

# Central charge of self-dual strings from holography

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Based on 1811.12375, 1812.00923

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**Science & Technology**  
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# Outline

## **Background and motivation**

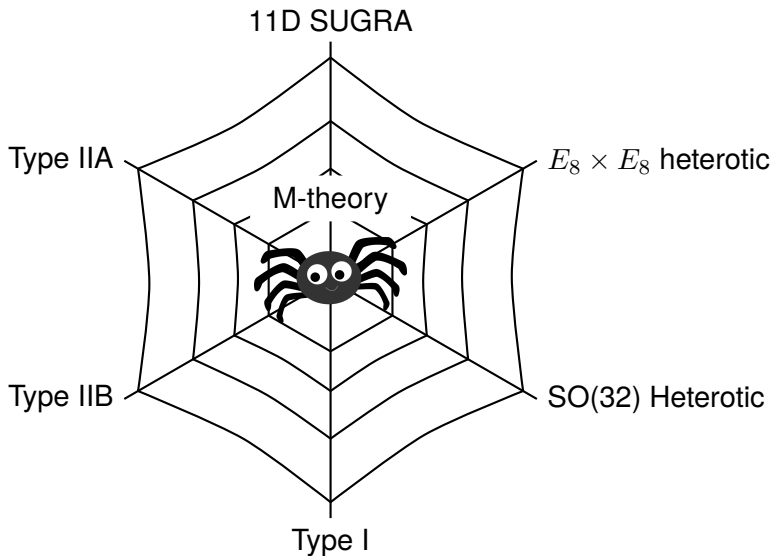
- M-theory
- Central charges
- Entanglement entropy

## Results

## Summary and outlook

# M-theory

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# 11D supergravity

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$$16\pi G_N S = \int d^{11}x \sqrt{-G} \left( R - \frac{1}{2} |F_4|^2 \right) - \frac{1}{6} \int C_3 \wedge F_4 \wedge F_4$$

$$F_4 = dC_3, \quad F_7 = \star F_4 = dC_6 - \frac{1}{2} dC_3 \wedge C_3$$

$C_3 \Rightarrow$  M2-branes

$C_6 \Rightarrow$  M5-branes

# 11D supergravity

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$$F_4 = dC_3, F_7 = \star F_4 = dC_6 - \frac{1}{2} dC_3 \wedge C_3$$

$C_3 \Rightarrow$  M2-branes

$C_6 \Rightarrow$  M5-branes

Stack of D-branes: Yang-Mills

What theory describes a stack of M-branes?

# Stacks of M2-branes: ABJM

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2+1 dimensions,  $\mathcal{N} = 6$  or 8 superconformal symmetry

$U(N)_k \times U(N)_{-k}$  Chern-Simons

[Aharony, Bergman, Jafferis, Maldacena, 0806.1218]

$N \gg k^5$ : dual to 11D SUGRA on  $AdS_4 \times S^7/\mathbb{Z}_k$

Degrees of freedom  $\sim N^{3/2}$

[Klebanov, Tseytlin, hep-th/9604089]

# Stacks of M5-branes

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5+1 dimensions,  $\mathcal{N} = (2, 0)$  superconformal symmetry

$\mathfrak{su}(M)$  gauge algebra

Tensor multiplet:

- Self-dual two-form gauge field  $A_2$ ,  $F_3 = \star F_3$
- One pair of symplectic Majorana-Weyl fermions
- Five real scalar fields

$M \gg 1$ : dual to 11D SUGRA on  $AdS_7 \times S^4$

Degrees of freedom  $\sim M^3$

[Klebanov, Tseytlin, [hep-th/9604089](https://arxiv.org/abs/hep-th/9604089)]

# Why study $\mathcal{N} = (2, 0)$ theory?

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Important part of M-theory – interesting for quantum gravity

6 dimensions = maximal number with superconformal symmetry

6D SCFTs can serve as “master theories”

4D  $\mathcal{N} = 2$  theories from compactification on Riemann surface

[Alday, Gaiotto, Tachikawa 0906.3219]

$\mathcal{N} = 4$  SYM obtained by compactifying  $\mathcal{N} = (2, 0)$  on  $T^2$

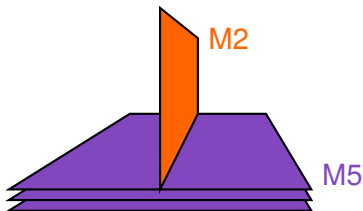
Modular transformations  $\Rightarrow$  S-duality in  $\mathcal{N} = 4$  SYM

[Vafa, hep-th/9707131]



# Self-dual strings

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Two-dimensional intersection: “string”

Electric charge = magnetic charge:

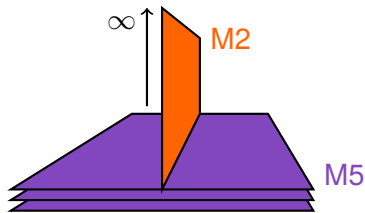
$$Q_E = \int_{S^3} \star F_3 = \int_{S^3} F_3 = Q_B, \quad \text{“self-dual”}$$

Tension  $\sim$  length of M2-brane

[Howe, Lambert, West, hep-th/9709014]

# Self-dual strings

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Infinite tension: length of M2 to infinity

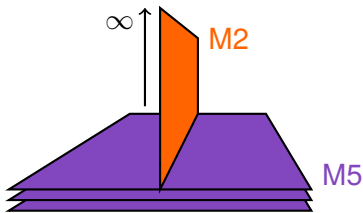
**M5-brane perspective:** Insertion of half-BPS “Wilson surface” into M5-brane theory

$$\langle W_{\mathcal{R}} \rangle = \text{tr}_{\mathcal{R}} \exp \left( i \int A_2 + \dots \right)$$

Depends on representation  $\mathcal{R}$  of Wilson surface

# Self-dual strings

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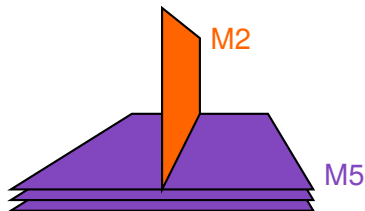


Infinite tension: length of M2 to infinity

**M2-brane perspective:** ABJM with a boundary

Boundary degrees of freedom: [\[Niarchos 1509.0767\]](#)

# Self-dual strings



Planar intersection

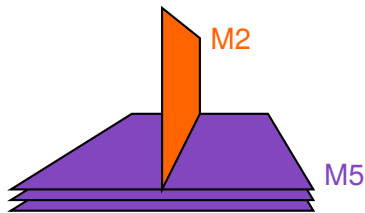
	0	1	2	3	4	5	6	7	8	9	10
M5	•	•	•	•	•	•	-	-	-	-	-
M2	•	•	-	-	-	-	•	-	-	-	-

Preserves 2D large  $\mathcal{N} = (4, 4)$  SUSY

**M5-brane perspective:**

$$SO(6, 2) \times SO(5)_R \rightarrow SO(2, 2) \times SO(4)_R \times SO(4)_R$$

# Self-dual strings



Planar intersection

	0	1	2	3	4	5	6	7	8	9	10
M5	•	•	•	•	•	•	-	-	-	-	-
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Preserves 2D large  $\mathcal{N} = (4, 4)$  SUSY

**M2-brane perspective:**

$$SO(3, 2) \times SO(8)_R \rightarrow SO(2, 2) \times SO(4)_R \times SO(4)_R$$

# Self-dual strings in holography

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M5-branes:  $M \gg 1$ , asymptotically  $AdS_7 \times S^4$

M2-branes:  $N \gg 1$ , asymptotically  $(AdS_4/\mathbb{Z}_2) \times S^7$

SUGRA solutions with flux

- Partition of M2-branes on M5-branes
- Young tableau

Representation depends on configuration of fluxes

$$\# \text{ boxes in Young tableau} = N$$

[Lunin, 0704.3442; D'Hoker, Estes, Gutperle, Krym, 0806.0605;...]

# Central charge: 2D CFT

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2D CFTs are characterised by their central charge  $c$ .

Virasoro algebra:

$$[L_m, L_n] = (m - n)L_{m+n} + \frac{c}{12}(m^3 - m)\delta_{m,-n}$$

# Central charge: 2D CFT

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$$[L_m, L_n] = (m - n)L_{m+n} + \frac{c}{12}(m^3 - m)\delta_{m,-n}$$

Conformal anomaly:  $\langle T^\mu{}_\mu \rangle = \frac{c}{24\pi}R$

Stress-tensor OPE:  $T(z)T(0) = \frac{c}{2z^4} + \dots$

Thermal entropy:  $S = \frac{\pi}{3}cLT$

Entanglement entropy:  $S_{\text{EE}} = \frac{c}{3} \log \left( \frac{2\ell}{\epsilon} \right) + \dots$

Central charge counts degrees of freedom.



# Central charge: 2D defects

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Higher dimensional CFT with 2D defect or boundary

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# Central charge: 2D defects

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Entanglement entropy:

$$S_{\text{EE}}^{\text{defect}} = \frac{b}{3} \log \left( \frac{2\ell}{\epsilon} \right) + \dots$$

# Central charge: 2D defects

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Entanglement entropy:

$$S_{\text{EE}}^{\text{defect}} = \frac{b}{3} \log \left( \frac{2\ell}{\epsilon} \right) + \dots$$

Conformal anomaly:

$$\langle T^\mu{}_\mu \rangle = (\text{bulk terms}) + \delta(x_\perp) \frac{1}{24\pi} \left( c\hat{R} + d_1\dot{\Pi}^2 - d_2 W_{ab}{}^{ab} \right)$$

## Plan:

- Build brane intersections in SUGRA
- Compute entanglement entropy
- Extract  $b$

# Outline

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- M-theory
- Central charges
- **Entanglement entropy**

## Results

## Summary and outlook

# Entanglement entropy

---

Factorize Hilbert space  $\mathcal{H} = \mathcal{A} \otimes \mathcal{B}$

Reduced density matrix on  $\mathcal{A}$ :

$$\rho_{\mathcal{A}} = \sum_{|\psi\rangle \in \mathcal{B}} \langle \psi | \rho | \psi \rangle$$

EE = von Neumann entropy of reduced density matrix:

$$S_{\text{EE}} = -\text{Tr} \rho_{\mathcal{A}} \log \rho_{\mathcal{A}}$$

# Entanglement entropy

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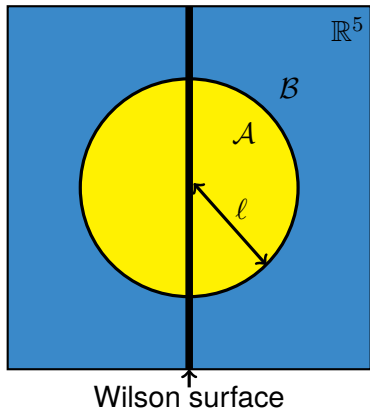
For a QFT, natural splitting of  $\mathcal{H}$  into spatial subregions.

2D CFT: [[Cardy, Calabrese, hep-th/0405152](#)]

$$S_{\text{EE}} = \frac{c}{3} \log \left( \frac{2\ell}{\epsilon} \right)$$

# Entanglement entropy with 2D defect

Entangling region: sphere of radius  $\ell$



Defect contribution to EE:

$$S_{\text{EE}} = S_{\text{EE}}^{d=6} + S_{\text{EE}}^{d=2}$$

Determined by central charge:

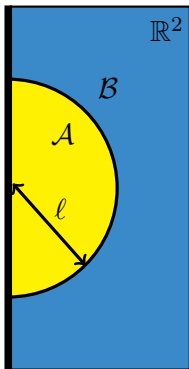
$$S_{\text{EE}}^{d=2} = \frac{b}{3} \log \left( \frac{2\ell}{\epsilon} \right) + \dots$$

$$b = 3\ell \frac{d}{d\ell} S_{\text{EE}}^{d=2}$$



# Entanglement entropy with boundary

Entangling region: semicircle of radius  $\ell$



Defect contribution to EE:

$$S_{\text{EE}} = \frac{1}{2} S_{\text{EE}}^{d=3} + S_{\text{EE}}^{d=2}$$

Determined by central charge:

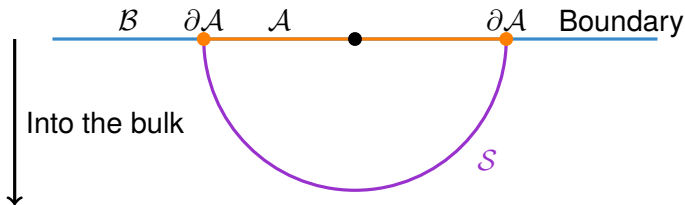
$$S_{\text{EE}}^{d=2} = \frac{b}{3} \log \left( \frac{2\ell}{\epsilon} \right) + \dots$$

$$b = 3\ell \frac{d}{d\ell} S_{\text{EE}}^{d=2}$$

# Holographic entanglement entropy

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Minimal surface extending into bulk:



$$S_{\text{EE}} = \frac{\text{Area}[S]}{4G_N}$$

[Ryu, Takayanagi, hep-th/0603001 & hep-th/0605073]

[Lewkowycz, Maldacena, 1304.4926]

# Outline

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- M-theory
- Central charges
- Entanglement entropy

## **Results**

## Summary and outlook

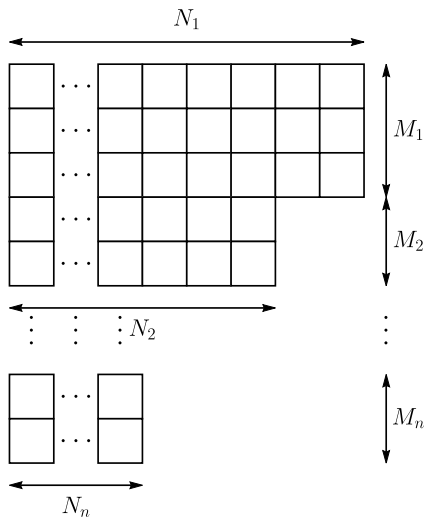
# Parameterisation of solutions

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$M_a$  M5-branes that have  $N_a$  M2-branes ending on them:

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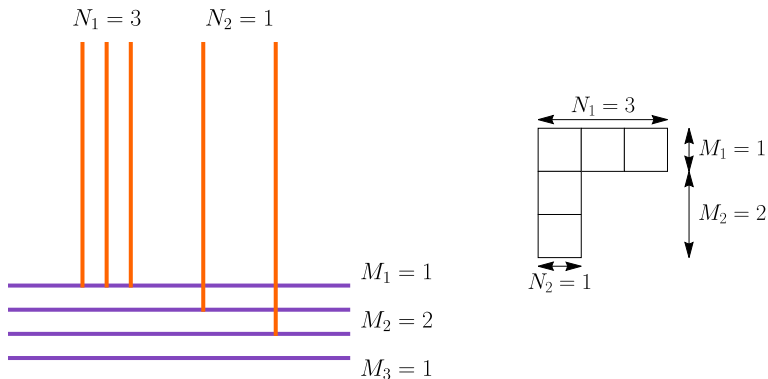
$$M = \sum_{a=1}^{n+1} M_a$$

$$N = \sum_{a=1}^n M_a N_a$$

$$N_{n+1} = 0$$

# Parameterisation of solutions

$M_a$  M5-branes that have  $N_a$  M2-branes ending on them:



$$M = \sum_{a=1}^{n+1} M_a, \quad N = \sum_{a=1}^n N_a M_a, \quad N_{n+1} = 0$$

# Central charge: M5-branes

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$$\begin{aligned} b_{M5} &= 3\ell \frac{d}{d\ell} \left[ S_{EE} - S_{EE}^{d=6} \right] \\ &= \frac{3}{5} \left[ 8MN + \frac{N^2}{M} + \sum_{a=1}^n \left( 8M_a^2 N_a - M_a N_a^2 - 16 \sum_{b=1}^a M_a M_b N_a \right) \right] \end{aligned}$$

In terms of highest weight  $\lambda$  and Weyl vector  $\rho$ ,

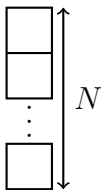
$$b_{M5} = \frac{48}{5}(\rho, \lambda) - \frac{3}{5}(\lambda, \lambda)$$

# Central charge: M5-branes

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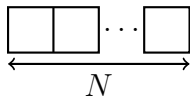
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## Antisymmetric representation



$$b_{M5} = \frac{24}{5}N(M - N)$$

## Symmetric representation



$$b_{M5} = \frac{24}{5}N(M - N/8)$$



# Central charge

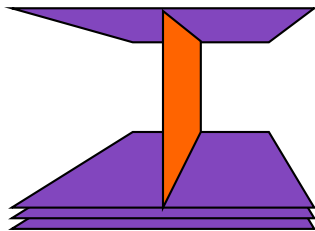
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Compare to result of [\[Berman, Harvey, hep-th/0408198\]](#)?

R-symmetry anomaly from inflow calculation

Coulomb branch:

$$SU(M + 1) \rightarrow SU(M) \times U(1)$$



Central charge:

$$b_{\text{BH}} = \frac{1}{4} N M$$

# Central charge

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Compare to result of [\[Niarchos, Siampos, 1206.2935\]](#)?

Blackfold approach: thermodynamic entropy

$$S = \frac{\pi}{3} b_{\text{NS}} L T + \dots$$

$$b_{\text{NS}} \approx 1.2 \frac{N^2}{M}$$

# Central charge: M2-branes

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$$\begin{aligned} b_{\text{M2}} &= 3\ell \frac{d}{d\ell} \left[ S_{\text{EE}} - \frac{1}{2} S_{\text{EE}}^{d=3} \right] \\ &= 3 \sum_{a=1}^n \left[ -M_a^2 N_a + \frac{1}{2} M_a N_a^2 + 2 \sum_{b=1}^a N_a M_a M_b \right] \end{aligned}$$

Symmetric partition:  $b_{\text{M2}} = 3 \left( N^2 + \frac{1}{2} N \right)$

Antisymmetric partition:  $b_{\text{M2}} = 3 \left( \frac{1}{2} N^2 + N \right)$

What is field theory interpretation of Young tableau?

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## Results

## **Summary and outlook**

# Summary

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Despite much progress on M5-brane physics, fundamental questions remain

We have studied Wilson surfaces in M5-brane theory holographically

Entanglement entropy provides a definition of a self-dual string central charge

Extends previous M5-brane central charge calculations, shows similar scaling with number of branes

Central charge in terms of group theory quantities

$$b_{M5} = \frac{48}{5}(\rho, \lambda) - \frac{3}{5}(\lambda, \lambda)$$

# Outlook

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Does this central charge show up in other quantities?

- Entanglement entropy related to  $\langle T^\mu{}_\mu \rangle$   
[Jensen, O'Bannon, Robinson, RR, 1812.xxxxx]
- Does this explain differences with other calculations?

Can we compute all of the trace anomaly coefficients?

- Different geometries for brane intersections?

How do these quantities behave along RG flows? Can we understand the results from field theory?

- AGT? [Alday, Gaiotto, Tachikawa, 0906.3219]
- Quantities in compactified theories that tell you about the central charge?

Thank you!

**Backup slides**



# Bubbling geometries

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[Lunin, 0704.3442; D'Hoker, Estes, Gutperle, Krym, 0806.0605;  
Estes, Feldman, Krym, 1209.1845; Bachas, D'Hoker, Estes, Krym, 1312.5477]

11D SUGRA solutions with  $AdS_3 \times S^3 \times S^3$  isometry

Remaining two coordinates on Riemann surface,  $\Sigma_2$

$$ds^2 = f_1^2 ds_{AdS_3}^2 + f_2^2 ds_{S_2^3}^2 + f_3^2 ds_{S_3^3}^2 + 2\rho_{w\bar{w}}^2 dw d\bar{w}$$

$$C_{(3)} = b_1 ds_{AdS_3} + b_2 ds_{S_2^3} + b_3 ds_{S_3^3}$$

# Bubbling geometries

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$$C_{(3)} = b_1 ds_{AdS_3} + b_2 ds_{S_2^3} + b_3 ds_{S_3^3}$$

$f_i, b_i, \rho_{w\bar{w}}$  real functions on  $\Sigma_2$ , express in terms of  $\{\gamma, h, G\}$

Super algebra  $D(2, 1; \gamma) \oplus D(2, 1; \gamma)$

[Sevrin, Troost, Van Proeyen, Phys. Lett. B, **208**, 447 (1988)]

$$\partial\bar{\partial}h = 0, \quad \partial G = \frac{1}{2}(G + \bar{G})\partial \log h, \quad h > 0, \quad |G| < 1$$

$$h = 0, \quad G = \pm i \quad \text{on } \partial\Sigma_2$$

# Bubbling geometries

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$$\partial\bar{\partial}h = 0, \quad \partial G = \frac{1}{2}(G + \bar{G})\partial \log h$$

$w$

$$h > 0 \quad |G| < 1$$

$$h = 0 \quad G = \pm i$$

# $AdS_7 \times S^4$ asymptotics

$$\gamma = -\frac{1}{2}, \quad h = -i(w - \bar{w}), \quad G = -i \left( 1 + \sum_{j=1}^{2n+2} (-1)^j \frac{w - \xi_j}{|w - \xi_j|} \right)$$

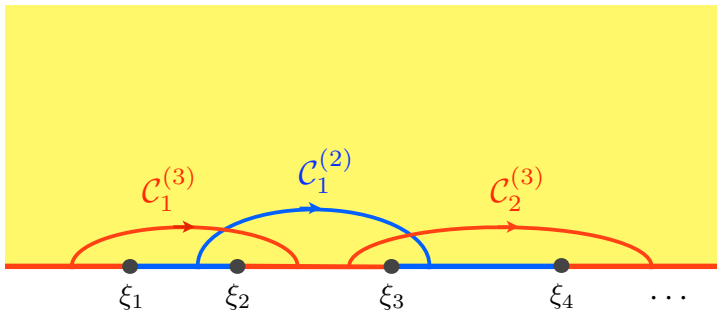


Figure from Bachas, D'Hoker, Estes, Krym 1312.5477

# $AdS_4 \times S^7$ asymptotics

$$\gamma < 0, \quad h = -i(w - \bar{w}), \quad G = -i \sum_{j=1}^{2n+1} (-1)^j \frac{w - \xi_j}{|w - \xi_j|}$$

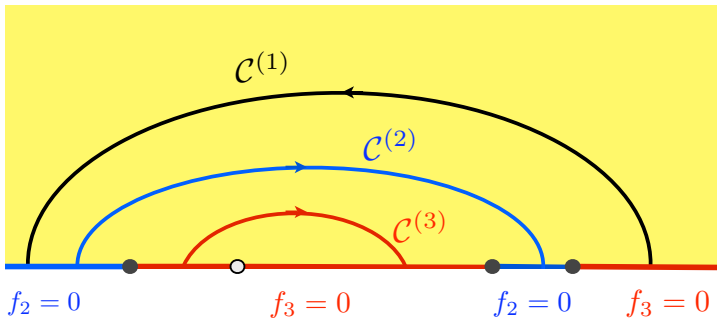


Figure from Bachas, D'Hoker, Estes, Krym 1312.5477

# Entanglement entropy

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Ryu-Takayanagi surface known for any geometry dual to a conformal defect: [\[Jensen, O'Bannon, 1309.4523\]](#)

Hemisphere in  $AdS_3$ , wrapping  $S_{2,3}^3$  and  $\Sigma_2$

$$ds^2 = f_1^2 ds_{AdS_3}^2 + f_2^2 ds_{S_2^2}^2 + f_3^2 ds_{S_3^2}^2 + 2\rho_{w\bar{w}}^2 dw d\bar{w}$$

Entanglement entropy

$$\begin{aligned} S_{EE} &= \frac{\text{Area}[\text{RT}]}{4G_N} \\ &= \frac{2(2\pi^2)^2}{4G_N} \int dw d\bar{w} (\rho_{w\bar{w}}^2 f_1 f_2^3 f_3^3) \underbrace{\int_{\epsilon}^{\ell} \frac{dz}{z} \frac{\ell}{\sqrt{\ell^2 - z^2}}}_{\log(2\ell/\epsilon) + \mathcal{O}(\epsilon)} \end{aligned}$$

[\[Gentle, Gutperle, Marasinou, 1506.00052\]](#)