Electroweak phenomena

Eleni Vryonidou



James Stirling Memorial Conference IPPP, 17/9/2019

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Eikonal regime of gravity-induced scattering at higher energy proton colliders

W.J. Stirling, E. Vryonidou, J.D. Wells

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Charm production in associati W	on with an electroweak g	auge boson at the LHC	Electroweak corrections and Bloch-Nordsied violations in 2-to-2 processes at the LHC	:k

W.J. Stirling, E. Vryonidou

W+charm as a probe of the strange PDF



Cross-section ratios are more sensitive:

$$R_c^{\pm} = \frac{\sigma(W^+ + \bar{c})}{\sigma(W^- + c)}$$

$$R_c^{\pm} = 1$$
 at the Tevatron $R_c^{\pm} < 1$ at the LHC

$$R_{c} = \frac{\sigma(W + c)}{\sigma(W + \text{jet})}$$

Sensitive to strange plus anti-strange

W. J. Stirling and EV Phys.Rev.Lett. 109 (2012) 082002

Comparisons of different PDF sets

Ratio	R_c^{\pm}	R_{c}
CT10	$0.953^{+0.009}_{-0.007}$	$0.124_{-0.012}^{+0.021}$
MSTW2008NLO	$0.921\substack{+0.022\\-0.033}$	$0.116\substack{+0.002\\-0.002}$
NNPDF2.1NLO	$0.944{\pm}0.008$	$0.104 {\pm} 0.005$
GJR08	$0.933 {\pm} 0.003$	$0.099 {\pm} 0.002$
ABKM09	$0.933 {\pm} 0.002$	$0.116 {\pm} 0.003$



d-contribution suppressed by factor of 20



- •NNPDF has a very small strange asymmetry
- •CT10 no asymmetry
- •MSTW larger asymmetry

W. J. Stirling and EV Phys.Rev.Lett. 109 (2012) 082002

Impact of LHC measurements



Precise 13TeV measurement more promising

What we can learn from Z plus charm

$$R_c^Z = \frac{\sigma(Z+c)}{\sigma(Z+jet)}$$

PDF set	R_c^Z
CT10	$0.0619\substack{+0.0032\\-0.0032}$
MSTW2008NLO	$0.0640\substack{+0.0014\\-0.0016}$
NNPDF2.1NLO	$0.0660 {\pm} 0.0013$
GJR08	$0.0611 {\pm} 0.0011$
ABKM09	$0.0605 {\pm} 0.0019$



NNPDF collaboration arXiv:1605.06515

$$R_c^{\pm}(Z) = \frac{\sigma(Z + \bar{c})}{\sigma(Z + c)}$$

A probe of intrinsic charm?

W. J. Stirling and EV Phys.Rev.Lett. 109 (2012) 082002

Gauge boson polarisation at the LHC

Electroweak gauge boson polarisation at the LHC



Only LH quarks and RH antiquarks couple to W (V-A) In the proton: quark more likely to have more momentum than the antiquark \longrightarrow W moves in the quark direction \longrightarrow LH W⁺(73%)

$$\frac{1}{\sigma}\frac{d\sigma}{d\cos\theta^*} = \frac{3}{8}(1-\cos\theta^*)^2 f_L + \frac{3}{8}(1+\cos\theta^*)^2 f_R + \frac{3}{4}\sin^2\theta^* f_0$$

$$f_0 = 2 - 5\langle \cos\theta^{*2} \rangle,$$

$$f_L = -\frac{1}{2} - \langle \cos\theta^{*} \rangle + \frac{5}{2} \langle \cos\theta^{*2} \rangle,$$

$$f_R = -\frac{1}{2} + \langle \cos\theta^{*} \rangle + \frac{5}{2} \langle \cos\theta^{*2} \rangle,$$

W polarisation fractions

 θ is the angle in the W rest frame between the charged lepton and the W flight direction in the lab frame

W. J. Stirling and EV JHEP 1207 (2012) 124

W+jets: polarisation fractions



Small differences between W⁺ and W⁻ due to PDF difference of u and d-quarks.

W. J. Stirling and EV JHEP 1207 (2012) 124



W bosons are predominantly lefthanded at high p_T independently of the number of jets as seen also by Bern et al. Phys.Rev.D84:034008,2011

W polarisation in top production



Polarisation is frame dependent and interaction dependent

Conclusions

Thesis submitted in September 2013 Title: Phenomenology of the Standard Model and beyond at hadron colliders Extremely lucky to have James

as my supervisor



