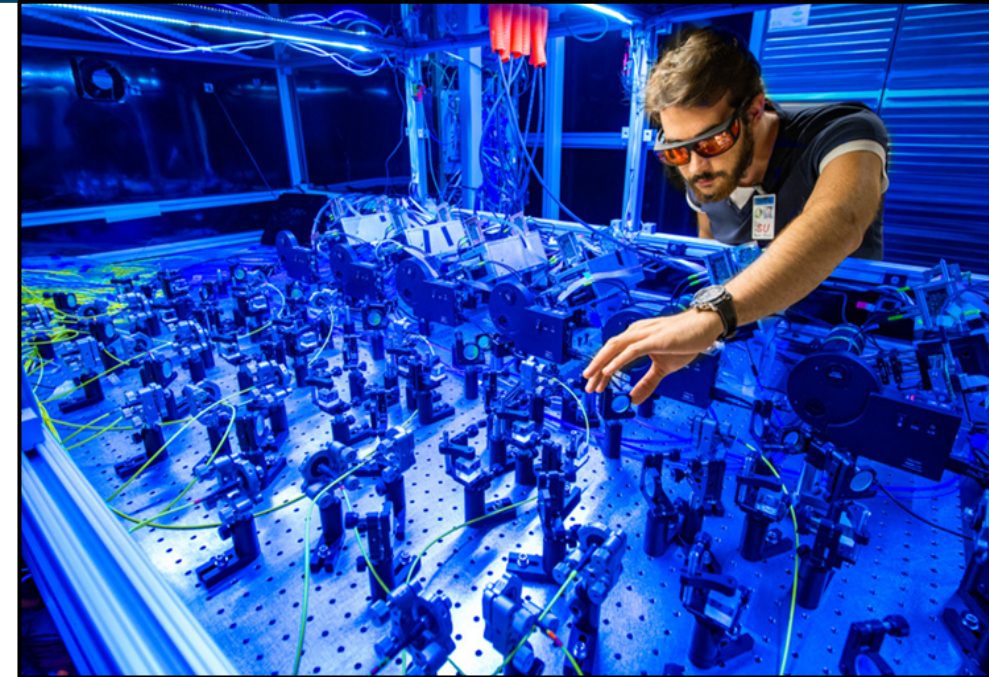


The calorimeters performed well during the run meeting all requirements



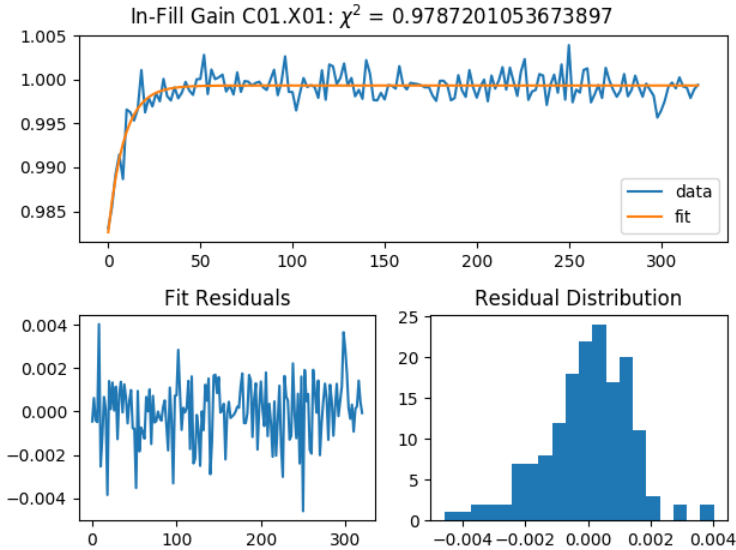
Laser calibration system

The laser calibration system allows any gain variations over time to be calibrated out



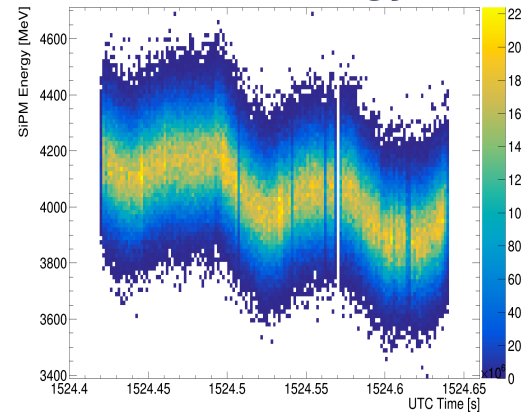
Sends laser pulses to every calorimeter both in and out of fill

Allows for both long and short term gain corrections

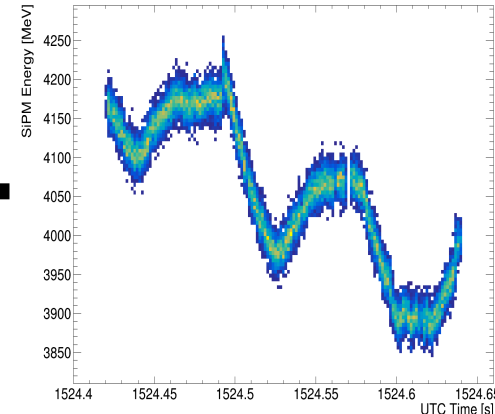


Performed well achieving gain stability of 0.04%

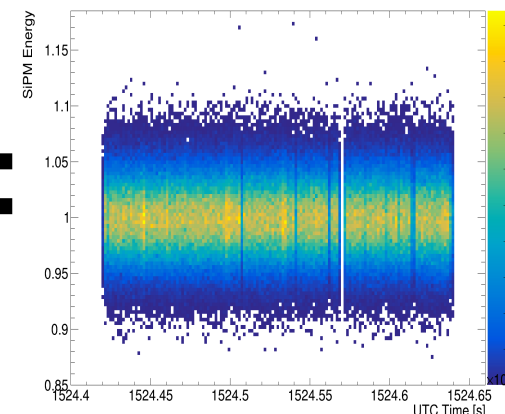
Raw SiPM energy



Out of fill correction



Corrected energy

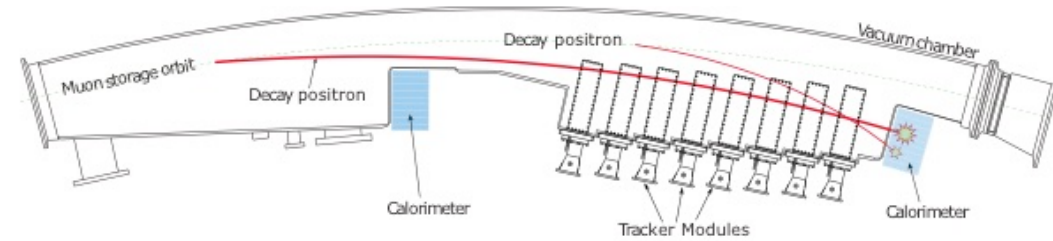


Tracking detectors

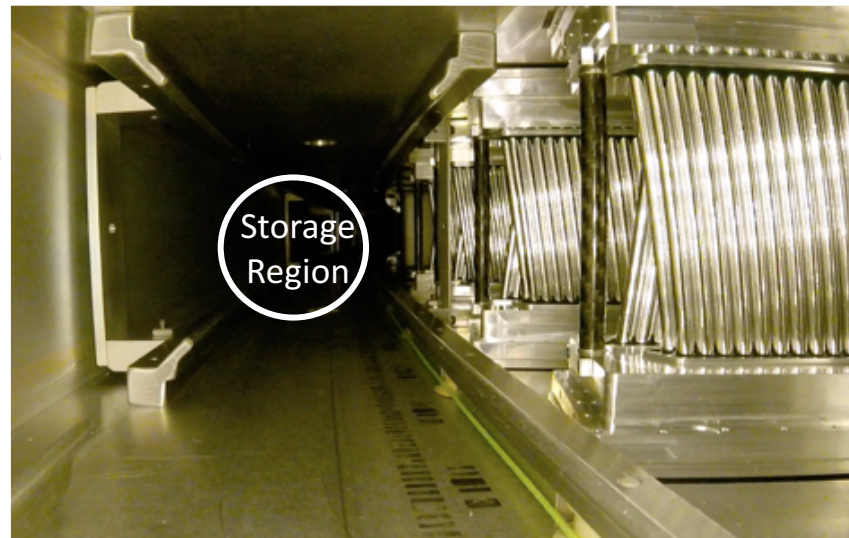
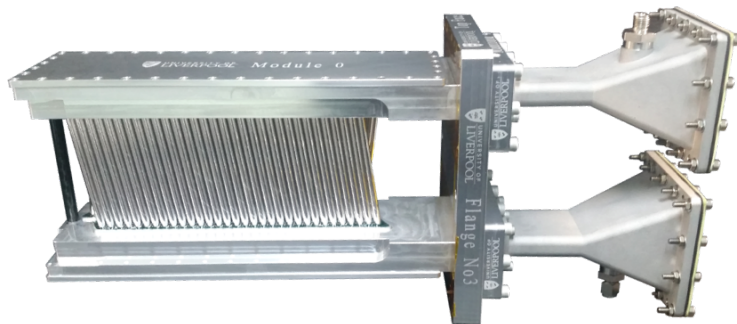
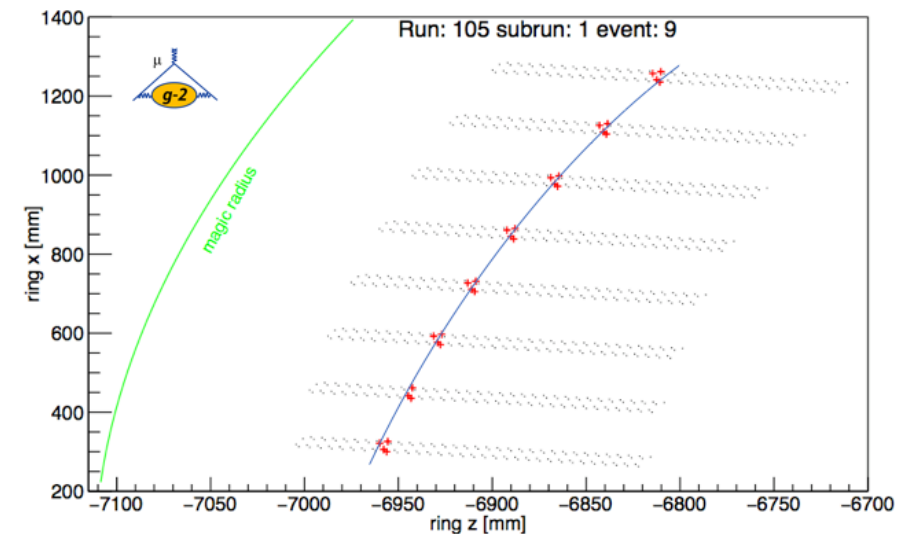
Tracking detectors are located at two locations around the ring allowing reconstruction and traceback of positron tracks

Aims :

- Measure the beam profile in multiple locations around the ring as a function of time
- Calibration and acceptance of the calorimeters
 - Pile up, gain, lost muons
- Measure or set a limit on a muon EDM



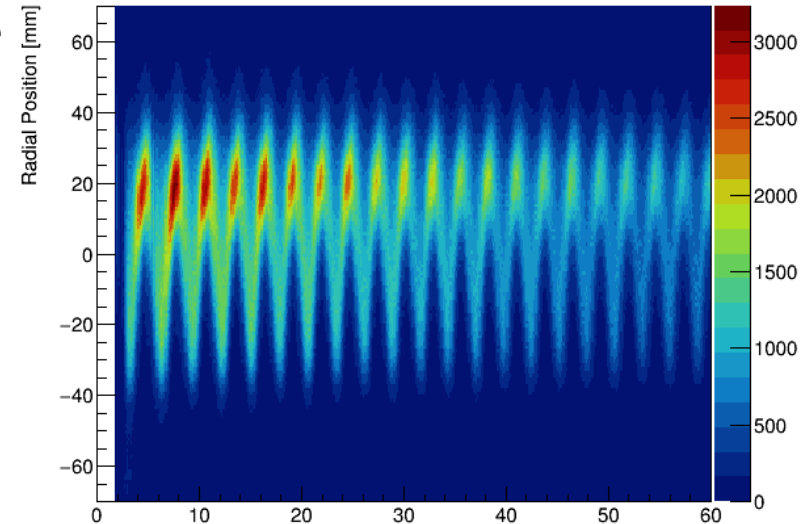
One of the first reconstructed tracks :



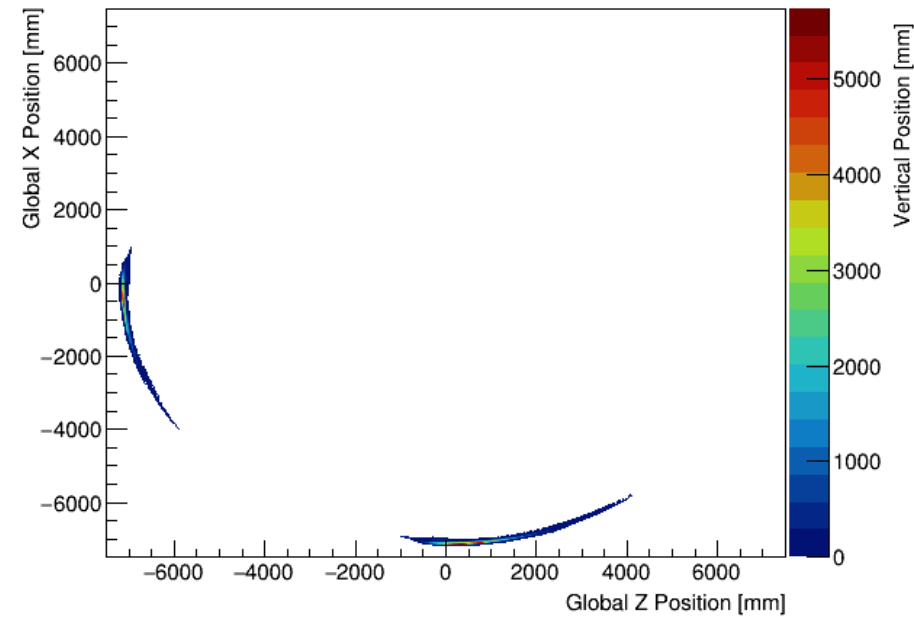
Tracker performance

The trackers performed very well providing a good insight into the beam movements

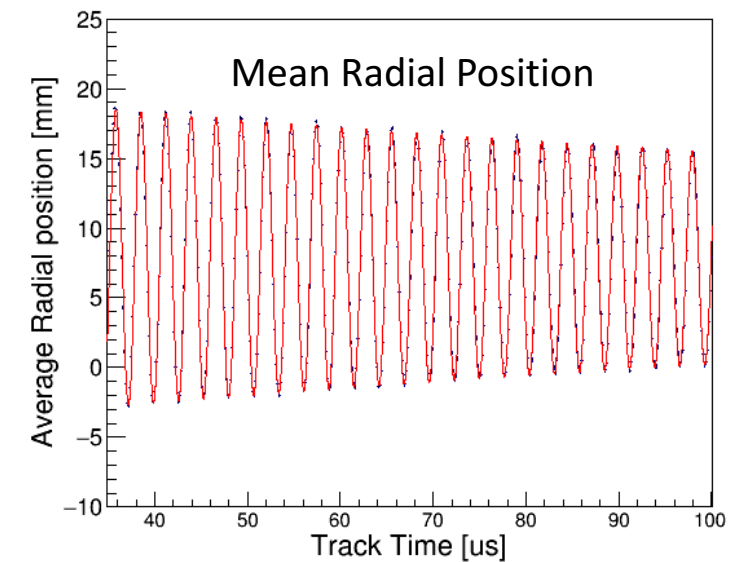
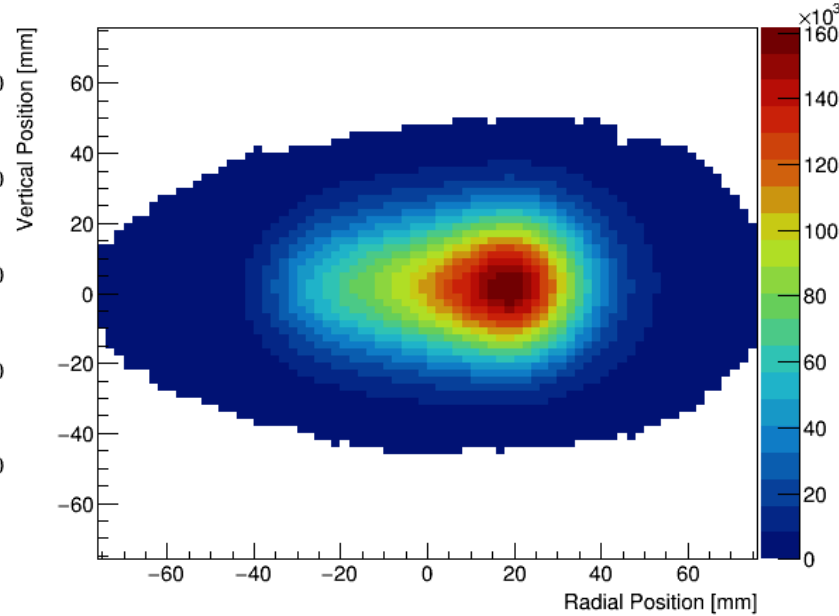
Radial position vs. time



Decay Vertices



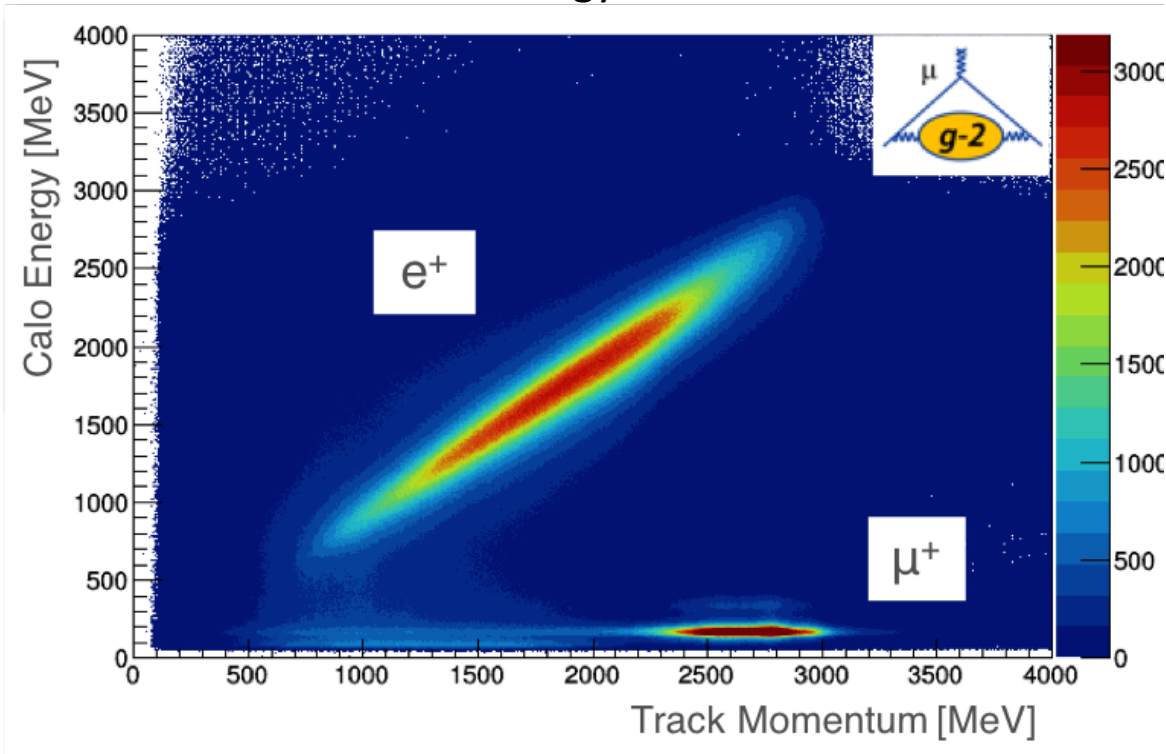
Reconstructed beam distribution



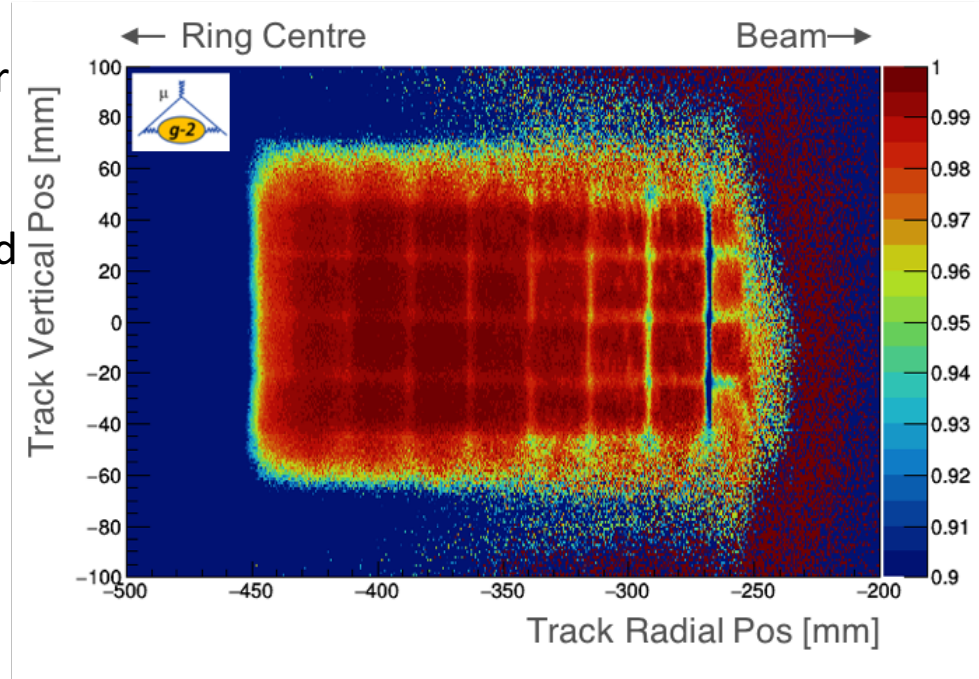
Tracker performance

The trackers are located in front of 2 of the calorimeters so can be used for systematic checks in terms of gain and pile up

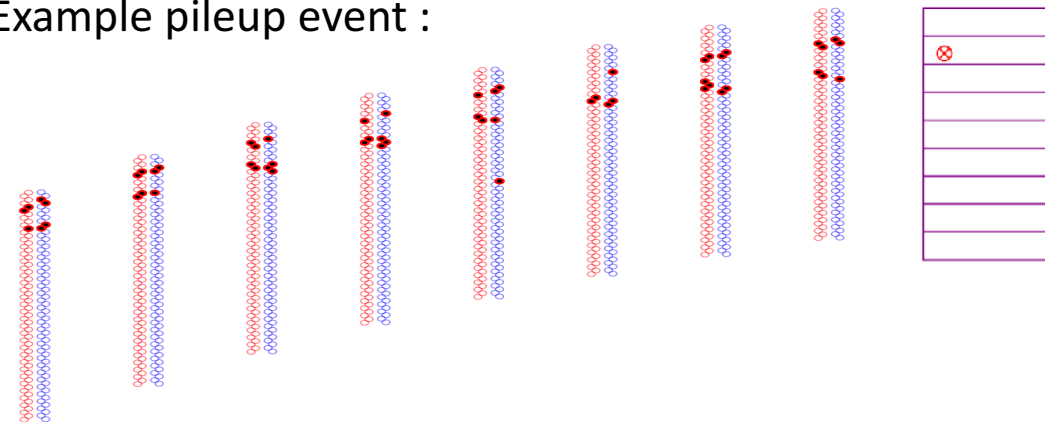
Calorimeter energy vs. track momentum



Calorimeter efficiency (based on extrapolated tracks)



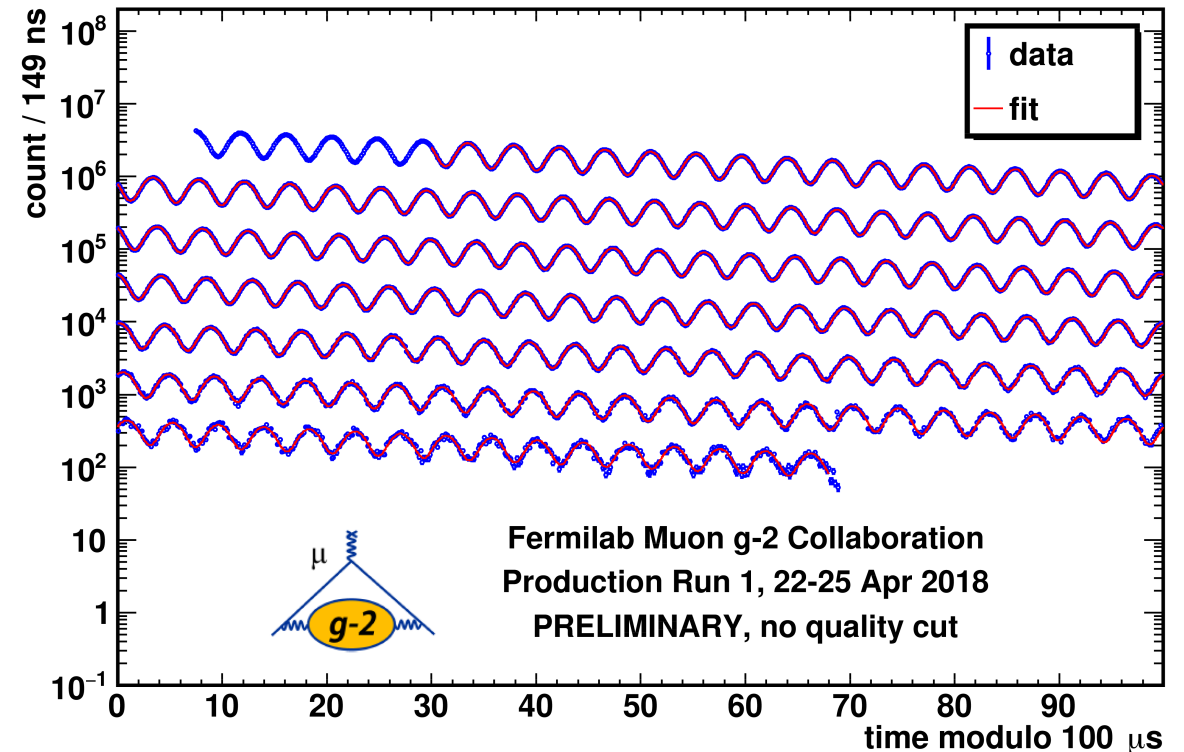
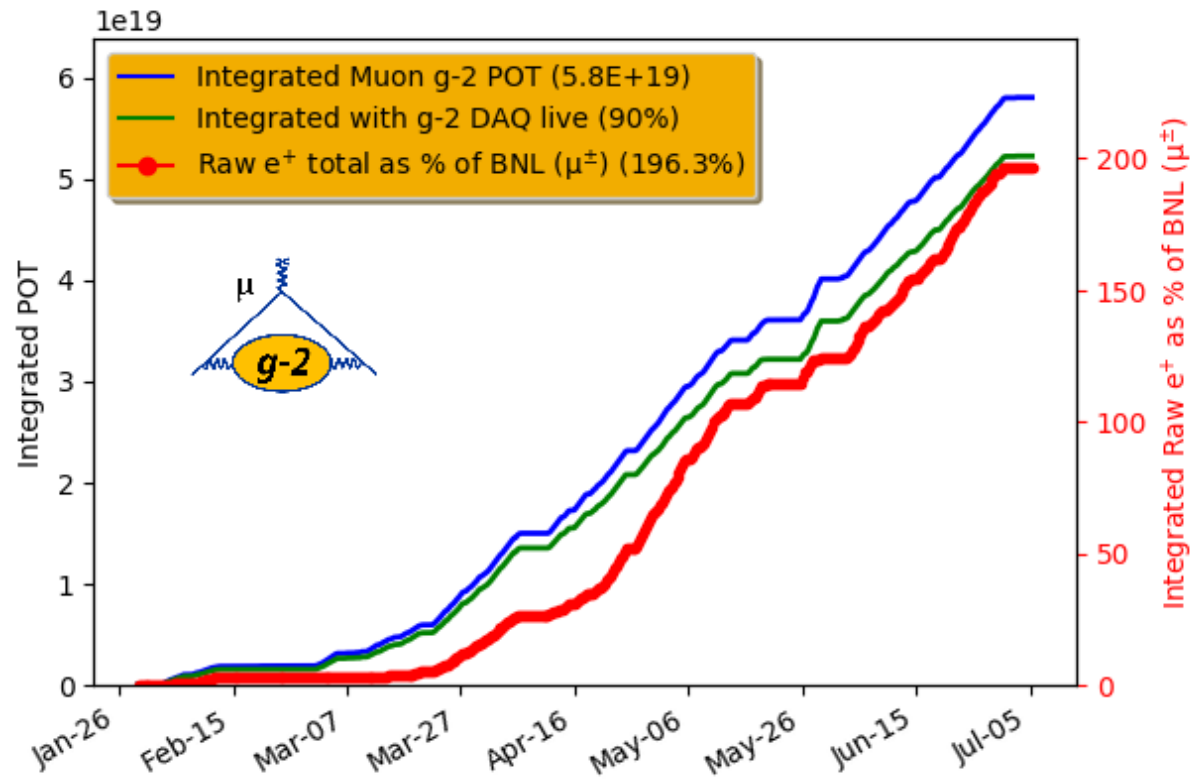
Example pileup event :



Where we are now

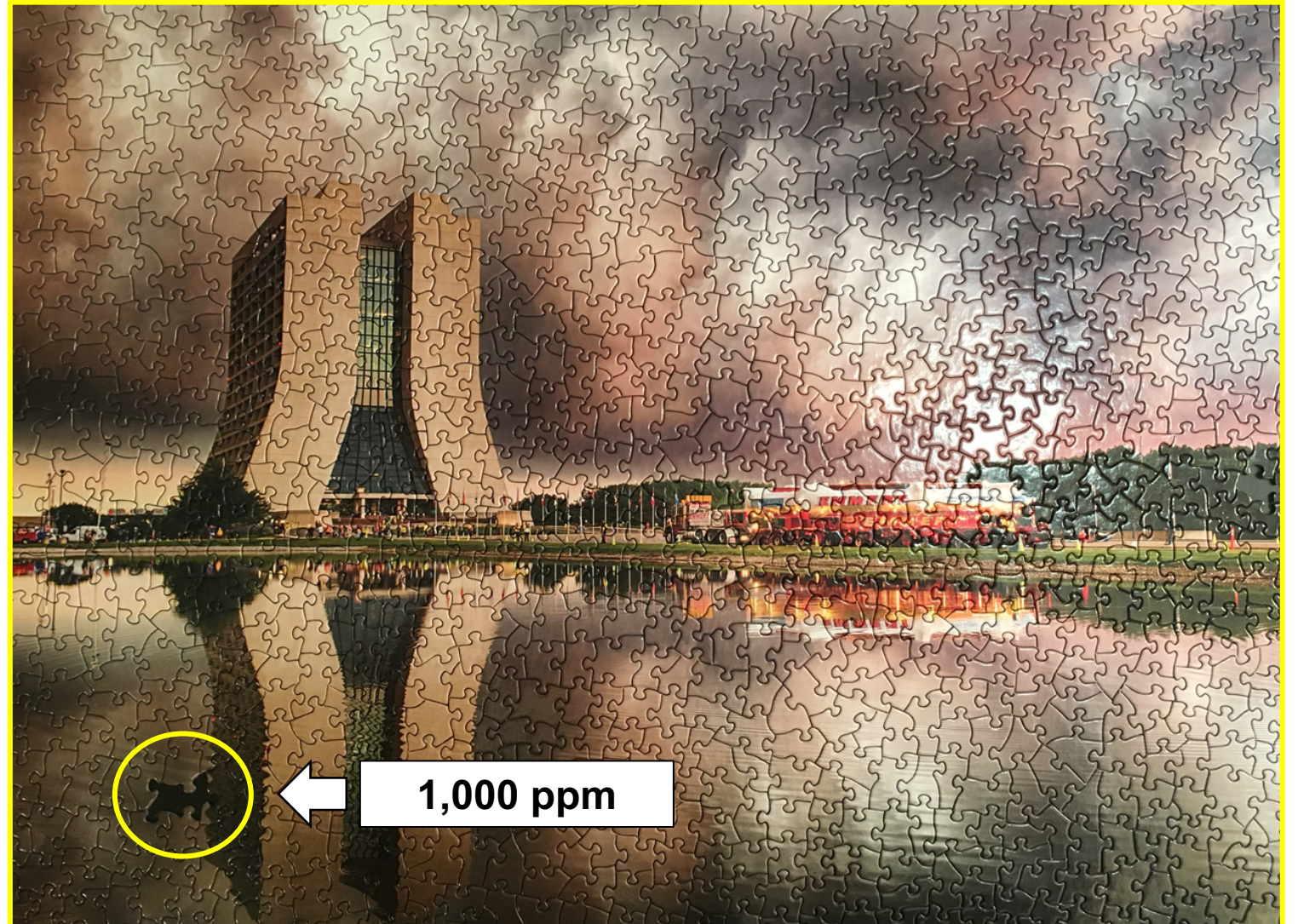
The first year of data taking finished in July collecting nearly 2 x BNL statistics

- First physics run complete, the analysis is underway expecting a physics result in spring next year
- Next run starts in October and will run through to July
- Currently systems are being upgraded to increase reliability, uptime and muon storage



What we are aiming for

To put the precision into context consider this 1000 piece jigsaw with 1 missing piece...



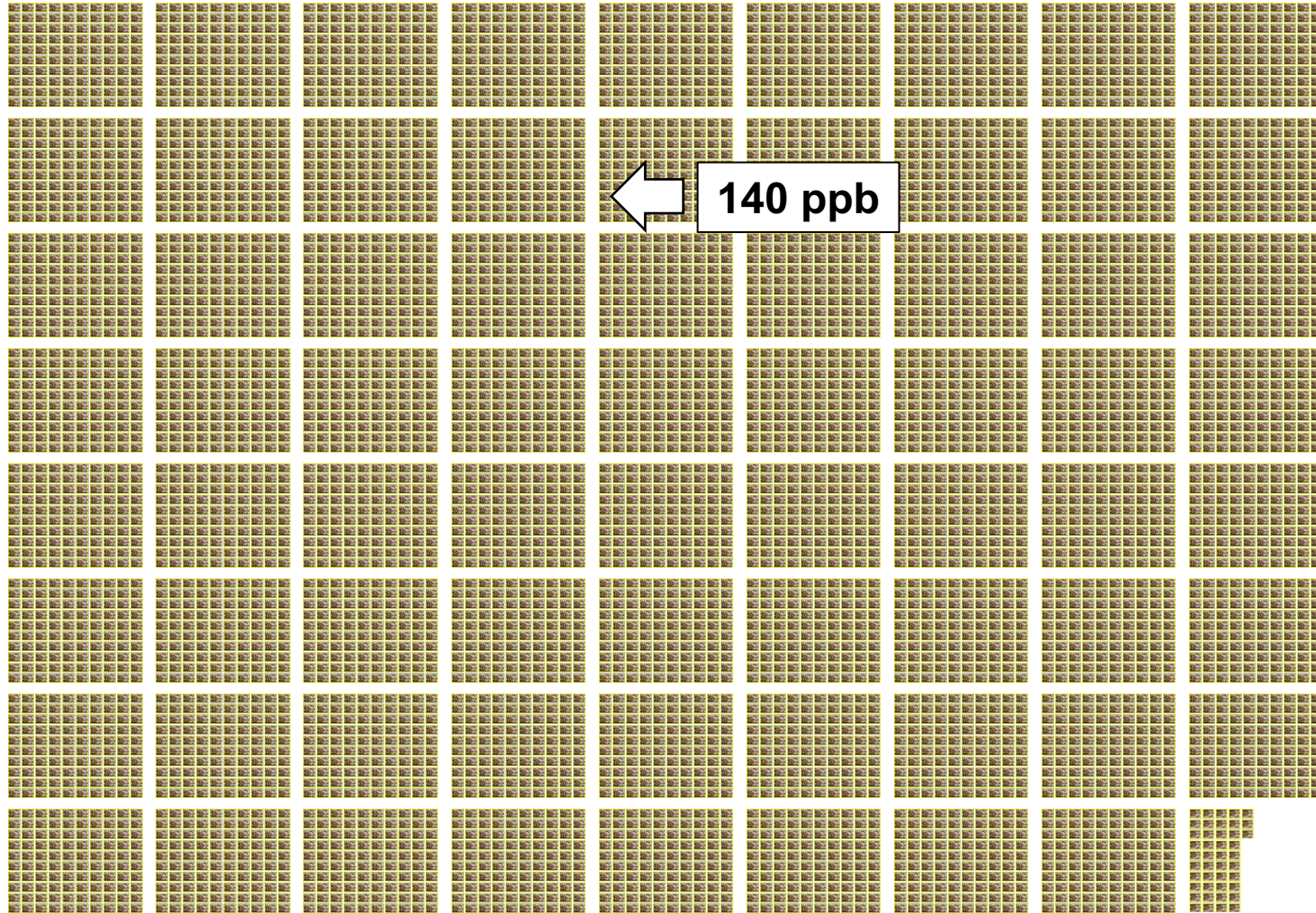
What we are aiming for



Consider 100 jigsaw puzzles with only one missing piece

Similar precision to the
CERN III experiment (1976)

What we are aiming for



7143 jigsaw puzzles with one missing piece

Lose one piece
→ 140ppb (Fermilab aim)

Every detail counts!

Aside : EDM

The g-2 experiment at Fermilab can also look for a potential muon EDM

Fundamental particles can also have an EDM defined by an equation similar to the MDM:

$$\vec{d} = \eta \frac{Qe}{2mc} \vec{s} \quad \vec{\mu} = g \frac{e}{2mc} \vec{s}$$

Defined by the Hamiltonian: $H = -\vec{\mu} \cdot \vec{B} - \vec{d} \cdot \vec{E}$

	E	B	μ or d
P	-	+	+
C	-	-	-
T	+	-	-

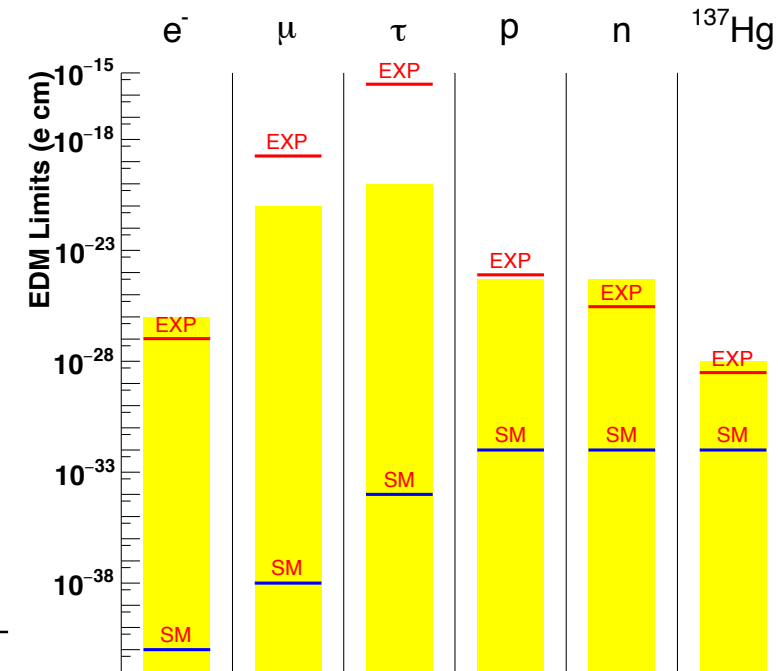
Provides an additional source of CP violation

Standard scaling : $\frac{d_\mu}{d_e} \sim \frac{m_\mu}{m_e}$

d_e limits imply d_μ scale of $10^{-25} \text{ e}\cdot\text{cm}$

But some BSM models predict non-standard scalings

(quadratic or even cubic)



The muon is a unique opportunity to search for an EDM in the 2nd generation

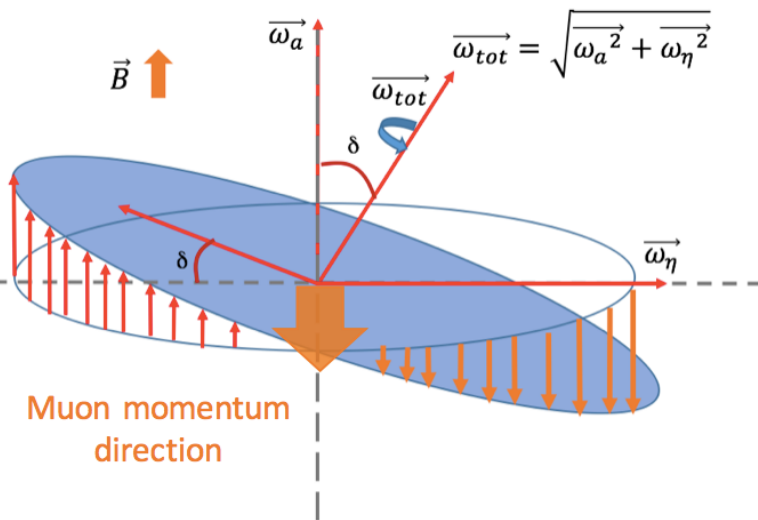
Aside : EDM

If an EDM is present the spin equation is modified to:

$$\vec{\omega}_{a\eta} = \vec{\omega}_a + \vec{\omega}_\eta = \underbrace{-\frac{Qe}{m} a \vec{B}}_{\text{MDM}} - \eta \underbrace{\frac{Qe}{2m} \left[\frac{\vec{E}}{c} + \vec{\beta} \times \vec{B} \right]}_{\text{EDM}}$$

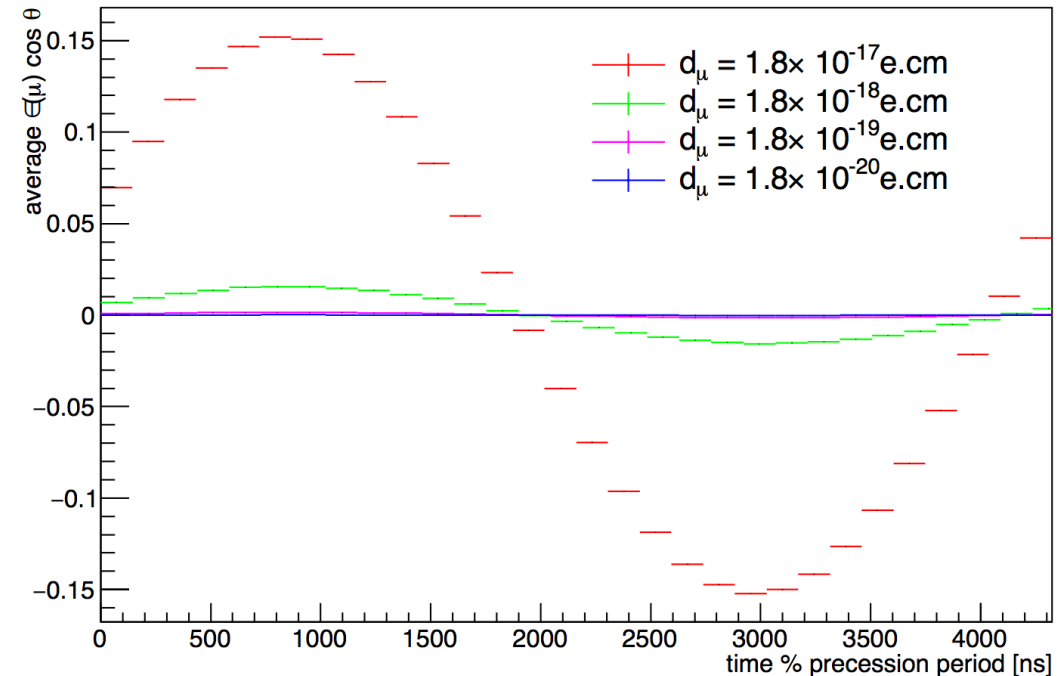
Dominant term

An EDM tilts the precession plane towards the centre of the ring
 → Vertical oscillation ($\pi/2$ out of phase)



Expect tilt of \sim mrad
 for $d_\mu \sim 10^{-19}$

An EDM also increases the
 precession frequency



Should reach BNL sensitivity in a few weeks (\sim 1 million tracks)
 Expect to reach 10^{-21} by the end of the experiment (several billion tracks)

Aside : JPARC g-2

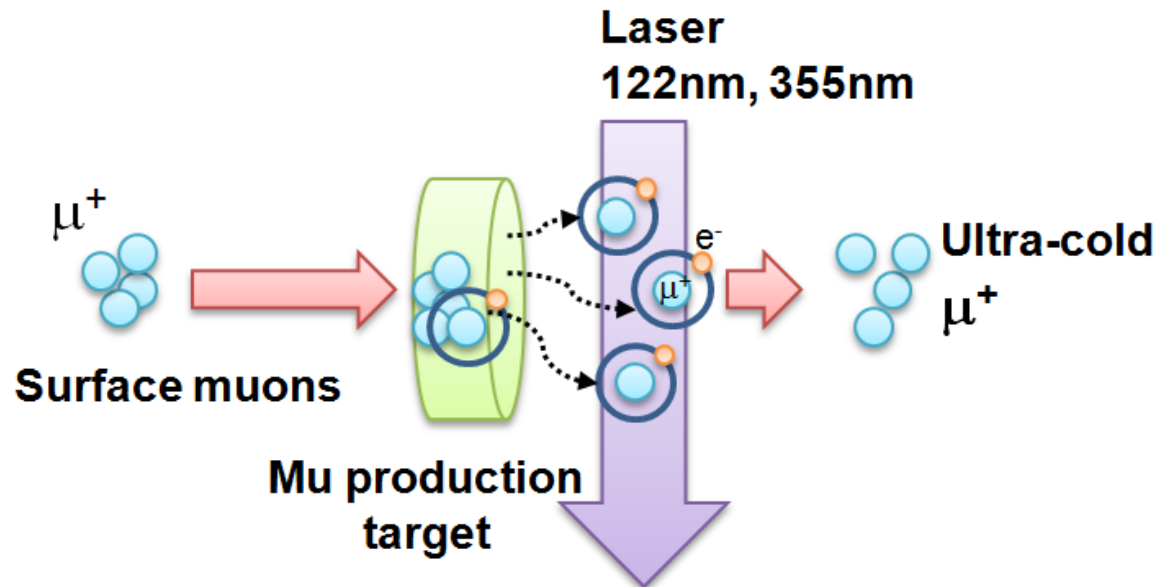
There is also an experiment to measure the muon g-2 at JPARC which uses a different technique

$$\vec{\omega}_a = \frac{e}{mc} \left[a_\mu \vec{B} - \left(a_\mu - \frac{1}{\gamma^2 - 1} \right) \vec{\beta} \times \vec{E} - a_\mu \left(\frac{\gamma}{\gamma + 1} \right) \left(\vec{\beta} \cdot \vec{B} \right) \vec{\beta} \right]$$

The Fermilab experiment cancels this term by running at the magic momentum

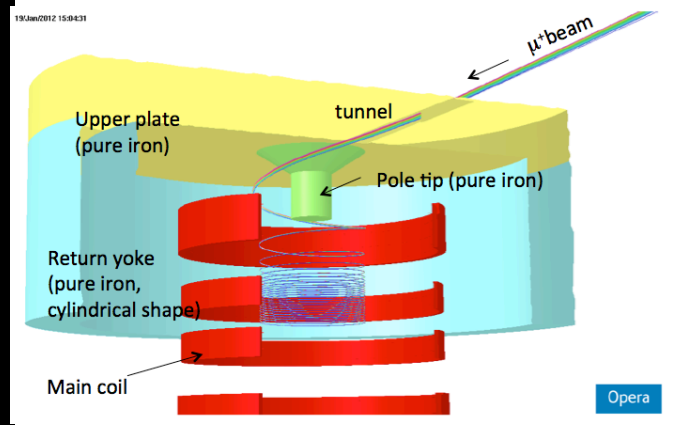
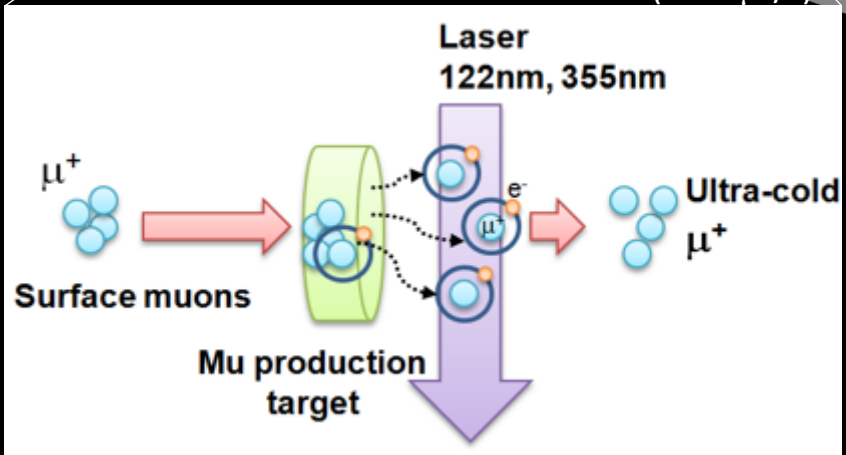
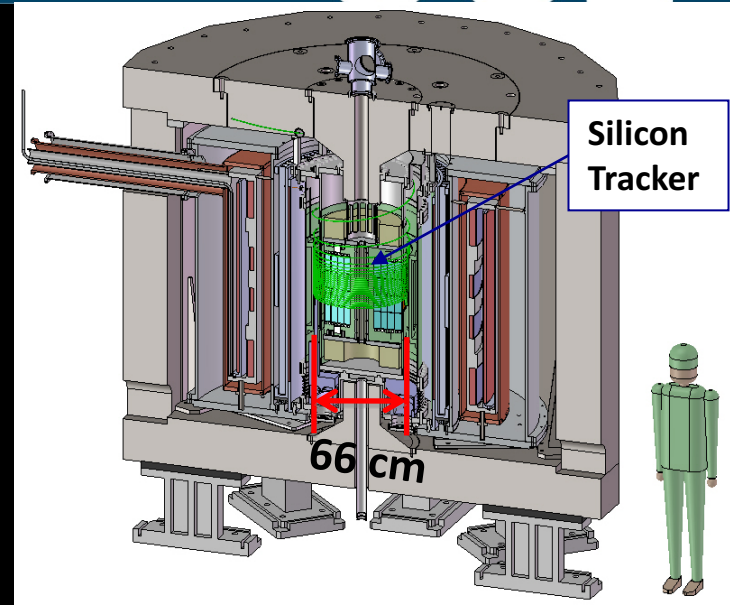
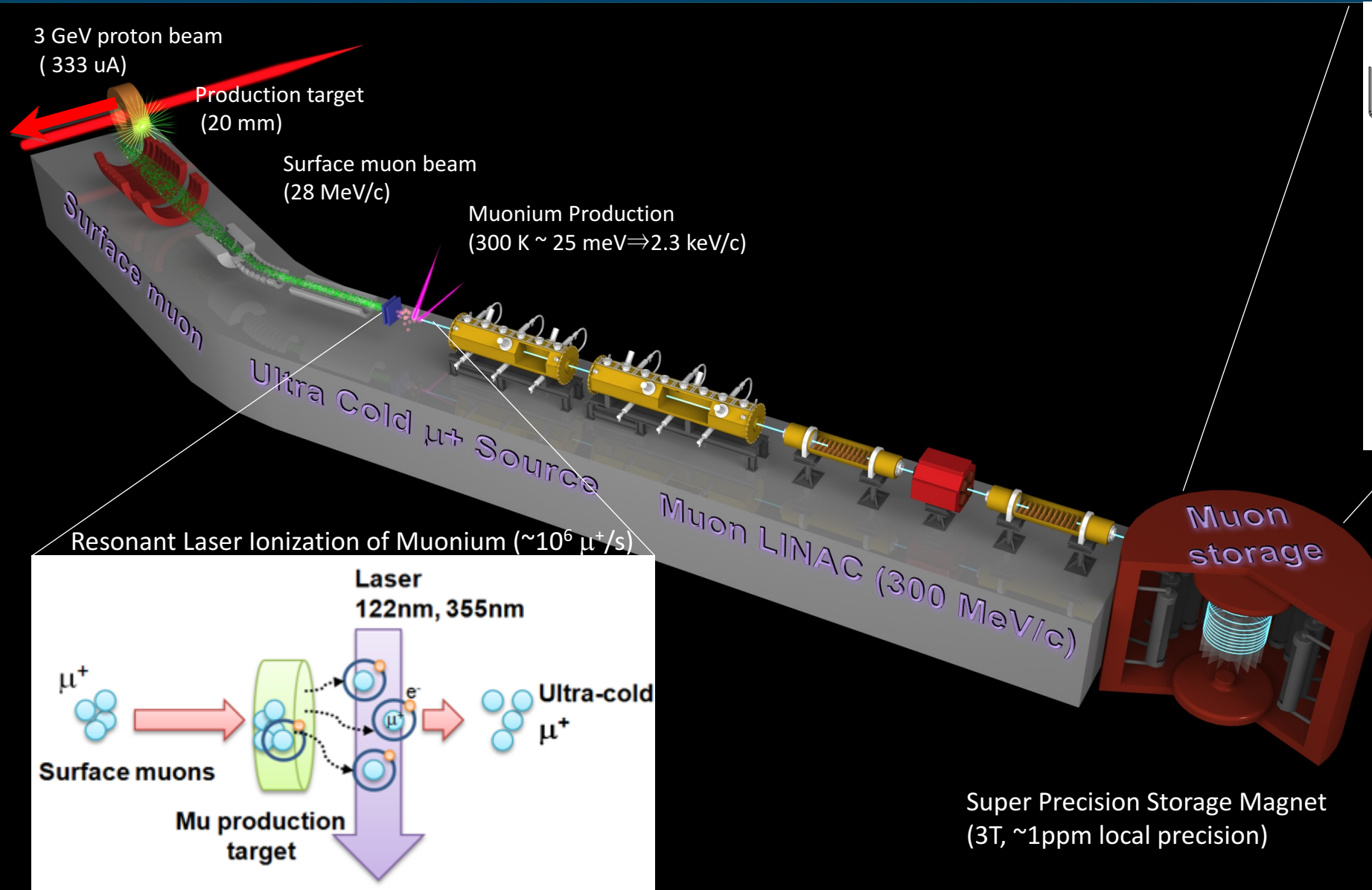
The JPARC experiment instead runs with no electric field
 → Use ultra cold muons to reduce the emittance

To produce the muons laser resonant ionization of muonium is used



The muons produced have very low transverse dispersion which are then accelerated

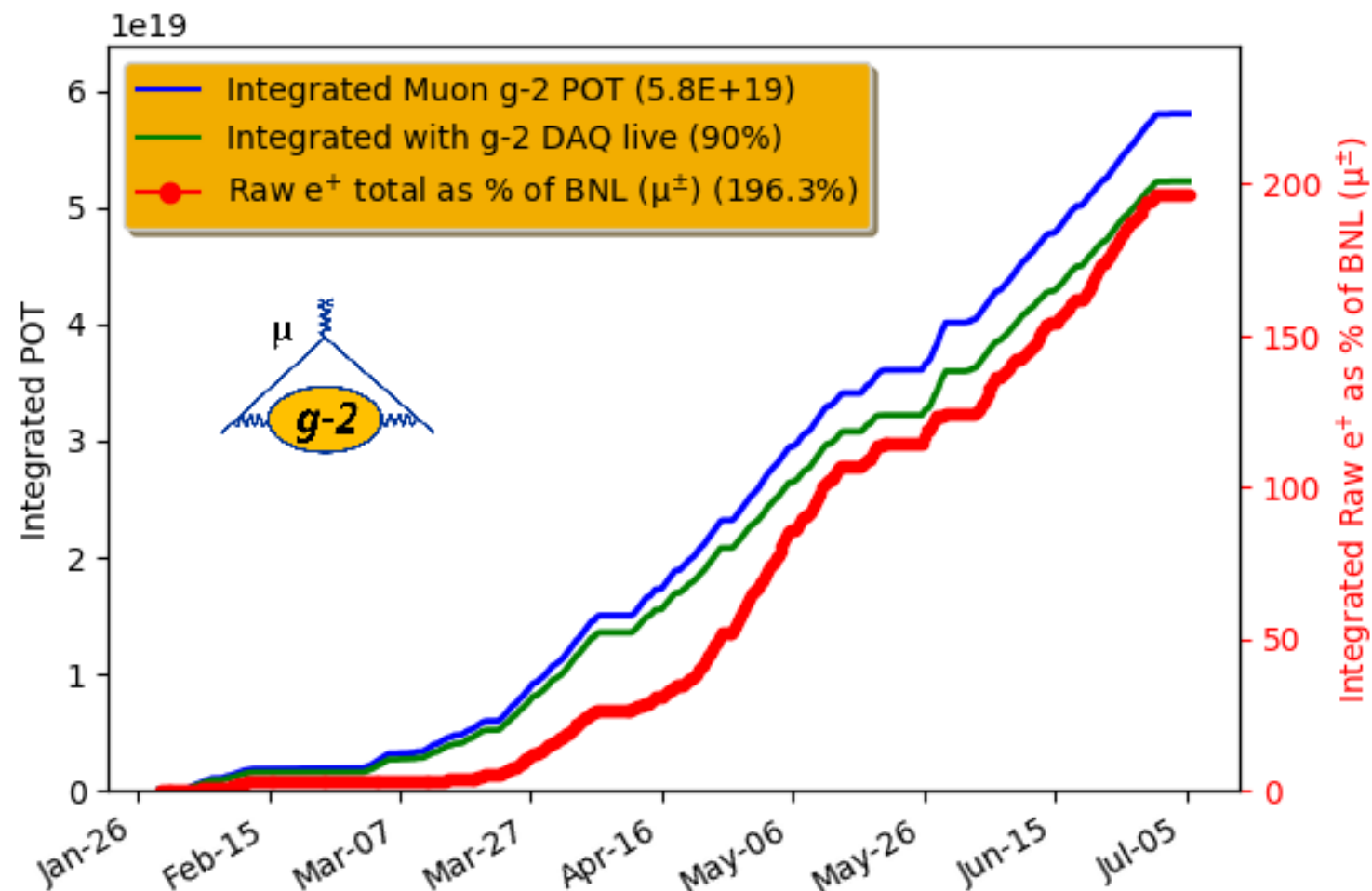
Aside : JPARC g-2



Summary

The new g-2 experiment at Fermilab has just finished the first year of physics data taking

- The new experiment aims to reduce the experimental uncertainty by a factor of 4 to investigate the current discrepancy between experiment and theory of ~ 3.5
- Expect to publish an early result with comparable to BNL precision by early 2019 (based on the data taken between now and the summer)
- An intermediate result will be published in 2020 and then the final full precision result in 2021



Back up

Fermilab Muon g-2 Collaboration ...



US Universities

- Boston
- Cornell
- Illinois
- James Madison
- Kentucky
- Massachusetts
- Michigan
- Michigan State
- Mississippi
- Northern Illinois
- Regis
- Texas
- Virginia
- Washington

National Labs

- Argonne
- Brookhaven
- Fermilab



Italy

- INFN
 - LNF Frascati,
 - Naples
 - Pisa
 - Roma 2
 - Trieste
 - Lecce
- Udine
- Naples
- Trieste
- Rijeka
- Molise
- SNS Pisa



China

- Shanghai



The Netherlands

- Groningen



Germany

- Dresden (thy)



England

- Cockcroft Institute
- Lancaster
- Liverpool
- University College London



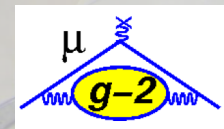
Korea

- KAIST
- CAPP



Russia

- Dubna
- Novosibirsk



Magnetic field systematics

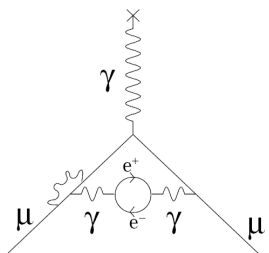
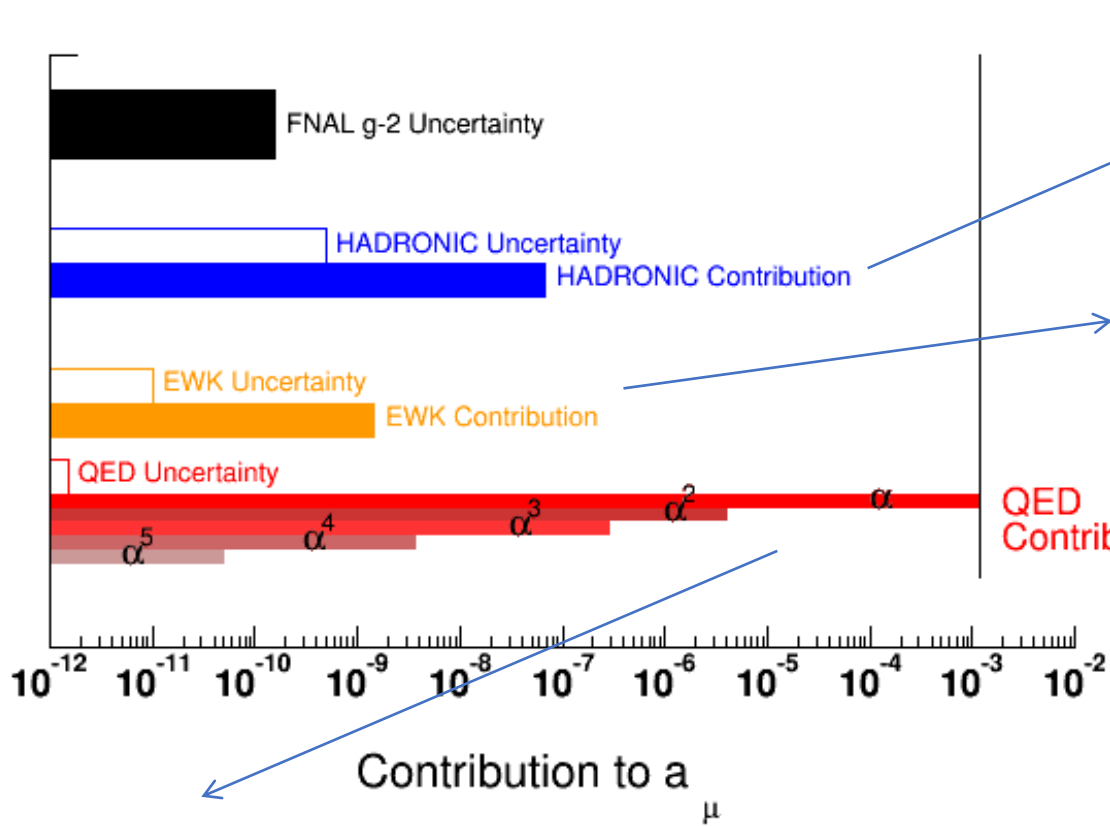
E821 Error	Size [ppm]	Plan for the E989 $g - 2$ Experiment	Goal [ppm]
Absolute field calibrations	0.05	Special 1.45 T calibration magnet with thermal enclosure; additional probes; better electronics	0.035
Trolley probe calibrations	0.09	Absolute cal probes that can calibrate off-central probes; better position accuracy by physical stops and/or optical survey; more frequent calibrations	0.03
Trolley measurements of B_0	0.05	Reduced rail irregularities; reduced position uncertainty by factor of 2; stabilized magnet field during measurements; smaller field gradients	0.03
Fixed probe interpolation	0.07	More frequent trolley runs; more fixed probes; better temperature stability of the magnet	0.03
Muon distribution	0.03	Additional probes at larger radii; improved field uniformity; improved muon tracking	0.01
Time-dependent external B fields	—	Direct measurement of external fields; simulations of impact; active feedback	0.005
Others	0.10	Improved trolley power supply; trolley probes extended to larger radii; reduced temperature effects on trolley; measure kicker field transients	0.05
Total	0.17		0.07

Spin precession systematics

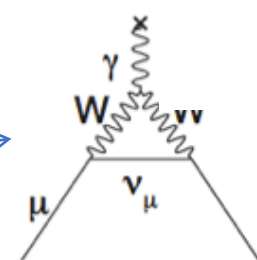
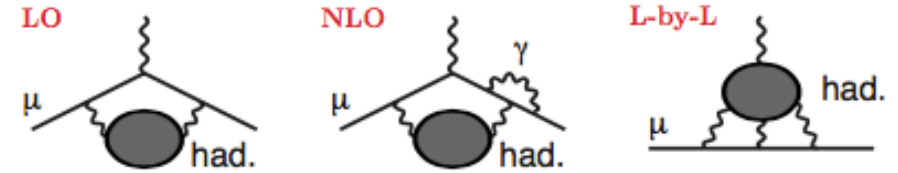
E821 Error	Size [ppm]	Plan for the E989 $g - 2$ Experiment	Goal [ppm]
Gain changes	0.12	Better laser calibration; low-energy threshold; temperature stability; segmentation to lower rates; no hadronic flash	0.02
Lost muons	0.09	Running at higher n -value to reduce losses; less scattering due to material at injection; muons reconstructed by calorimeters; tracking simulation	0.02
Pileup	0.08	Low-energy samples recorded; calorimeter segmentation; Cherenkov; improved analysis techniques; straw trackers cross-calibrate pileup efficiency	0.04
CBO	0.07	Higher n -value; straw trackers determine parameters	0.03
E-Field/Pitch	0.06	Straw trackers reconstruct muon distribution; better collimator alignment; tracking simulation; better kick	0.03
Diff. Decay	0.05 ¹	better kicker; tracking simulation; apply correction	0.02
Total	0.20		0.07

The Magnetic Moment

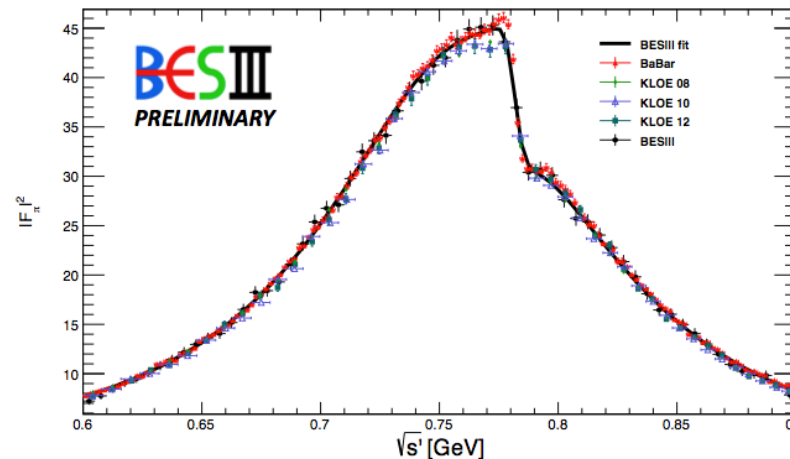
The hadronic uncertainty dominates in the theoretical calculation



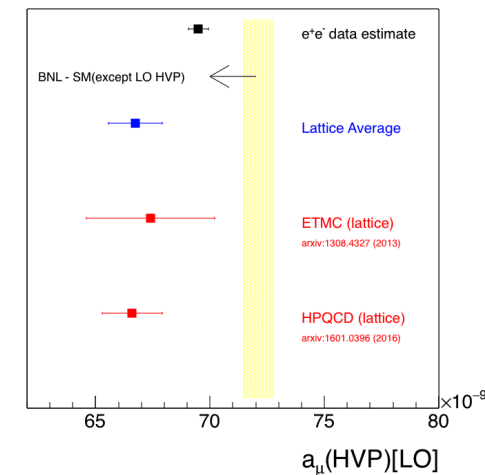
$$a_\mu^{\text{had}} = a_\mu^{\text{had,VP LO}} + a_\mu^{\text{had,VP NLO}} + a_\mu^{\text{had,Light-by-Light}}$$



Improved data should bring down the uncertainty :



Lattice calculations are also catching up :

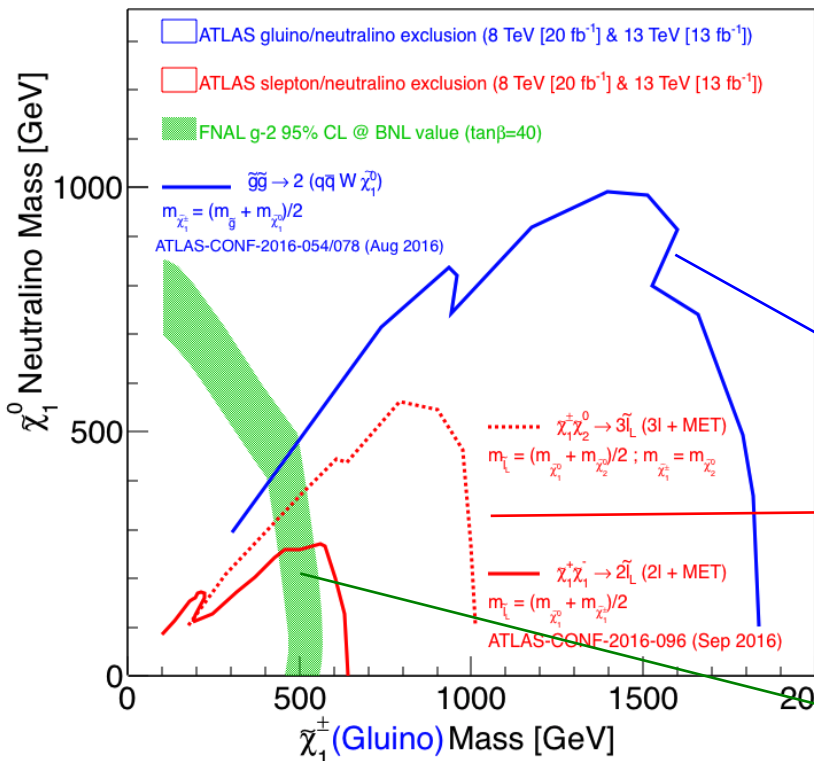
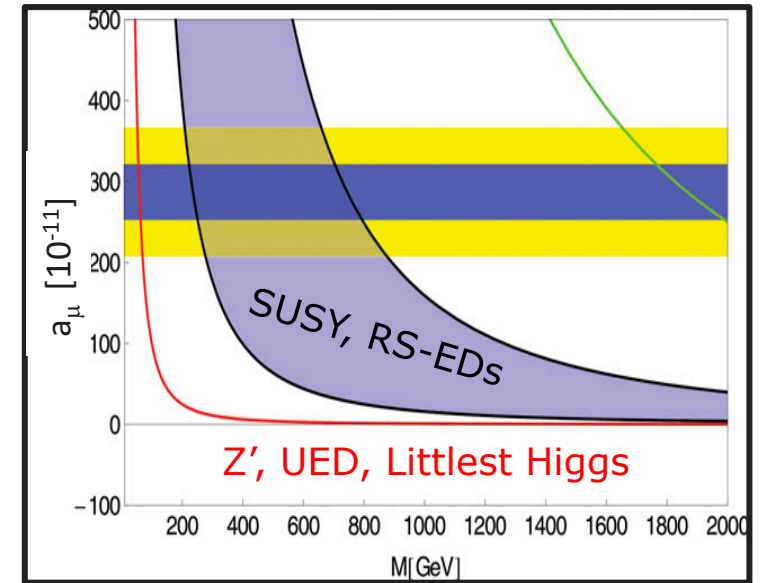


The muon g-2 can probe new physics at TeV scales – complementary to the LHC

The value of the muon g-2 can help set limits on models of new physics

The g-2 interactions flip the chirality of the muon but conserve flavour and CP

Radiative muon mass / technicolor



The LHC has good sensitivity to strongly interacting new physics (SUSY)

But is less sensitive to weakly interacting new physics

Muon g-2 is probing similar phase space as the LHC with more sensitivity in some areas

More μ per proton

Lower inst. rate

Fewer pions

Unique capabilities
of FNAL accelerators

Improved detectors

Improved stored muon
beam dynamics

Improved field uniformity, field
measurement & calibration

Improved modeling of beam
& detectors

BNL \rightarrow FNAL

$[54 \text{ (stat.)} \oplus 33 \text{ (syst.)} \rightarrow 11 \text{ (stat.)} \oplus 11 \text{ (syst.)}] \times 10^{-11}$

0.54 ppm \rightarrow 0.14 ppm

New / improved technologies

Additional collaborators

Building on wealth of experience
from BNL E821 & other expts