



# EFT Constraints from Higgs measurements

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Nicholas Wardle – Imperial College London

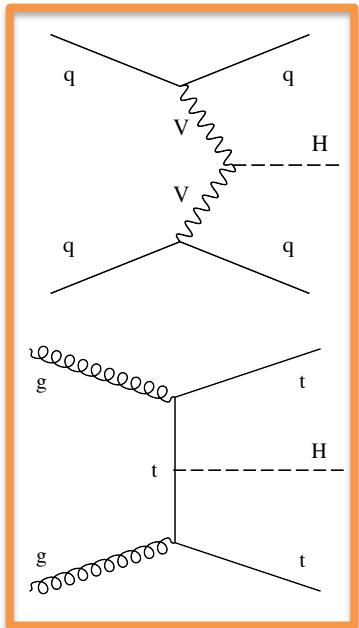
**Pushing the Boundaries - Standard Model and Beyond at LHC**  
**Durham, UK**

19/09/2019

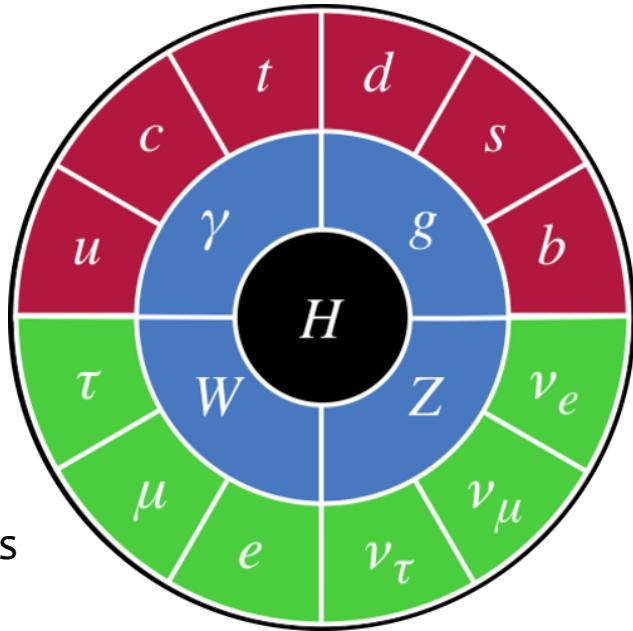
# Completing the SM

Discovery of the Higgs boson completed the Standard Model

- Great success of LHC @ Run-1
- Coupling structure in SM allows for rich program of Higgs boson studies @ LHC



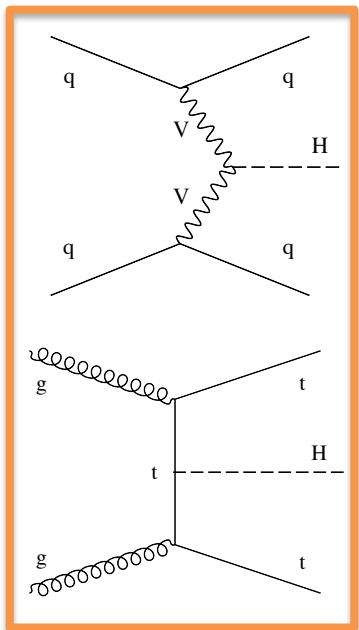
**Tree level** coupling to vector bosons and heavy fermions



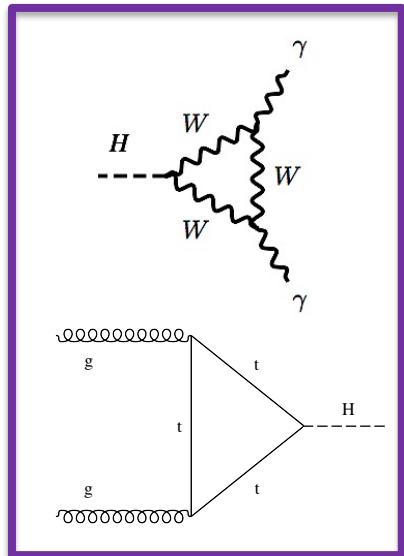
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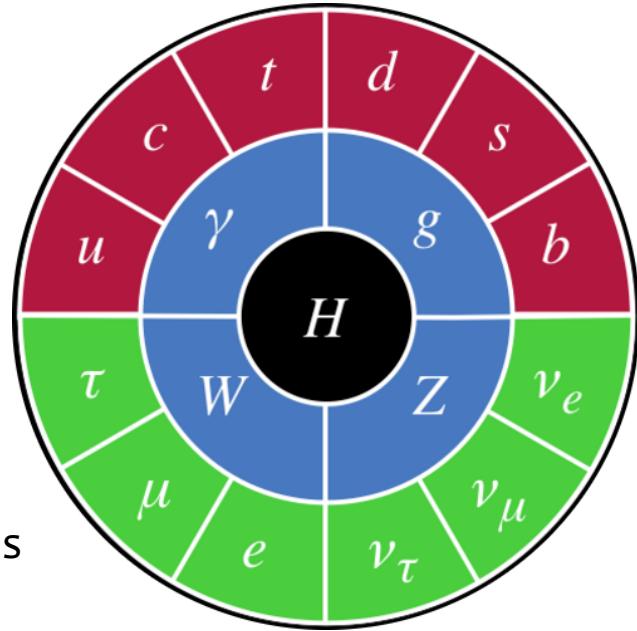
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**Tree level** coupling to vector bosons and heavy fermions



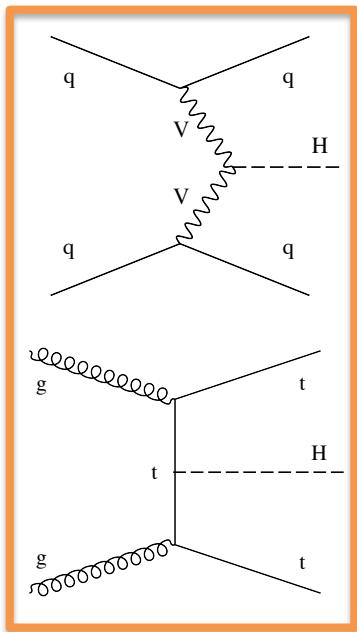
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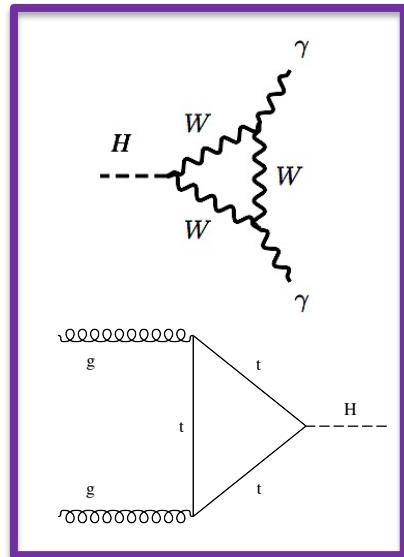
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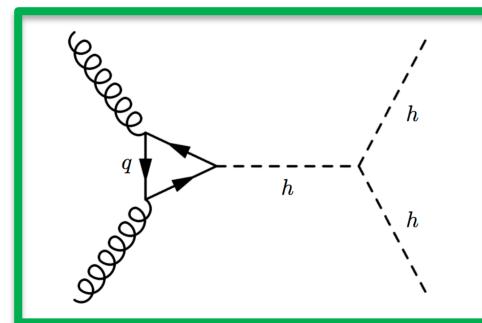
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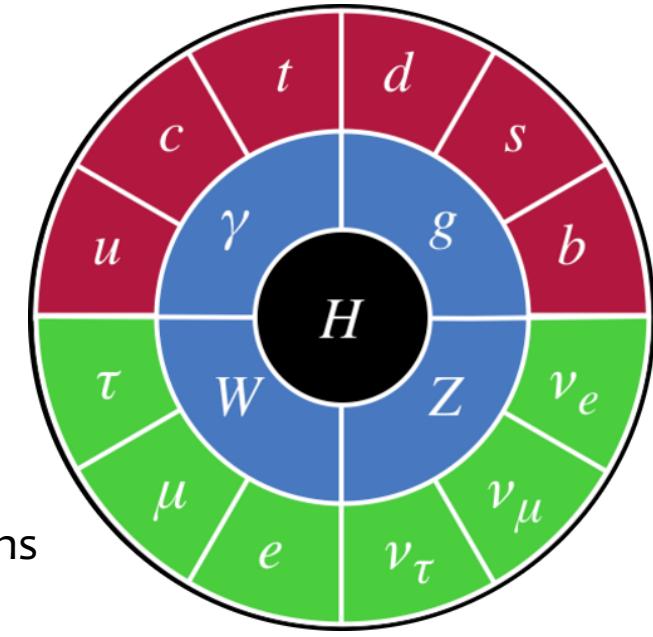
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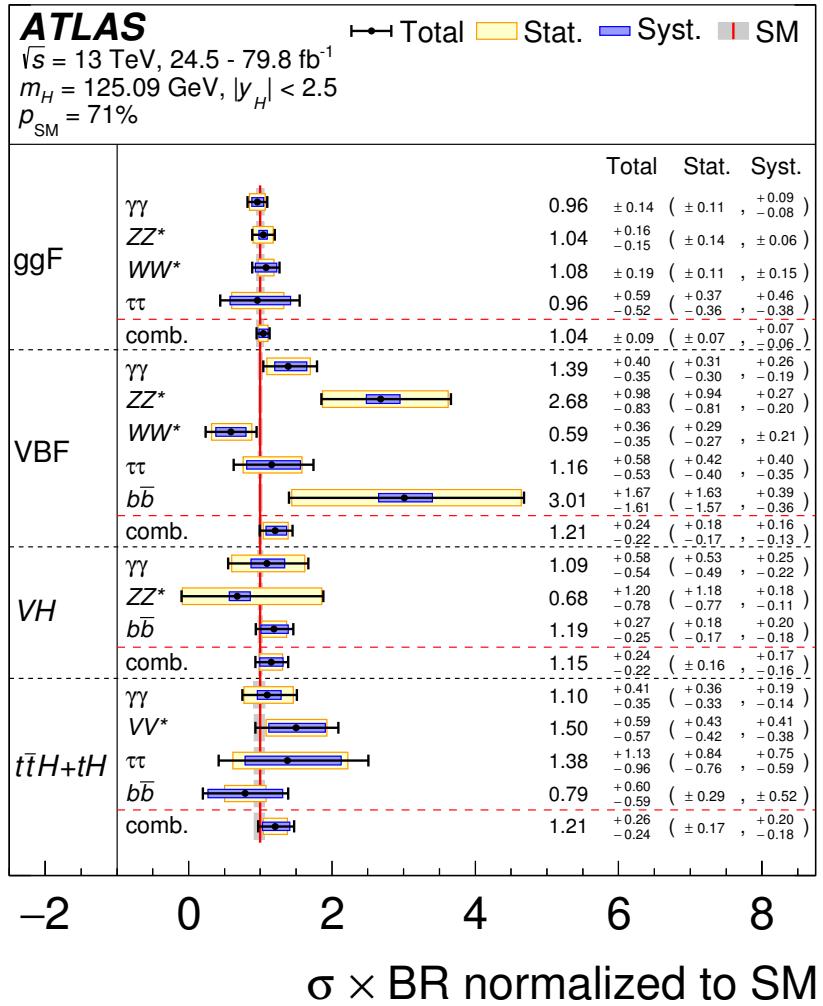


Higgs boson  
the only particle that **couples to itself!**

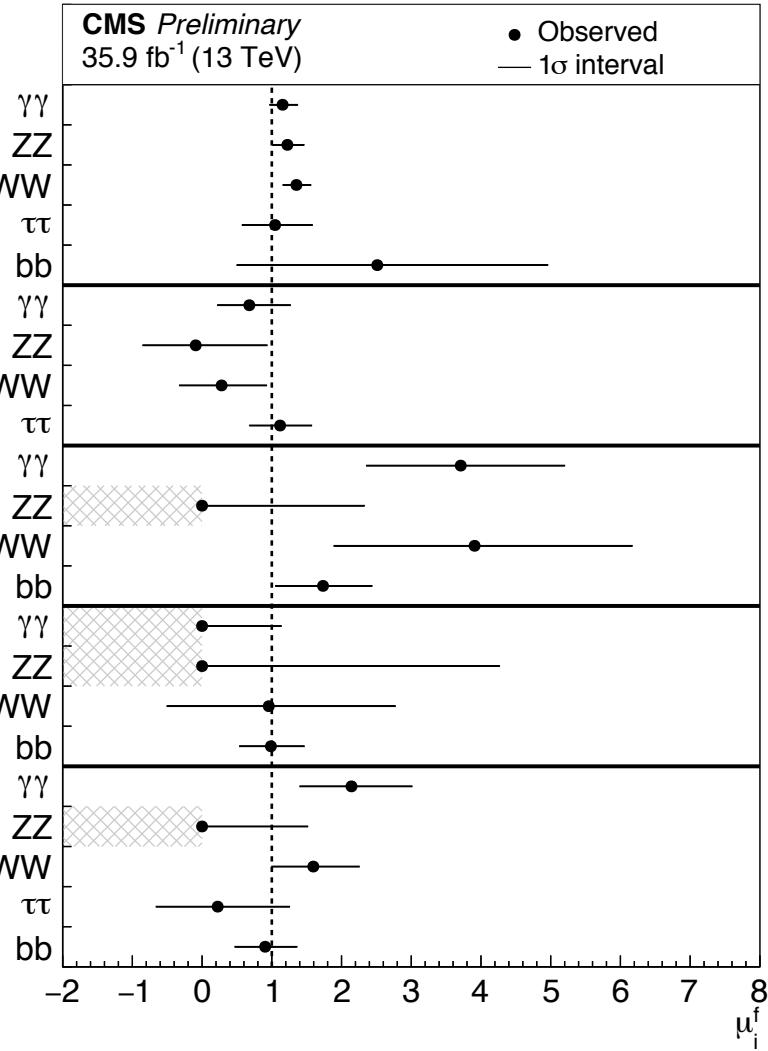


# Great for experiments!

arXiv:1909.02845 (sub to Phys Rev D)



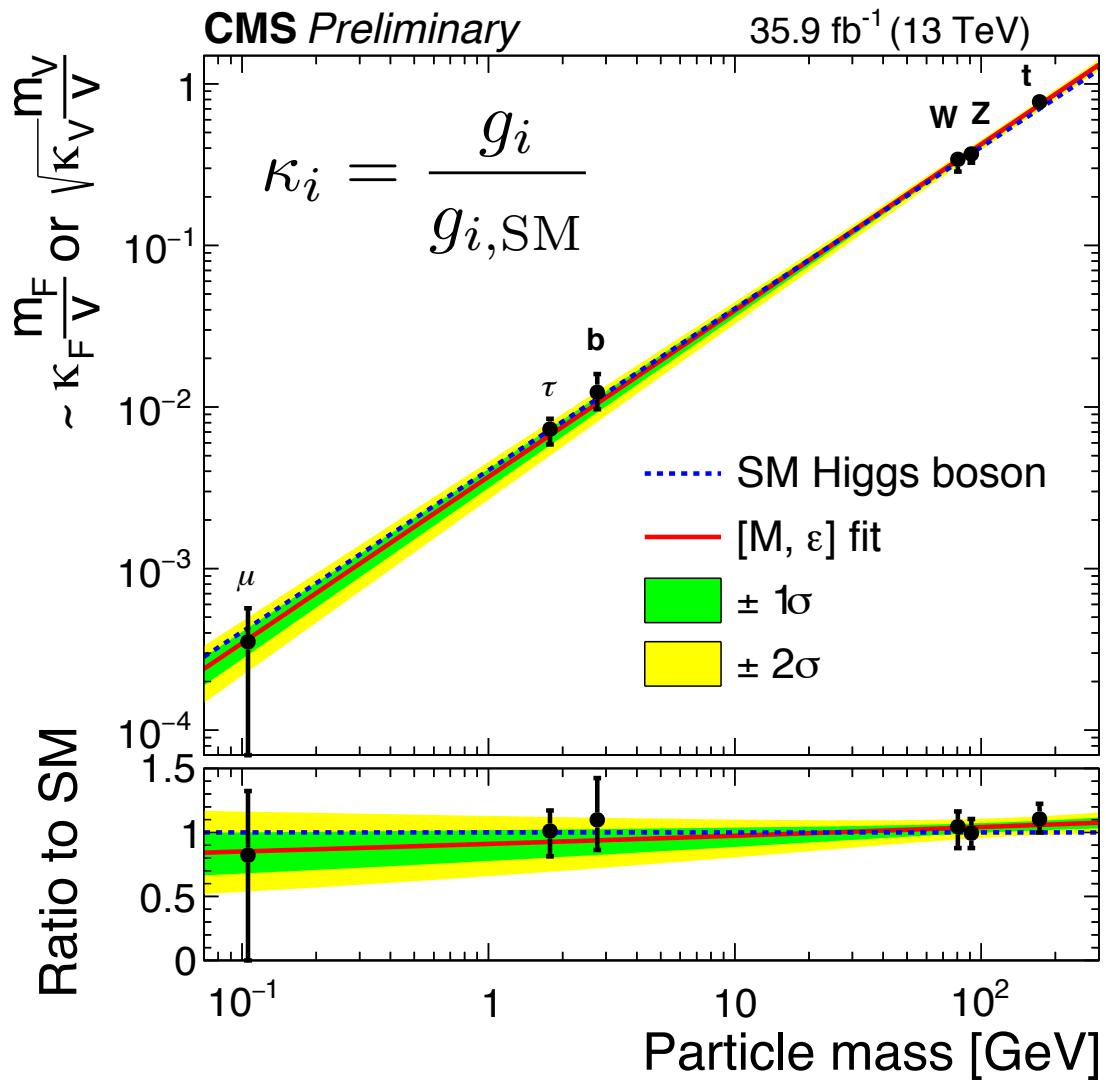
Eur. Phys. J. C 79 (2019)



Wide range of measurements possible  
to study Higgs boson properties

# What do we know?

Couplings to heavy fermions and bosons well established!

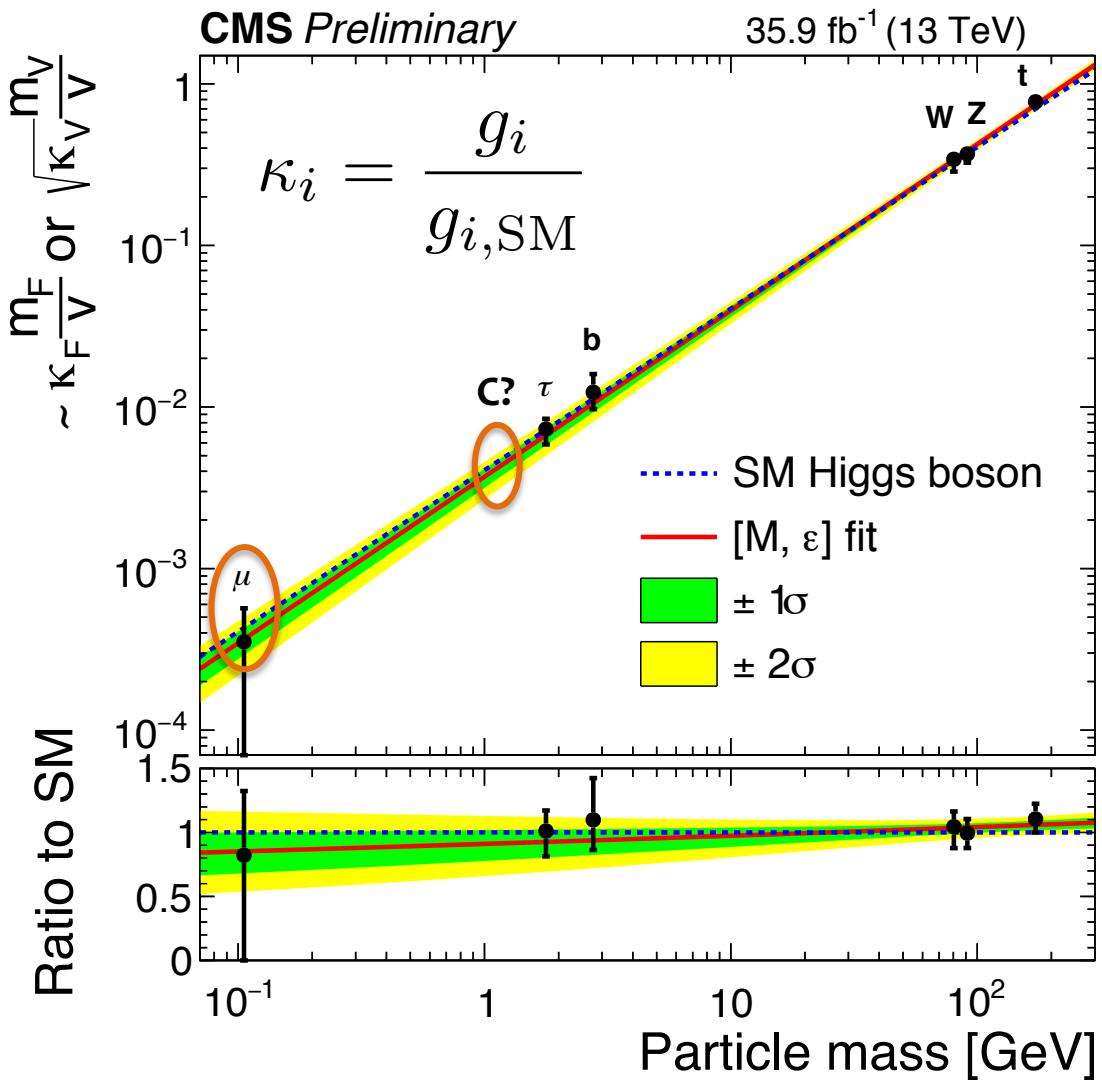


# What don't we know?

Couplings to heavy fermions and bosons well established!

Yet to confirm **coupling to 2<sup>nd</sup> Generation**

- $B(H \rightarrow \mu\mu) < 1.7 \times \text{SM}$  @ 95% CL  
(ATLAS-CONF-2019-028)
- $B(H \rightarrow cc) < 70 \times \text{SM}$  @ 95 % CL  
(CMS-PAS-HIG-18-031)



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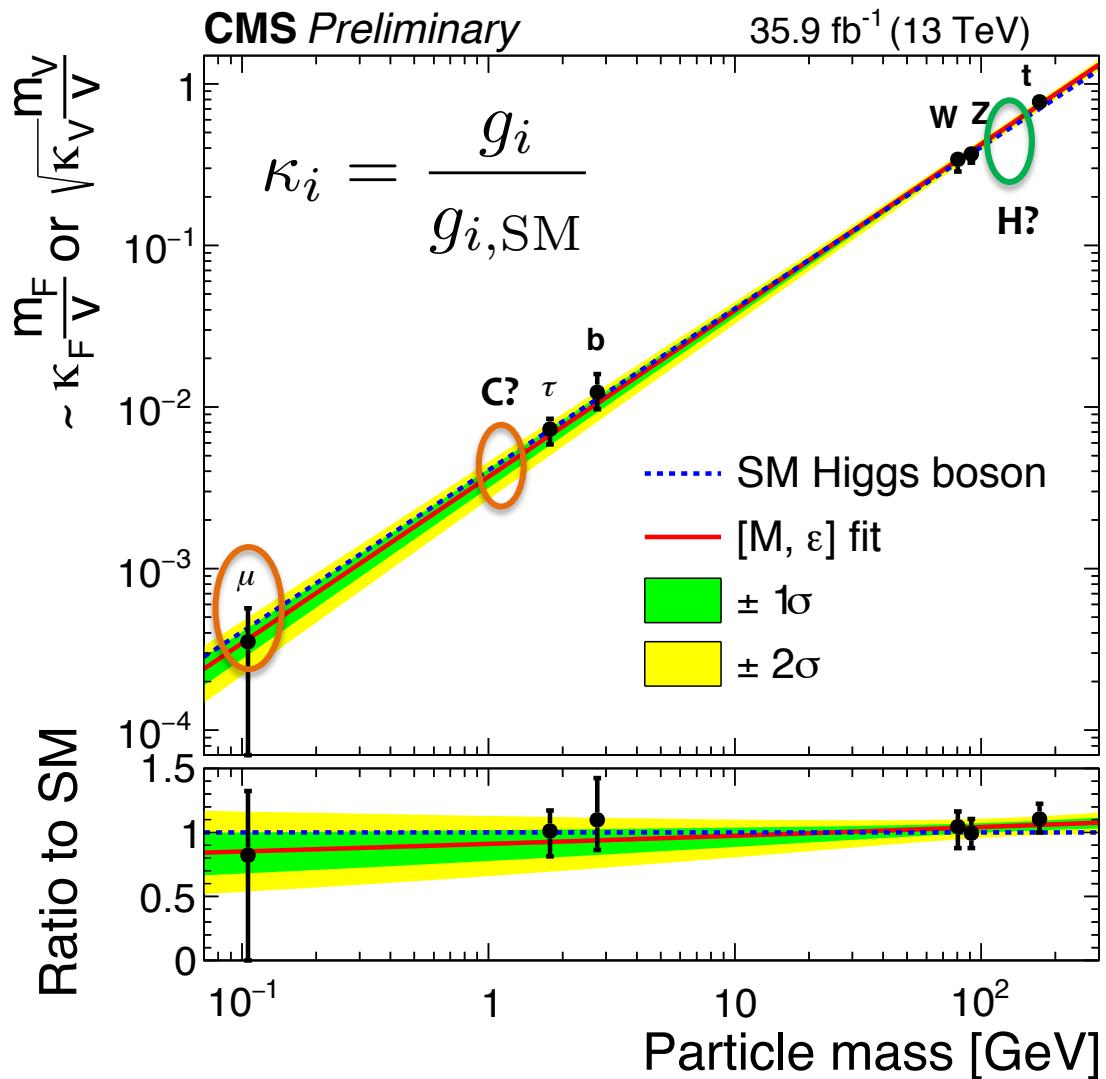
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Haven't observed **self-coupling** of the Higgs boson

- $\text{pp} \rightarrow \text{HH}$  searches severely limited by statistics @LHC



# What don't we know?

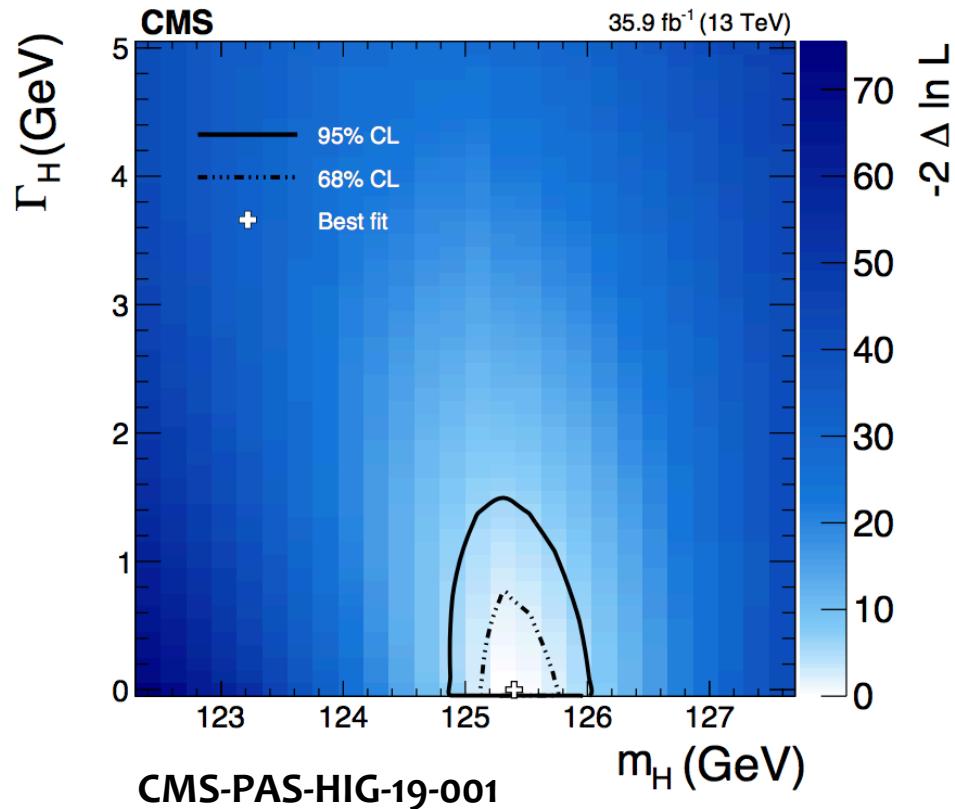
$$\Gamma_{\text{SM}} = 4 \text{ MeV}$$

Relatively large modifications possible from new physics

- Modifications to  $B(\rightarrow \text{SM})$  values
- Additional decays present if NP includes light objects

Width (directly) constrained from  $H \rightarrow 4l$  line-shape

- Large uncertainty with Run-2 measurements
- Plenty of available room from BSM contributions with current data



$\Gamma_H < 1.1 \text{ GeV} @ 95\% \text{ CL}$

# Effective field theory

$$\mathcal{L} = L_{SM} + \sum \frac{c_i}{\Lambda^2} \mathcal{O}_i^{d=6} + \sum \frac{c_i}{\Lambda^4} \mathcal{O}_i^{d=8} + \dots$$

Well-defined theoretical approach to explore physics beyond the SM

- Assume new states are Heavy → only include light (SM) states in effective Lagrangian
- BSM effects show up as deviations of expansion coefficients from 0 in the data

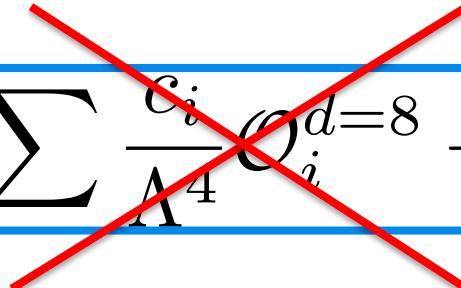
Not specific to Higgs interactions ...

→ Great framework to combine electroweak, top, Higgs measurements!

Matching to UV complete theories ~straightforward ...

→ Complimentary to direct searches for heavy states

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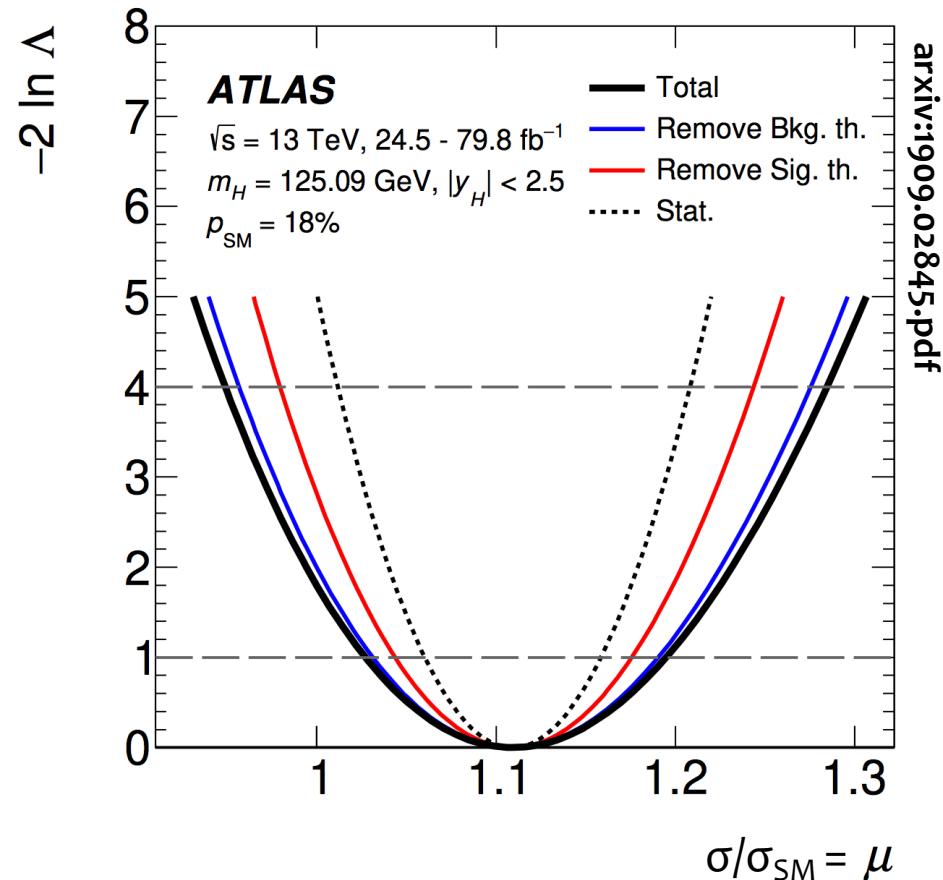
In Higgs, limited sensitivity with Run-2 → Focus on dim-6 operators in experiments

# Overall signal rate

Measuring the inclusive Higgs rate compared to that expected under the SM ( $\mu$ )...

Uncertainty in overall rate < 10%  
with latest combinations

- Dominated by systematic component
- Theoretical uncertainty on signal cross-sections are major component of measurement



$$\mu = 1.11^{+0.09}_{-0.08} = 1.11 \pm 0.05 \text{ (stat.)} {}^{+0.05}_{-0.04} \text{ (exp.)} {}^{+0.05}_{-0.04} \text{ (sig. th.)} \pm 0.03 \text{ (bkg. th.)}$$

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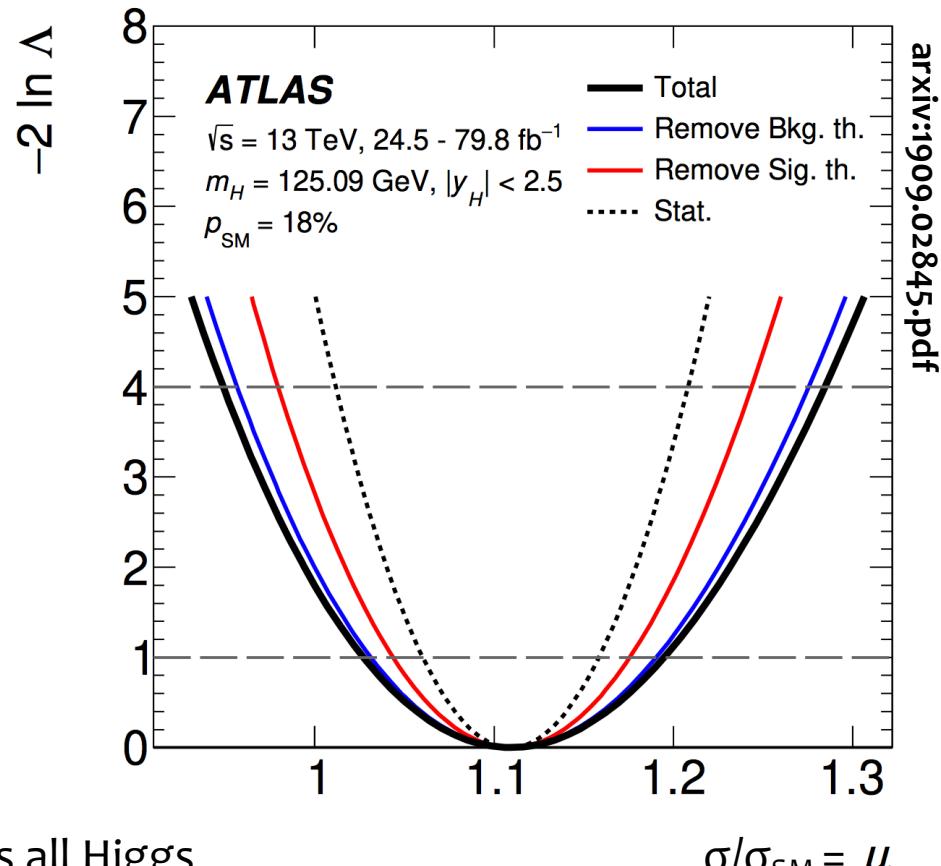
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$$\mathcal{O}_H \sim \frac{c_H}{\Lambda^2} \partial^\mu (\Phi^\dagger \Phi) \partial_\mu (\Phi^\dagger \Phi)$$

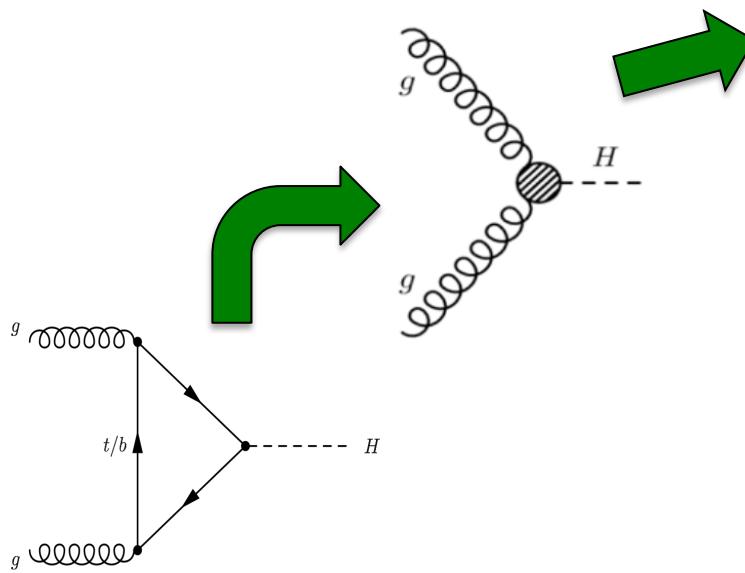

Related to constraint on  $c_H \rightarrow O_H$  rescales all Higgs  
prod/decay processes

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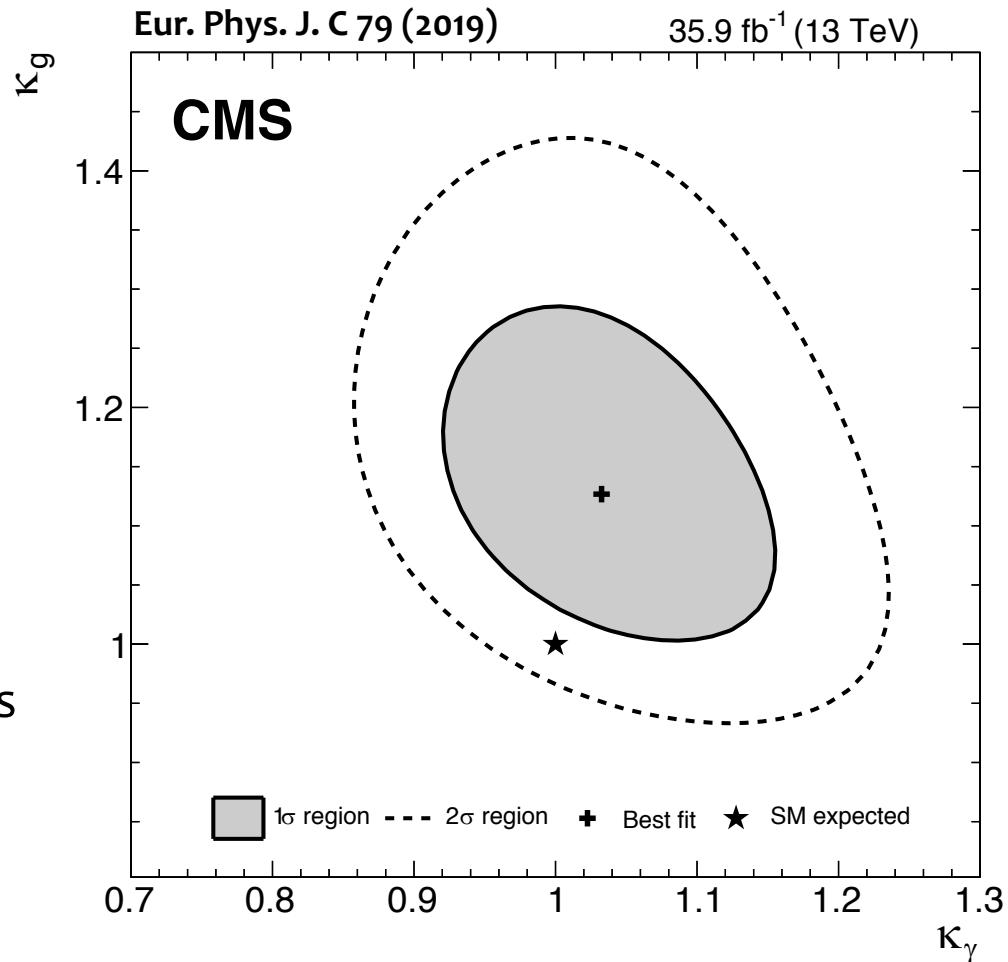


# Effective Couplings ( $\kappa$ -framework)

In LO  $\kappa$ -framework from Run-1, we also relax assumption on loop-induced processes and treat as “**effective couplings**”

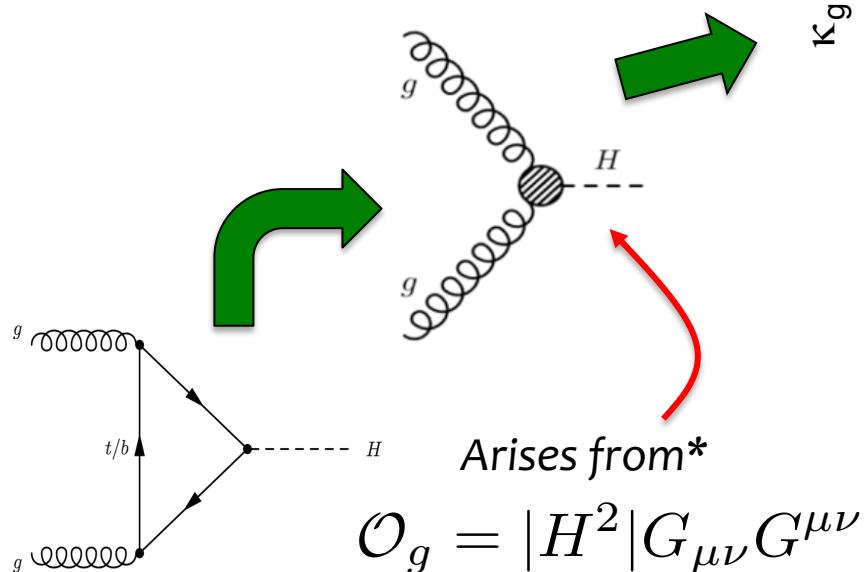


Assume BSM physics at Heavy scales  
→ Accessible only in loops  
→ No additional direct decays  
(  $\text{BR}_{\text{BSM}}=0$  )



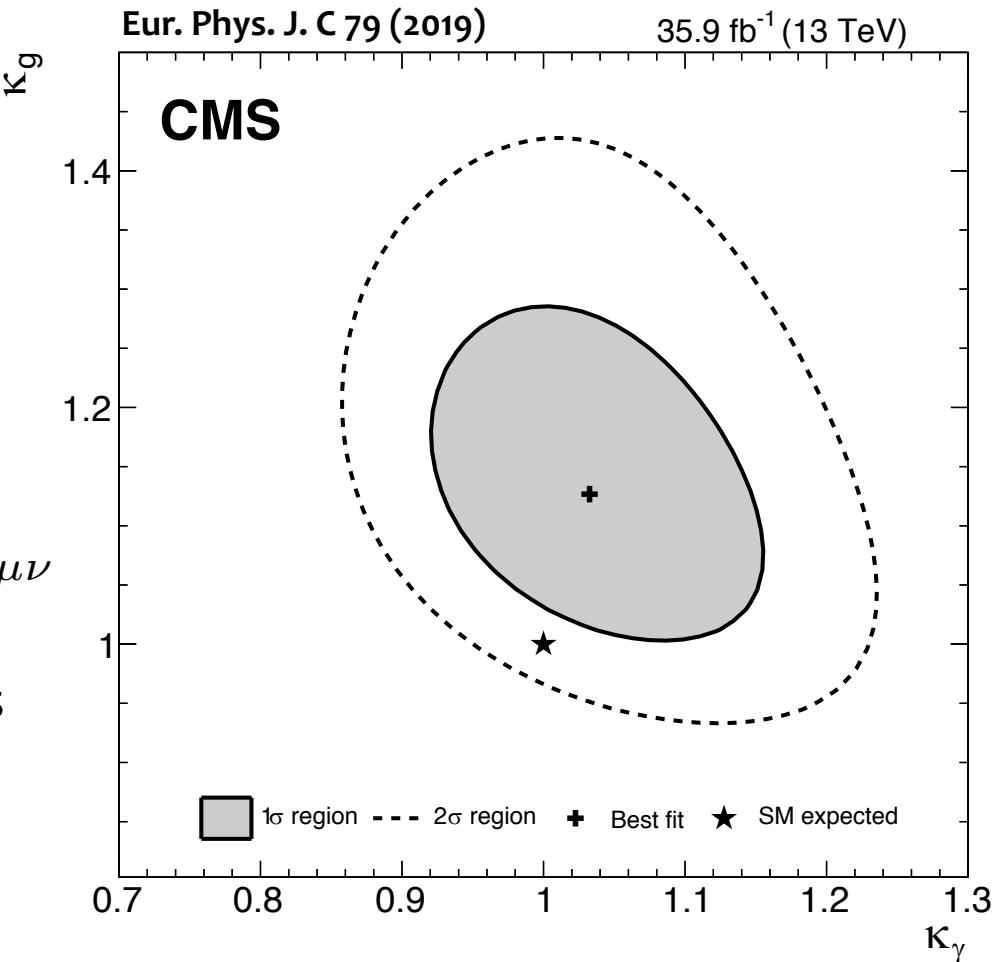
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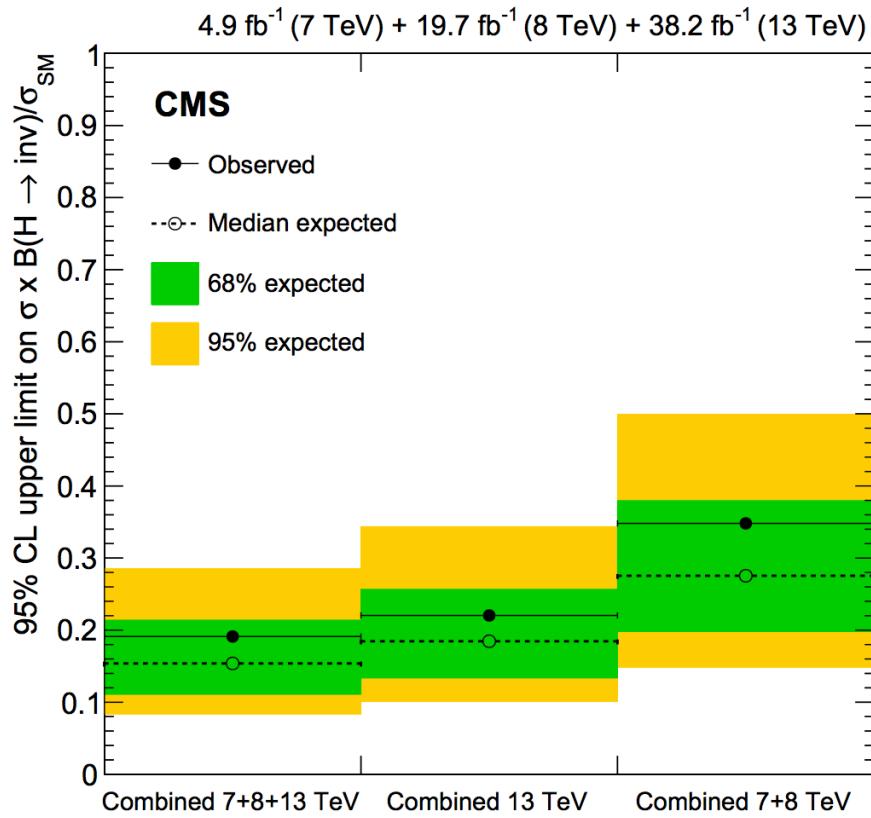
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(\*In SILH basis)



# Effective theories to compare experiments

Searches for  $H \rightarrow$ invisibles can be interpreted under EFTs for  $DM^* \rightarrow$  Higgs portal models [1]

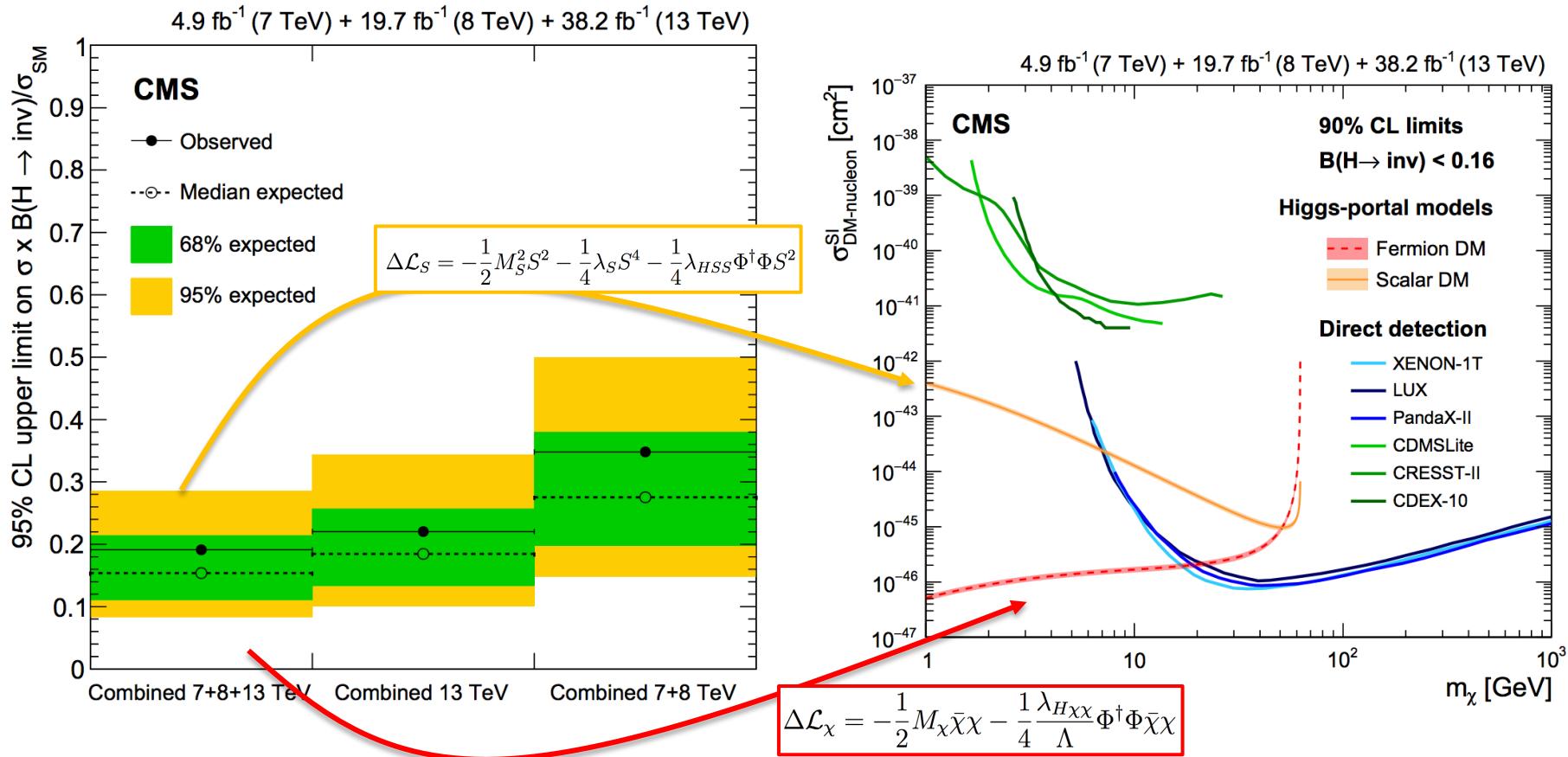


[1] Review in <https://arxiv.org/abs/1903.03616>

\*note this **does** require addition of new light states!

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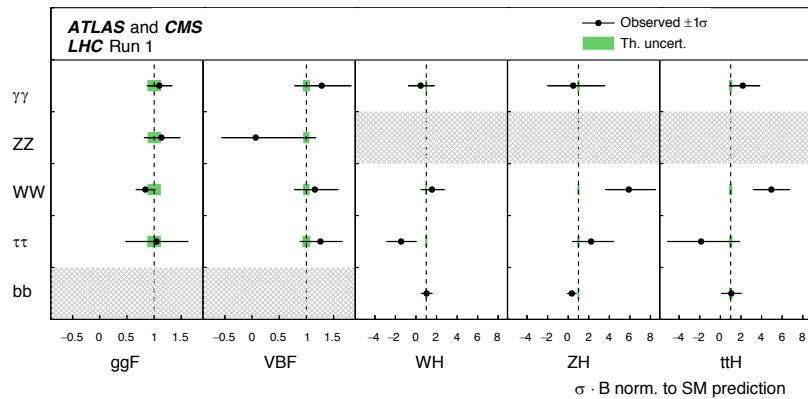
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$$\mathcal{L} = \mathcal{L}_{SM} + \sum \frac{c_i}{\Lambda^2} \mathcal{O}_i^{d=6}$$

On-shell



$$\delta\mu \approx \left(\frac{v}{\Lambda}\right)^2$$

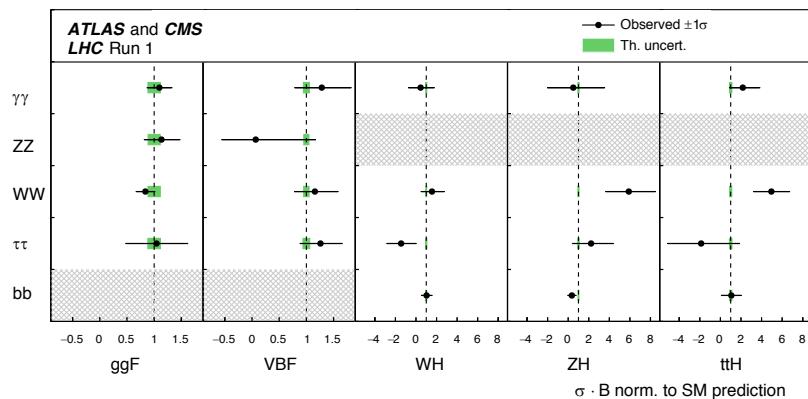
Inclusive  $\mu/\kappa$  : high-precision yields precision on new physics scale

$$\delta\mu = 1\% \rightarrow \Lambda \sim 2.5 \text{ TeV}$$

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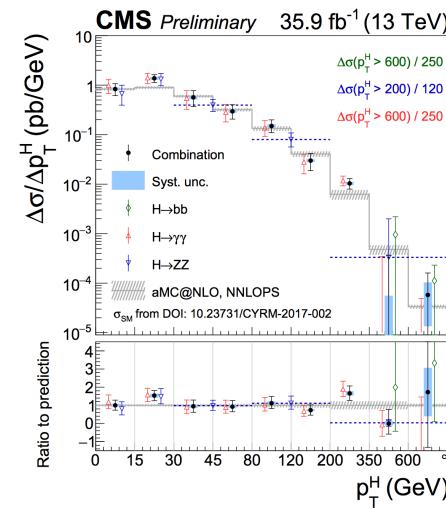


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Off-shell / large  $q^2$



$$\delta\mu \approx \left(\frac{q}{\Lambda}\right)^2$$

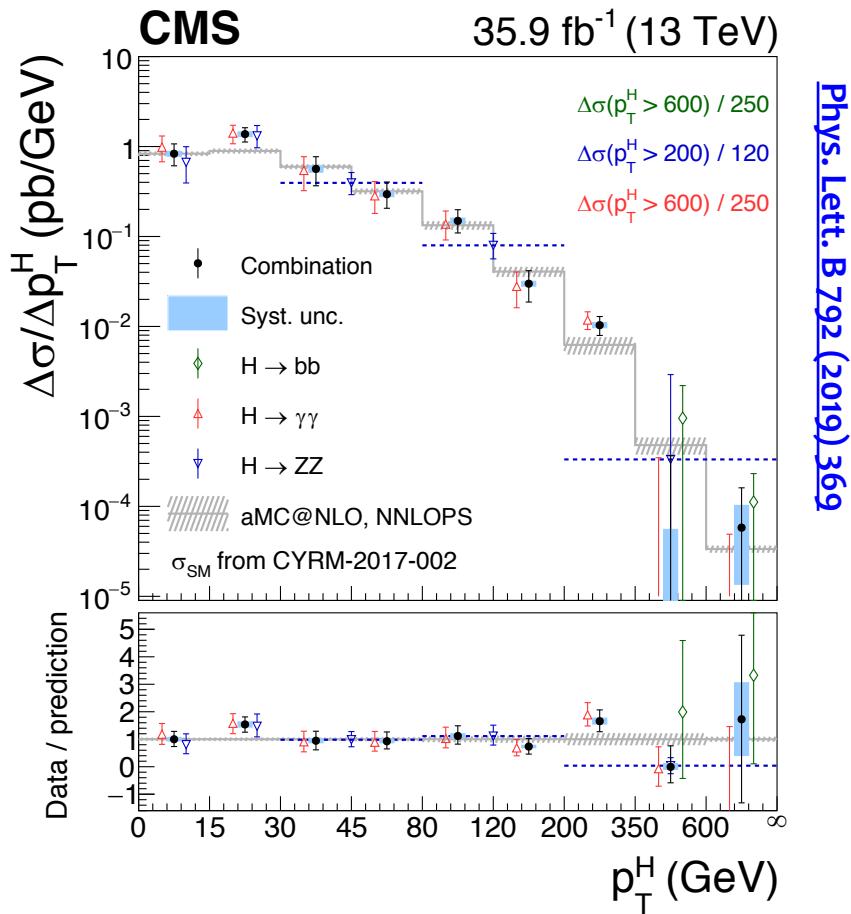
Differential: High momentum production sensitive to new physics

$$\delta\sigma = 15\% (q=1\text{TeV}) \rightarrow \Lambda \sim 2.5 \text{ TeV}$$

Need to use differential measurements to exploit sensitivity at LHC!

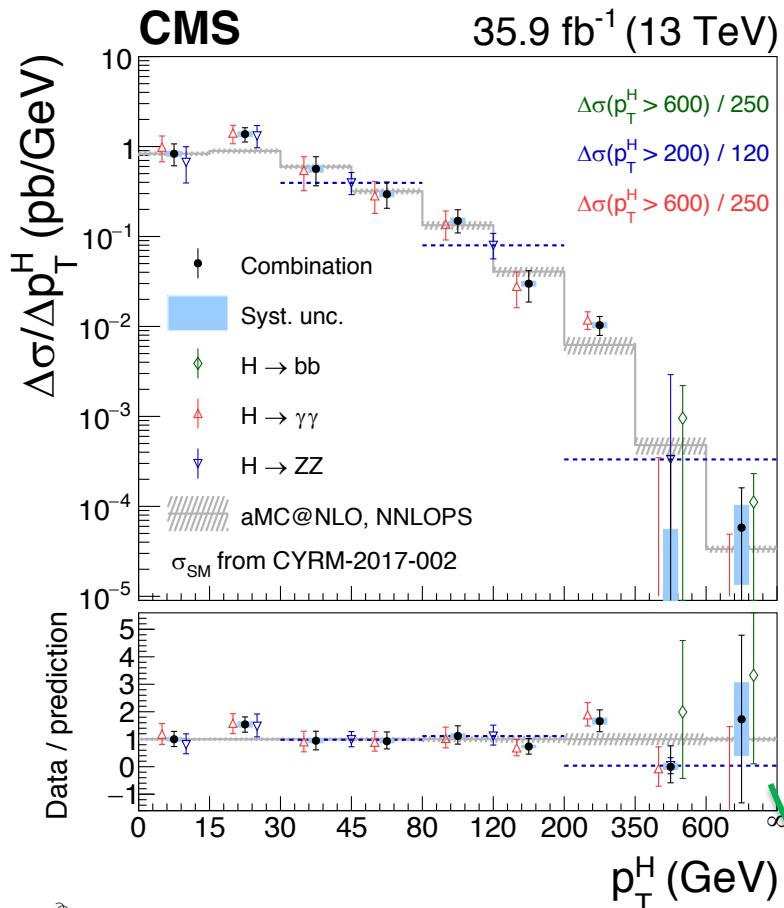
# Differential Higgs

Differential  $p_T(H)$  Combination in  $gg \rightarrow H$   
mode:  $H \rightarrow ZZ + H \rightarrow \gamma\gamma$  (+boosted  $H \rightarrow bb$ )



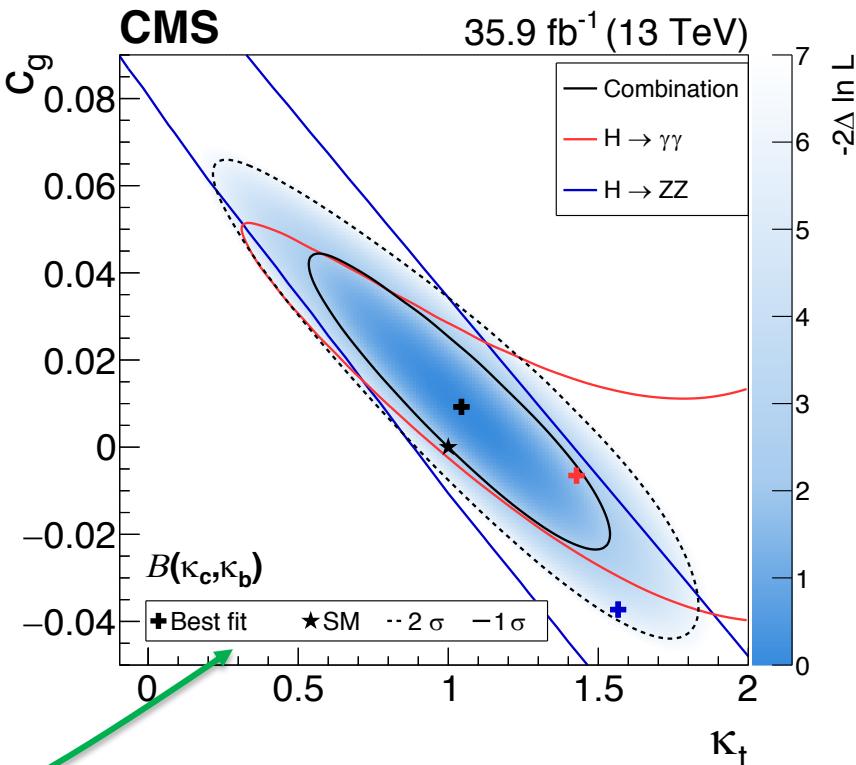
[Phys. Lett. B 792 \(2019\) 369](#)

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Use measurements to constrain Yukawas/Wilson coeffs (following [1])

$\curvearrowleft \mathcal{O}_1 = |H|^2 G_{\mu\nu}^a G^{a,\mu\nu}, \quad \mathcal{O}_2 = |H|^2 \bar{Q}_L H^c u_R + h.c., \quad \mathcal{O}_3 = |H|^2 \bar{Q}_L H d_R + h.c.$

**Yukawas**

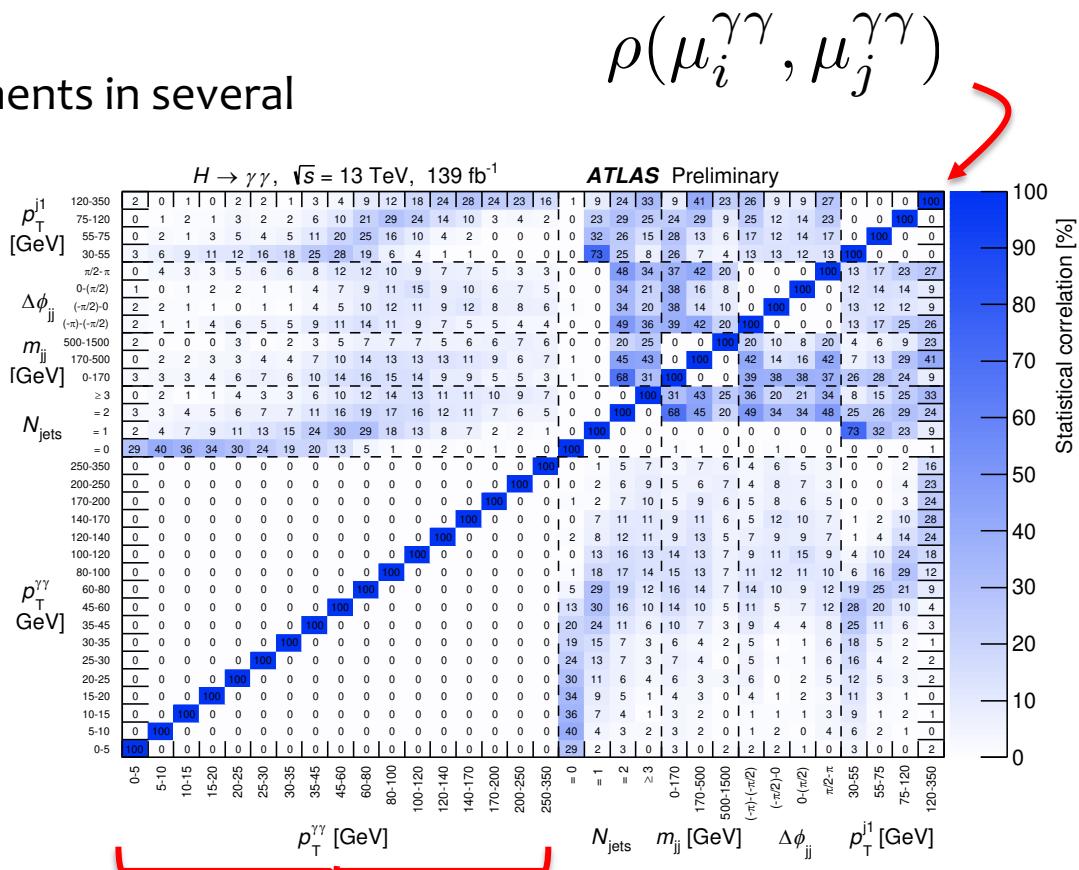
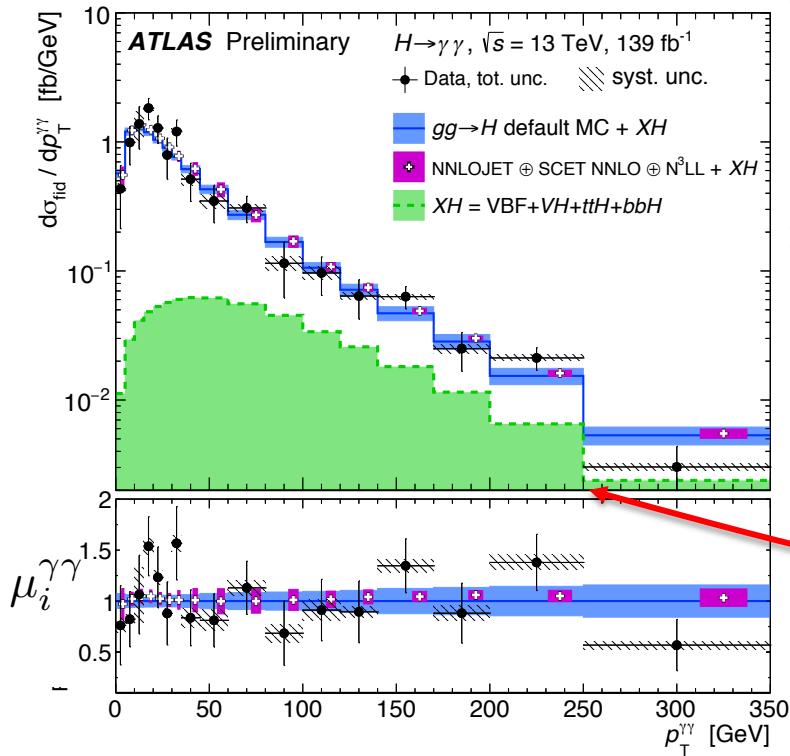
[1] <https://arxiv.org/abs/1705.05143>

# Differential Higgs

Fiducial differential Higgs measurements in several (1D) variables with  $H \rightarrow \gamma\gamma$

- $p_T^{\gamma\gamma}, \Delta\phi_{jj}, p_T^{j1}, m_{jj}, N_{\text{jets}}$

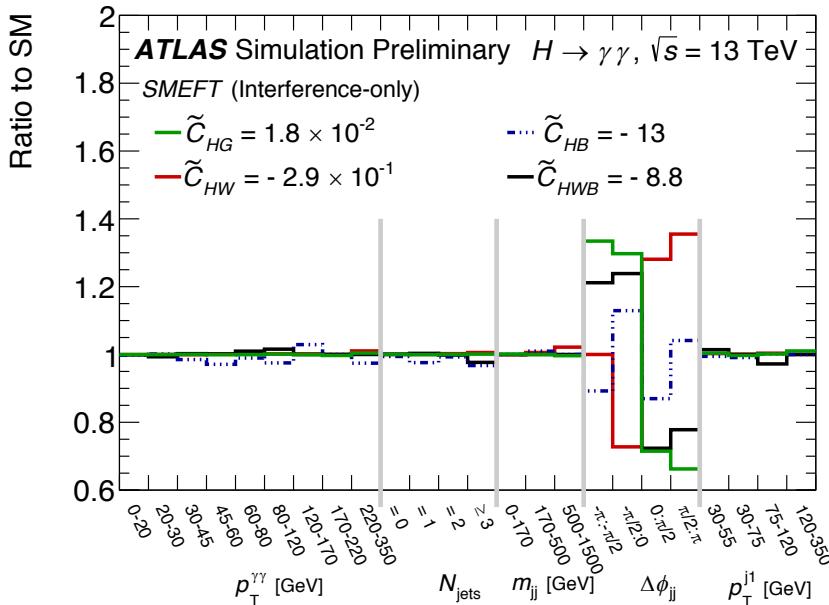
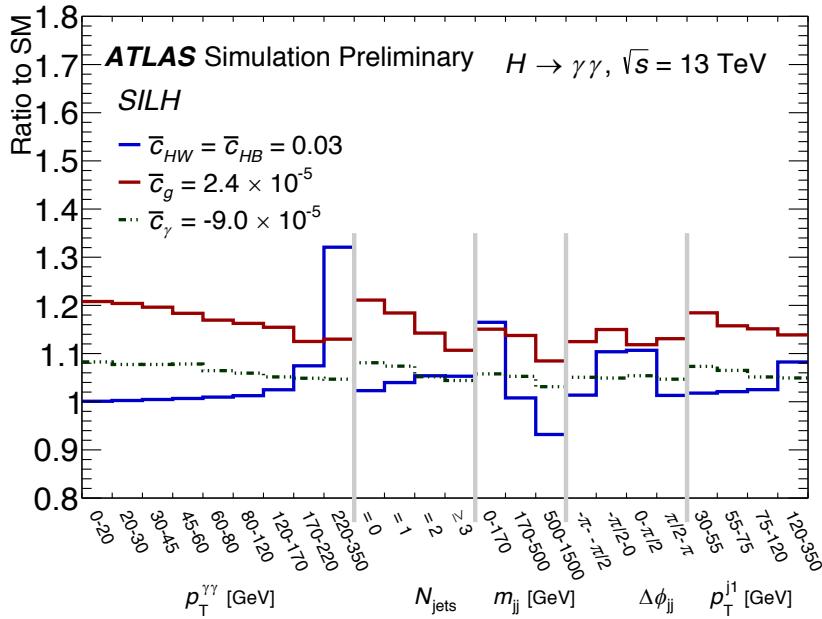
ATLAS-CONF-2019-029



Correlations ( $\rho$ ) between different bins in different distributions calculated using bootstrapping

# Differential Higgs

Deviation from SM expressed in terms of Wilson coefficients  
(using SILH / SMEFT basis @ LO)



Build parameterization of  $\mu_i^{\gamma\gamma}(\bar{c}, \tilde{c})$  using many MC points + professor algorithm [1] and construct log-likelihood function ...

$$-2 \log L = (\hat{\mu}^{\gamma\gamma} - \mu^{\gamma\gamma}(\bar{c}, \tilde{c}))^T V^{-1} (\hat{\mu}^{\gamma\gamma} - \mu^{\gamma\gamma}(\bar{c}, \tilde{c}))$$

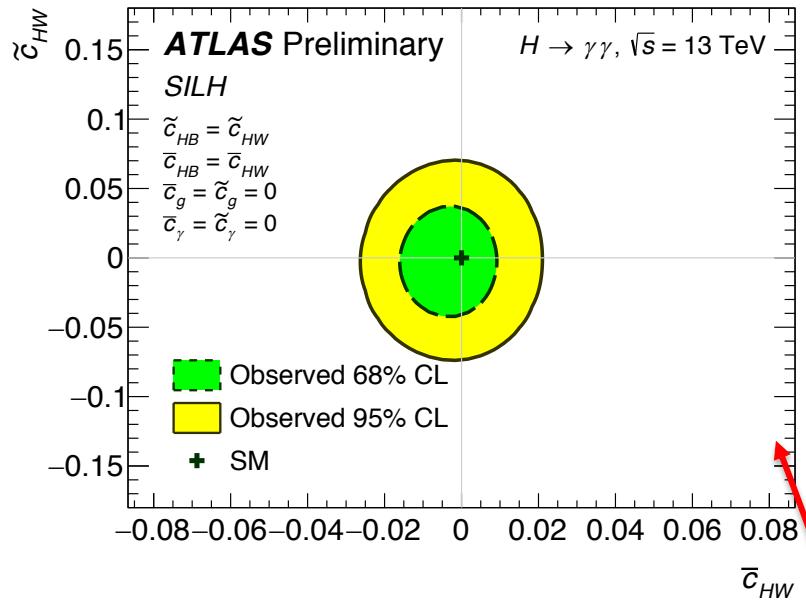
[1] <https://arxiv.org/pdf/0907.2973.pdf>

$$V_{i,j} = \rho(\mu_i^{\gamma\gamma}, \mu_i^{\gamma\gamma}) \sigma_{\mu_i^{\gamma\gamma}} \sigma_{\mu_j^{\gamma\gamma}}$$

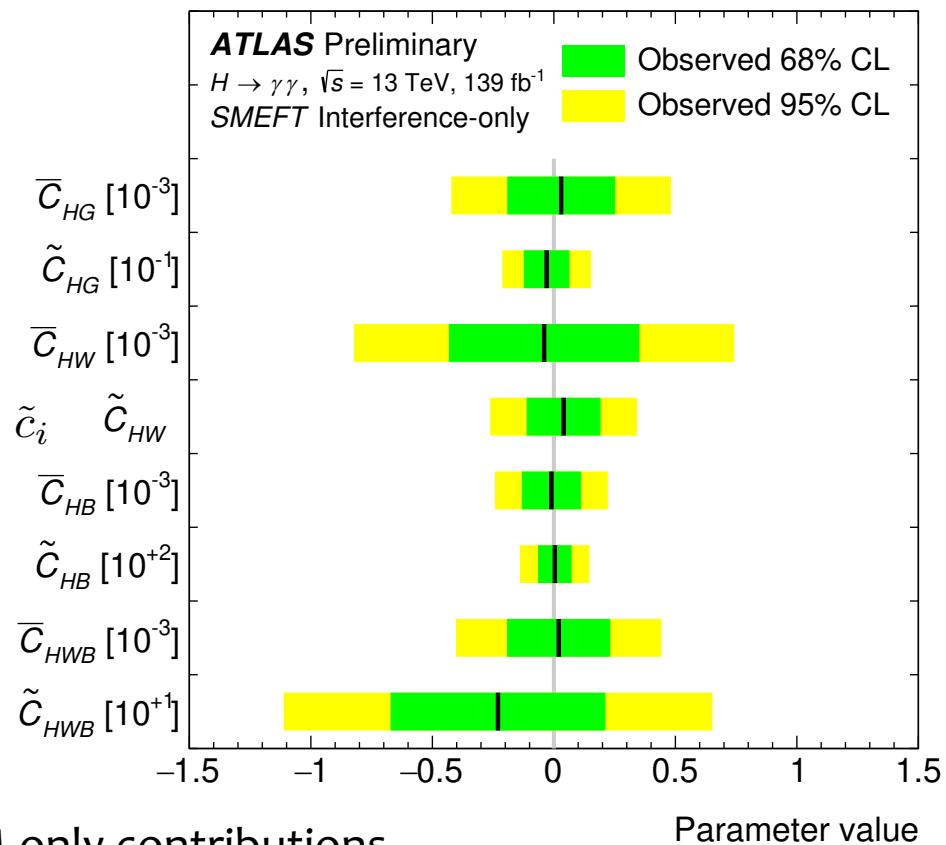
# Differential Higgs

Minimize log-likelihood to extract constraints on Wilson coefficients

→ Limits extracted fixing other coeffs to 0



Sensitivity to both CP-even and CP-odd contributions through inclusion of distributions like  $\Delta\varphi_{jj}$



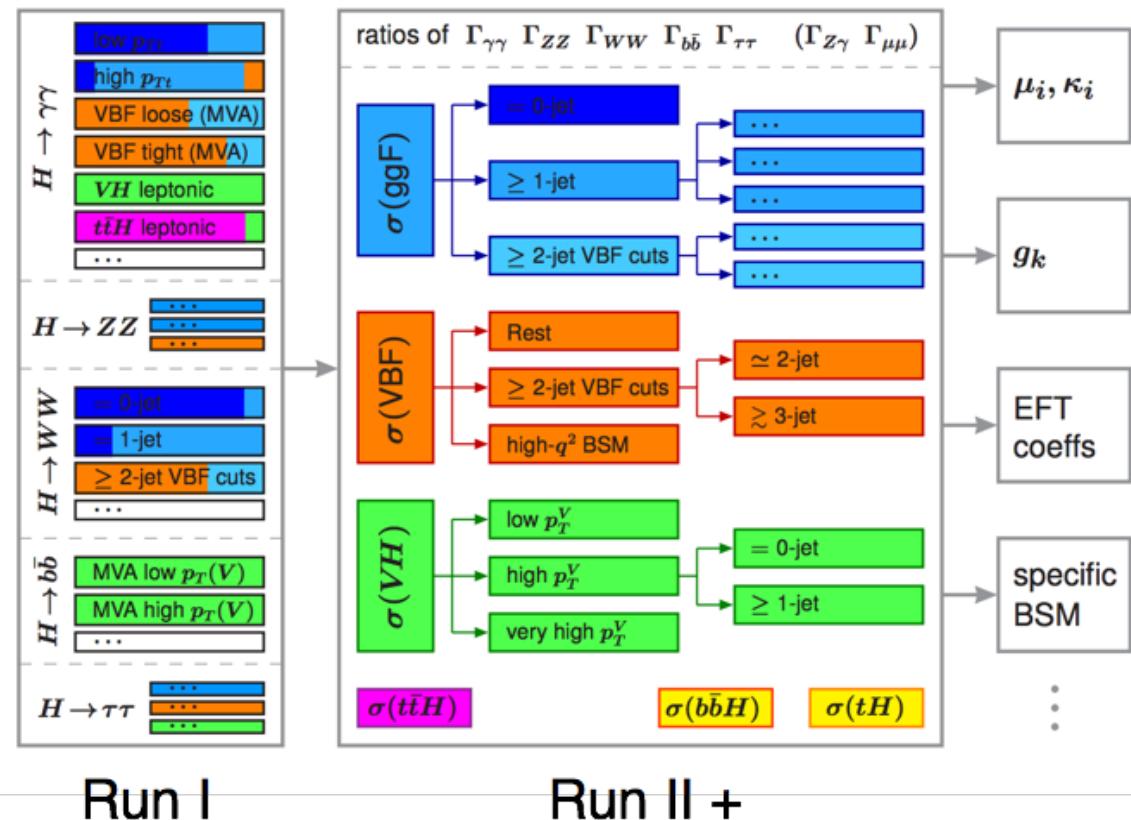
Also studied effects of including or not BSM-only contributions

→ Constraints much greater when including BSM only contributions for CP-odd operators  $\sim \tilde{C}_i$

# Simplified Template X-Sections (STXS)

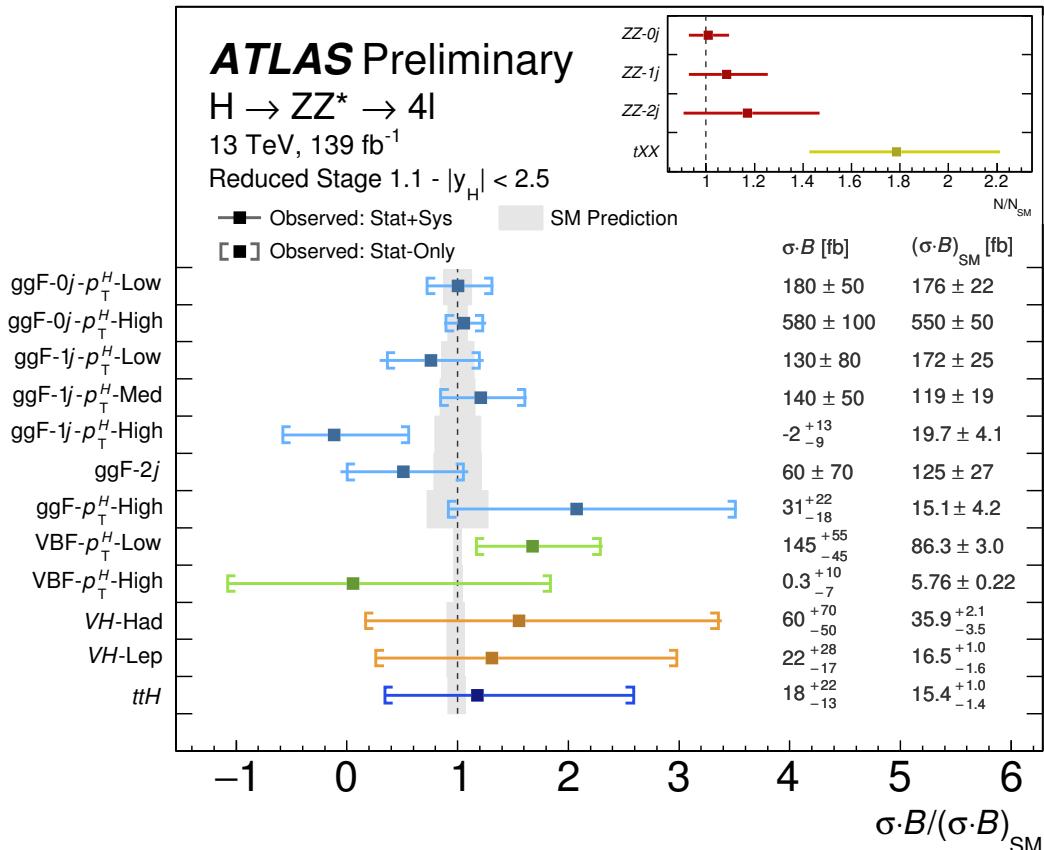
For Run-2 ATLAS/CMS have additional produced Higgs results in terms of STXS measurements

- Introduce differential binning but maintain process distinction
- Binning chosen to decouple theory uncertainty from measurement and maintain experimental sensitivity
- In principle, measurements can be (re)interpreted under generic BSM scenarios  $\rightarrow \kappa$ -framework, EFTs ...



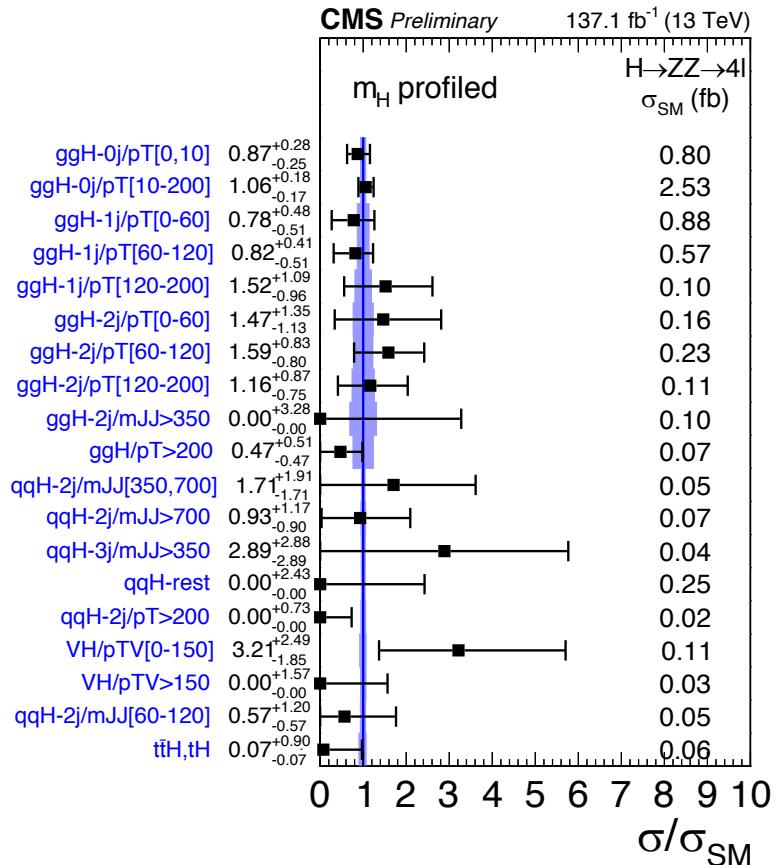
# Simplified Template X-Sections (STXS)

ATLAS-CONF-2019-005



ATLAS/CMS have results using full Run-2 datasets (several more STXS measurements with partial Run-2 also)

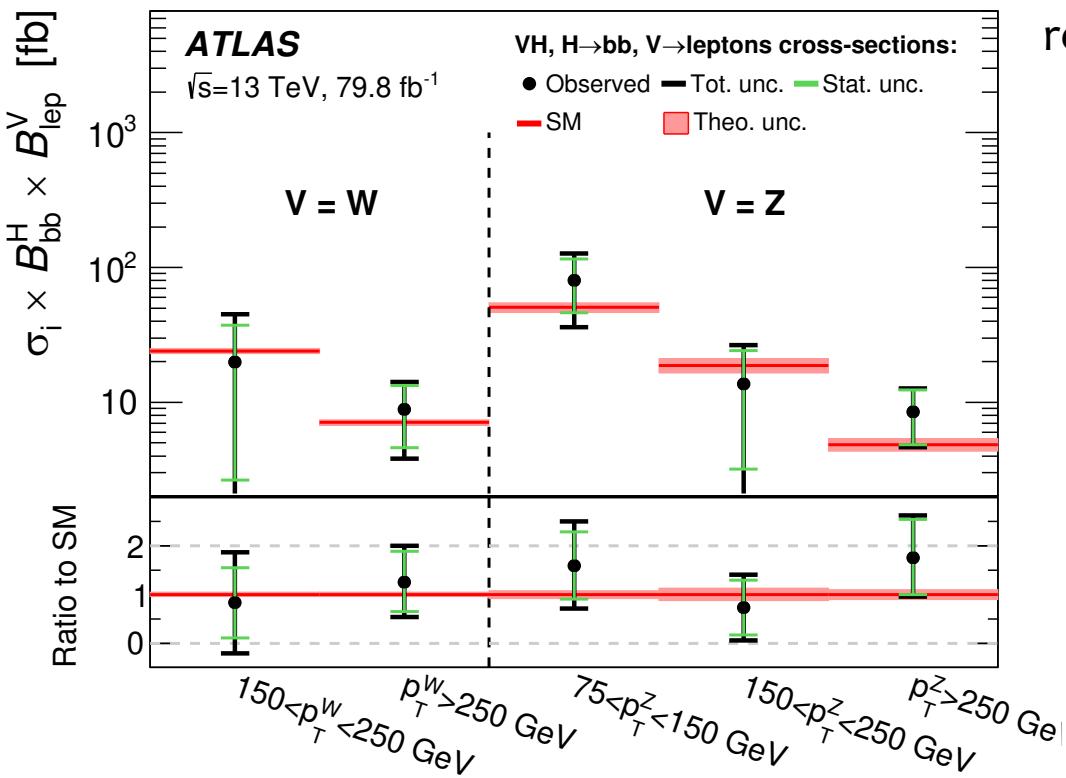
CMS-PAS-HIG-19-001



# STXS VH( $\rightarrow$ bb)

ATLAS V( $\rightarrow$ lep)H( $\rightarrow$ bb) cross-sections measured in bins of  $p_T(V)$

[JHEP 05 \(2019\) 141](#)



Include  $O_d = y_d |H|^2 \bar{Q}_L H d_R$  for H $\rightarrow$ bb branching ratio

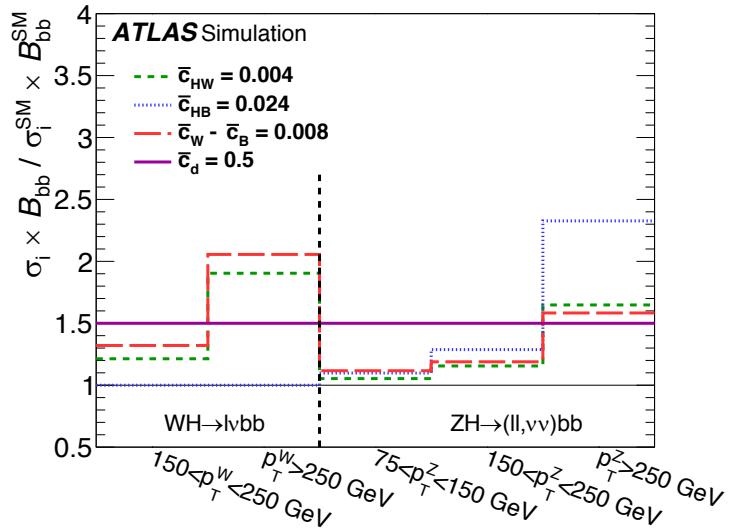
Parameterise\* VH cross-sections, relative to SM, in terms of

$$O_{HW} = i (D^\mu H)^\dagger \sigma^a (D^\nu H) W_{\mu\nu}^a,$$

$$O_{HB} = i (D^\mu H)^\dagger (D^\nu H) B_{\mu\nu},$$

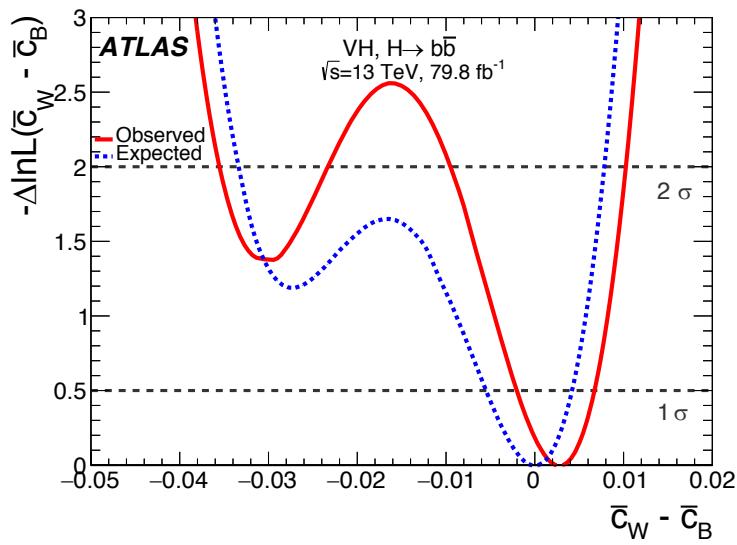
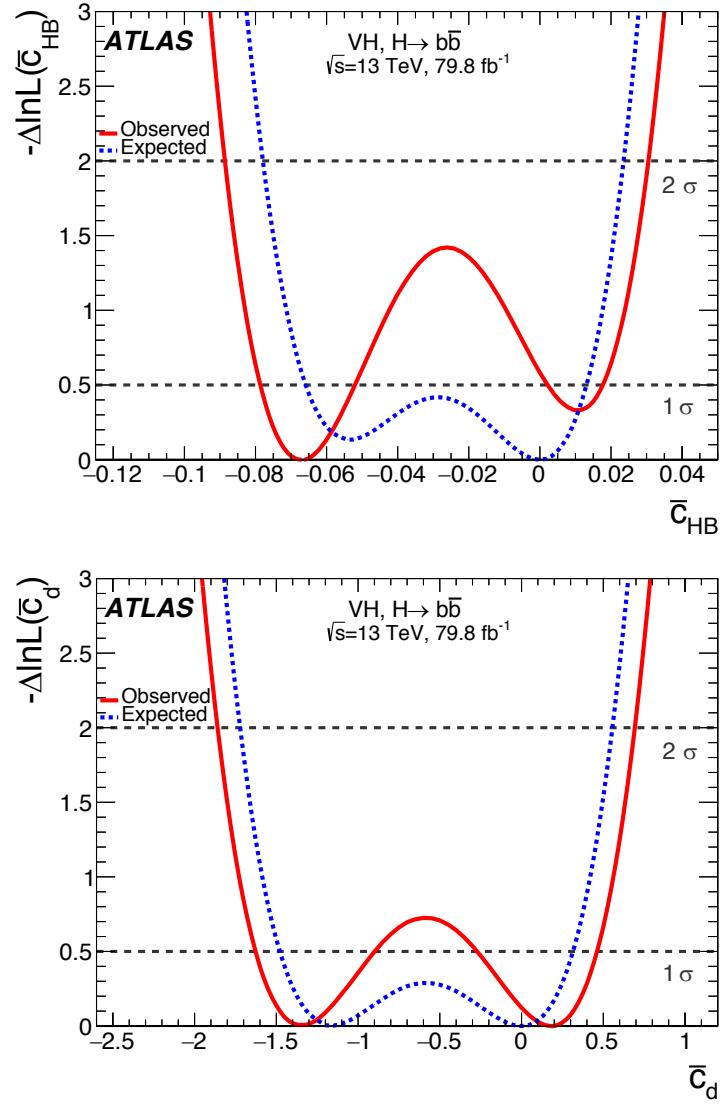
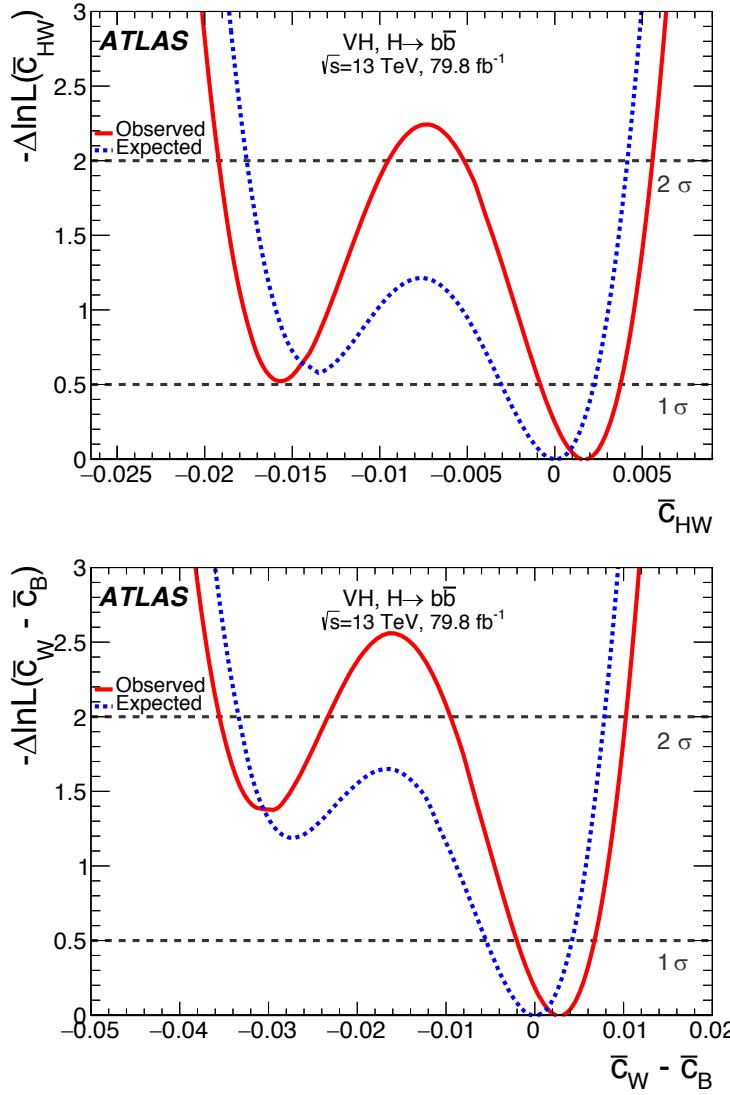
$$O_W = \frac{i}{2} \left( H^\dagger \sigma^a \overset{\leftrightarrow}{D}{}^\mu H \right) D^\nu W_{\mu\nu}^a,$$

$$O_B = \frac{i}{2} \left( H^\dagger \overset{\leftrightarrow}{D}{}^\mu H \right) \partial^\nu B_{\mu\nu}.$$



\*Based on LHCHXSWG-INT-2017-001

# STXS VH( $\rightarrow b\bar{b}$ )



# STXS combinations

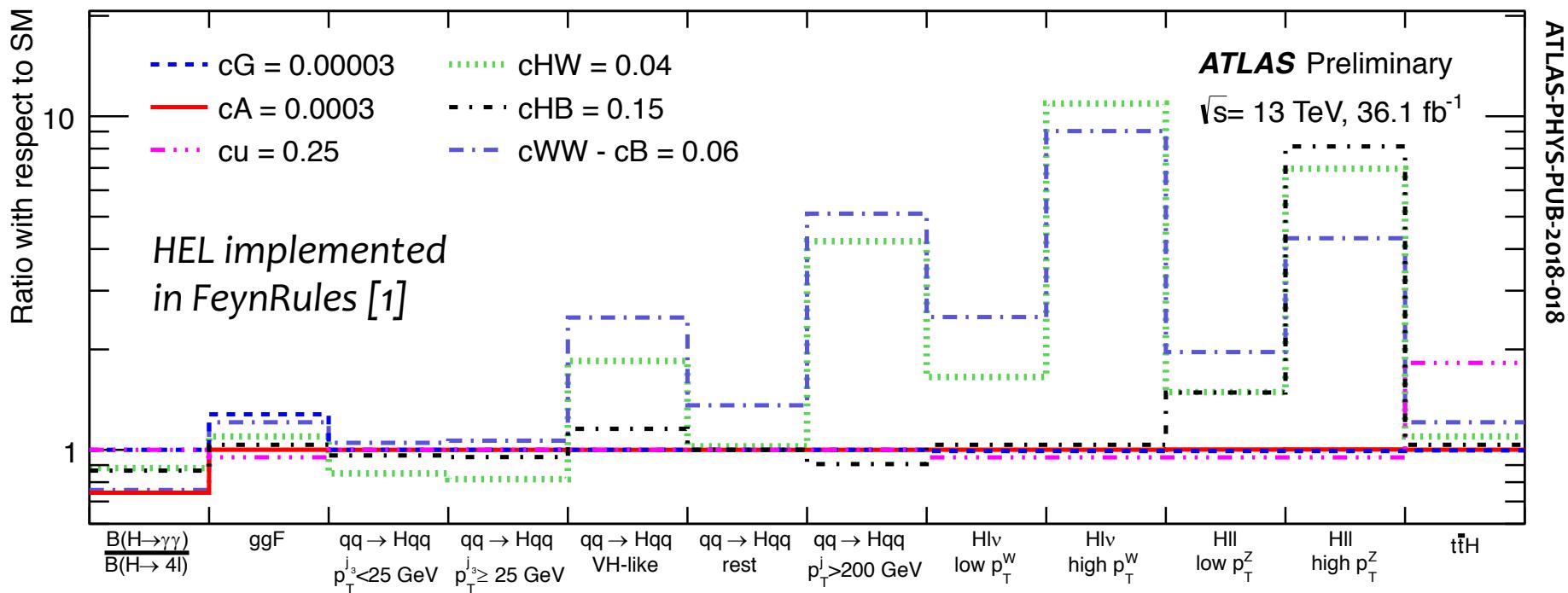
Common "fiducial" definitions for each STXS bin

→ Parametrise STXS scaling as,

$$\mu = \frac{\sigma}{\sigma_{\text{SM}}} = 1 + \sum_i A_i \bar{c}_i + \sum_{ij} B_{ij} \bar{c}_i \bar{c}_j$$

Calculated from MC  
with MadGraph+STXS  
Rivet routine

→ STXS includes ratios of Branching ratios to allow combinations of decay channels



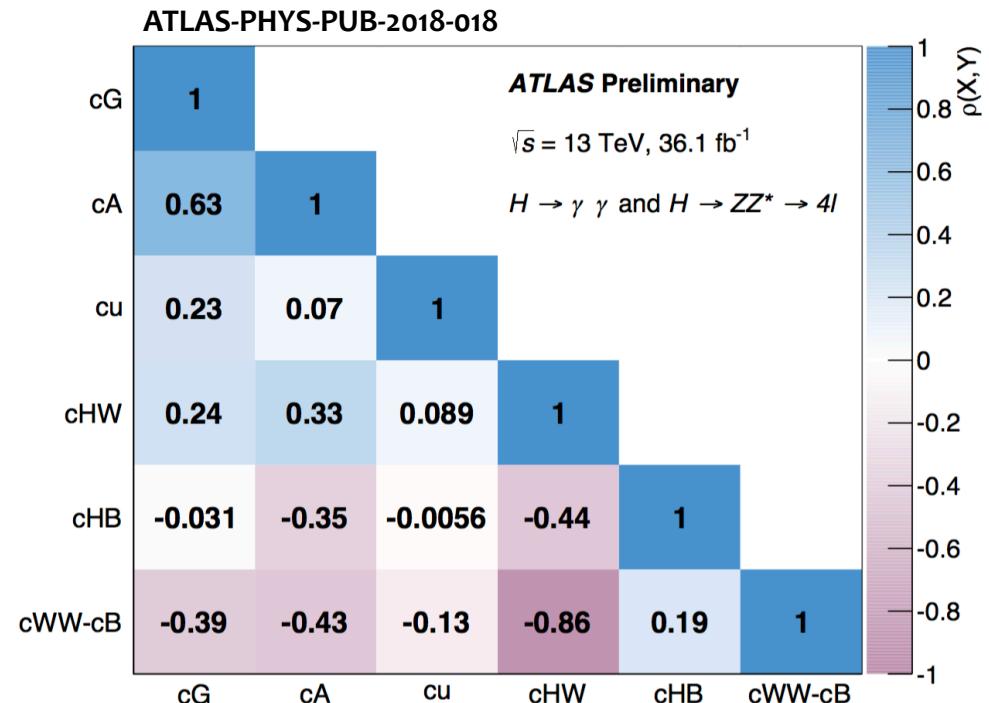
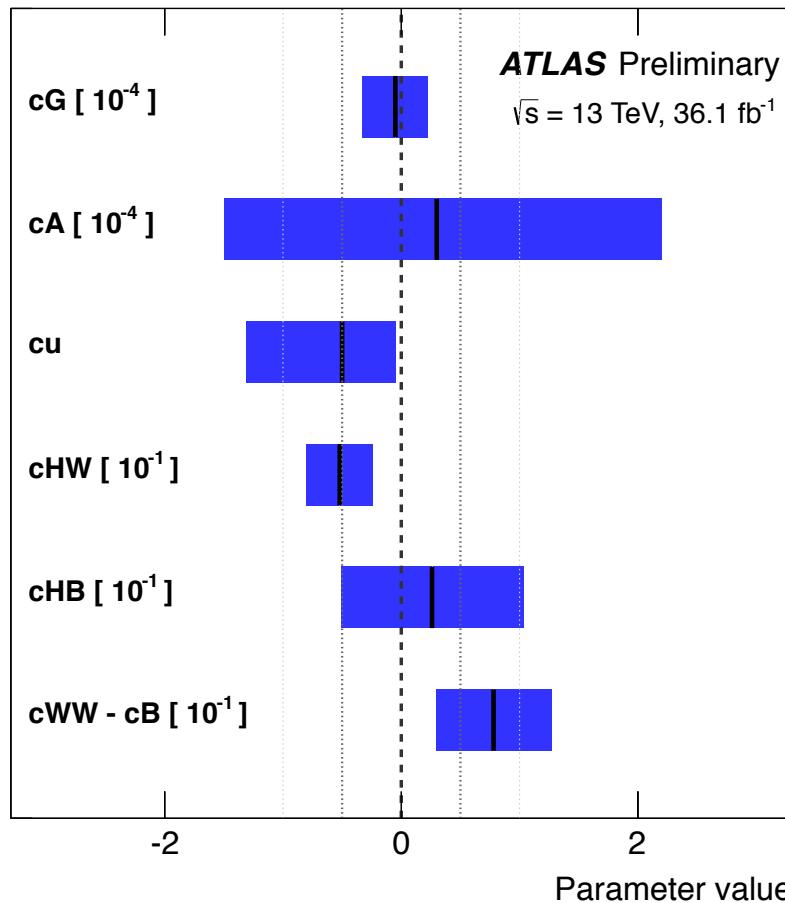
Similar studies @ CMS (See talk from J. Langford)

[1] [JHEP 1404 \(2014\) 110](#)

# STXS combinations

Combinations allow for exploring more EFT operators, simultaneously, without introducing (strong) degeneracies between coefficients

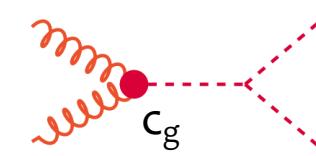
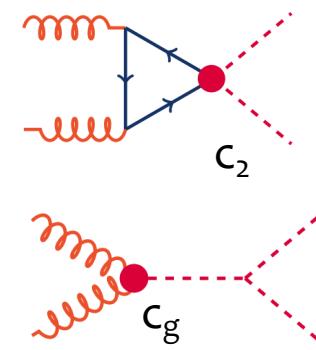
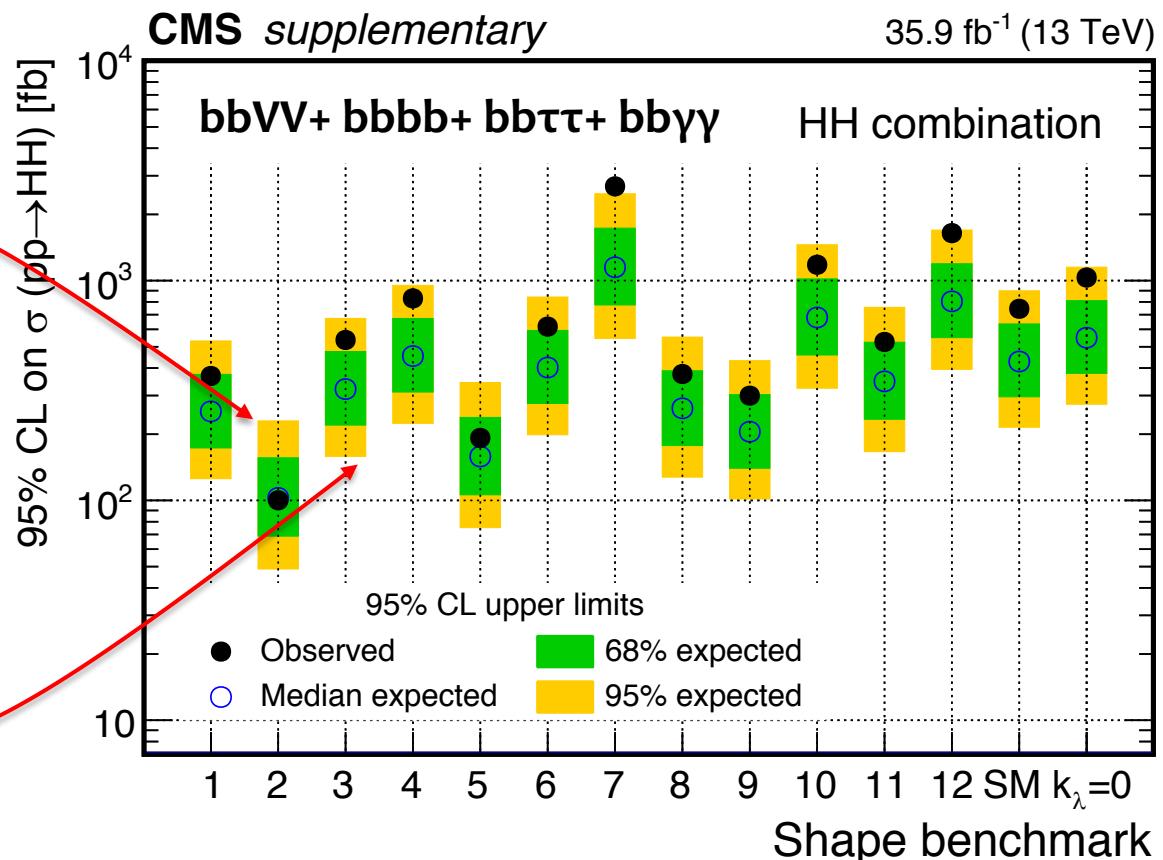
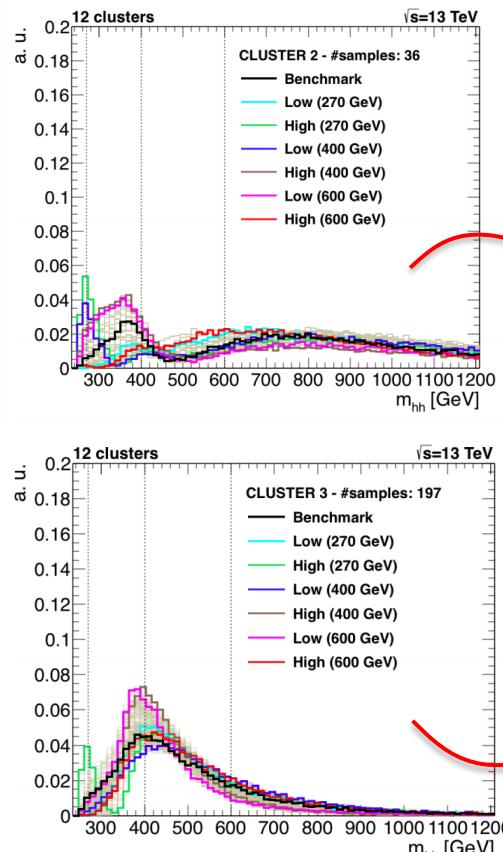
Observed HEL constraints with  $H \rightarrow ZZ^*$  and  $H \rightarrow \gamma\gamma$



# Double Higgs

Limits on HH production expressed under 12 different EFT scenarios varying 5 effective couplings ...  $c_g$ ,  $c_2$ ,  $c_{2g}$ ,  $y_t$ ,  $k_\lambda$

→ Benchmarks picked from clusters of scenarios [1] which share kinematic properties (rather than full 5D exploration)



# Self-coupling

“Higgs potential has only been observed as a formula in text-books” → no direct experimental confirmation ...

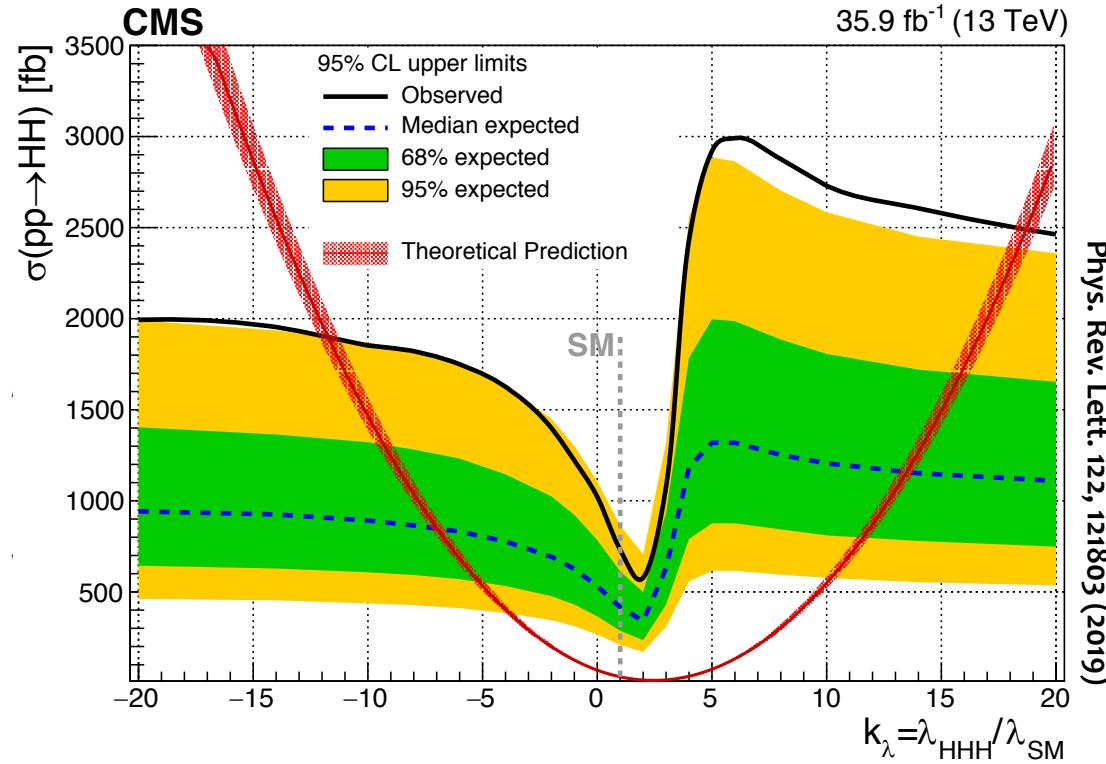
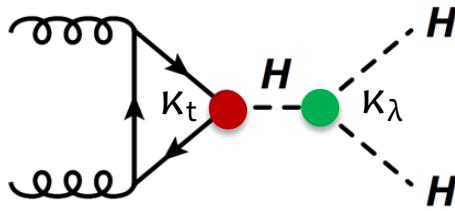
$$V(H) = \frac{m_H^2}{2} H^2 + \lambda_3 v H^3 + \lambda_4 H^4$$

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Double Higgs production  
sensitive to Higgs self-  
coupling  $K_\lambda$

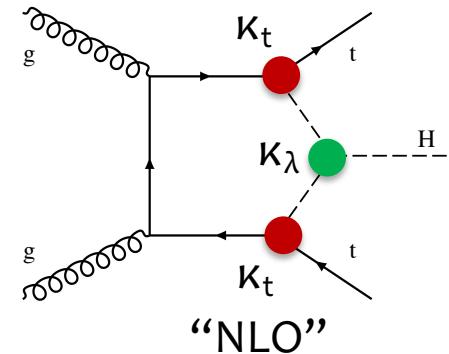
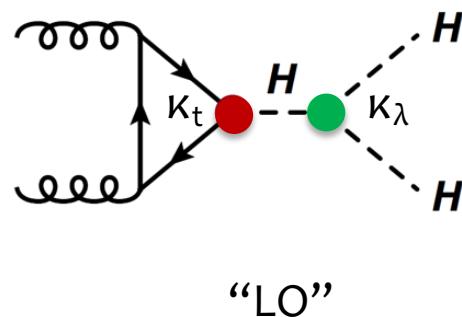


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$$V(H) = \frac{m_H^2}{2} H^2 + \lambda_3 v H^3 + \lambda_4 H^4$$

Double Higgs **and single Higgs** processes can be expressed in terms of effective coupling modifier  $K_\lambda$

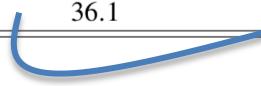


# Self-coupling

ATLAS-PHYS-PUB-2019-009

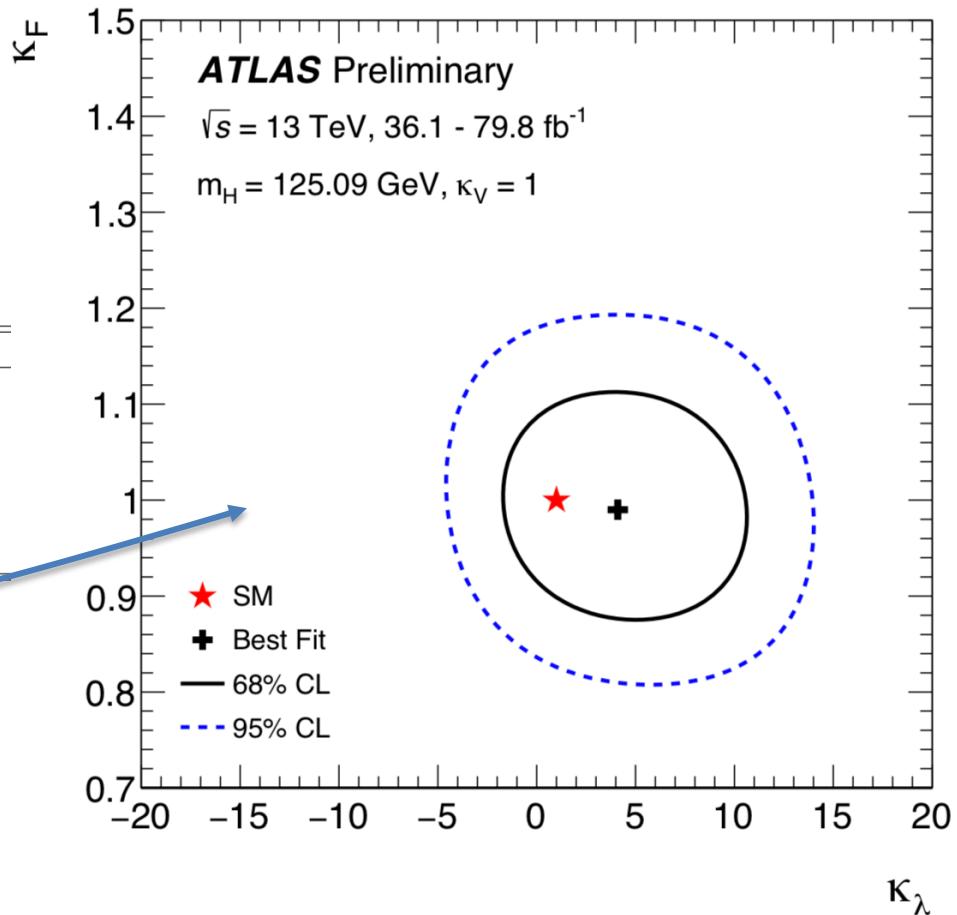
Re-interpretation of single Higgs production measurements in terms of Higgs self-coupling and H-fermion coupling (following [1-2])

| Analysis                                                                                                 | Integrated luminosity ( $\text{fb}^{-1}$ ) |
|----------------------------------------------------------------------------------------------------------|--------------------------------------------|
| $H \rightarrow \gamma\gamma$ (including $t\bar{t}H$ , $H \rightarrow \gamma\gamma$ )                     | 79.8                                       |
| $H \rightarrow ZZ^* \rightarrow 4\ell$ (including $t\bar{t}H$ , $H \rightarrow ZZ^* \rightarrow 4\ell$ ) | 79.8                                       |
| $H \rightarrow WW^* \rightarrow e\nu\mu\nu$                                                              | 36.1                                       |
| $H \rightarrow \tau\tau$                                                                                 | 36.1                                       |
| $VH$ , $H \rightarrow b\bar{b}$                                                                          | 79.8                                       |
| $t\bar{t}H$ , $H \rightarrow b\bar{b}$ and $t\bar{t}H$ multilepton                                       | 36.1                                       |



Introducing additional degrees of freedom leads to large degeneracies in the fit.

Can we combine single + double Higgs measurements under EFT?



[1] Eur. Phys. J. C77 (2017) 887, [2] JHEP 12 (2016) 080

# Self-coupling

ATLAS-PHYS-PUB-2019-009

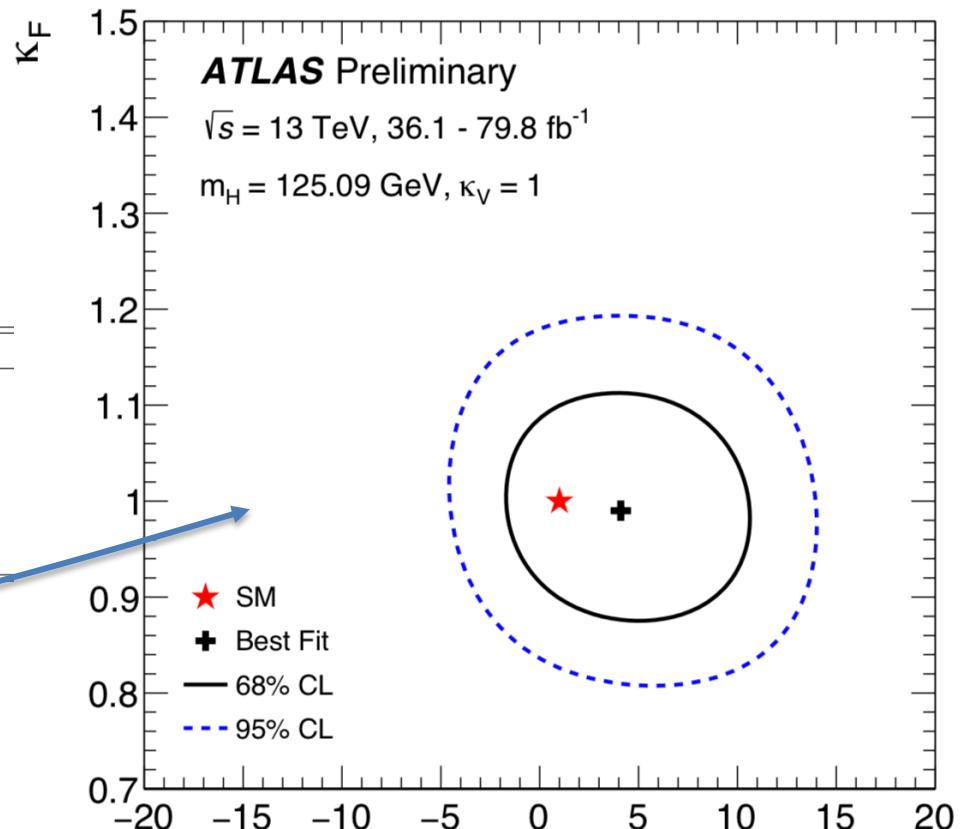
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Introducing additional degrees of freedom leads to large degeneracies in the fit.

Can we combine single + double Higgs measurements under EFT?

→ Yes but we need NLO implementation  
(SMEFTsim @ NLO?)



$$\frac{\bar{c}_6 \lambda}{v^2} [\Phi^\dagger \Phi]^3$$

[1] Eur. Phys. J. C77 (2017) 887, [2] JHEP 12 (2016) 080

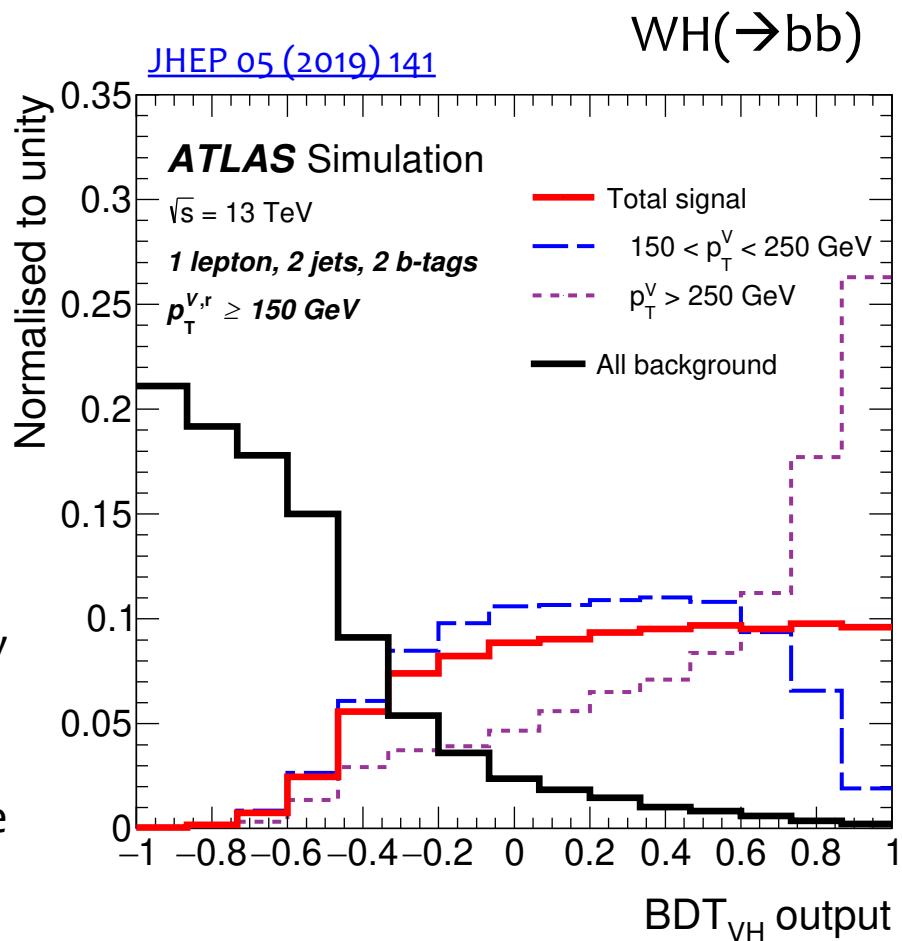
# EFT Caveats 1.

STXS (and differential measurements) are great for EFT interpretations ...

However, need to be careful when extrapolating to “fiducial regions”

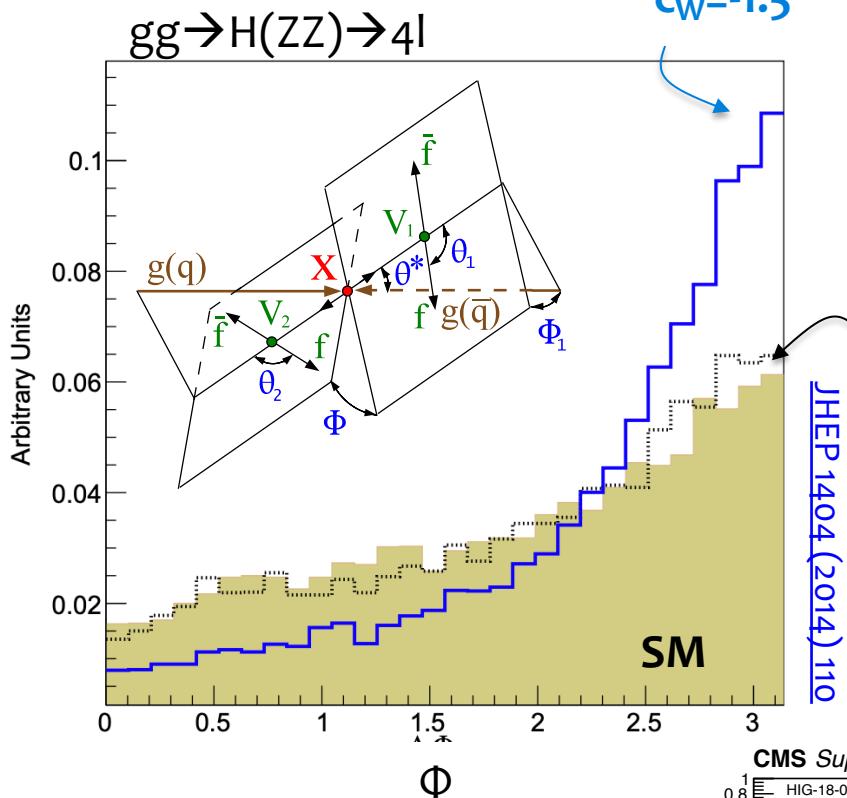
→ Often experiments rely on SM templates\* to determine efficiencies/acceptances etc

→ Furthermore, experiments often employ sophisticated categorization methods (eg BDTs) to perform measurements. What happens to  $\text{WH}_{[150 < p_T < 250]}(\text{BDT})$  if we include dim-6 operators?



\* The “T” in STXS is exactly this!

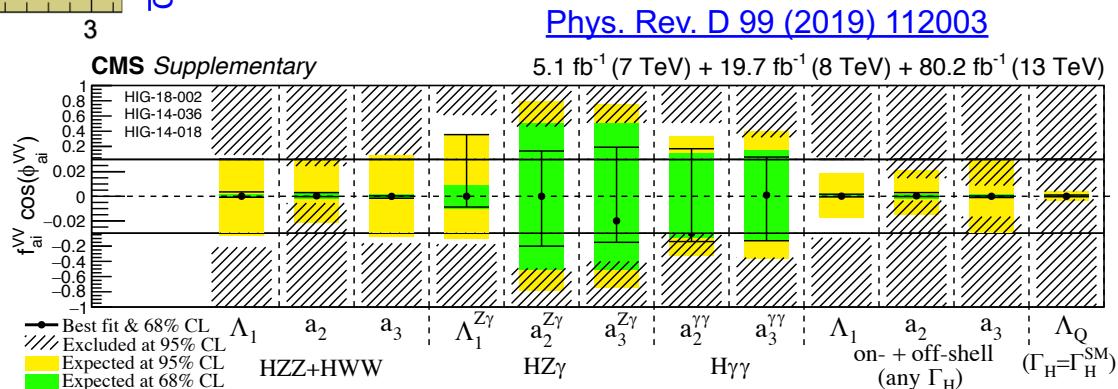
# EFT Caveats 2.



Can STXS be extended - eg with Pseudo Observables?\*

STXS (and combinations of differential measurements) don't include relevant information about decay of the Higgs

- Angular information (eg in 4l final state) sensitive to BSM effects
- ALTAS/CMS use MELA/BDT to exploit this information – with SM templates



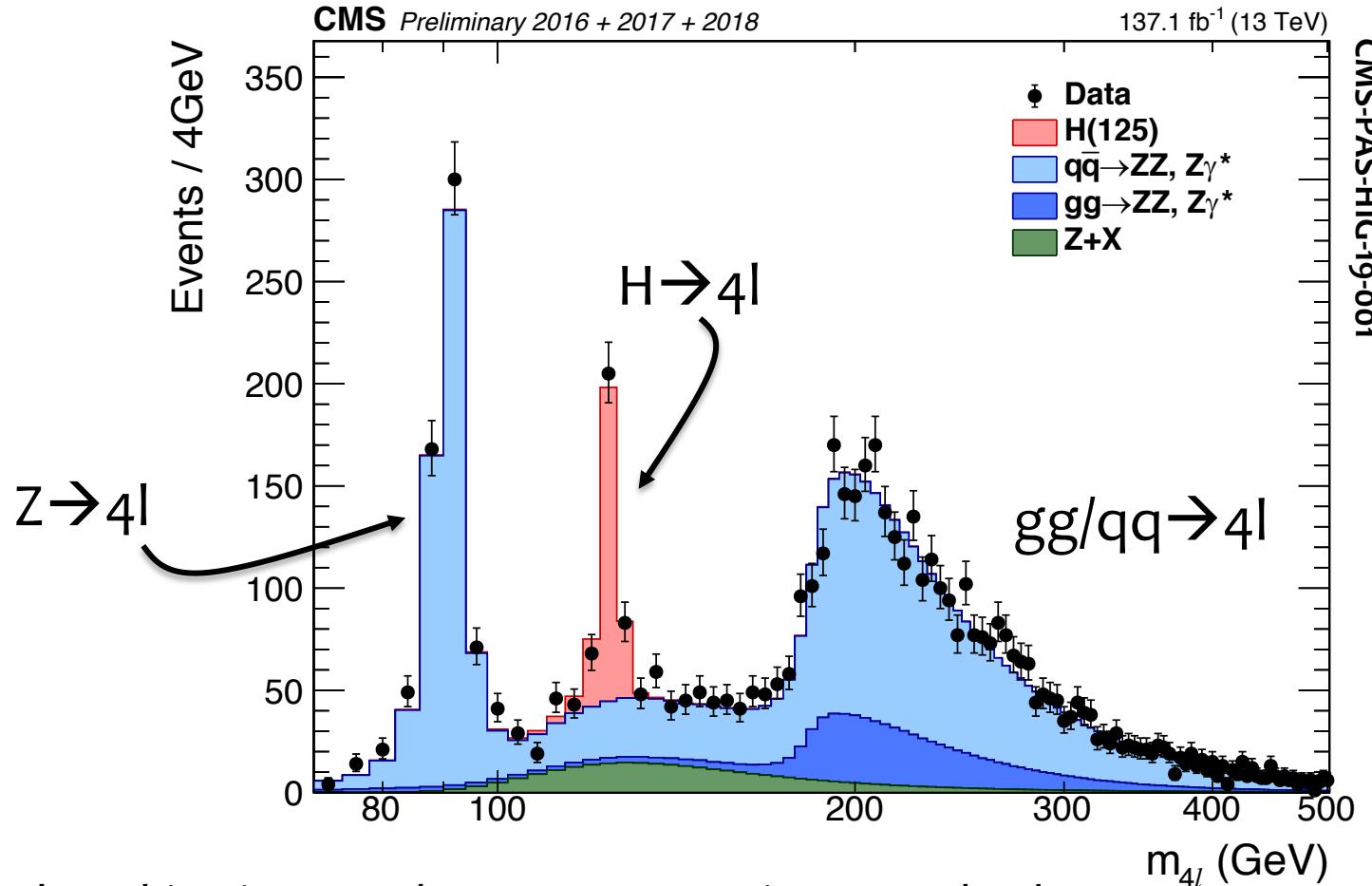
\*Discussions at LesHouches this year : <https://phystev.cnrs.fr/lh2019>

# EFT Caveats 3.

See talk from Z. Grout

In CMS/ATLAS we are used to thinking of **Signal / Background**

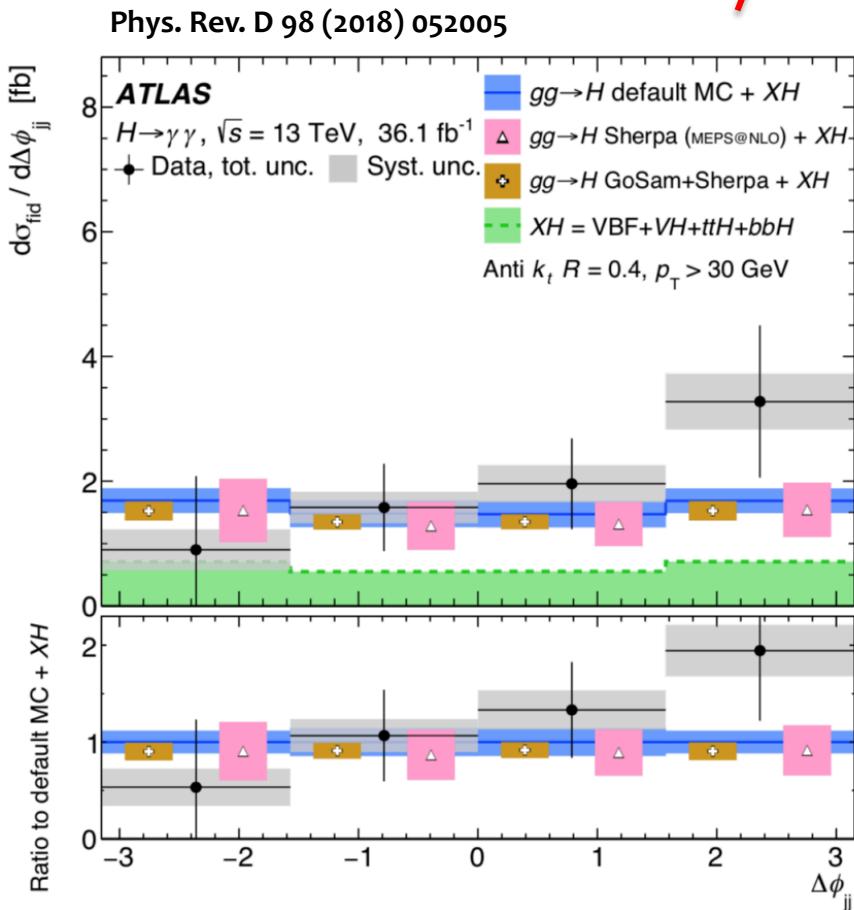
→ But EFT is a global approach!



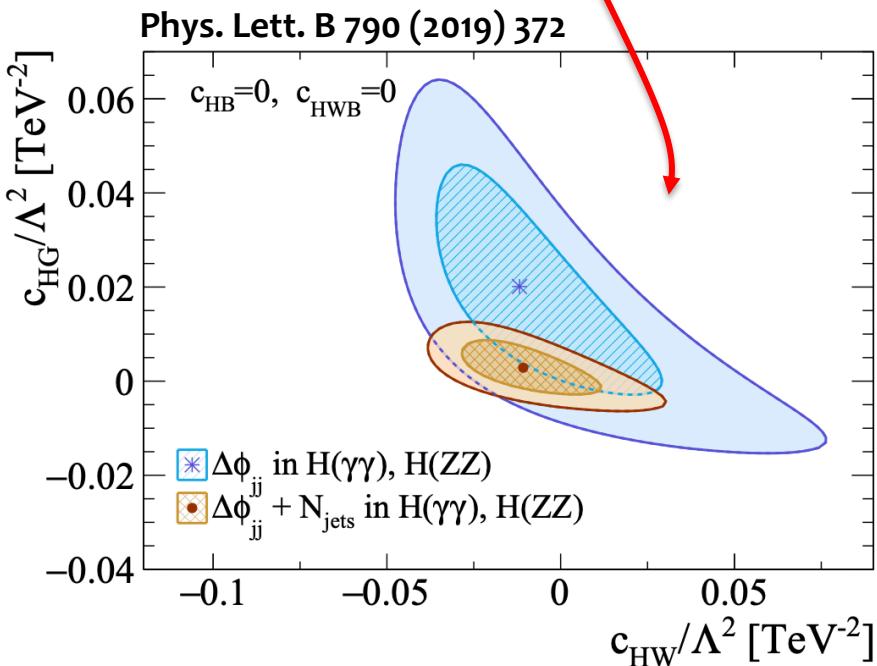
Full  $\text{pp} \rightarrow 4l$  combinations are the correct way to interpret the data  
→ Need to consider all contributions together

# (Re)Interpretations

Providing differential measurements allows theorists to produce EFT interpretations



$$L_i(\mathbf{c}/\Lambda^2) = \frac{1}{\sqrt{2\pi\sigma^2}} \exp\left(-\frac{(\mathbf{o}_i - \boldsymbol{\tau}_i)^2}{2\sigma^2}\right)$$



Aside from the caveats mentioned, what other effects (approximations) could yield inaccuracies?

→ Experiments should help validate these approaches by including EFT interpretations in papers – **how much do correlations / HO moments matter?**

# Speaking the same language

Outcome of LesHouches-2019 :

<https://github.com/ajgilbert/EFT2Obs-Demo>

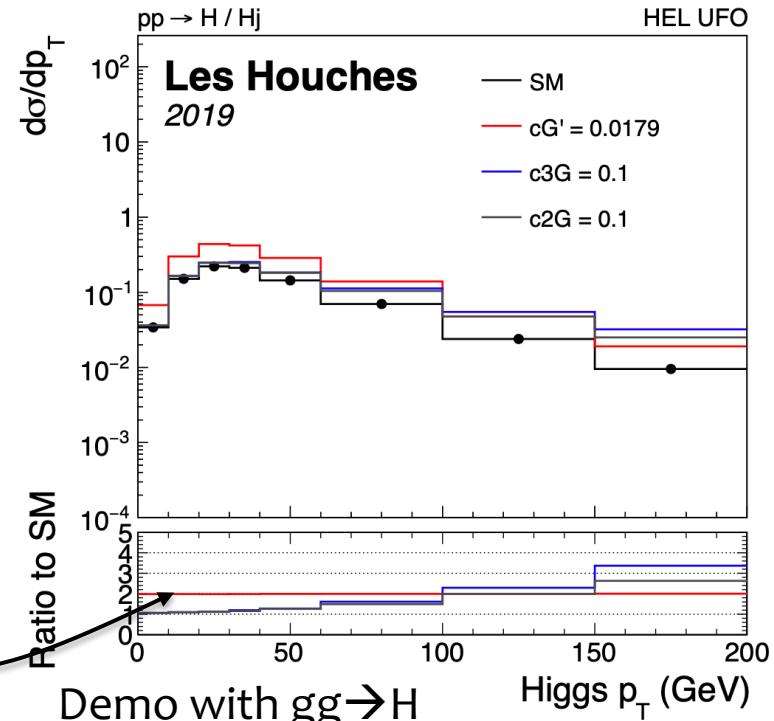
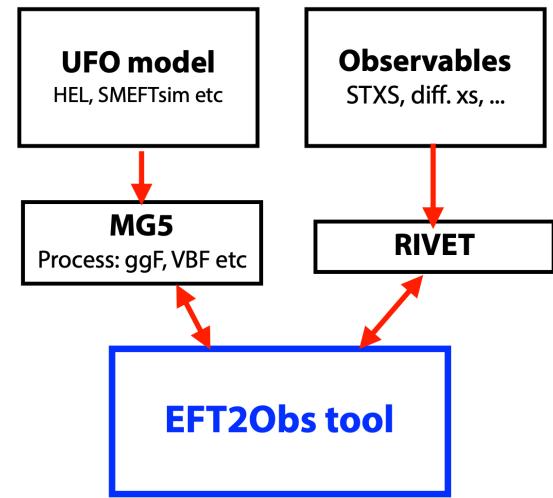
- Usable by both inside & outside CMS, therefore built on public tools (+ theorists prefer not to use ROOT\*?)
- Agnostic to specific EFT implementation, easy to implement new models (currently extensively tested with HEL, ongoing validations with SMEFTsim too)

Reweighting steering with simple configs

```
[  
  {  
    "block": "newcoup",  
    "index": 12,  
    "name": "cg",  
    "step": 0.0001  
  },  
  {  
    "block": "newcoup",  
    "index": 30,  
    "name": "c3g",  
    "step": 0.0001  
  },  
  ...  
]
```

$$\mu_i = 1 + \sum_j c_j A_{i,j} + \sum_{jk} c_j c_k B_{i,jk}$$

\*Outputs in YODA (also supports ROOT)

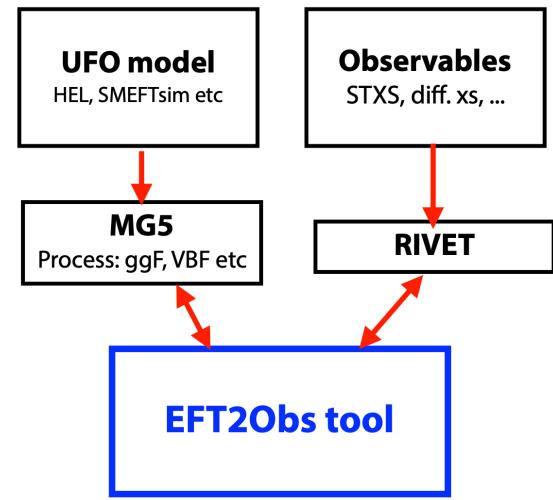


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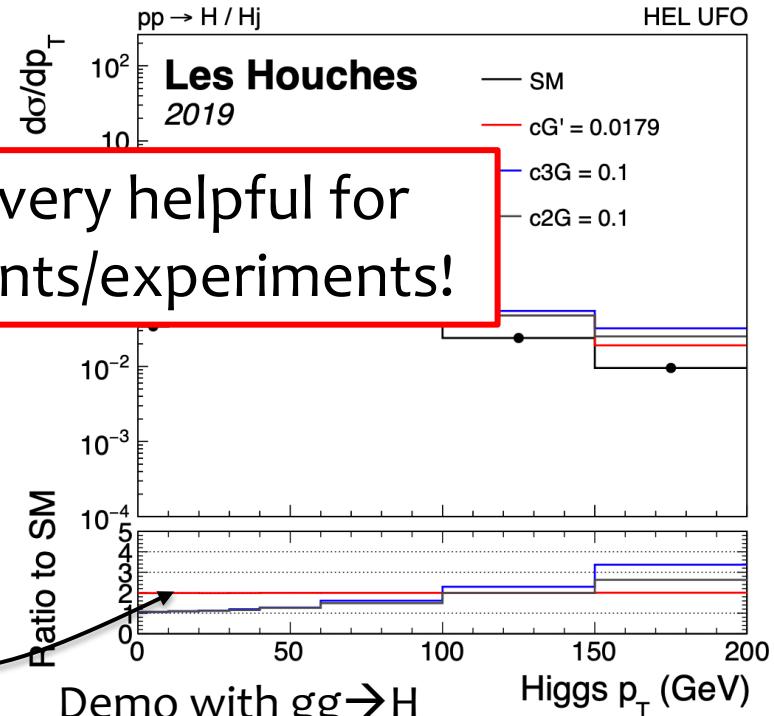
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```

Common tool will be very helpful for combining measurements/experiments!

$$\mu_i = 1 + \sum_j c_j A_{i,j} + \sum_{jk} c_j c_k B_{i,jk}$$

\*Outputs in YODA (also supports ROOT)



# EFT Roadmap from H-perspective?

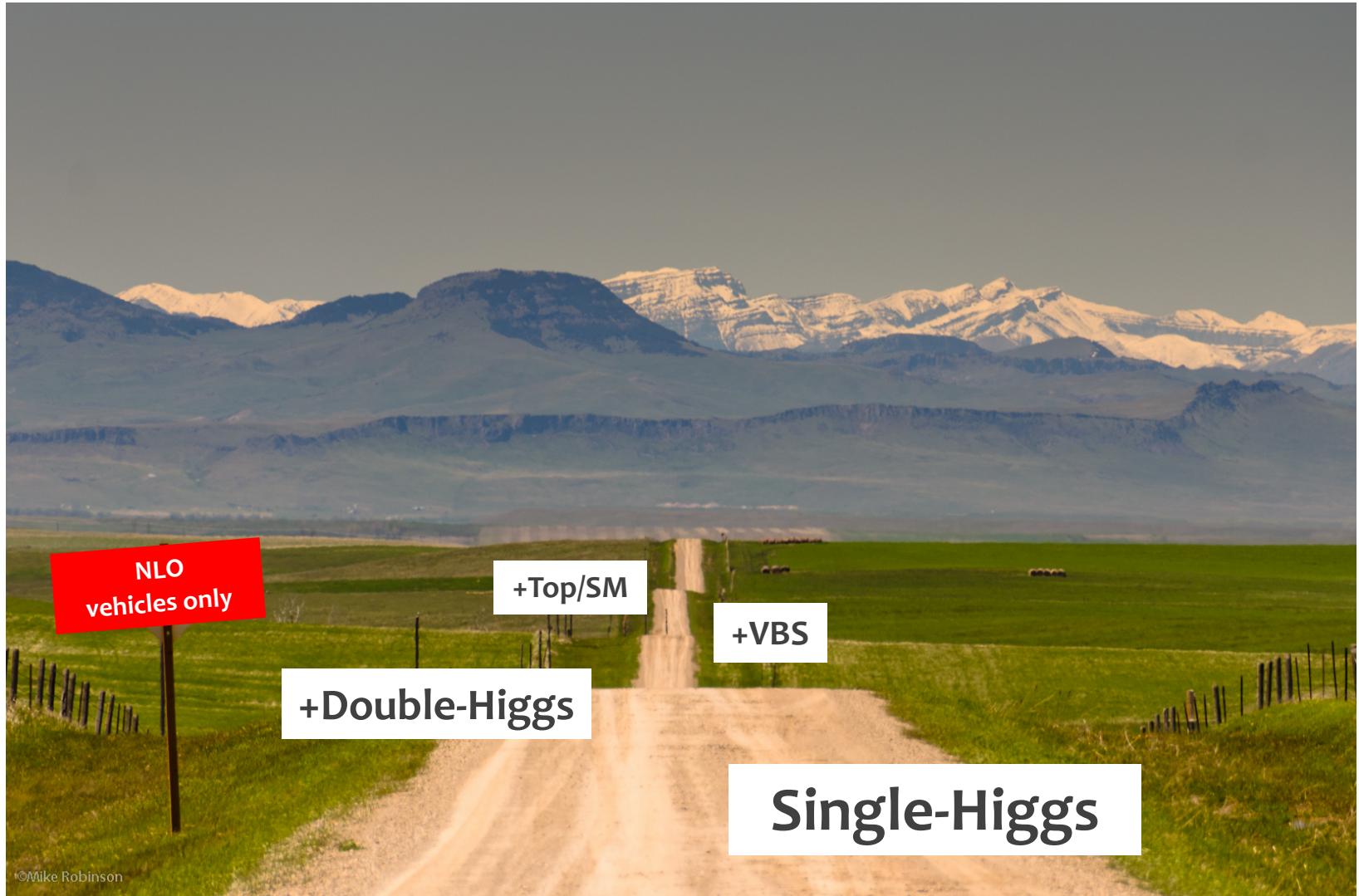


Image from <https://tauo.wordpress.com/2011/06/12/down-the-road-again/>

# EFT Roadmap from H-perspective?

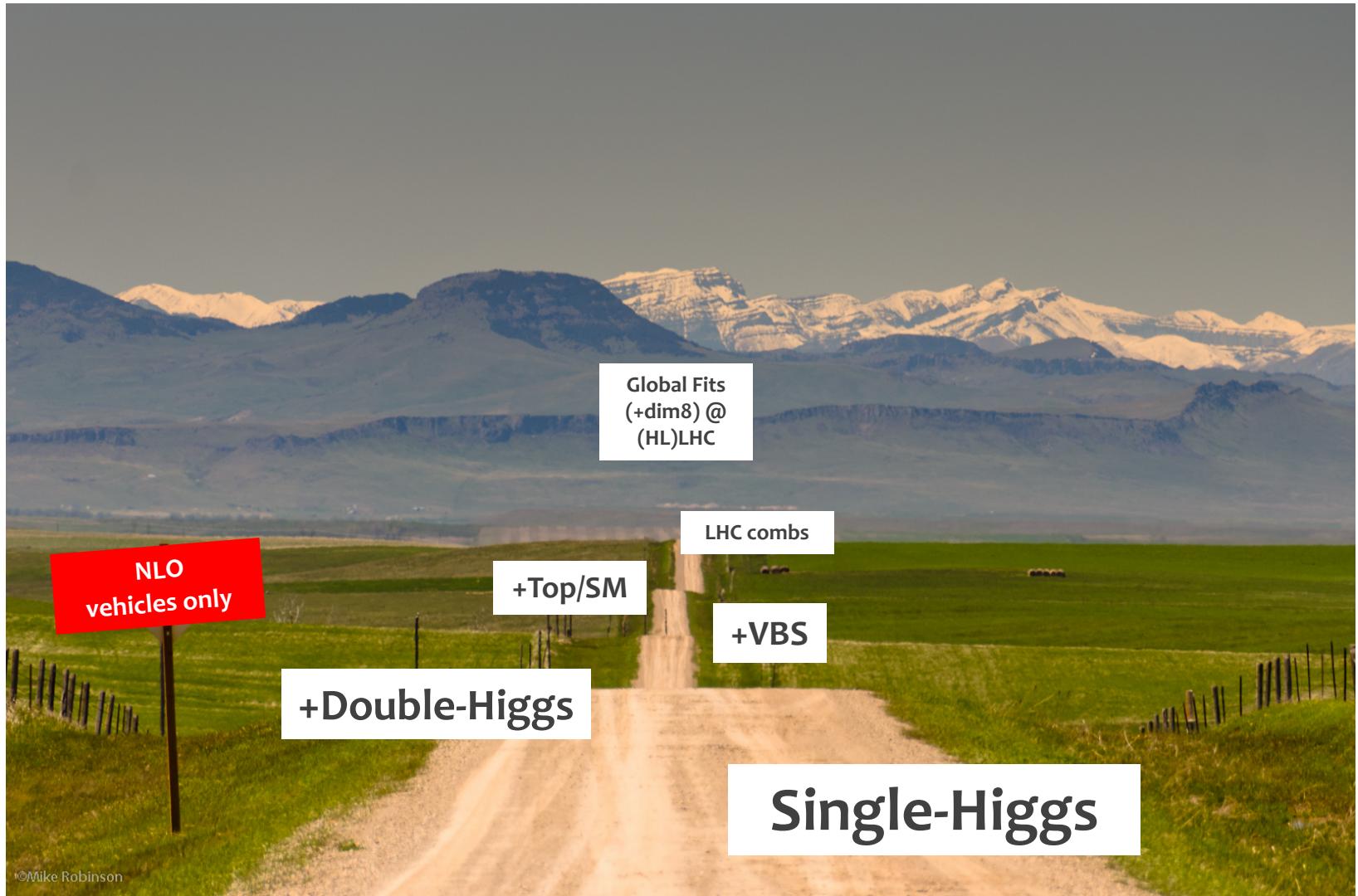


Image from <https://tauo.wordpress.com/2011/06/12/down-the-road-again/>

# Summary

EFT is a great tool to explore potential BSM physics in the Higgs sector

- A number of interpretations already from ATLAS/CMS
- Differential / STXS measurements already being used to explore EFT scenario

Experiments ultimately provide **measurements** (cross-sections/rates etc)

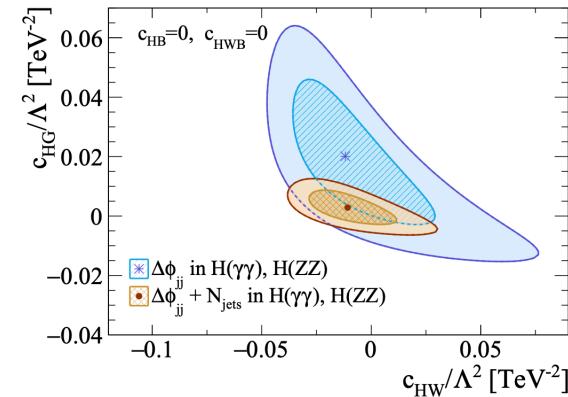
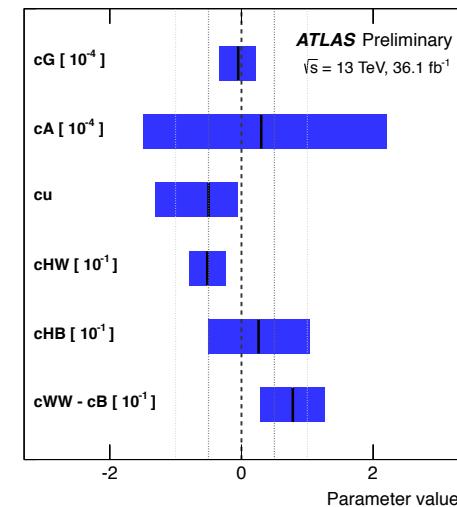
- EFT Interpretations within the experiments exciting
- Can provide additional info if needed for accurate interpretations *outside of the experiment*

Global fits/combinations are the challenging but exciting way forward

- Common tools (eg EFT implementations) help experiments+theorists
- Need to continue synergy among similar efforts (LHC HXS WG, Les Houches, HiggsTools, HiggsDays, this workshop...)

**LHC is not done yet → More results with EFT to come, but let's keep in touch!**

Observed HEL constraints with  $H \rightarrow ZZ^*$  and  $H \rightarrow \gamma\gamma$



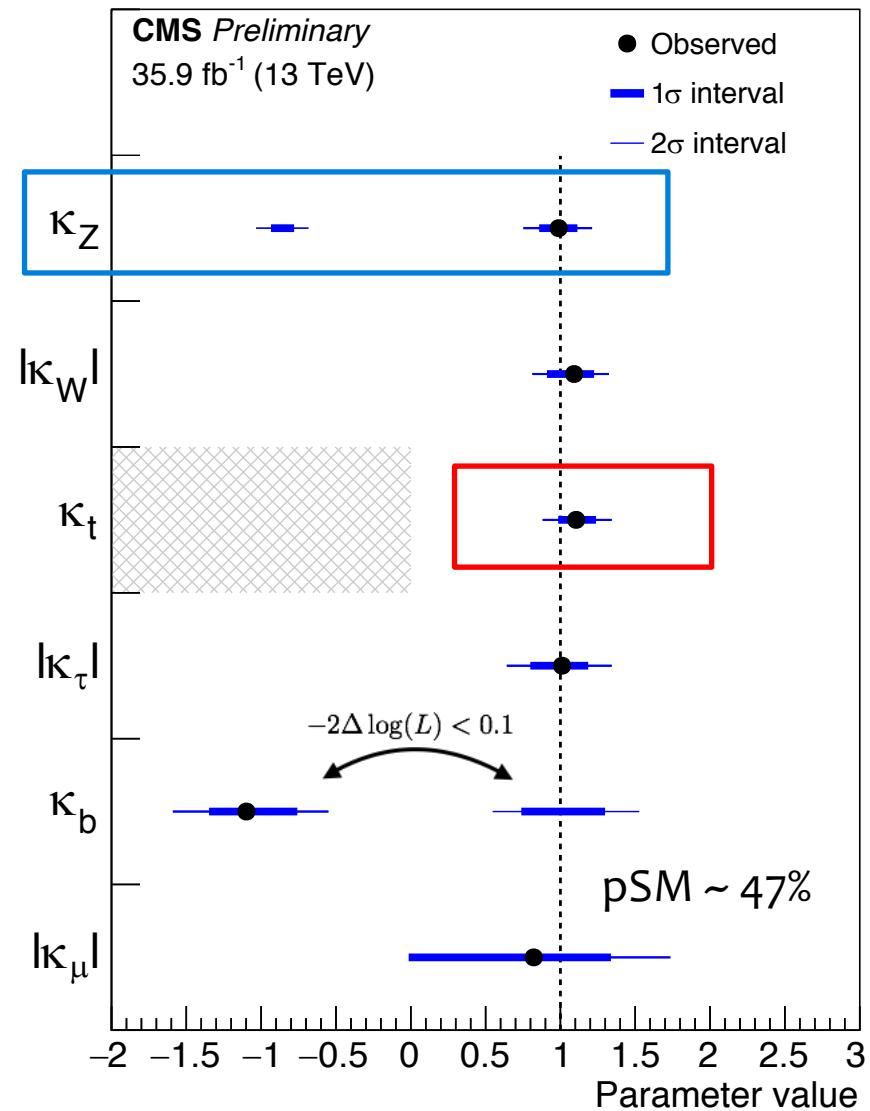
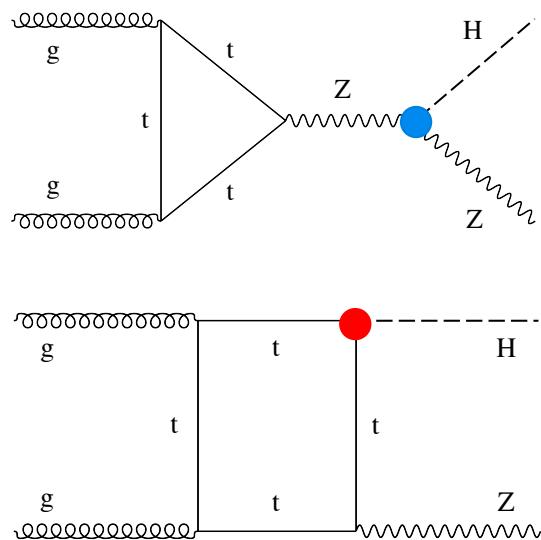
# THANKS!



# SM Coupling Constraints

Leading order coupling modifier framework used to correlate prod/decay rates

Interference between diagrams helps constrain degeneracies EG:  
 $\kappa_z$  sign degeneracy broken in  $gg \rightarrow ZH$  production

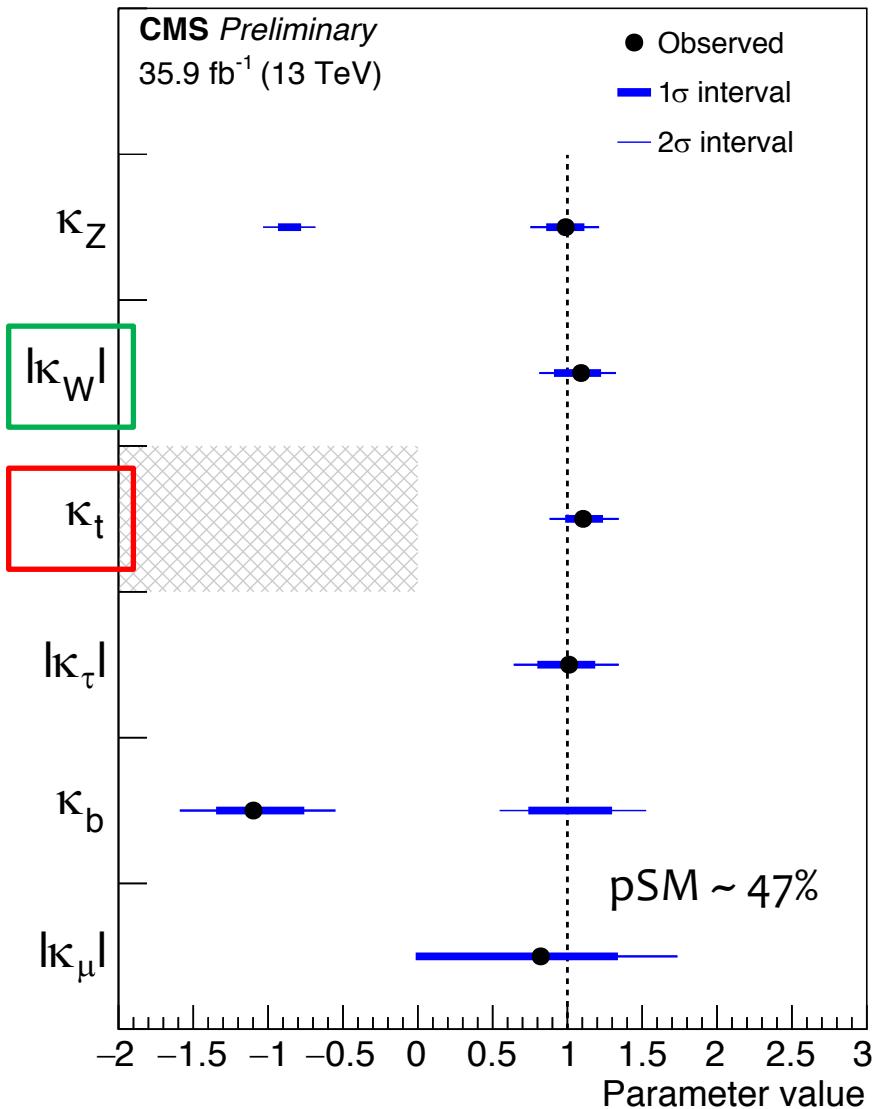
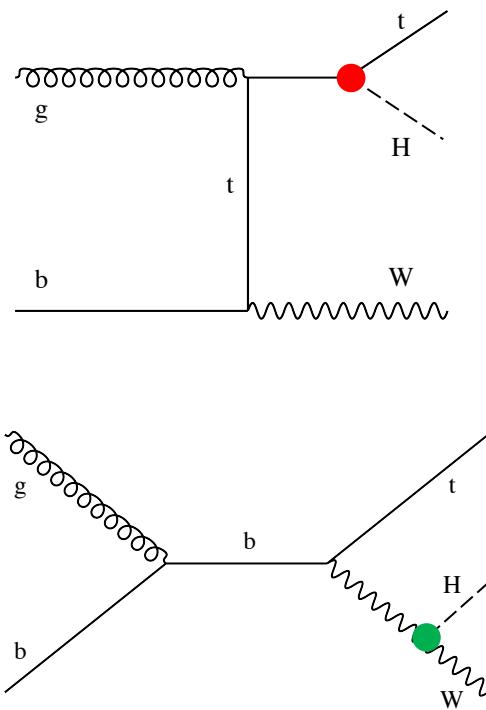


-ve sign for  $\kappa_b$  slightly preferred due to small excess in  $ggH$

# SM Coupling Constraints

Overall sign degeneracy – here we fix  $\kappa_t > 0$  and only look at +ve combination of  $\kappa_W \times \kappa_t$

Relative sign probed in single-top Higgs production Eg: tHW



Contribution from +ve sign combination included in relevant categories

| Operator             | Expression                                            | HEL coefficient                          | Vertices             |
|----------------------|-------------------------------------------------------|------------------------------------------|----------------------|
| $\mathcal{O}_g$      | $ H ^2 G_{\mu\nu}^A G^{A\mu\nu}$                      | $cG = \frac{m_W^2}{g_s^2} \bar{c}_g$     | $Hgg$                |
| $\mathcal{O}_\gamma$ | $ H ^2 B_{\mu\nu} B^{\mu\nu}$                         | $cA = \frac{m_W^2}{g'^2} \bar{c}_\gamma$ | $H\gamma\gamma, HZZ$ |
| $\mathcal{O}_u$      | $y_u  H ^2 \bar{u}_l H u_R + \text{h.c.}$             | $cu = v^2 \bar{c}_u$                     | $Ht\bar{t}$          |
| $\mathcal{O}_{HW}$   | $i (D^\mu H)^\dagger \sigma^a (D^\nu H) W_{\mu\nu}^a$ | $cHW = \frac{m_W^2}{g} \bar{c}_{HW}$     | $HWW, HZZ$           |
| $\mathcal{O}_{HB}$   | $i (D^\mu H)^\dagger (D^\nu H) B_{\mu\nu}$            | $cHB = \frac{m_W^2}{g'} \bar{c}_{HB}$    | $HZZ$                |
| $\mathcal{O}_W$      | $i (H^\dagger \sigma^a D^\mu H) D^\nu W_{\mu\nu}^a$   | $cWW = \frac{m_W^2}{g} \bar{c}_W$        | $HWW, HZZ$           |
| $\mathcal{O}_B$      | $i (H^\dagger D^\mu H) \partial^\nu B_{\mu\nu}$       | $cB = \frac{m_W^2}{g'} \bar{c}_B$        | $HZZ$                |

# Why Higgs Couplings?

Measuring Higgs bosons couplings remain a key goal of Future Higgs measurements

- Higgs coupling measurements are good to test SM compatibility
  - Synergies with EFT approaches

BSM contributions require O(%) level Higgs property measurements

arXiv:1310.8361 M<sub>NP</sub> ~ 1 TeV

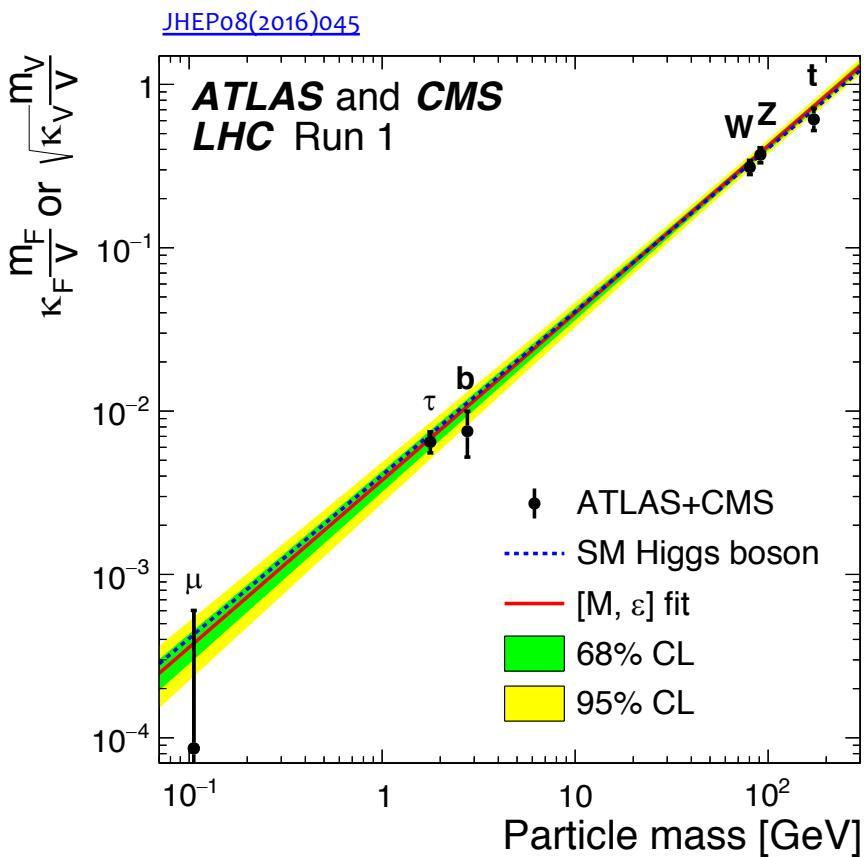
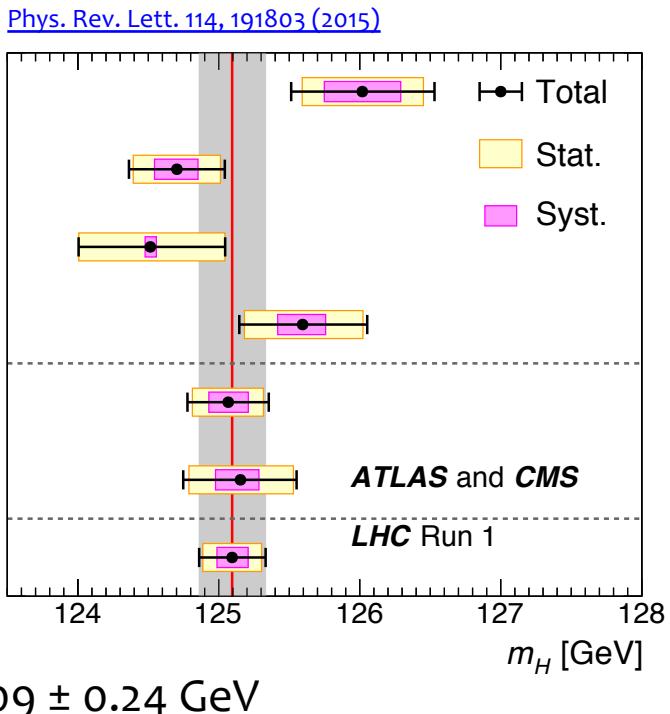
| Model           | $\kappa_V$ | $\kappa_b$  | $\kappa_\gamma$ |
|-----------------|------------|-------------|-----------------|
| Singlet Mixing  | ~ 6%       | ~ 6%        | ~ 6%            |
| 2HDM            | ~ 1%       | ~ 10%       | ~ 1%            |
| Decoupling MSSM | ~ -0.0013% | ~ 1.6%      | ~ -.4%          |
| Composite       | ~ -3%      | ~ -(3 - 9)% | ~ -9%           |
| Top Partner     | ~ -2%      | ~ -2%       | ~ +1%           |

# Completing the SM

Discovery of the Higgs boson completed the Standard Model

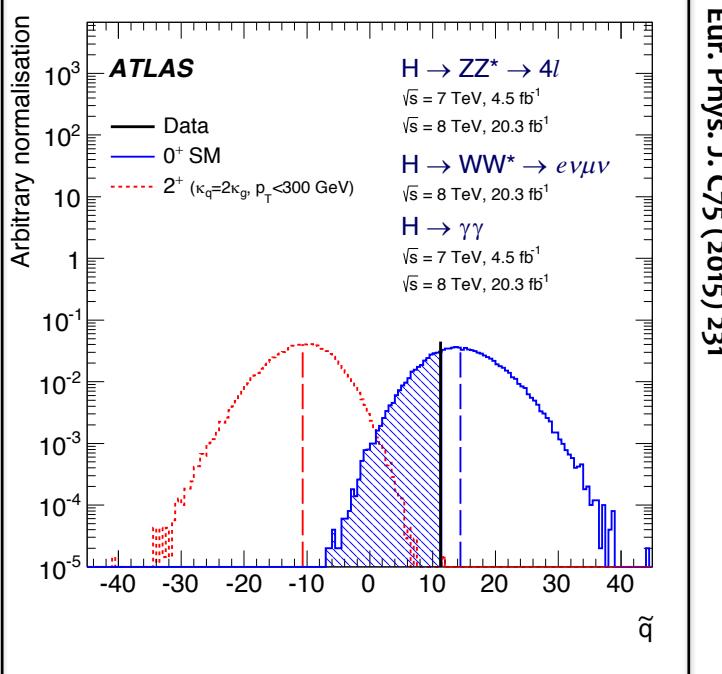
→ Great success of LHC @ Run-1

Precision measurements of Higgs boson properties through its couplings to SM particles ...



... and remarkable precision on mass measurement (< 0.2% uncertainty)

# Spin-Parity

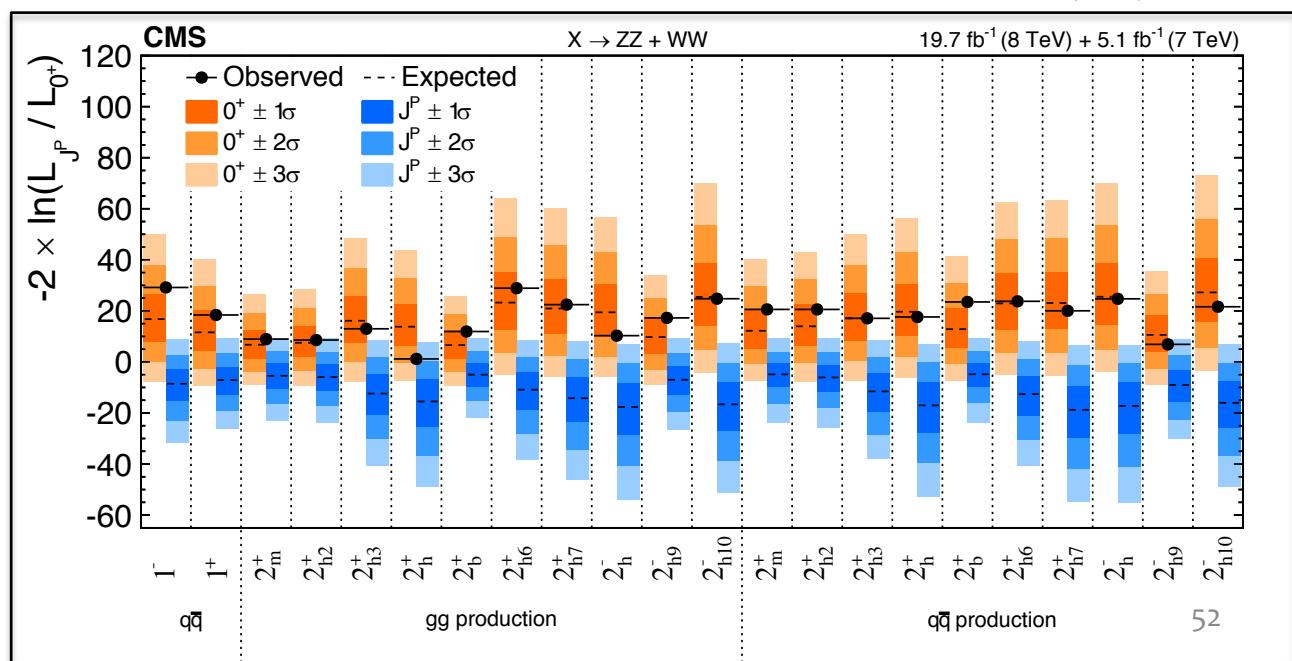
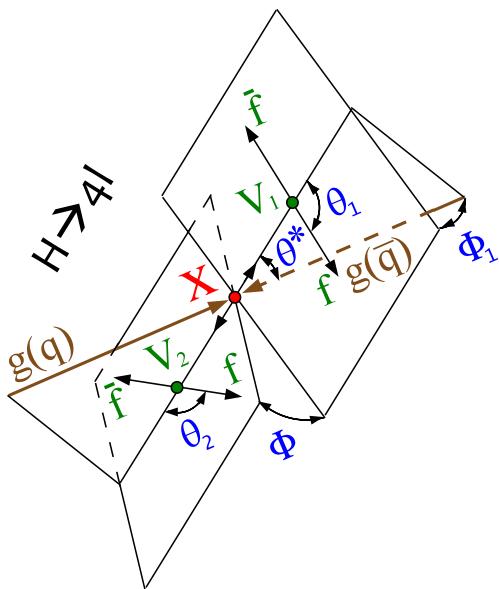


Eur. Phys. J. C75 (2015) 231

ATLAS and CMS  $H \rightarrow WW, ZZ, \gamma\gamma$  modes use angular information to distinguish between various  $J^P$  hypotheses...

Run-1 data is enough to rule out spin-2 (and many other  $J^P$  states) at  $> 99.9\%$  confidence level

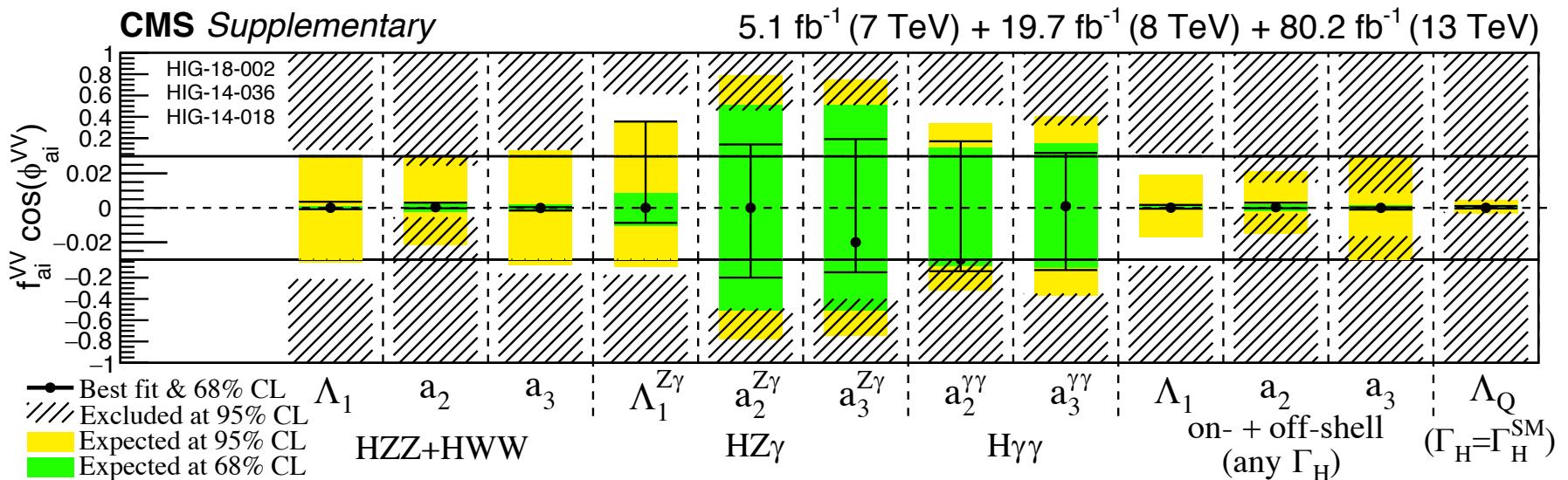
PRD 92 (2015) 012004



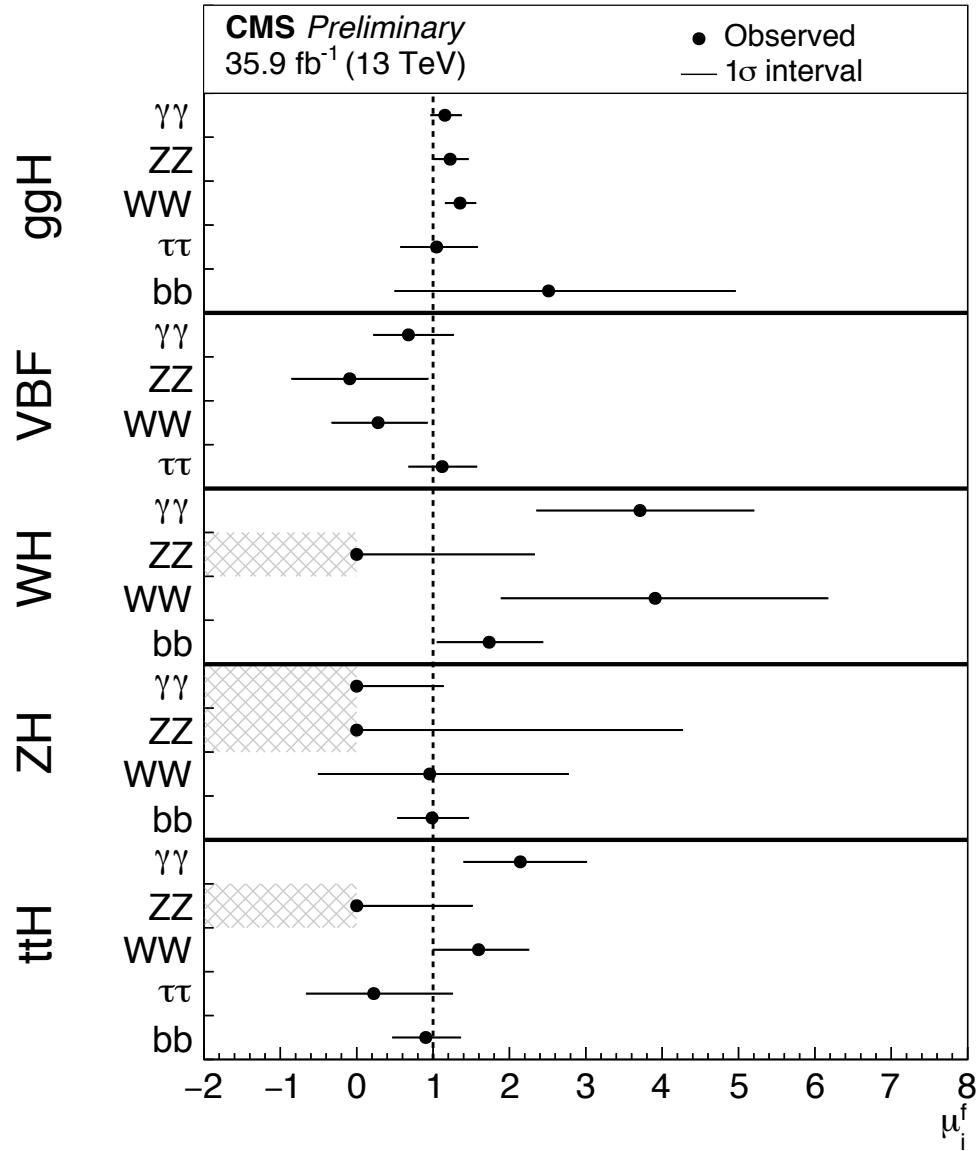
$$A(HVV) \sim \left[ a_1^{VV} + \frac{\kappa_1^{VV} q_1^2 + \kappa_2^{VV} q_2^2}{(\Lambda_1^{VV})^2} \right] m_{V1}^2 \epsilon_{V1}^* \epsilon_{V2}^* + a_2^{VV} f_{\mu\nu}^{*(1)} f^{*(2)\mu\nu} + a_3^{VV} f_{\mu\nu}^{*(1)} \tilde{f}^{*(2)\mu\nu},$$

$$f_{ai} = \frac{|a_i|^2 \sigma_i}{\sum_{j=1,2,3\dots} |a_j|^2 \sigma_j},$$

$$\phi_{ai} = \arg \left( \frac{a_i}{a_1} \right),$$



# What do we actually measure?



# What do we actually measure?

**Likelihood** to interpret the combined datasets from across Higgs channels ....

$$L(D|\mu, \theta) = \prod_n Prob \left( d_n | \sum_{i,f} \boxed{\mu_i \mu^f} S_{i,n}^f(\theta) + \sum_k B_k(\theta) \right) \times Gauss(\tilde{\theta}|\theta)$$


Extract “**signal strengths**” from  
Maximum likelihood estimators

Re-parameterize strengths in terms  
of “coupling modifiers”  $\kappa$

$$\mu_i = \frac{\sigma_i}{(\sigma_i)_{\text{SM}}} \quad \text{and} \quad \mu^f = \frac{\text{BR}^f}{(\text{BR}^f)_{\text{SM}}}.$$

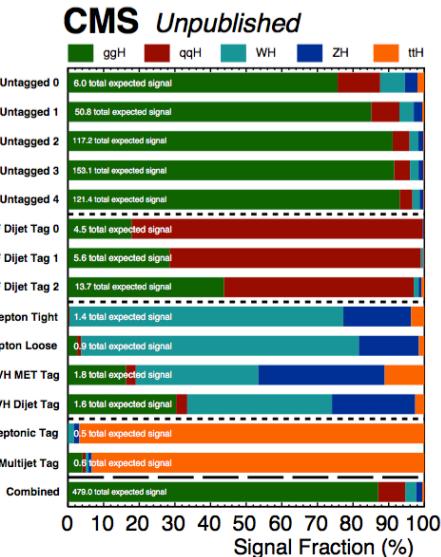
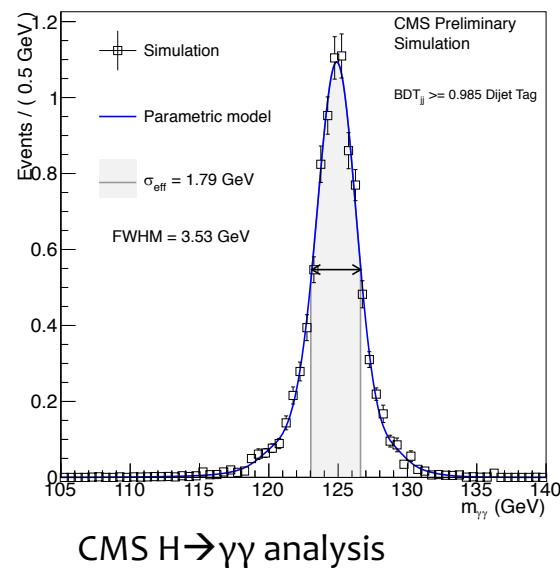
$$\mu \rightarrow \mu(\kappa)$$

Standard model defined by:  $\mu_i = \mu^f = 1$  or  $\kappa = 1$

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**Likelihood** to interpret the combined datasets from across Higgs channels ....

$$L(D|\mu, \theta) = \prod_n Prob \left( d_n | \sum_{i,f} \mu_i \mu^f S_{i,n}^f(\theta) + \sum_k B_k(\theta) \right) \times Gauss(\tilde{\theta}|\theta)$$



$$\times \mathcal{L} \times \varepsilon \times A$$

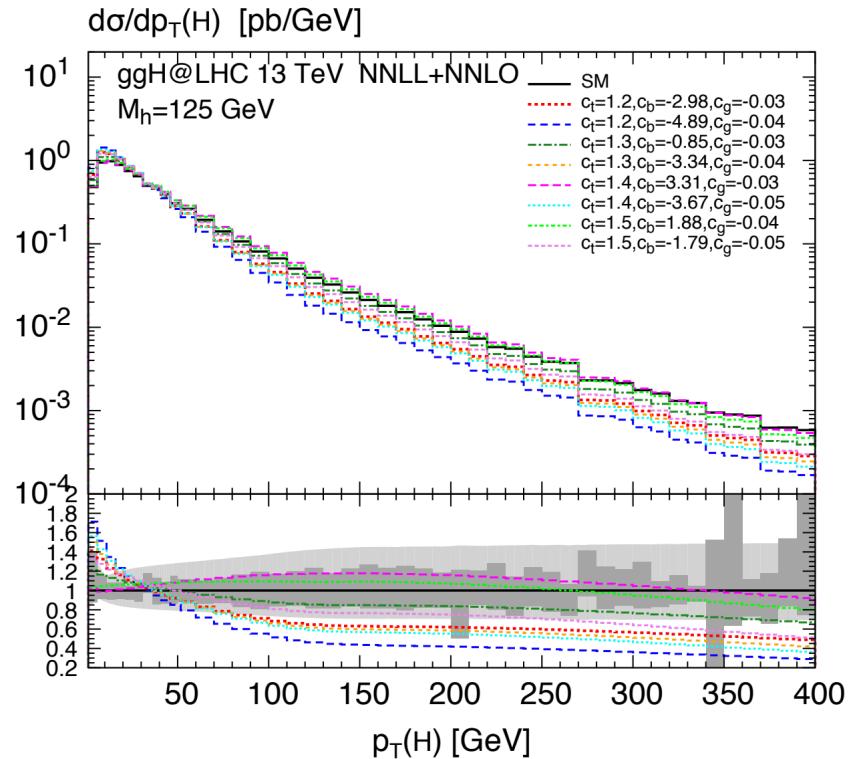
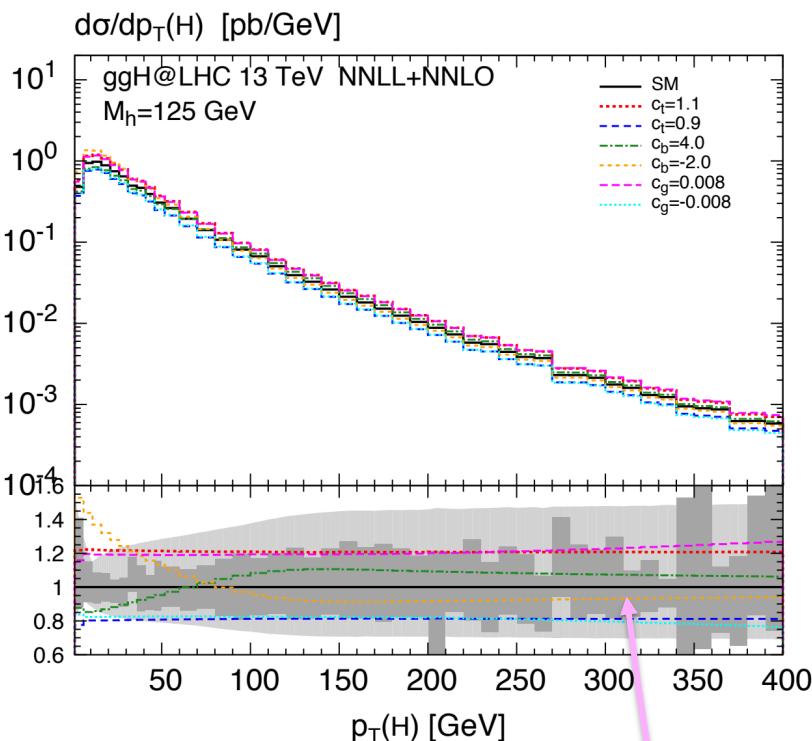
Signal model, accounts for “**shape**” of signal processes

- Efficiency \* acceptance
- Relative composition across signal regions (analysis bins, BDT output ...)

**Rely on SM Higgs Predictions to calculate in each channel (V-p<sub>T</sub>, n-jets etc)**

# Differential Higgs @ NLL+NLO

$$\frac{c_1}{\Lambda^2} \mathcal{O}_1 \rightarrow \frac{\alpha_S}{\pi v} c_g h G_{\mu\nu}^a G^{a,\mu\nu}, \quad \frac{c_2}{\Lambda^2} \mathcal{O}_2 \rightarrow \frac{m_t}{v} c_t h \bar{t} t, \quad \frac{c_3}{\Lambda^2} \mathcal{O}_3 \rightarrow \frac{m_b}{v} c_b h \bar{b} b,$$

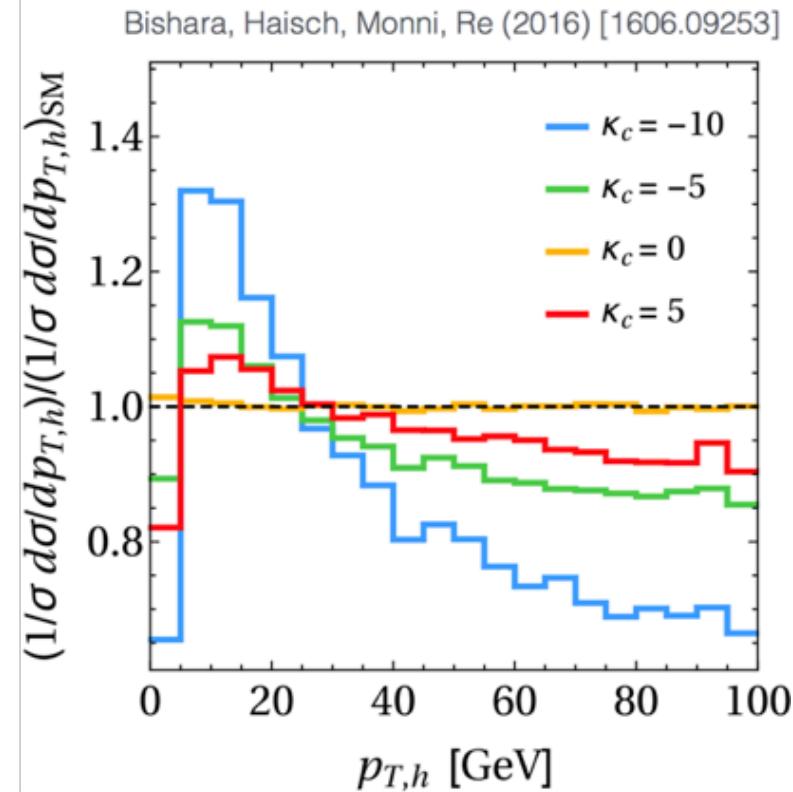
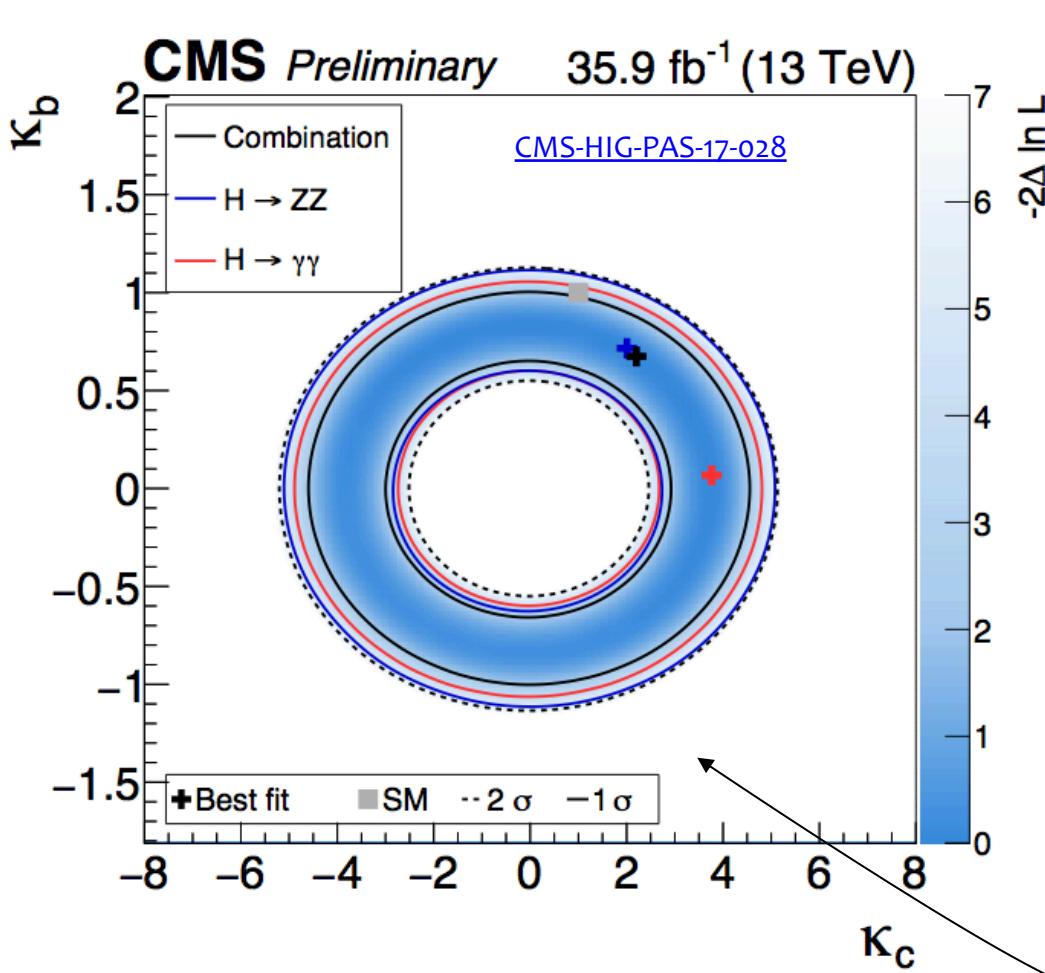


<https://arxiv.org/abs/1705.05143>

Relatively small variation from  $c_g$  vs  $p_T(H)$

# Differential Measurements

Constrain b & c Yukawa couplings from  
Differential measurements of inclusive Higgs  
production



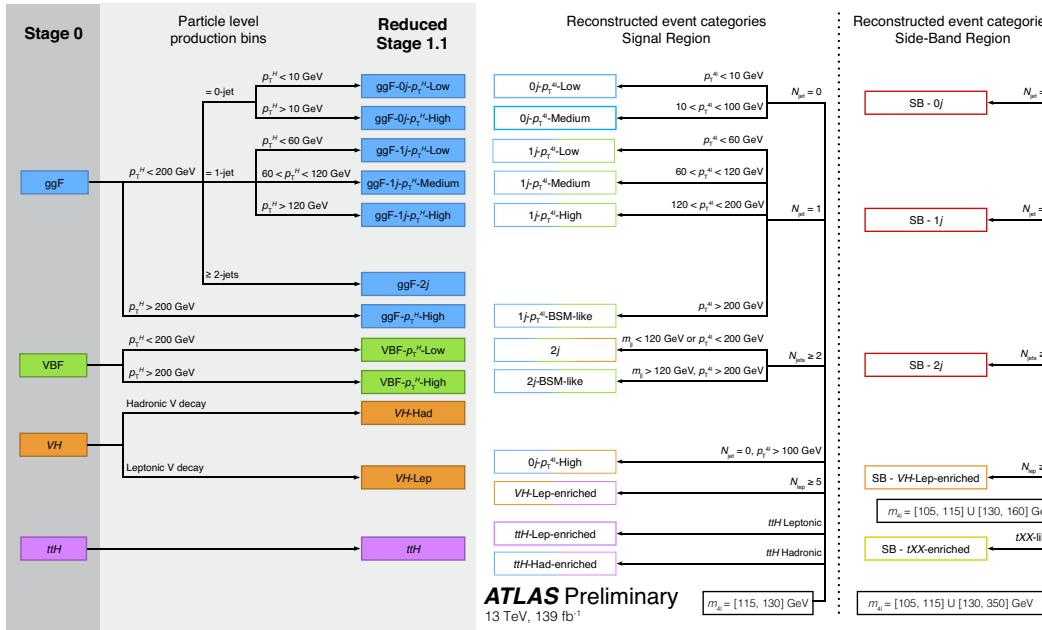
Large constraint from overall  
normalization (as in inclusive  
couplings) but also access more  
info via “shape”

# ATLAS differential diphoton SMEFT

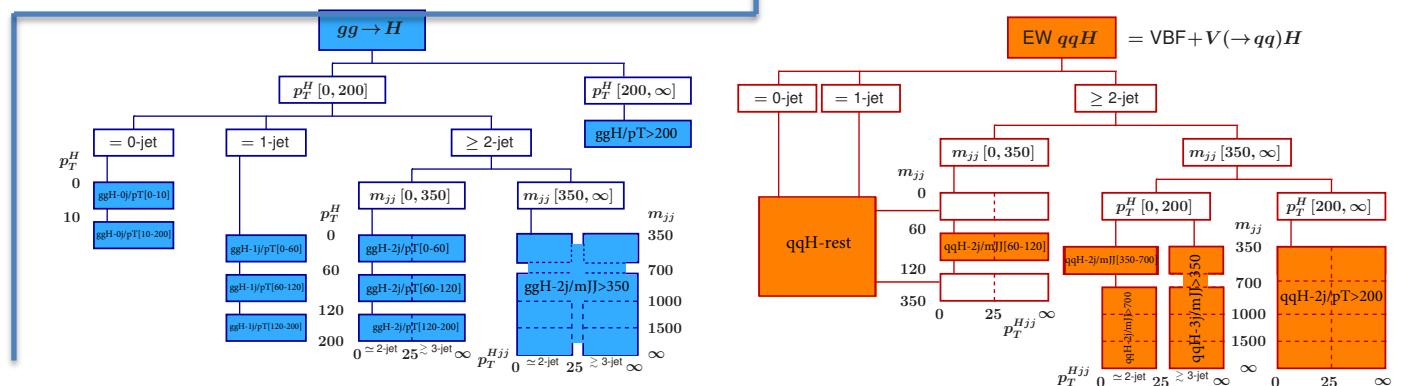
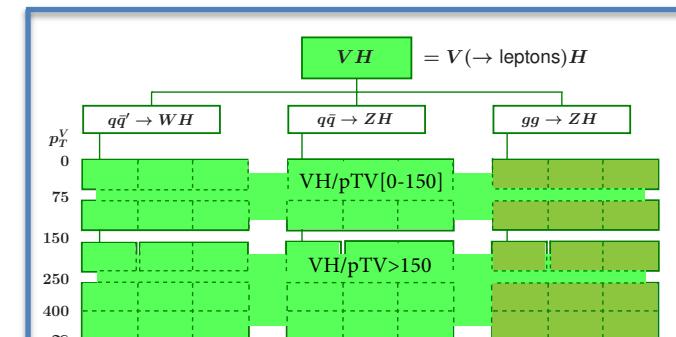
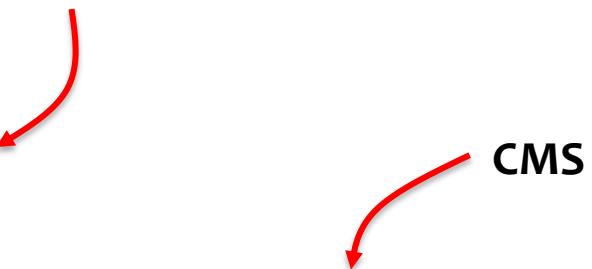
Table 6: The 95% CL observed limits on the  $\bar{C}_{HG}$ ,  $\bar{C}_{HW}$ ,  $\bar{C}_{HB}$ ,  $\bar{C}_{HWB}$  Wilson coefficients of the SMEFT basis and their CP-odd counterparts using interference-only terms and using both interference and quadratic terms. Limits are derived fitting one Wilson coefficient at a time while setting the other coefficients to zero.

| Coefficient       | 95% CL, interference-only terms | 95% CL, interference and quadratic terms |
|-------------------|---------------------------------|------------------------------------------|
| $\bar{C}_{HG}$    | $[-4.2, 4.8] \times 10^{-4}$    | $[-6.1, 4.7] \times 10^{-4}$             |
| $\tilde{C}_{HG}$  | $[-2.1, 1.6] \times 10^{-2}$    | $[-1.5, 1.4] \times 10^{-3}$             |
| $\bar{C}_{HW}$    | $[-8, 2, 7.4] \times 10^{-4}$   | $[-8.3, 8.3] \times 10^{-4}$             |
| $\tilde{C}_{HW}$  | $[-0.26, 0.33]$                 | $[-3.7, 3.7] \times 10^{-3}$             |
| $\bar{C}_{HB}$    | $[-2.4, 2.3] \times 10^{-4}$    | $[-2.4, 2.4] \times 10^{-4}$             |
| $\tilde{C}_{HB}$  | $[-13.0, 14.0]$                 | $[-1.2, 1.1] \times 10^{-3}$             |
| $\bar{C}_{HWB}$   | $[-4.0, 4.4] \times 10^{-4}$    | $[-4.2, 4.2] \times 10^{-4}$             |
| $\tilde{C}_{HWB}$ | $[-11.1, 6.5]$                  | $[-2.0, 2.0] \times 10^{-3}$             |

# CMS/ATLAS H $\rightarrow$ ZZ STXS defs

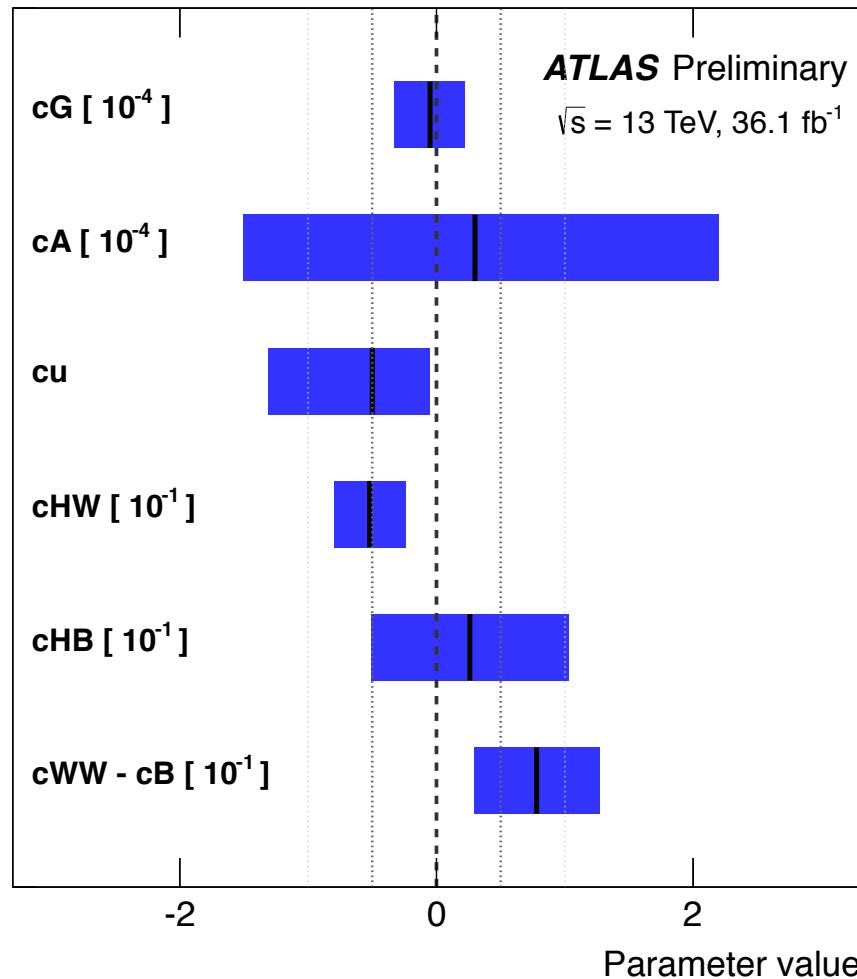


ATLAS

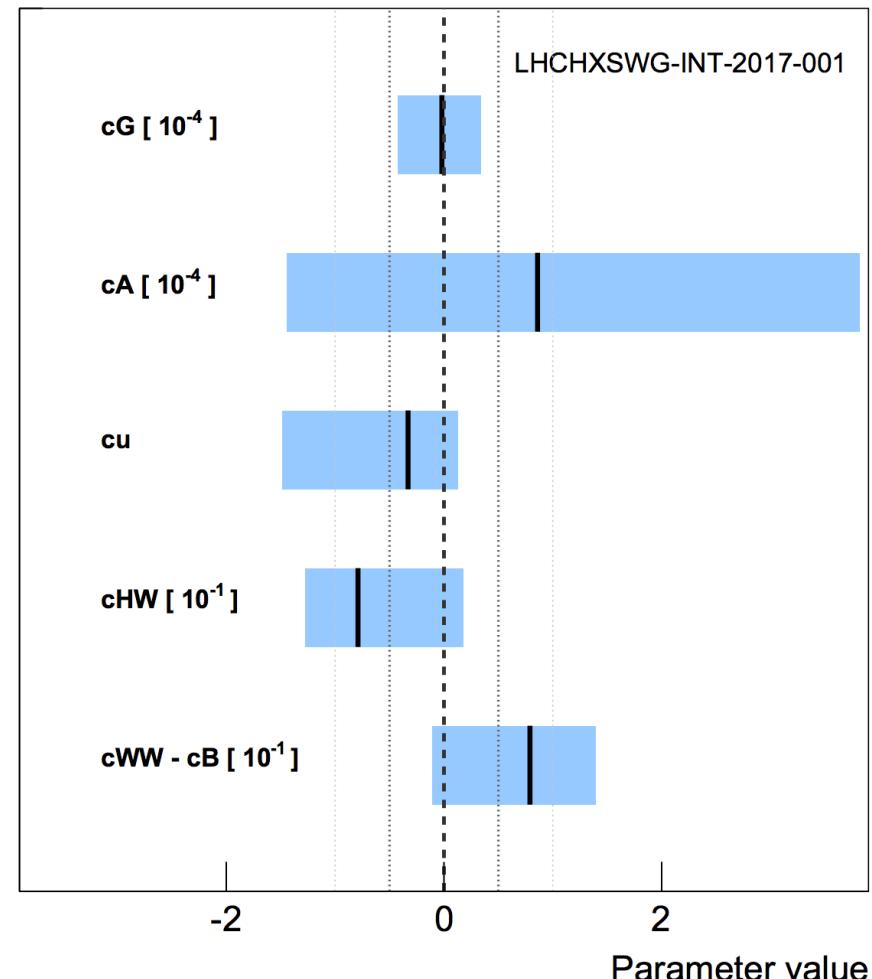


# Compare ~full LH vs $\chi^2$

Observed HEL constraints with  $H \rightarrow ZZ^*$  and  $H \rightarrow \gamma\gamma$



Fit to ATLAS STXS measurements (ATLAS-CONF-2017-047)



# Clusters for double Higgs

| Benchmark | $\kappa_\lambda$ | $\kappa_t$ | $c_2$ | $c_g$ | $c_{2g}$ |
|-----------|------------------|------------|-------|-------|----------|
| 1         | 7.5              | 1.0        | -1.0  | 0.0   | 0.0      |
| 2         | 1.0              | 1.0        | 0.5   | -0.8  | 0.6      |
| 3         | 1.0              | 1.0        | -1.5  | 0.0   | -0.8     |
| 4         | -3.5             | 1.5        | -3.0  | 0.0   | 0.0      |
| 5         | 1.0              | 1.0        | 0.0   | 0.8   | -1.0     |
| 6         | 2.4              | 1.0        | 0.0   | 0.2   | -0.2     |
| 7         | 5.0              | 1.0        | 0.0   | 0.2   | -0.2     |
| 8         | 15.0             | 1.0        | 0.0   | -1.0  | 1.0      |
| 9         | 1.0              | 1.0        | 1.0   | -0.6  | 0.6      |
| 10        | 10.0             | 1.5        | -1.0  | 0.0   | 0.0      |
| 11        | 2.4              | 1.0        | 0.0   | 1.0   | -1.0     |
| 12        | 15.0             | 1.0        | 1.0   | 0.0   | 0.0      |
| SM        | 1.0              | 1.0        | 0.0   | 0.0   | 0.0      |

